## 14 Review of Fourier Analysis, Partial Differential Equations

### 14.1 Fourier Series/Transform/Integrals

- orthogonality of eigenfunctions of the S-L problem $\rightarrow$ orthogonality of trigonometric system $\rightarrow$ Euler formulas for the Fourier series
- odd/even $\rightarrow$ Fourier sine/cosine series, Half-range expansion
- extension to nonperiodic functions $\rightarrow$ Fourier integral
- Fourier transform (sine/cosine)
- Recall that Fourier series are deeply connected with periodic phenomena involving ordinary differential equations (ODEs). Now we will see they are extremely useful to find the solutions for more complex differential equations called partial differential equations (PDEs).


### 14.2 Basic Concepts of Partial Differential Equations

- PDE : diff eqn involving partial derivatives of two or more independent variables
- The second order differential equation

$$
A u_{x x}+2 B u_{x y}+C u_{y y}=F
$$

Discriminant $B^{2}-A C$ :

$$
\begin{array}{ll}
B^{2}-A C=0: & \text { Parabolic } \\
B^{2}-A C>0: & \text { Hyperbolic } \\
B^{2}-A C<0: & \text { Elliptic }
\end{array}
$$

Example 1. Important linear partial differential equations of the second order

- 1-dimensional wave equation

$$
\frac{\partial^{2} u}{\partial^{2} t}=c^{2} \frac{\partial^{2} u}{\partial x^{2}}
$$

Discriminant: $B^{2}-A C=0-1 \cdot\left(-c^{2}\right)=c^{2}>0$ : hyperbolic

- 1-dimensional heat equation

$$
\frac{\partial T}{\partial t}=c^{2} \frac{\partial^{2} T}{\partial x^{2}}
$$

Discriminant: $B^{2}-A C=0-0 \cdot c^{2}=0$ : parabolic

- 2-dimensional Laplace equation (steady state heat equation with no heat generation)

$$
\frac{\partial^{2} u}{\partial x^{2}}+\frac{\partial^{2} u}{\partial y^{2}}=0
$$

Discriminant: $B^{2}-A C=0-1 \cdot 1=-1,0$ : elliptic

- 2-dimensional Poisson equation (steady state heat equation with heat generation)

$$
\frac{\partial^{2} T}{\partial x^{2}}+\frac{\partial^{2} T}{\partial y^{2}}=f(x, y)
$$

- 2-dimensional wave equation

$$
\frac{\partial^{2} u}{\partial^{2} t}=c^{2}\left(\frac{\partial^{2} u}{\partial x^{2}}+\frac{\partial^{2} u}{\partial y^{2}}\right)
$$

- 3- dimensional Laplace equation

$$
\frac{\partial^{2} u}{\partial x^{2}}+\frac{\partial^{2} u}{\partial y^{2}}+\frac{\partial^{2} u}{\partial z^{2}}=0
$$

Theorem 1 Fundamental Theorem (Superposition or linearity principle)
If $u_{1}$ and $u_{2}$ are any solutions of a linear homogeneous partial differential equation in some region $R$, then

$$
u=c_{1} u_{1}+c_{2} u_{2}
$$

with any constants $c_{1}$ and $c_{2}$ is also a solution of that equation in $R$.
Example 2 Find a solution $u(x, y)$ of the partial differential equation

$$
u_{x x}-u=0
$$

solution)

$$
u(x, y)=A(y) \cdot e^{x}+B(y) \cdot e^{-x}
$$

Example 3 Solve the partial differential equation

$$
u_{x y}=-u_{x}
$$

solution)

$$
\begin{gathered}
u_{x}=P, \quad \Rightarrow \quad P_{y}=-P \\
\frac{P_{y}}{P}=-1 \quad \Longrightarrow \quad \ln P=-y+\tilde{c}(x) \\
\therefore P=c(x) \cdot e^{-y} \\
u(x, y)=f(x) e^{-y}+g(y) \quad \text { where } \quad f(x)=\int c(x) d x
\end{gathered}
$$

### 14.3 Modeling; Vibrating String; Wave Equation

## Physical Assumptions

1. Homogeneous, perfectly elastic string.
2. Gravitational force is negligible, compared to the tension.
3. String performs small transverse motions in a vertical plane.

## Derivation of the Differential Equation from Forces

Force balance

$$
\begin{gathered}
x: \quad T_{1} \cos \alpha=T_{2} \cos \beta=T=\mathrm{const} \\
y: \quad T_{2} \cdot \cos \beta-T_{1} \cdot \sin \alpha=\rho \Delta x \frac{\partial^{2} u}{\partial t^{2}} \\
\frac{T_{2} \sin \beta}{T_{2} \cos \beta}-\frac{T_{1} \sin \alpha}{T_{1} \cos \alpha}=\tan \beta-\tan \alpha=\frac{\rho \Delta x}{T} \cdot \frac{\partial^{2} u}{\partial t^{2}} \\
\tan \alpha=\left.\left(\frac{\partial u}{\partial x}\right)\right|_{x} \quad \text { and } \quad \tan \beta=\left.\left(\frac{\partial u}{\partial x}\right)\right|_{x+\Delta x}
\end{gathered}
$$

- By Taylor expansion around $x$,

$$
\begin{aligned}
\left.\frac{\partial u}{\partial x}\right|_{x+\Delta x} & =\left.\left(\frac{\partial u}{\partial x}\right)\right|_{x}+\left.\frac{\partial^{2} u}{\partial x^{2}}\right|_{x} \Delta x+\cdots \\
{\left[\left.\left(\frac{\partial u}{\partial x}\right)\right|_{x+\Delta x}-\left.\left(\frac{\partial u}{\partial x}\right)\right|_{x}\right] } & =\left[\left.\frac{\partial u}{\partial x}\right|_{x}+\left.\frac{\partial^{2} u}{\partial x^{2}}\right|_{x} \Delta x-\left.\frac{\partial u}{\partial x}\right|_{x}\right]=\frac{\rho \Delta x}{T} \cdot \frac{\partial^{2} u}{\partial t^{2}} \\
& \therefore \frac{\partial^{2} u}{\partial t^{2}}=\frac{T}{\rho} \cdot \frac{\partial^{2} u}{\partial x^{2}}
\end{aligned}
$$

- One-dimensional wave equation

$$
\frac{\partial^{2} u}{\partial t^{2}}=c^{2} \frac{\partial^{2} u}{\partial x^{2}} \quad\left(c^{2}=\frac{T}{\rho}\right)
$$

- Unit of $c$

$$
\begin{aligned}
& {\left[c^{2}\right]=\left[\frac{T}{\rho}\right]=\left[\frac{\mathrm{kg} \cdot \mathrm{~m} / \mathrm{s}^{2}}{\mathrm{~kg} / \mathrm{m}}\right]=[\mathrm{m} / \mathrm{s}]^{2} } \\
\therefore & {[c]=[\mathrm{m} / \mathrm{s}] \quad \text { (propagation velocity) } }
\end{aligned}
$$

