

§ Line patterns to visualize flows

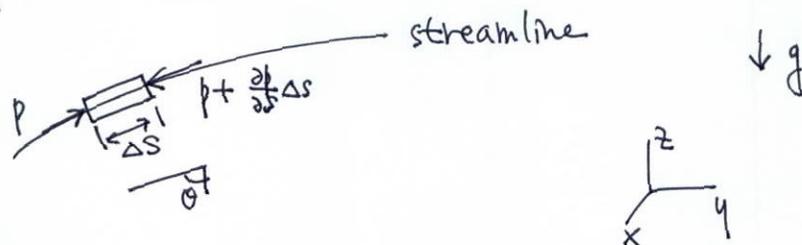
- streamline
- pathline
- streakline
- timeline

Chap. 4. Pressure and Momentum

§ Bernoulli's principle

- Euler's equation in curved streamline (steady)

(1) pressure gradient along streamline



$$\rightarrow \sum F = pA - (p + \frac{\partial p}{\partial s} \Delta s) A - \rho A \Delta s g \cos \theta = \rho A \Delta s \frac{dv}{dt}$$

$$-\frac{dp}{ds} - \rho g \cos \theta = \rho \frac{dv}{dt}$$

$$\frac{dp}{ds} + \rho \frac{dv}{dt} + \rho g \cos \theta = 0$$

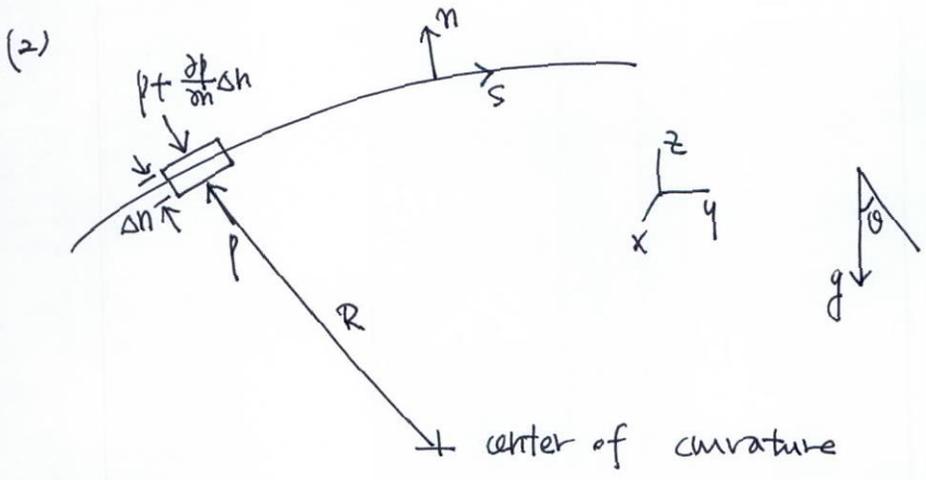
$$dp + \rho \left(\frac{ds}{dt} \right) dv + \rho g ds \cos \theta = 0$$

$$dp + \rho v dv + \rho g dh = 0$$

$$d \left(p + \frac{1}{2} \rho v^2 + \rho g h \right) = 0$$

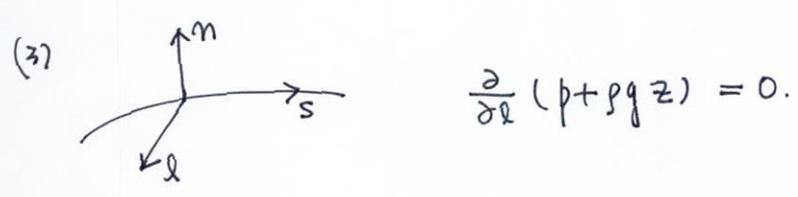
OR $\frac{d}{ds} \left(p + \frac{1}{2} \rho v^2 + \rho g z \right) = 0$

Bernoulli's eq.

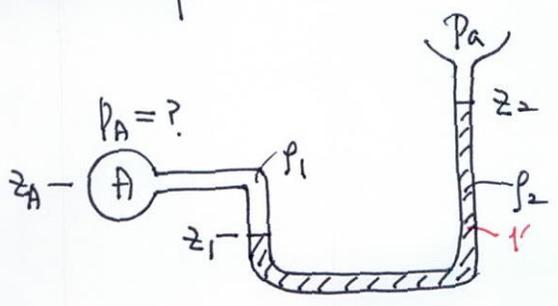


$$\begin{aligned} \sum F &= -pA + (p + \frac{\partial p}{\partial n} \Delta n) A + \rho g \Delta n A \cos \theta \\ &= (\frac{\partial p}{\partial n} + \rho g \cos \theta) A \Delta n \\ &= \rho A \Delta n \frac{v^2}{R} \end{aligned}$$

$$\frac{\partial}{\partial n} (p + \rho g z) = \frac{\rho v^2}{R}$$



§ Manometry



$p_1 = p_1'$ (Pascal's law)

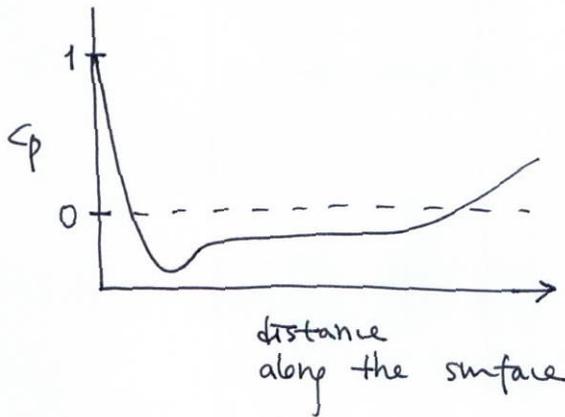
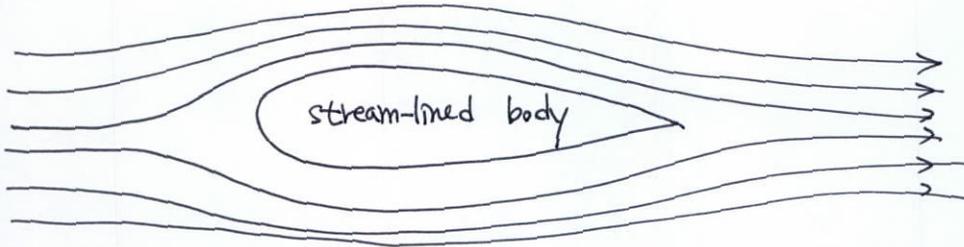
$p_1' = p_a + \rho_2 g (z_2 - z_1)$

$p_1 = p_A + \rho_1 g (z_A - z_1)$

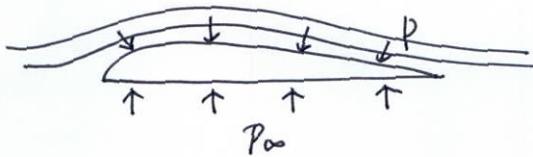
$\therefore p_A = p_a + \rho_2 g (z_2 - z_1) - \rho_1 g (z_A - z_1)$

§ pressure coefficients

$$C_p = \frac{P - P_{\infty}}{\frac{1}{2} \rho U^2}$$

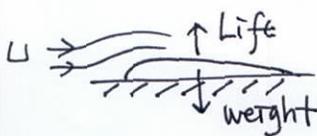


• Lift



$$L = \int_A (P_{\infty} - P) dA_{\perp} = \int_A \left(\frac{1}{2} C_p \rho U^2 \right) dA = \frac{1}{2} \rho U^2 \int_A C_p(x) dA$$

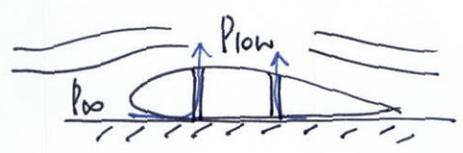
e.g. plane (가자비), ray (가오리)



as $U \uparrow$ - Lift \uparrow

	(stip speed)	(lift-off speed)
$U: 0$	0.2 m/s	0.5 m/s
	quiet	fin
		beating
		to stay
		on
		bottom
		actively
		dig in
		or
		take-off

e.g. slotted sand dollar (성하)

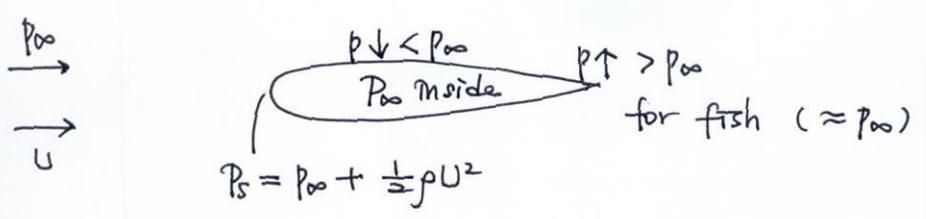


presence of slots - reduces lift

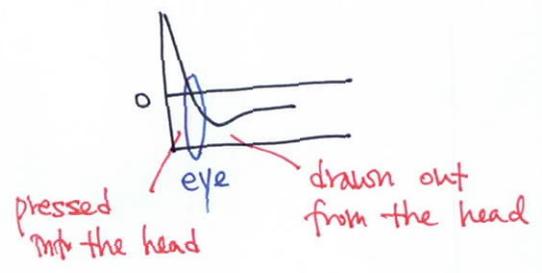
- raise the dislodgement speed by ~ 20%

§ Pressure distributions around flexible organisms

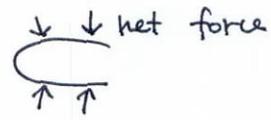
e.g. the form of fishes



fish eye location : $C_p = 0$ or $p = P_{\infty}$

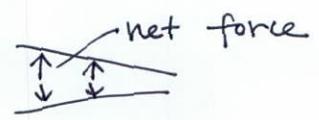


where $p > P_{\infty}$



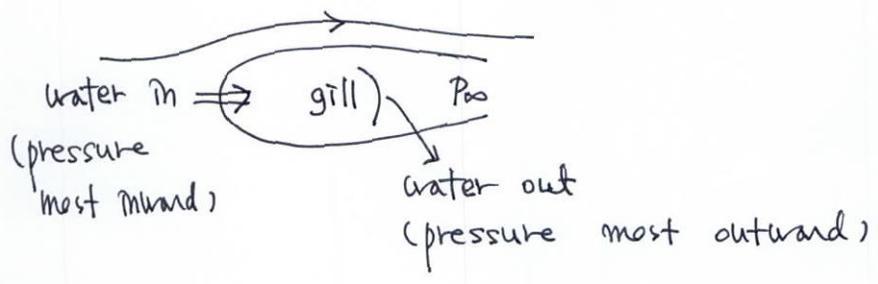
~ compression - resistant skull

where $p < P_{\infty}$



~ flexible but minimally extensible skin to maintain normal shape

ram ventilation (augmentation of ventilation by the motion of fish)



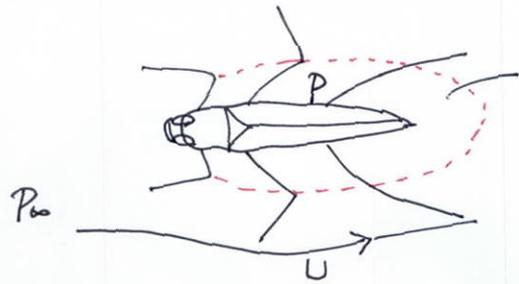
e.g. engulfing whales (Fig. 4.11)



$$\begin{aligned} \text{stretching force} &= \int (p_{\infty} - p) dA \\ &= \frac{1}{2} \rho U^2 \int_A C_p(x) dA \end{aligned}$$

sufficient force to expand cavity if $U = 3\text{m/s}$
 \approx speed of feeding fin whales

e.g. beetle's bubble



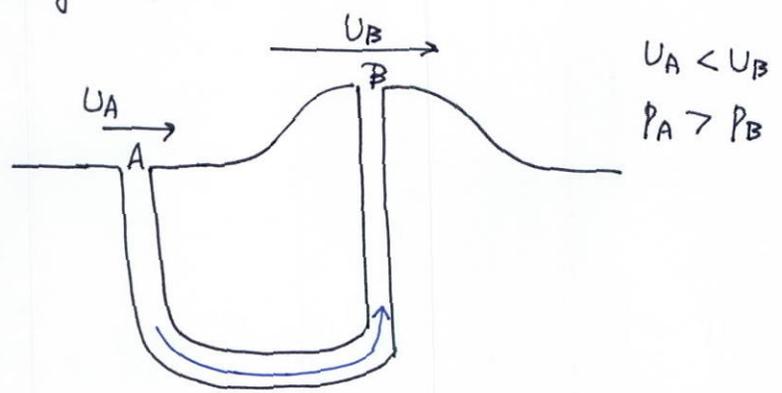
to maintain the bubble, dissolved gas must diffuse inward

$$\begin{aligned} P &< P_{\infty} \\ P &< P_{\text{atm}} \end{aligned}$$

bubble = physical gill
 (permanent only in moving water saturated with air at P_{atm})

* to be subatmospheric, live in very shallow and very rapid water
 $P_{\infty} = P_{\text{atm}} + \rho gh$

§ Forcing further flows



e.g. prairie dog's burrows : ventilation, olfactory assistance

e.g. African termites' mounds : ventilation

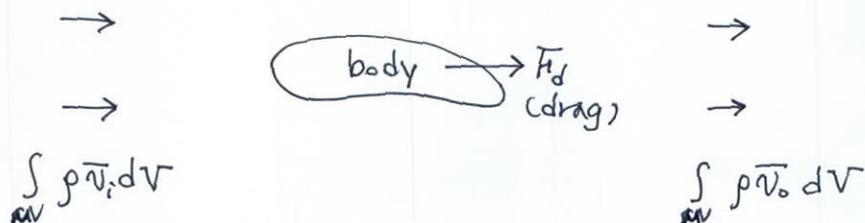
e.g. seawater moving across ripples of sand



e.g. breathing volcano

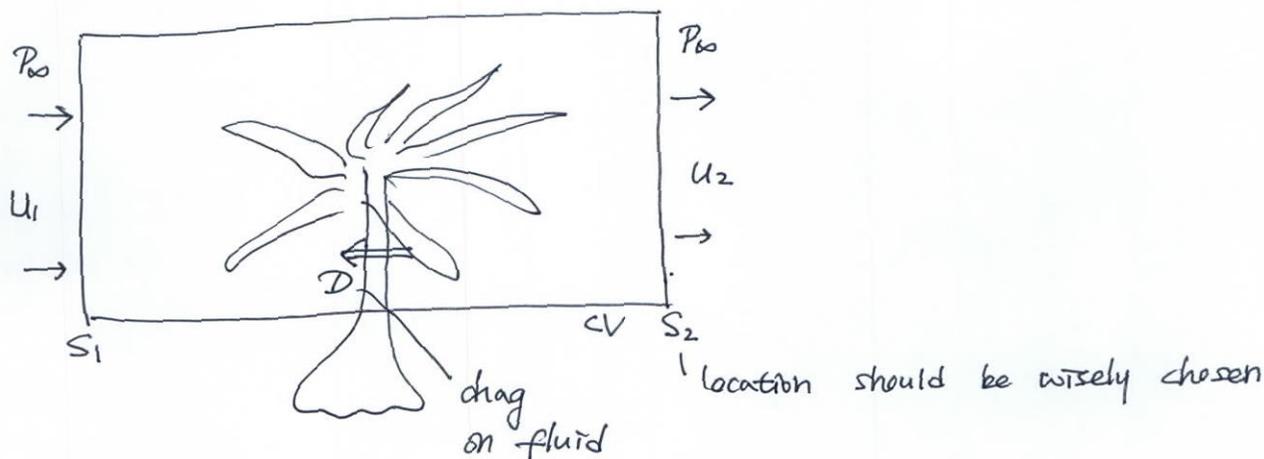
§ Momentum

§ Indirect force measurements



$$\frac{D}{Dt} \int_{CV} \rho \bar{u} dV = F_D$$

If the moving fluid exerts a drag on a body, then the body must remove momentum from the fluid at a rate that just balances its drag.



$$\rightarrow -D = \int_{CS} \rho u v_{rn} dA$$

$$-D = \int_{S_1} \rho u_1 (-u_1) dA + \int_{S_2} \rho u_2 \cdot u_2 dA$$

if $S_1 = S_2 = S$

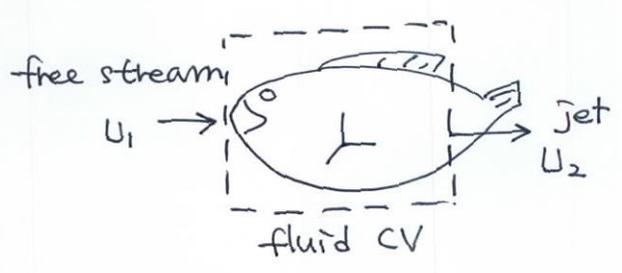
$$D = \int_S \rho (u_1^2 - u_2^2) dA$$

continuity: $\int_{S_1} \rho u_1 dA = \int_{S_2} \rho u_2 dA$

unknowns: u_2, D , # eqs: 2

§ Jet propulsion

Fig. 4.14 examples



$$\rightarrow T = \int_{CS_1} \rho U_1 (-U_1) dA + \int_{CS_2} \rho U_2 (U_2) dA$$

U_1, U_2 uniform across CS_1 and CS_2 , respectively.

$$T = \dot{m} (U_2 - U_1) > 0 \quad \text{: force on fluid} \rightarrow$$

$$\text{where } \dot{m} = \int_{CS_1} \rho U_1 dA = \int_{CS_2} \rho U_2 dA$$

by continuity

\therefore force on the swimmer \leftarrow : thrust

- power output : $P_{out} = T \cdot U_1 = \dot{m} U_1 (U_2 - U_1)$

- power input : $P_{in} = \frac{1}{2} \dot{m} (U_2^2 - U_1^2)$

- Froude propulsion efficiency

$$\eta_f = \frac{P_{out}}{P_{in}} = \frac{2U_1}{U_2 + U_1} = 1 \quad \text{if } U_1 = U_2$$

\downarrow as $U_2 \uparrow (> U_1)$

: propellers or fins achieve high η_f with only a small incremental velocity.

- real η_f of jet propulsion $> \frac{2U_1}{U_2 + U_1}$

\therefore water entrainment as rolling up into toroidal vortices