

if $U_{\text{say}} = 0.1 \text{ m/s}$

$d_{\text{xylem}} = 0.25 \text{ mm}$

$$\frac{\Delta P_{\text{ideal}}}{L} = \frac{8\mu U}{a^2} = \frac{8(10^{-3})(0.1)}{(0.125 \times 10^{-3})^2} = 51.2 \text{ kPa/m}$$

$$\frac{\Delta P_{\text{real}}}{L} = (0.7)^{-1} \frac{\Delta P_{\text{ideal}}}{L} = 73 \text{ kPa/m} \neq 7.3 \text{ kPa/m} \quad (\text{in textbook})$$

Chap. 15. Flow at very low Reynolds numbers

- creeping flow

N.-S. eq.

$$0 = -\nabla p + \mu \nabla^2 \bar{u} : \text{ Stokes eq}$$

e.g. microscopic organisms
proteins, cells, DNA, ...

- characteristics

- ① reversible
- ② mixing extremely difficult
- ③ significant wall effect

- propulsion slow

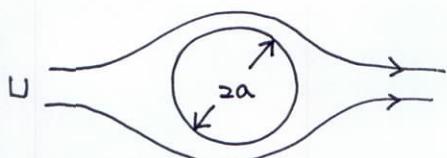
$$\text{Drag} \sim L^2$$

$$\text{Thrust} \sim \text{volume} \sim L^3 \quad (\text{engine size})$$

$$\frac{T}{D} \sim L \quad \downarrow \text{as } L \downarrow$$

§ Drag

(1) Spheres



$$D = 6\pi\mu U a \quad (Re < 1)$$

Stokes' law

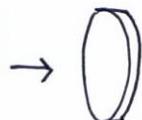
$$C_{Df} = \frac{24}{Re} + \frac{6}{1+Re^{1/2}} + 0.4 \quad (1 < Re < 2 \times 10^5)$$

$$= \frac{D}{\frac{1}{2} \rho U^2 S}$$

$$Re = \frac{U(2a)}{\nu}$$

Stokes' drag = skin friction + fore-and-aft pressure drop
 $(\frac{2}{3})$ $(\frac{1}{3})$

(2) circular disks



$$D = 16 \mu U a$$

} difference not large

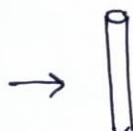


$$D = 10.67 \mu U a$$

(3) cylinders



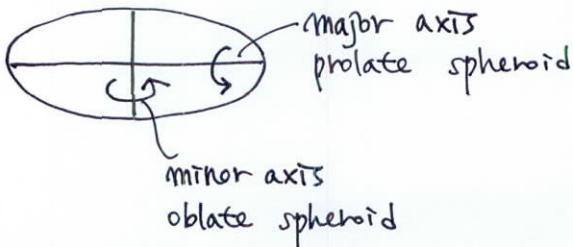
$$D = \frac{2\pi \mu U l}{\ln(\frac{l}{d}) - 0.807}$$



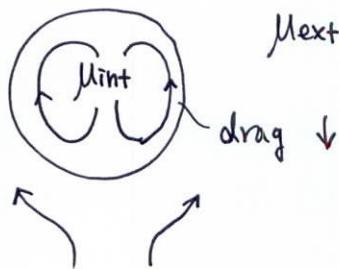
$$D = \frac{4\pi \mu U l}{\ln(\frac{l}{d}) + 0.193}$$

(4) Spheroids

ellipse

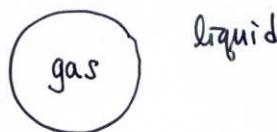


(5) Fluid spheres



$$D = 6\pi\mu_e U a \frac{1 + \frac{2}{3}(\mu_e/\mu_i)}{1 + \mu_e/\mu_i}$$

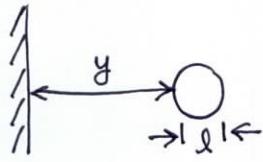
for gas bubbles



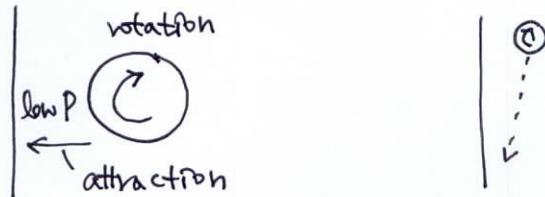
$$\frac{\mu_e}{\mu_i} \rightarrow \infty$$

$$D = 4\pi\mu_e U a$$

(6) Wall effects



Drag near wall > Drag isolated

to neglect wall effect, $\frac{y}{d} > \frac{20}{Re}$ $(Re < 1)$ 

§ Terminal velocity

force balance for a sphere falling at a constant vel.

$B = \frac{4}{3}\pi a^3 \rho_0 g$

ρ_0

$D = 6\pi\mu U a$

$W = \frac{4}{3}\pi a^3 \rho g$

$$W = B + D$$

$$\frac{4}{3}\pi a^3 (\rho - \rho_0) g = 6\pi\mu U a$$

terminal vel.

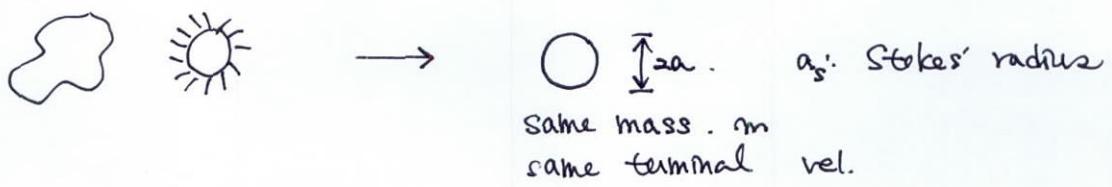
$$U = \frac{2a^2 g (\rho - \rho_0)}{9\mu} \quad (Re < 0.5)$$

Note: $U_t \sim a^2$

$$\therefore W - B \sim a^3$$

$$D \sim a$$

(1) Stokes' radius

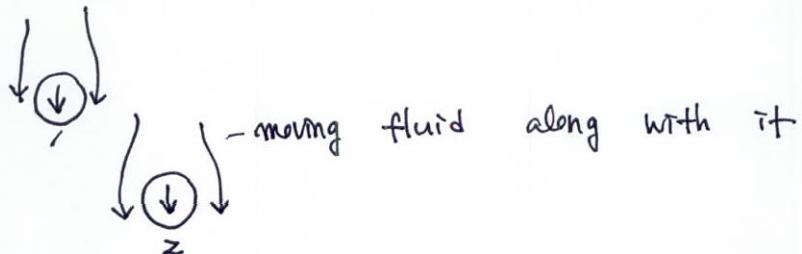


$$mg - \frac{4}{3} \rho_0 g \pi a_s^3 = 6\pi \mu U a_s$$

for air, ρ_0 term negligible

$$a_s = \frac{mg}{6\pi \mu U}$$

(2) Interactions among falling objects



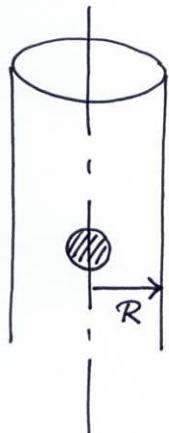
$$U_{1, \text{pair}} > U_{1, \text{isolated}}$$

$$U_{2, \text{pair}} < U_{2, \text{isolated}}$$



(3) Measuring terminal velocity

- wall effect !



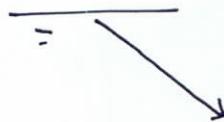
$$U_{\text{true}} = \frac{U_{\text{apparent}}}{1 + 2.4 \frac{a}{R}}$$

(4) When terminal velocity matters

- small sinkers - reduce sinking rate to travel longer distance

ways

- ① swim upward (negative geotaxis)
- ② mucus thread
- ③ enveloped in mucilage
- ④ irregular shapes



§ Propulsion at low Re

(circulation-based lift not efficient (not working)
drag-based scheme

- velocity gradient : wide, gentle
 - ~ to push against free stream, propulsive appendage should protrude through / beyond the gradient region



(1) changing the area of appendages : paddles and bristles

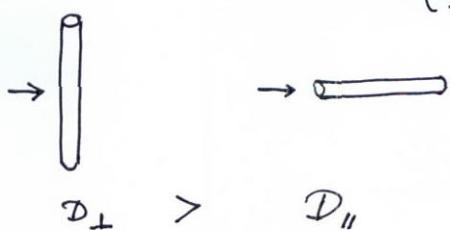
bristles - [power stroke : spread
recovery stroke : fold back]

Fig. 15.3

(2) changing the orientation of appendages : cilia and flagella : cylinders !

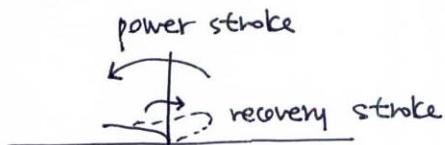
* cilium : reciprocating beat

(short: 5~10 μm in length)



$$D_{\perp} > D_{\parallel}$$

Fig. 15.4



* flagellum

planar, undulating waves

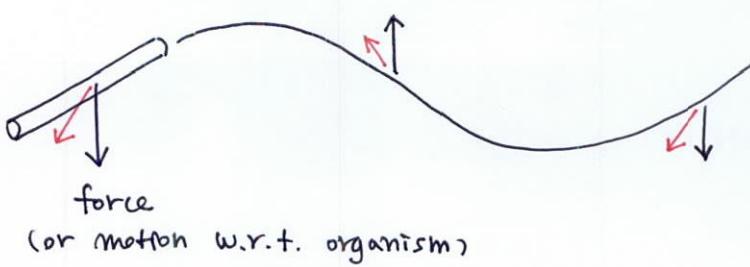
helical waves from base to tip

spiral waves

· obliquely falling cylinder



· moving oblique portion of appendage sideways



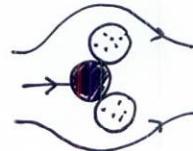
for cylinders $(Re < 1)$

$$\frac{D_{\perp}}{D_{\parallel}} = \begin{cases} 1.75 & l/a = 100 \\ 1.80 & l/a = 300 \end{cases}$$

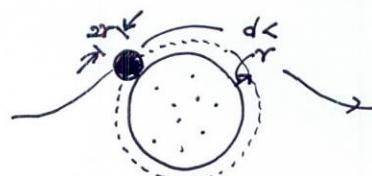
§ Filtration (suspension feeding)

- 5 mechanisms

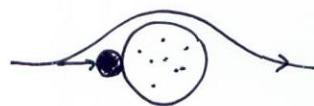
① sieving



② direct interception



③ inertial impaction



④ gravitational deposition



⑤ diffusional deposition

