

2018 Fall

“Phase Transformation *in* Materials”

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Contents for previous class

- Ternary Equilibrium: Ternary Phase Diagram**

1) Gibbs Triangle

Used to determine the overall composition

$$X_A + X_B + X_C = 1$$

2) Isothermal section $\rightarrow F = C - P$ 2 phases region

Tie line: A straight line joining any two ternary compositions

Amount of each phase present is determined by using the **Inverse Lever Rule**

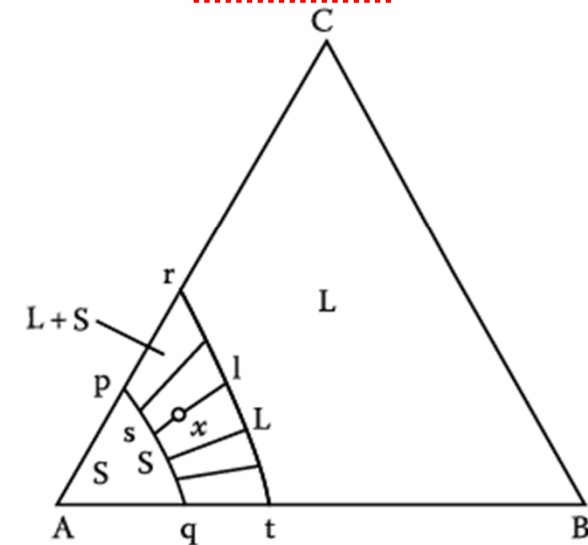
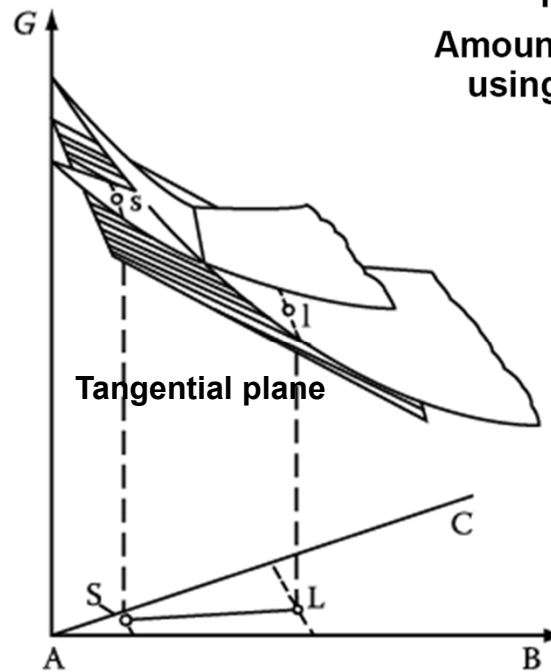
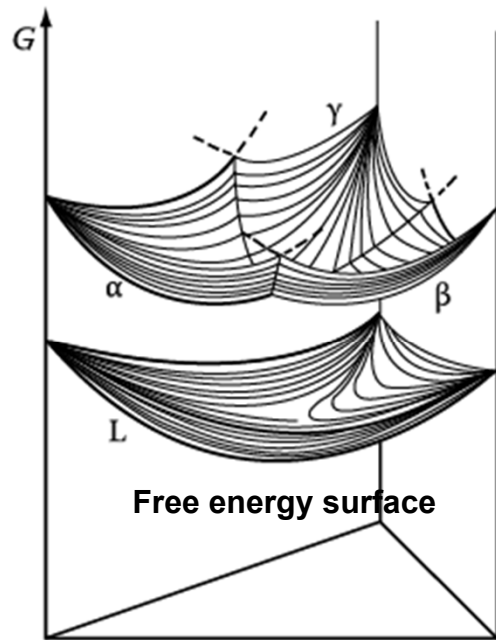
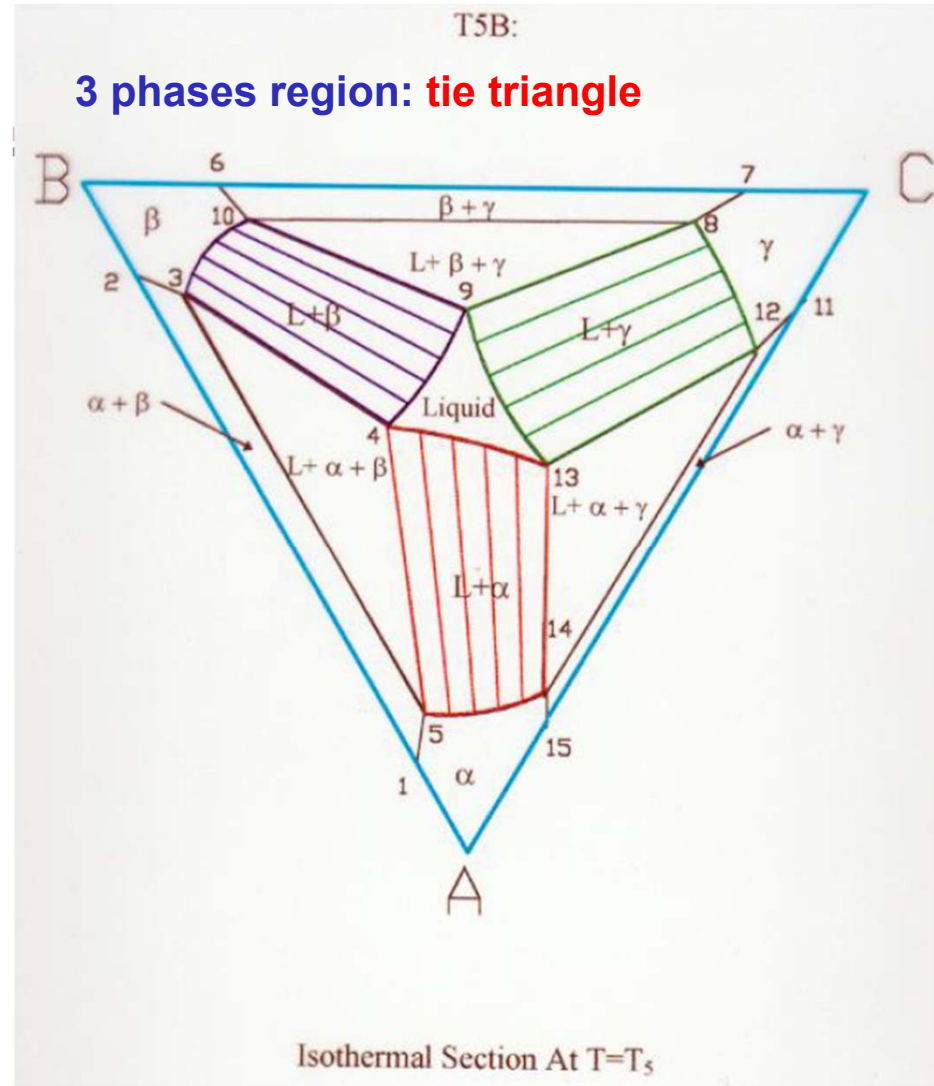
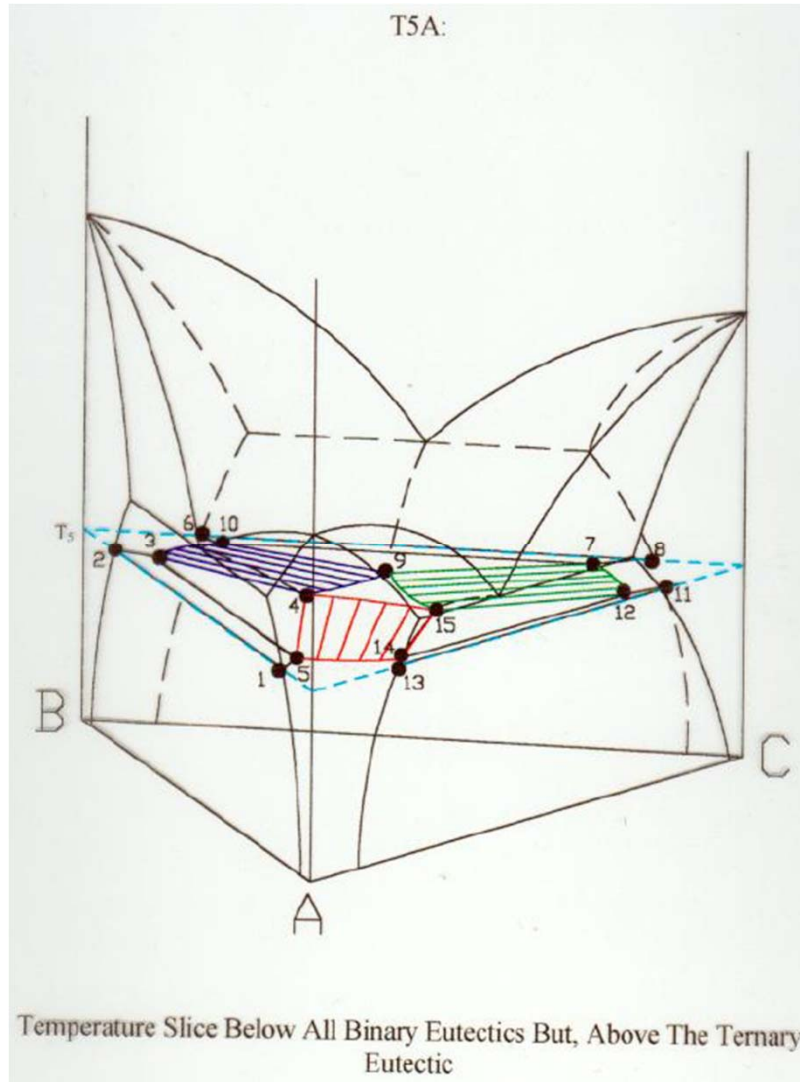


Fig. 1.41 (a) Free energy surface of a liquid and three solid phases of a ternary system.

(b) A tangential plane construction to the free energy surfaces defined equilibrium between s and l in the ternary system

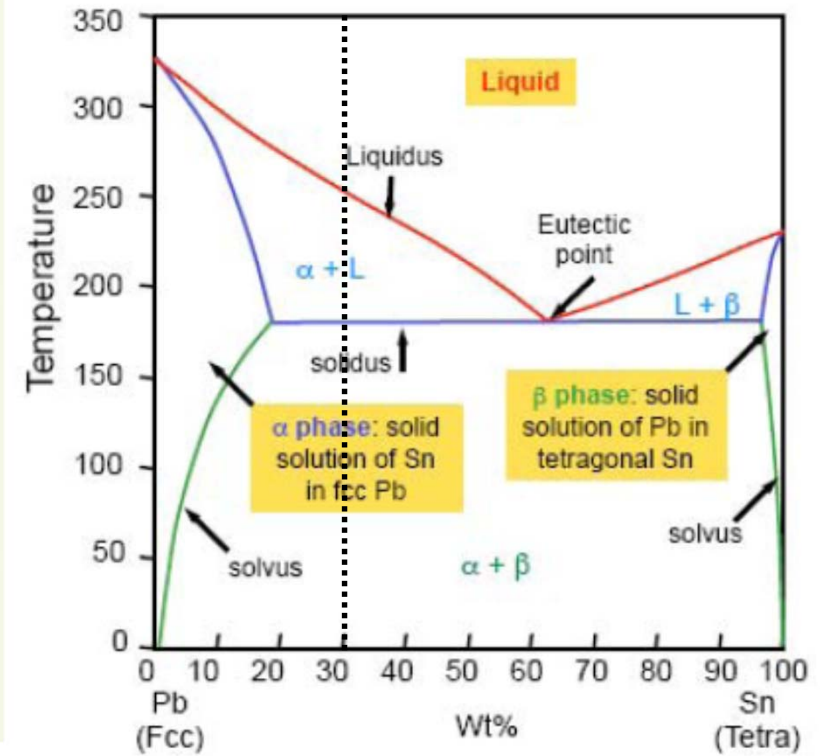
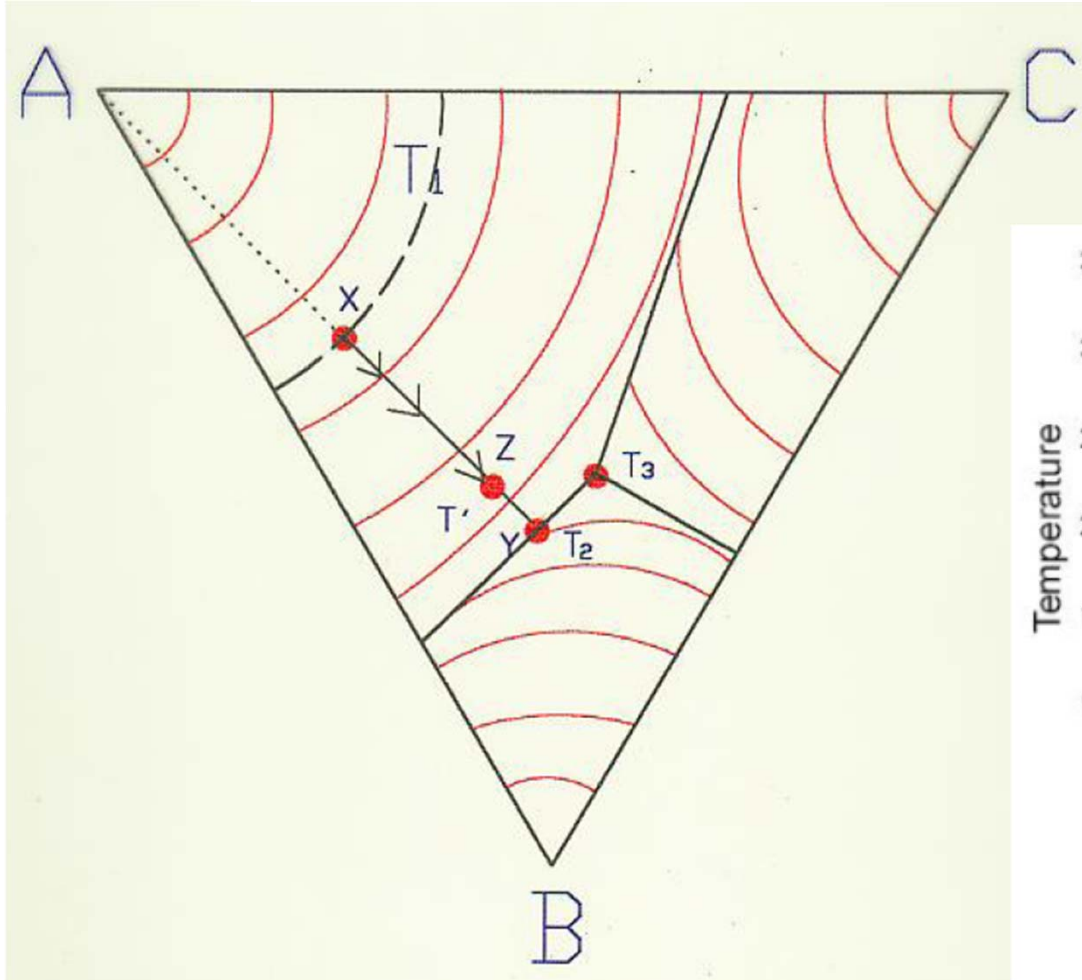
(c) Isothermal section through a ternary phase diagram

Ternary Eutectic System (with Solid Solubility)



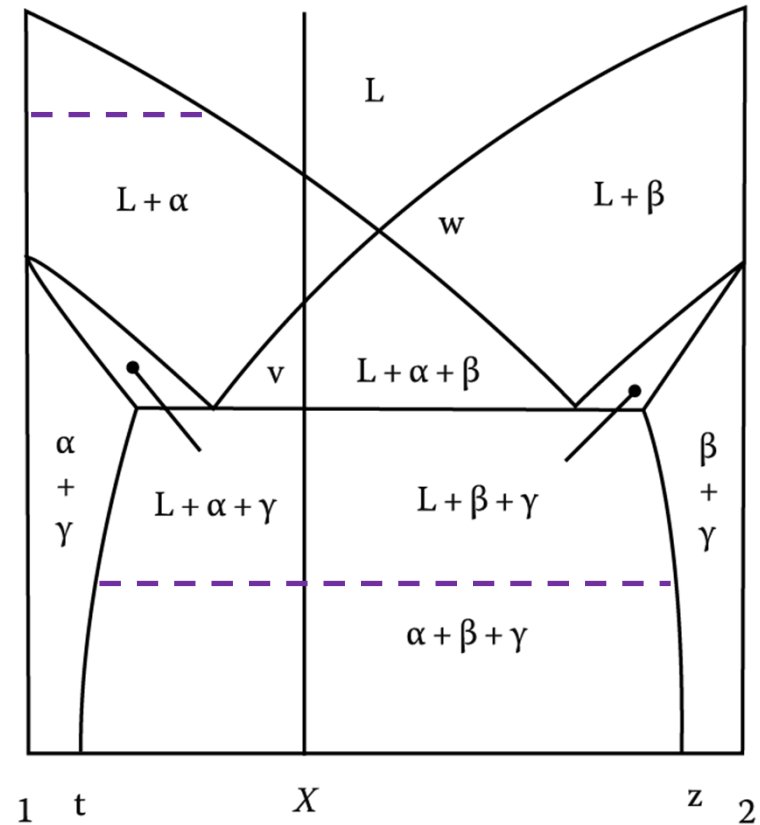
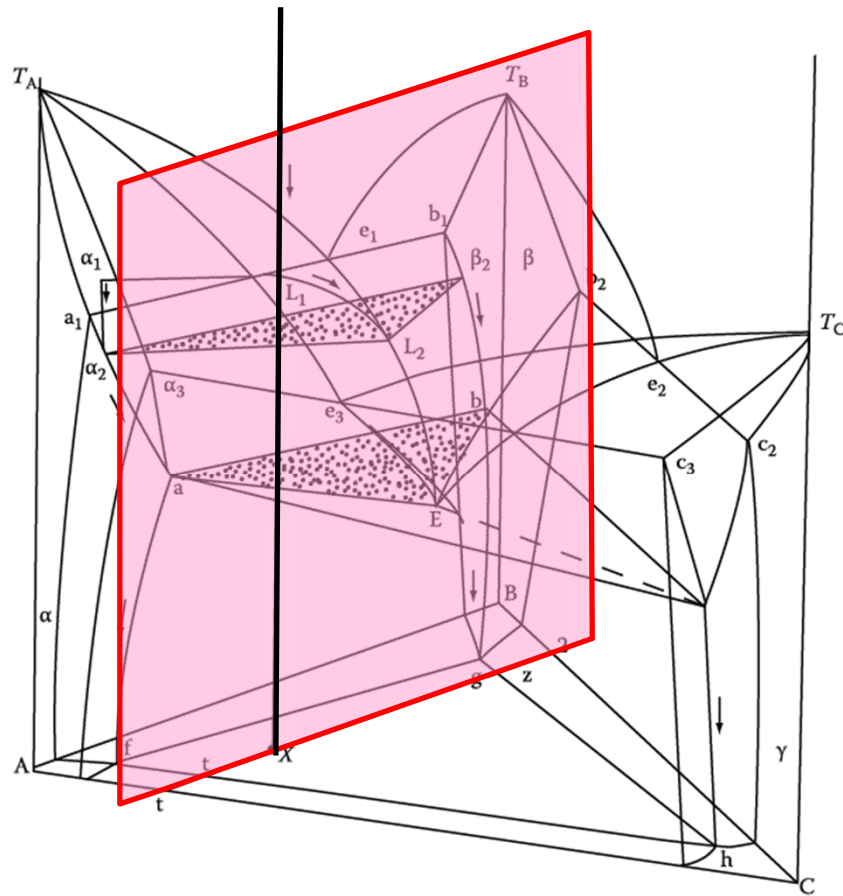
Ternary Eutectic System

3) Solidification Sequence: liquidus surface



Ternary Eutectic System

* Vertical section



* The horizontal lines are not tie lines.
(no compositional information)

* Information for equilibrium phases
at different temperatures

< Quaternary phase Diagrams >

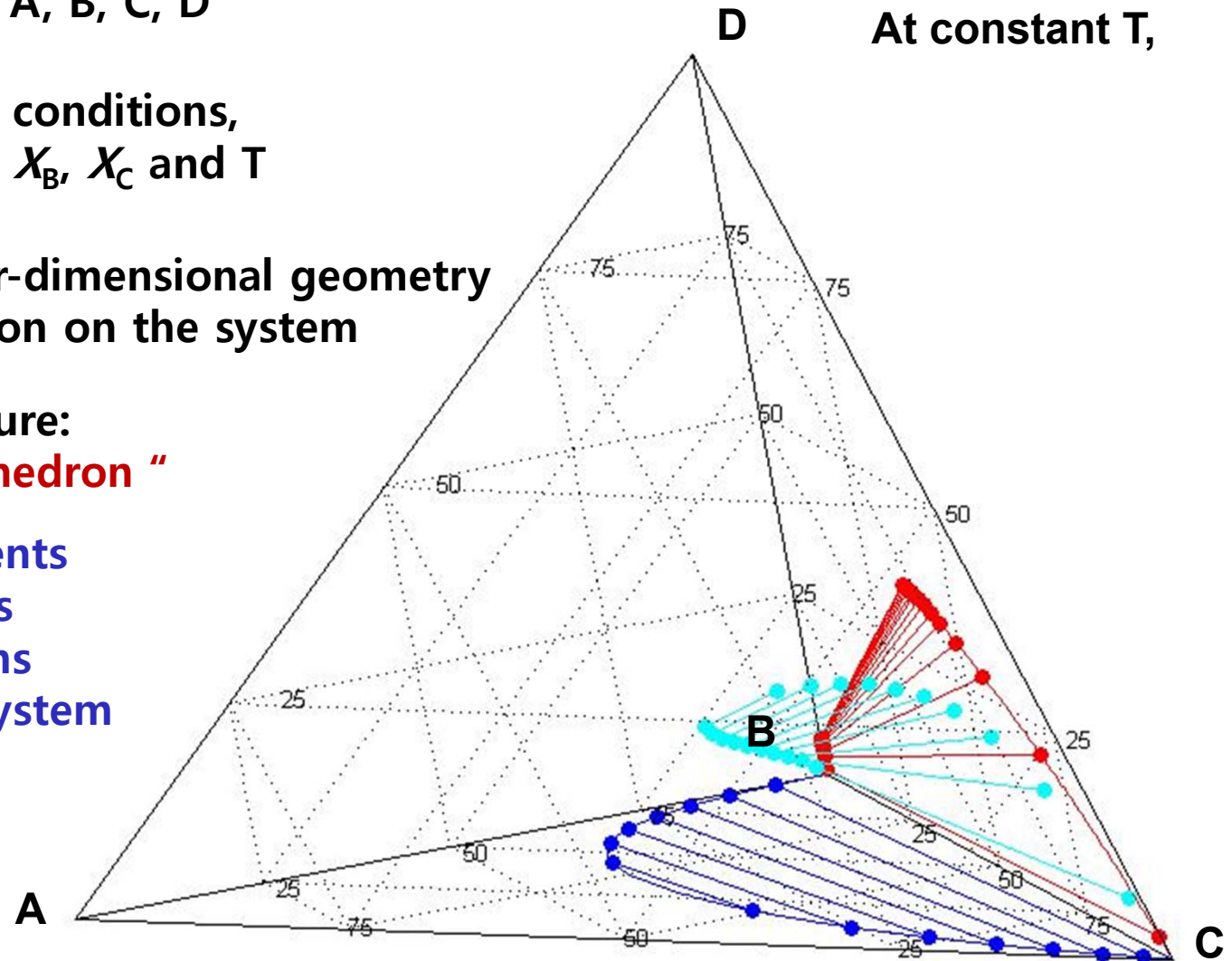
Four components: A, B, C, D

Assuming isobaric conditions,
Four variables: X_A , X_B , X_C and T

A difficulty of four-dimensional geometry
→ further restriction on the system

Most common figure:
" equilateral tetrahedron "

- 4 pure components
- 6 binary systems
- 4 ternary systems
- A quaternary system



* Draw four small equilateral tetrahedron
 → formed with edge lengths of a, b, c, d

$$a + b + c + d = 100$$

- %A = Pt = c,
- %B = Pr = a,
- %C = Pu = d,
- %D = Ps = b

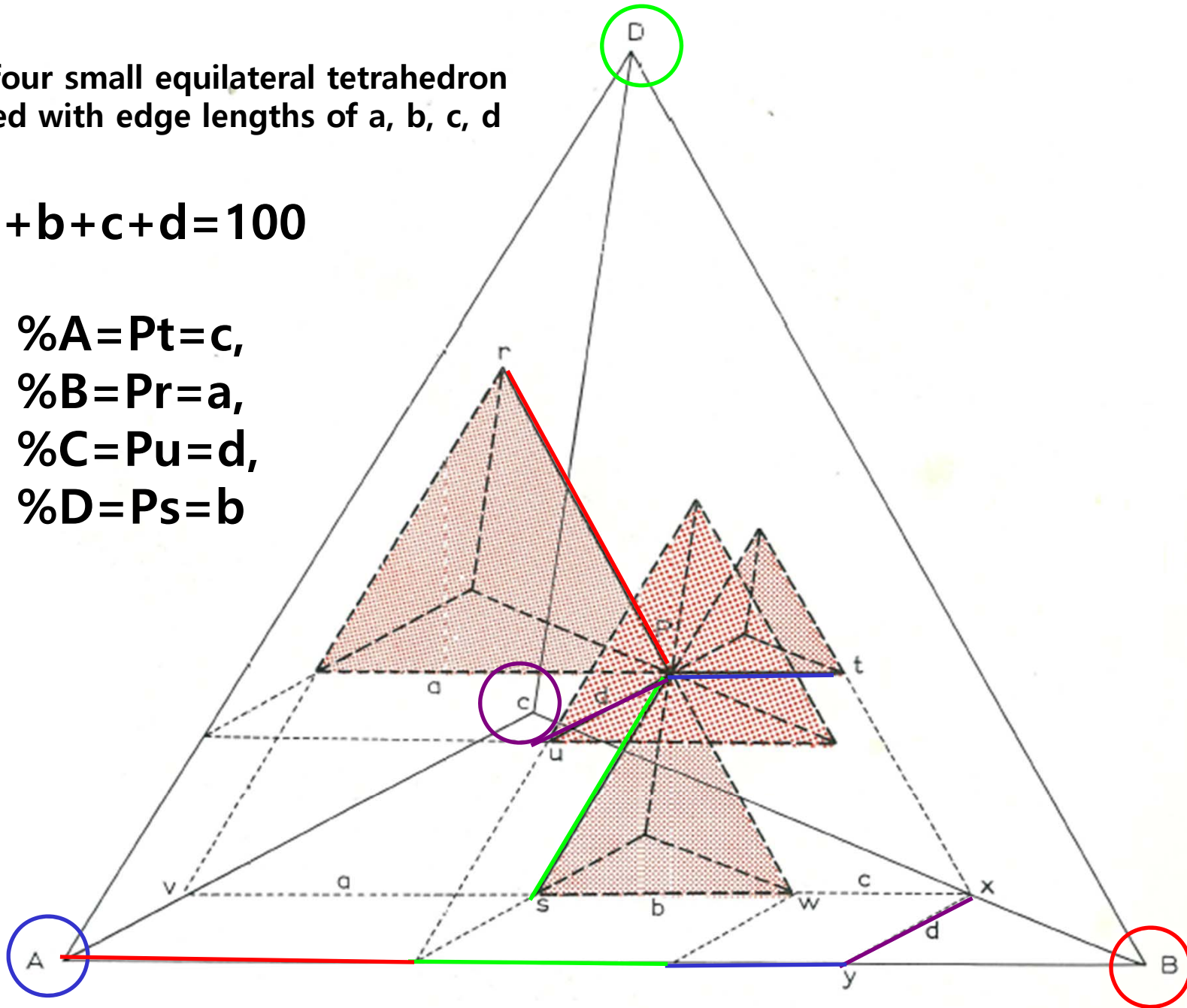
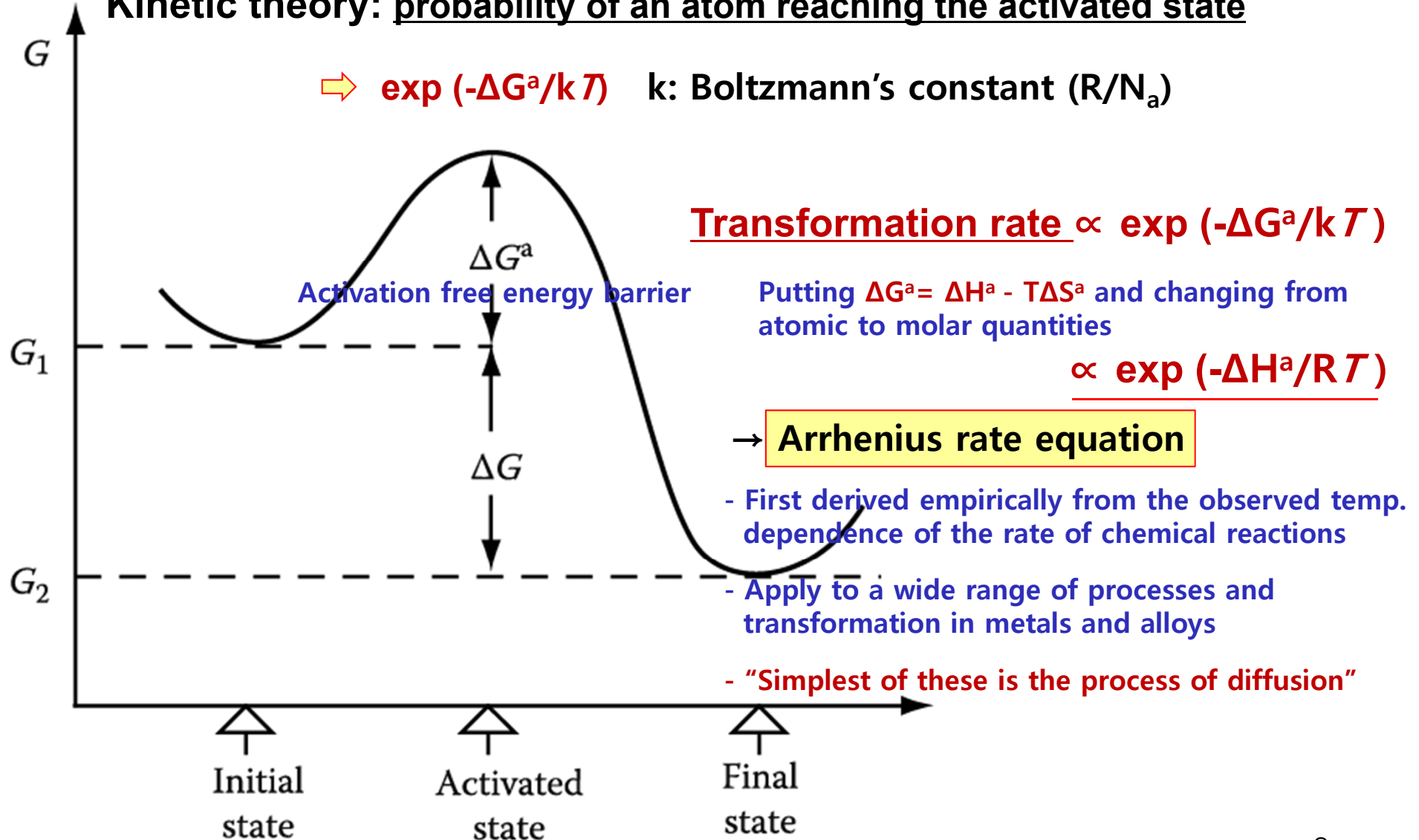


Fig. 247. Representation of a quaternary system by an equilateral tetrahedron.

1.10 The kinetics of phase transformations

Kinetic theory: probability of an atom reaching the activated state



Contents in Phase Transformation

Background
to understand
phase
transformation

(Ch1) Thermodynamics and Phase Diagrams

(Ch2) Diffusion: Kinetics

(Ch3) Crystal Interface and Microstructure

Representative
Phase
transformation

(Ch4) Solidification: Liquid \rightarrow Solid

(Ch5) Diffusional Transformations in Solid: Solid \rightarrow Solid

(Ch6) Diffusionless Transformations: Solid \rightarrow Solid

Contents for today's class

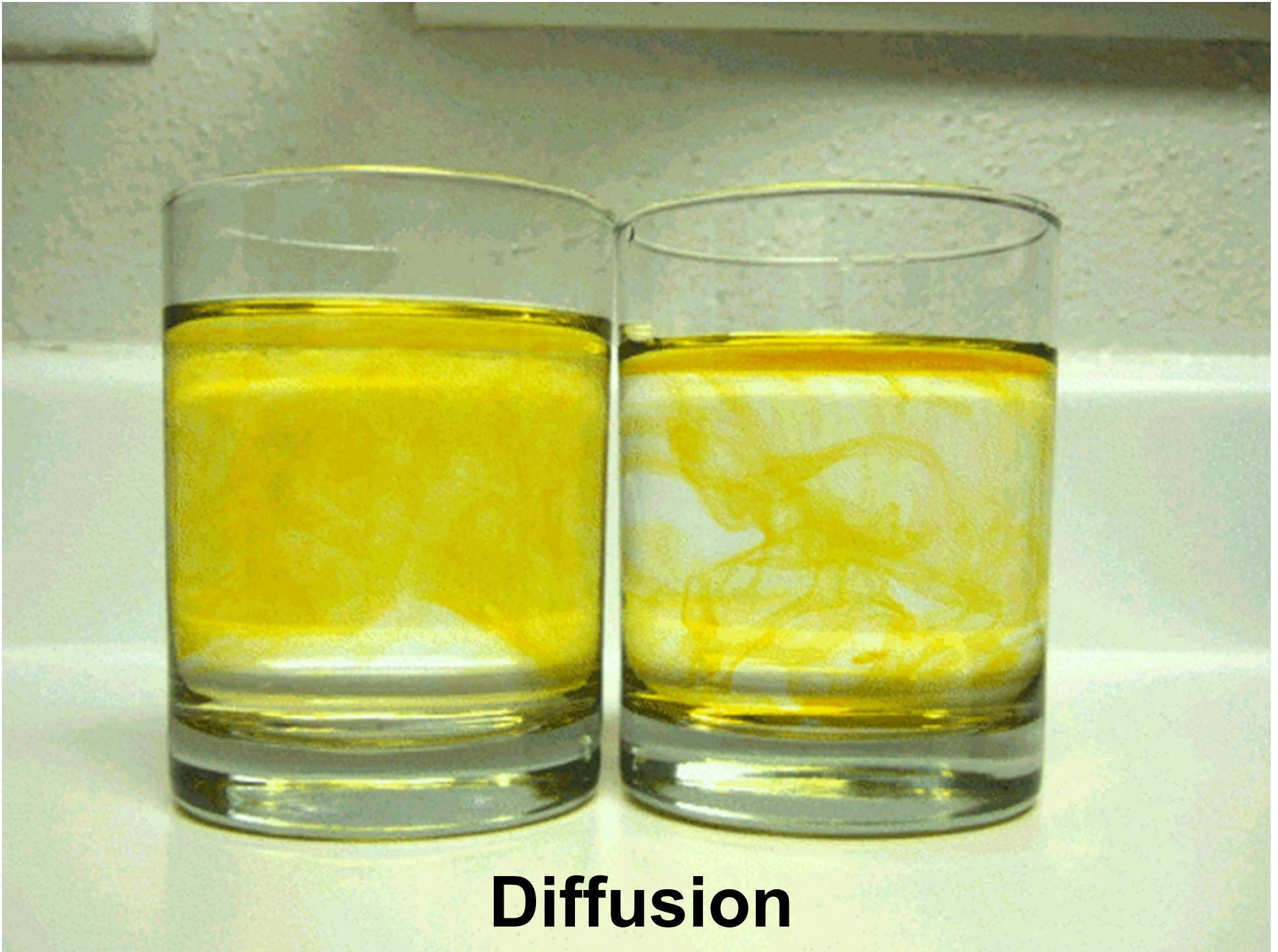
- **Diffusion**
- **Interstitial Diffusion – Fick's First Law**
 - Effect of Temperature on Diffusivity
 - Nonsteady-state diffusion – **Fick's Second Law**
- **Solutions to the diffusion equations**
- **Substitution Diffusion**
 1. Self diffusion in pure material
 2. Vacancy diffusion
 3. Diffusion in substitutional alloys

Q. What is the driving force for diffusion?

⇒ a concentration gradient (x)

⇒ a chemical potential gradient (o)

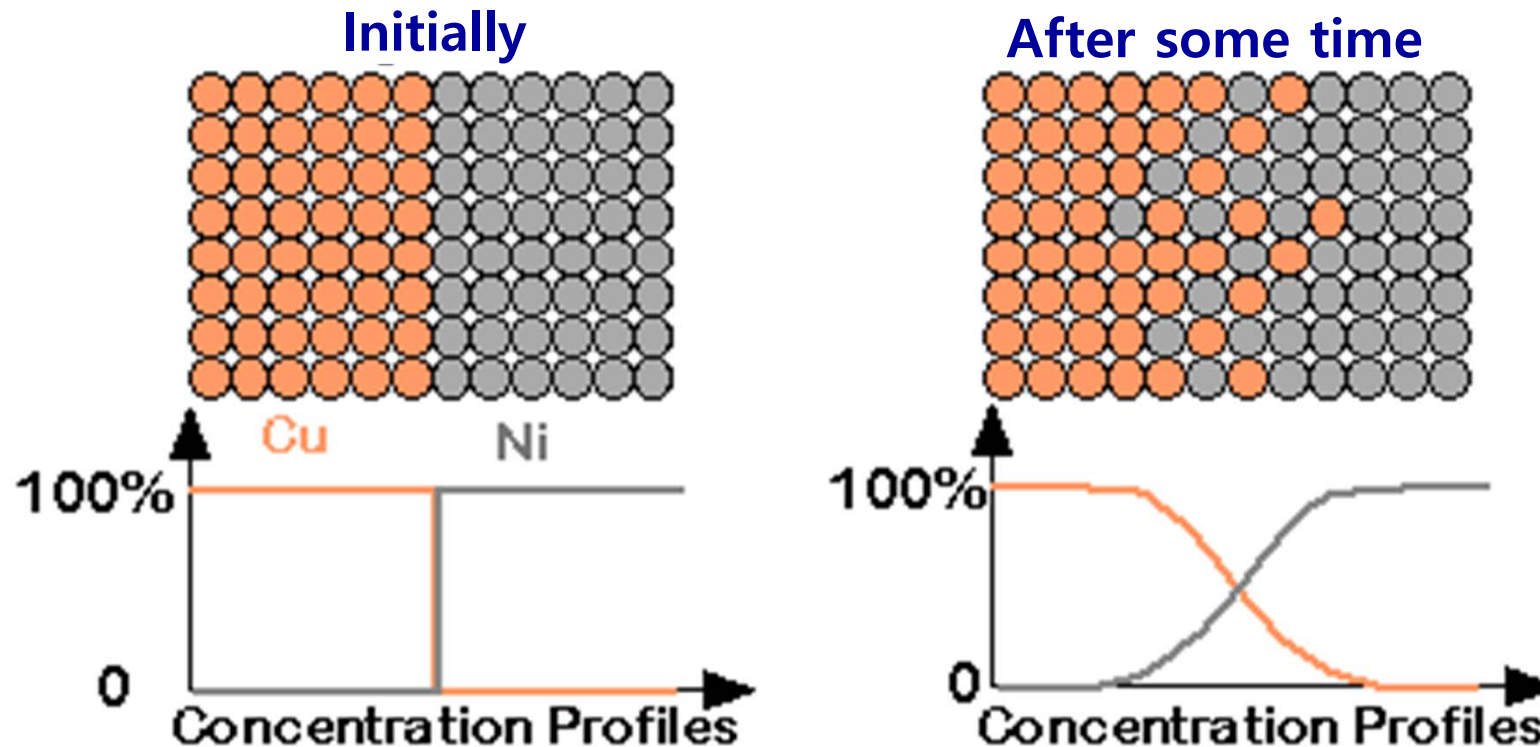
Diffusion ⇒ Movement of atoms to reduce its chemical potential μ .



Diffusion

Diffusion: THE PHENOMENON

- **Interdiffusion:** in a solid with more than one type of element (an alloy), atoms tend to migrate from regions of large concentration.



< Diffusion couple between Cu and Ni >

Mechanism of Solid-State Diffusion

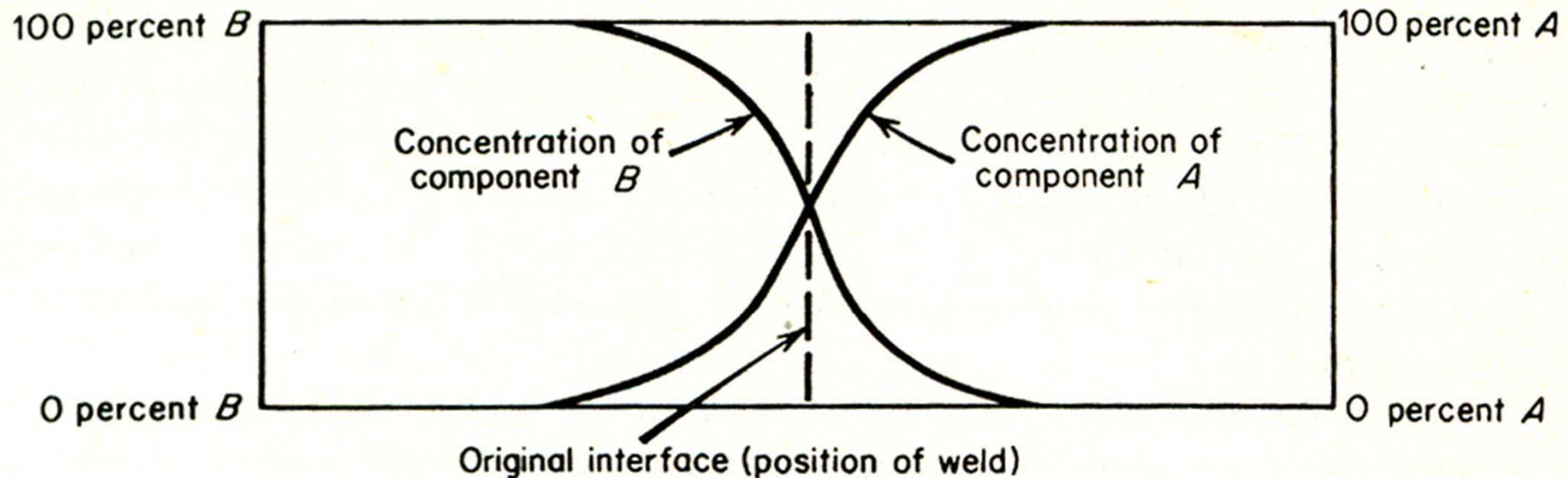


Fig. 10.6 Curves showing concentration as a function of distance along a diffusion couple. Curves of this type are usually called penetration curves.

Substitutional diffusion ↔ **Interstitial diffusion**

R.E. Reed-Hill, Physical Metallurgy Principles

Diffusion : Movement of atoms to reduce its chemical potential μ .

Driving force: Reduction of G

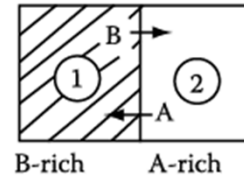
Down-hill diffusion

movement of atoms from a high concentration region to low concentration region.

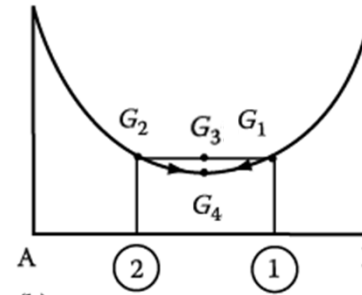
Up-hill diffusion

movement of atoms from a low concentration region to high concentration region

'down-hill' diffusion

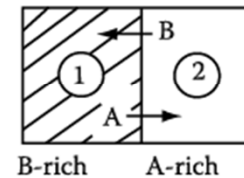


(a)

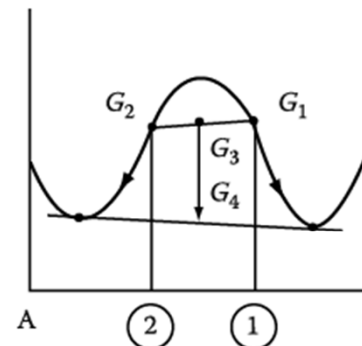


(b)

'up-hill' diffusion

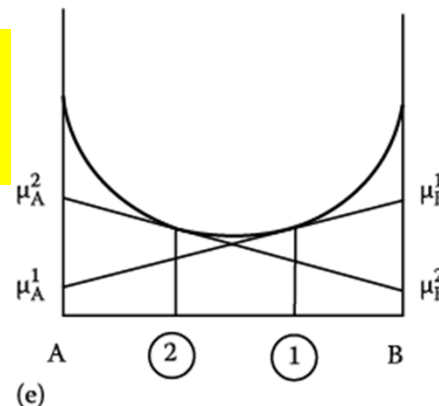


(c)

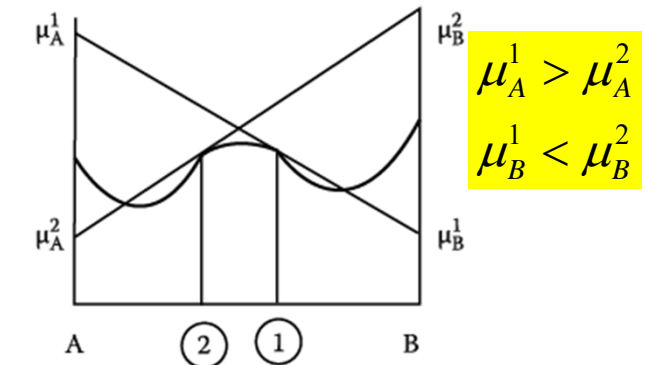


(d)

$$\begin{aligned} \mu_A^1 &< \mu_A^2 \\ \mu_B^1 &> \mu_B^2 \end{aligned}$$



(e)

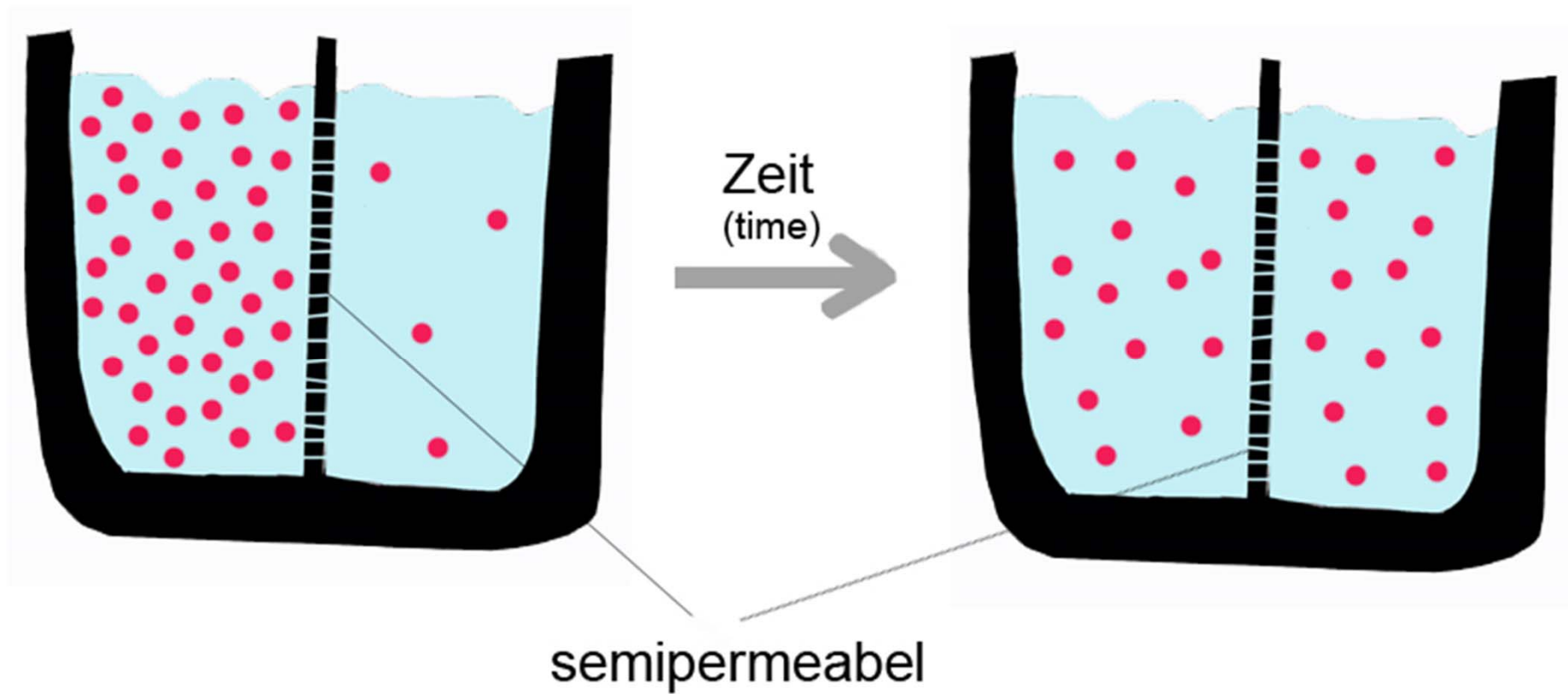


(f)

$$\begin{aligned} \mu_A^1 &> \mu_A^2 \\ \mu_B^1 &< \mu_B^2 \end{aligned}$$

Down-hill diffusion

movement of atoms from a high concentration region to low concentration region.



Diffusion

Movement of atoms to reduce its chemical potential μ .

Driving force: Reduction of G

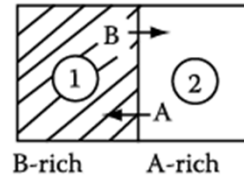
Down-hill diffusion

movement of atoms from a high concentration region to low concentration region.

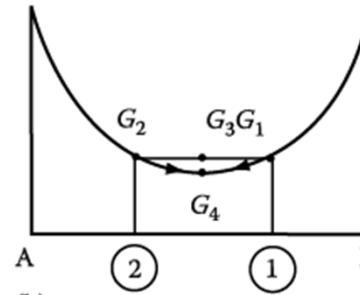
Up-hill diffusion

movement of atoms from a low concentration region to high concentration region

'down-hill' diffusion

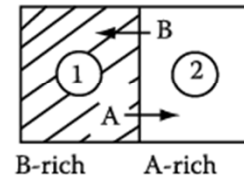


(a)

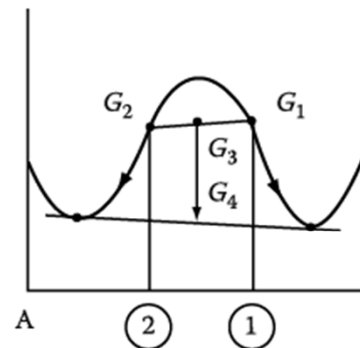


(b)

'up-hill' diffusion

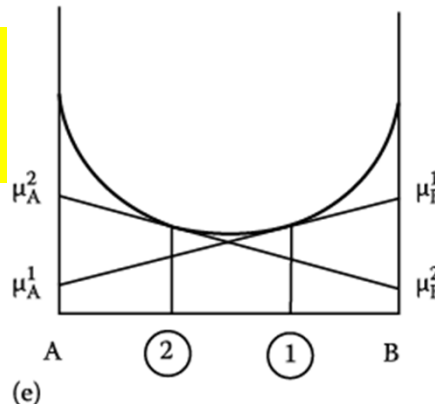


(c)

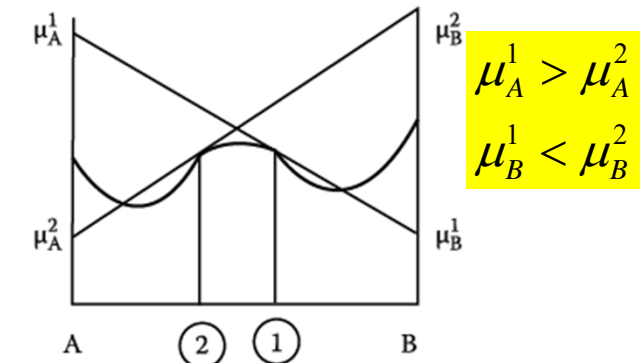


(d)

$$\begin{aligned} \mu_A^1 &< \mu_A^2 \\ \mu_B^1 &> \mu_B^2 \end{aligned}$$



(e)



(f)

$$\begin{aligned} \mu_A^1 &> \mu_A^2 \\ \mu_B^1 &< \mu_B^2 \end{aligned}$$

5.5.5 Spinodal Decomposition

Spinodal mode of transformation has no barrier to nucleation

: describing the transformation of a system of two or more components in a metastable phase into two stable phases

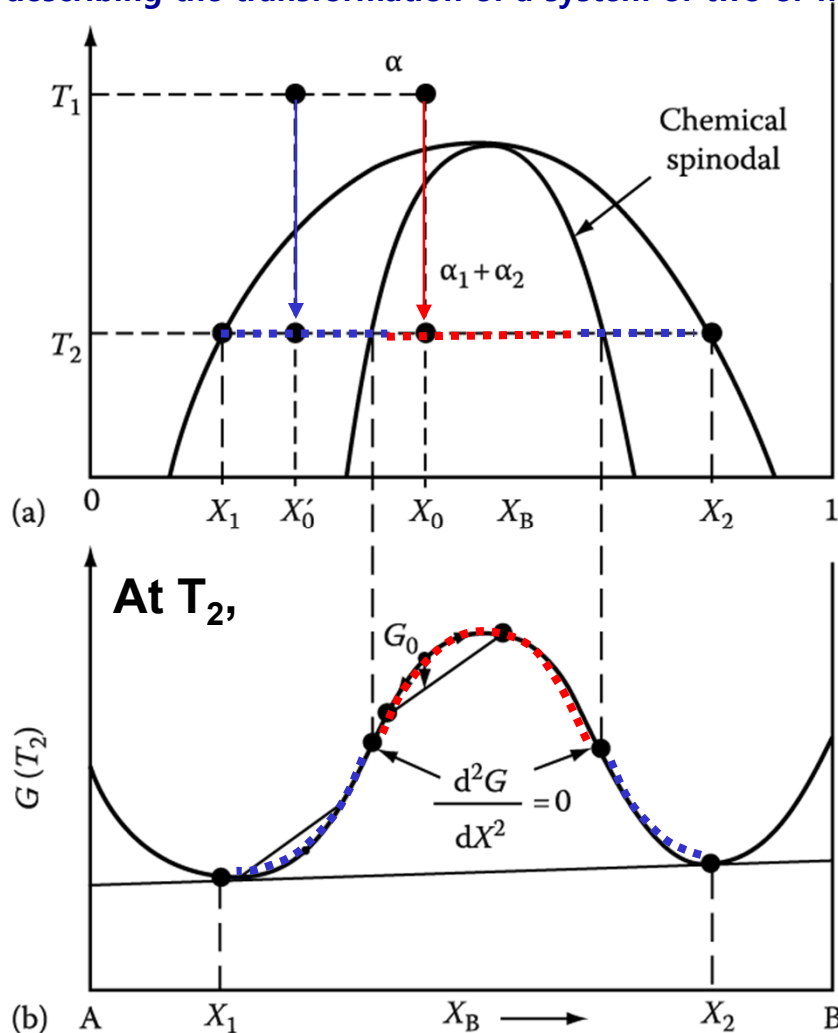


Fig. 5.38 Alloys between the spinodal points are unstable and can decompose into two coherent phases α_1 and α_2 without overcoming an activation energy barrier. Alloys between the coherent miscibility gaps and the spinodal are metastable and can decompose only after nucleation of the other phase.

How does it differ between **inside** and **outside the inflection point** of Gibbs free energy curve?

1) **Within the spinodal** $\frac{d^2G}{dX^2} < 0$

: phase separation by small fluctuations in composition/
"up-hill diffusion"

2) If the alloy lies **outside the spinodal**, small variation in composition leads to an increase in free energy and the alloy is therefore **metastable**.

The free energy can only be decreased if nuclei are formed with a composition very different from the matrix.

→ **nucleation and growth**

: "down-hill diffusion"

a) Composition fluctuations within the spinodal

b) Normal down-hill diffusion outside the spinodal

up-hill diffusion

down-hill diffusion

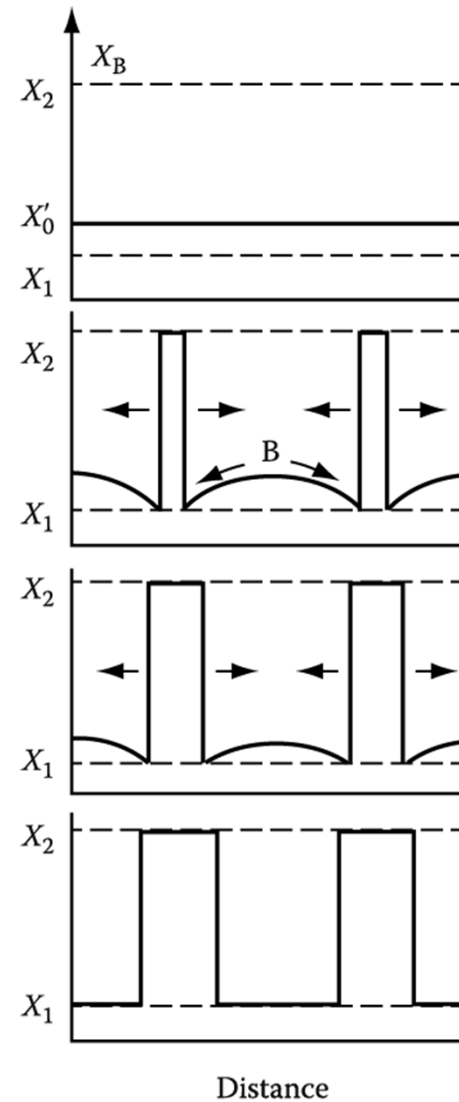
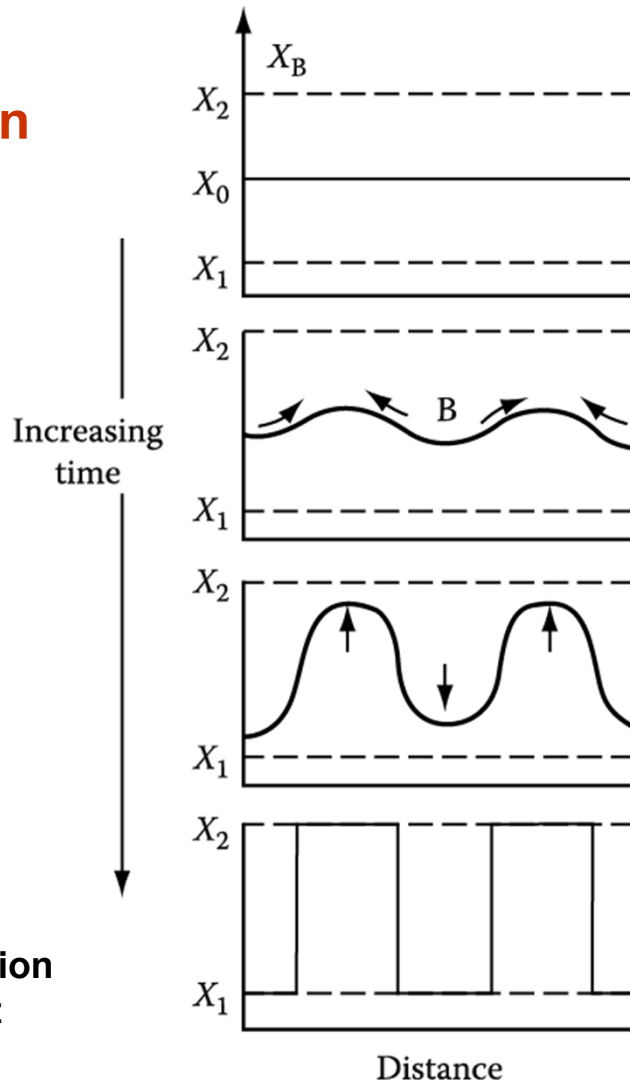


Fig. 5.39 & 5.40 schematic composition profiles at increasing times in (a) an alloy quenched into the spinodal region (X_0 in Figure 5.38) and (b) an alloy outside the spinodal points (X'_0 in Figure 5.38)

Diffusion

Diffusion : Mechanism by which matter transported through matter

What is the driving force for diffusion?

⇒ a concentration gradient (x)

⇒ a chemical potential gradient (o)

But this chapter will explain with

“concentration gradients for a convenience”.

(Down-hill diffusion)

Q. Interstitial diffusion vs Substitutional diffusion

Fick's First Law of Diffusion

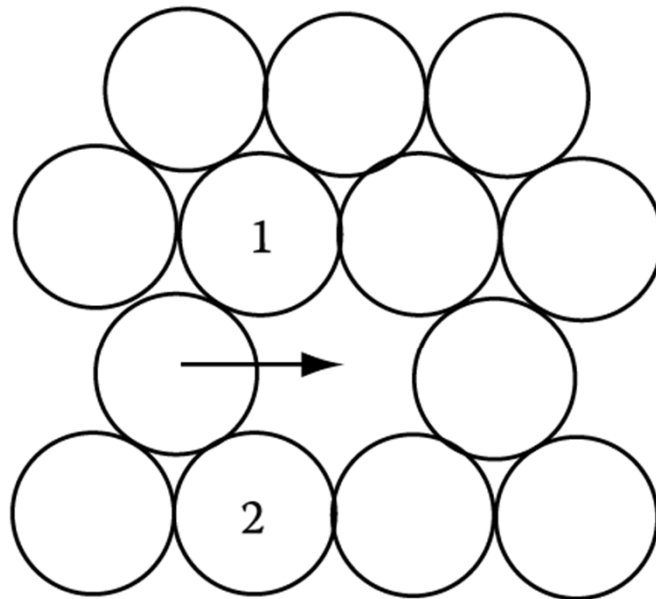
$$J_B = -D_B \frac{\partial C_B}{\partial X}$$

Atomic mechanisms of diffusion

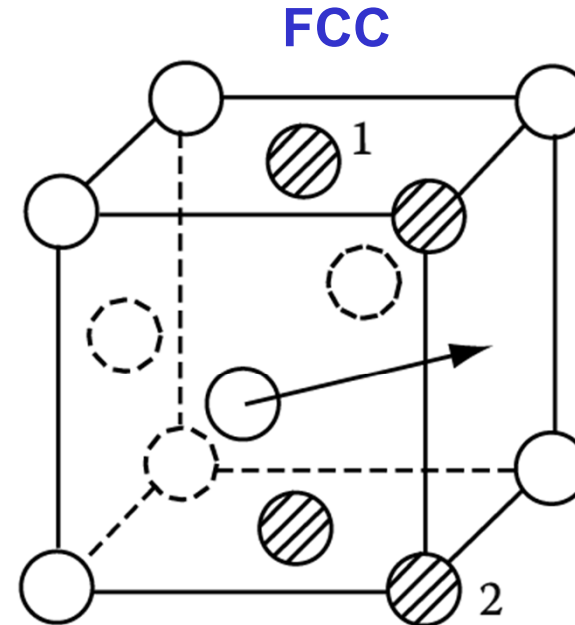
The mean vibrational energy is $3kT$, therefore the amplitude of thermal oscillation of an atom increases markedly with **temperature**. If a neighboring atom site is vacant, then the oscillation atom can change site with a certain probability (**atom jumping**). The probability consists of **creation of the vacant site** and movement of a neighboring atom into the site.

Substitutional vs. Interstitial diffusion

Substitutional diffusion



(a) # of vacant site



(b)

Interstitial diffusion

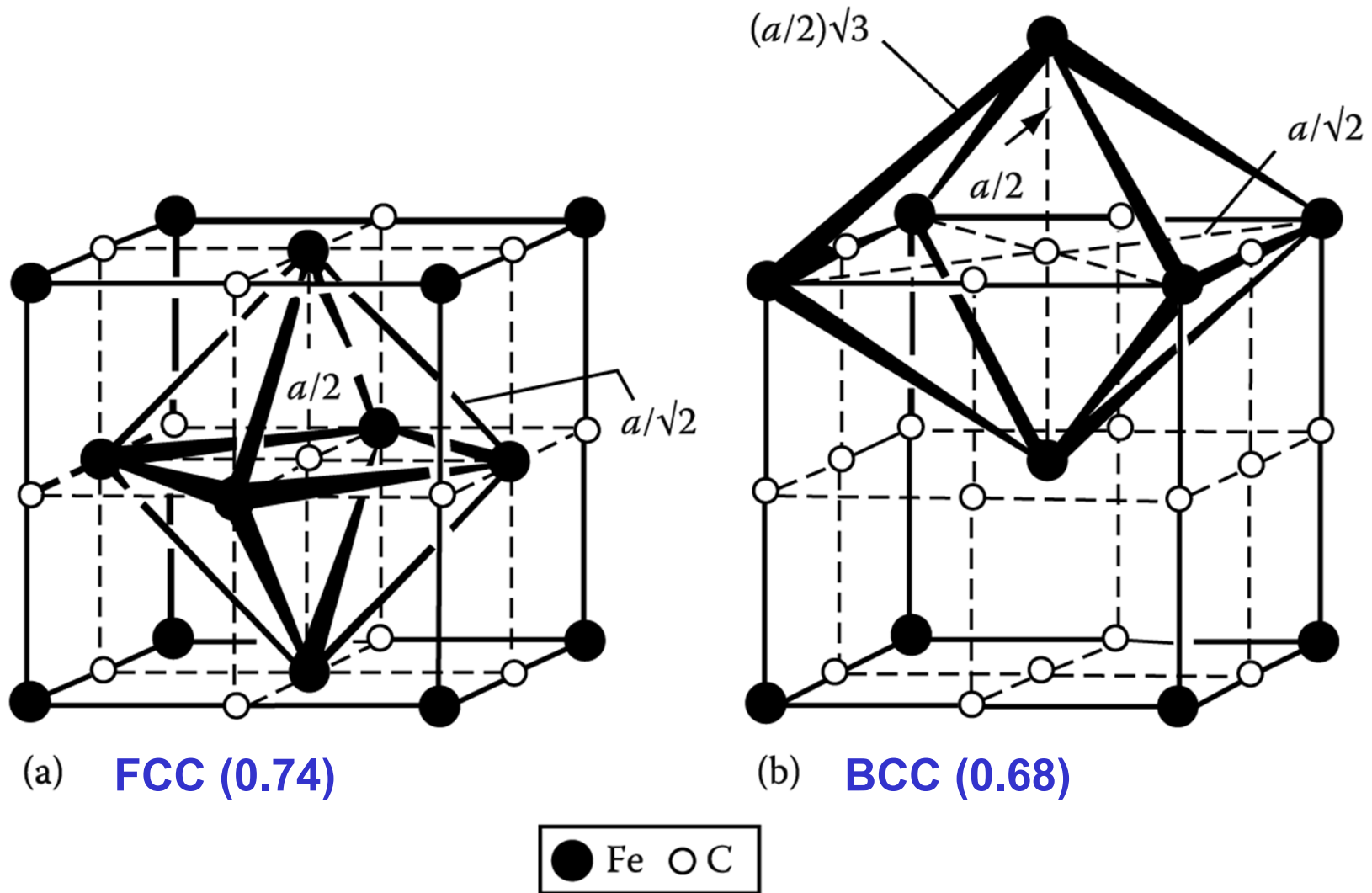
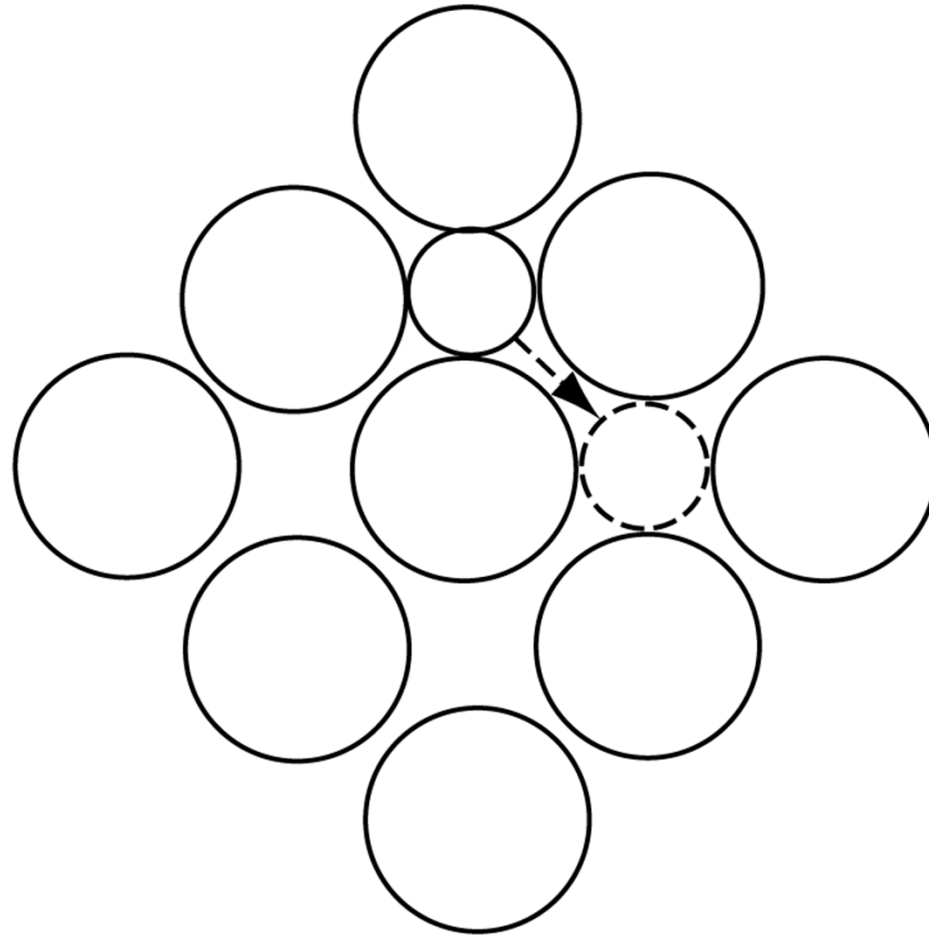


Fig. 2.3 (a) Octahedral interstices (O) in an fcc crystal. (b) Octahedral interstices in a bcc crystal.

Interstitial diffusion



How interstitial diffusion differs from substitutional diffusion?

Interstitial diffusion

Random jump of solute B atoms in a dilute solid solution of a simple cubic lattice

Assume that there is **no lattice distortion** and also that there are **always six vacant sites around the diffusion atom**.

$$J_B = \frac{1}{6}\Gamma_B n_1 - \frac{1}{6}\Gamma_B n_2$$



J_B : Net flux of B atom

Γ_B : Average jump rate of B atoms

n_1 : # of atoms per unit area of plane 1

n_2 : # of atoms per unit area of plane 2

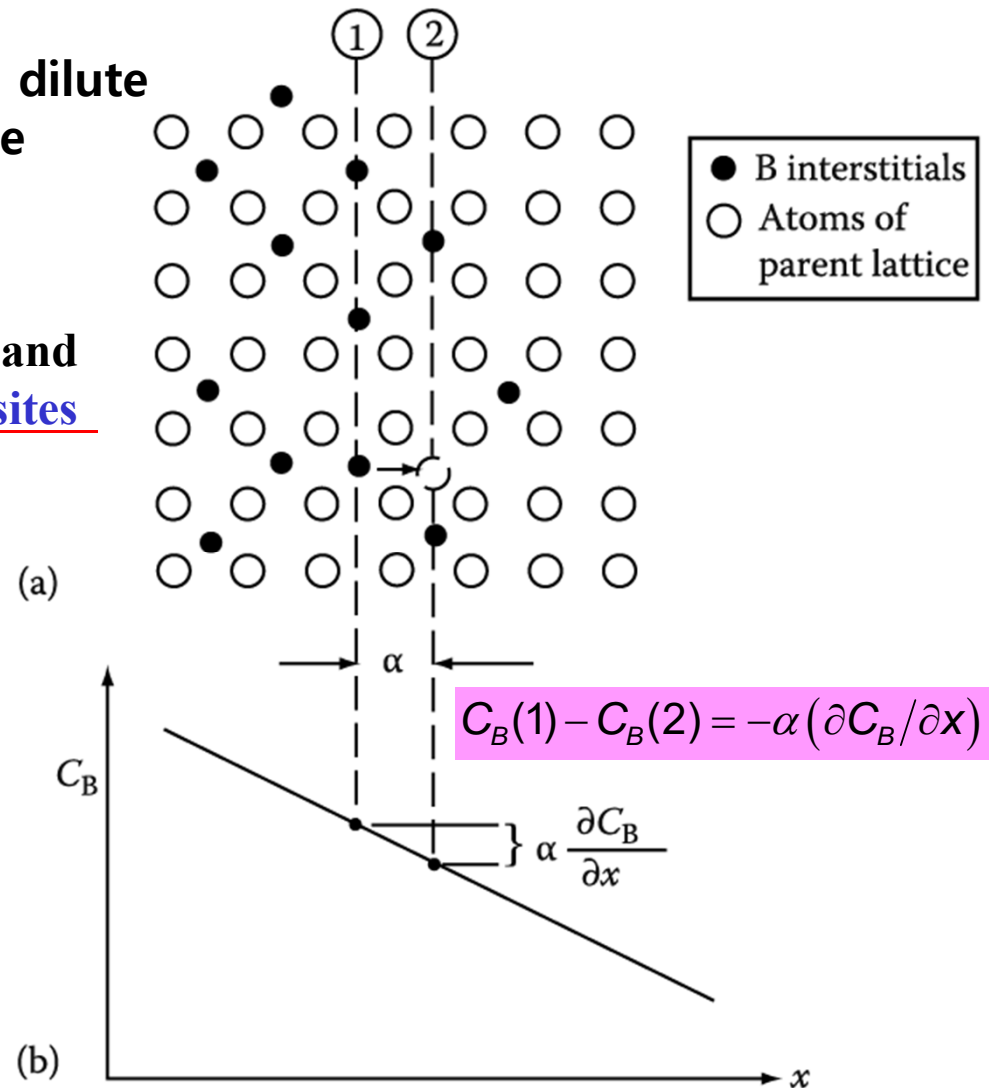


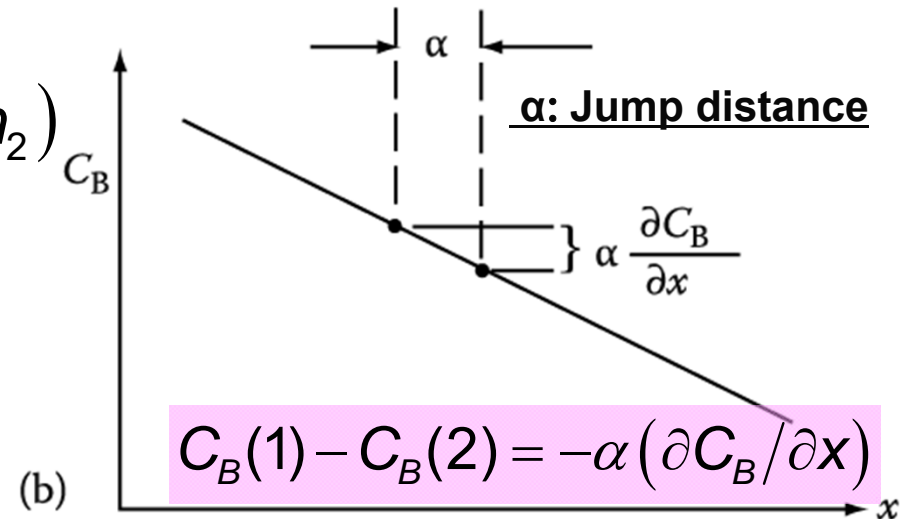
Fig. 2.5 Interstitial diffusion by random jumps in a concentration gradient.

Fick's First Law of Diffusion

$$J_B = \frac{1}{6} \Gamma_B n_1 - \frac{1}{6} \Gamma_B n_2 = \frac{1}{6} \Gamma_B (n_1 - n_2)$$

$$C_B(1) = n_1/\alpha, \quad C_B(2) = n_2/\alpha$$

$$(n_1 - n_2) = \alpha (C_B(1) - C_B(2))$$



$$J_B = \frac{1}{6} \Gamma_B (n_1 - n_2) = - \left(\frac{1}{6} \Gamma_B \alpha^2 \right) \frac{\partial C_B}{\partial x} = -D_B \frac{\partial C_B}{\partial x}$$

➔ Concentration varies with position. (atoms $\text{m}^{-2} \text{s}^{-1}$)

D_B : Intrinsic diffusivity or
Diffusion coefficient of B
⇒ depends on microstructure
of materials

Magnitude of D in various media

Gas : $D \approx 10^{-1} \text{ cm}^2/\text{s}$

Liquid : $D \approx 10^{-4} \sim 10^{-5} \text{ cm}^2/\text{s}$

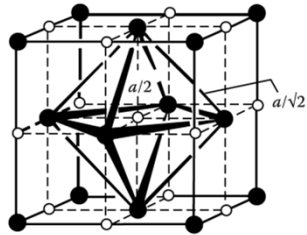
Solid : Materials near melting temp. $D \approx 10^{-8} \text{ cm}^2/\text{s}$

Elemental semiconductor (Si, Ge) $D \approx 10^{-12} \text{ cm}^2/\text{s}$

* Estimate the jump frequency of a carbon atom in γ -Fe (FCC) at 1000 °C.

$$D_B = \frac{1}{6} \Gamma_B \alpha^2$$

lattice parameter of γ -Fe : ~ 0.37 nm



$$\alpha = 0.37 / \sqrt{2} = 0.26 \text{ nm}$$

$$D_C = 2.5 \times 10^{-11} \text{ m}^2 \text{ s}^{-1}$$

$$\Gamma = 2 \times 10^9 \text{ jumps s}^{-1}$$

the vibration frequency of carbon : $\sim 10^{13}$

Only about one attempt in 10^4 results in a jump from one site to another.

* If the crystal structure is not cubic, then the probability of jump is anisotropic (different α). For example the probability of jumps along the basal direction and the axial direction of hcp crystal are different.

D_B , in general, is concentration-dependent.

For example, in γ -Fe at 1000 °C,

$$D_C = 2.5 \times 10^{-11} \text{ m}^2 \text{ s}^{-1}, \text{ when } C_C = 0.15 \text{ wt\%}$$

$$D_C = 7.7 \times 10^{-11} \text{ m}^2 \text{ s}^{-1}, \text{ when } C_C = 1.4 \text{ wt\%}$$

$$C_C \uparrow \rightarrow D_C \uparrow$$

∴ C atoms stain the Fe lattice thereby making diffusion easier as amount of strain increases.

Q. What is the radial distance, r from the origin in random work?

$$r = \alpha \sqrt{(\Gamma t)} = f(D) = 2.4 \sqrt{Dt}$$

For random walk in 3 dimensions,

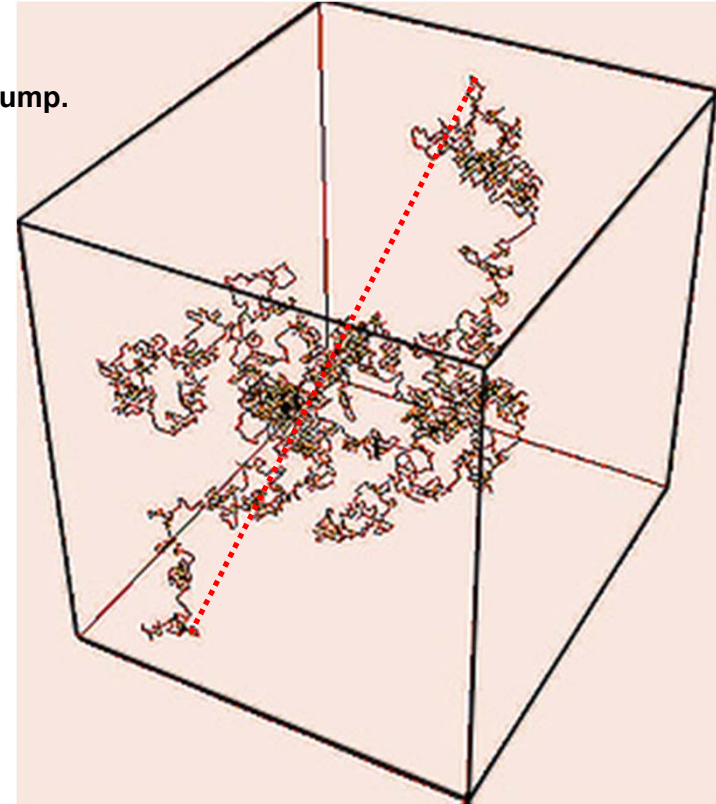
The direction of each new jump is independent of the direction of the previous jump.

after n steps of length α

$$\rightarrow \alpha\sqrt{n}$$

Net distance from its original position

During random walk,
an atom will move in time (t)
a radial distance (r) from the origin



$$r = \alpha\sqrt{(\Gamma t)} = f(D) = 2.4\sqrt{Dt} \quad \because \Gamma = \frac{6D}{\alpha^2} \rightarrow r = \sqrt{6}\sqrt{Dt}$$

$$D_B = \frac{1}{6}\Gamma_B\alpha^2$$

In 1 s, each carbon atom will move a total distance of **~ 0.5 m**
→ a net displacement : **~ 10 μm .**

⇒ **Very few of the atom jumps provide a useful contribution to the total diffusion distance.**

Q. What is the effect of the temperature on diffusivity?

$$D_B = D_{B0} \exp \frac{-Q_{ID}}{RT} \quad \text{(Arrhenius-type equation)}$$

EFFECT OF TEMPERATURE on Diffusivity

Thermal Activation

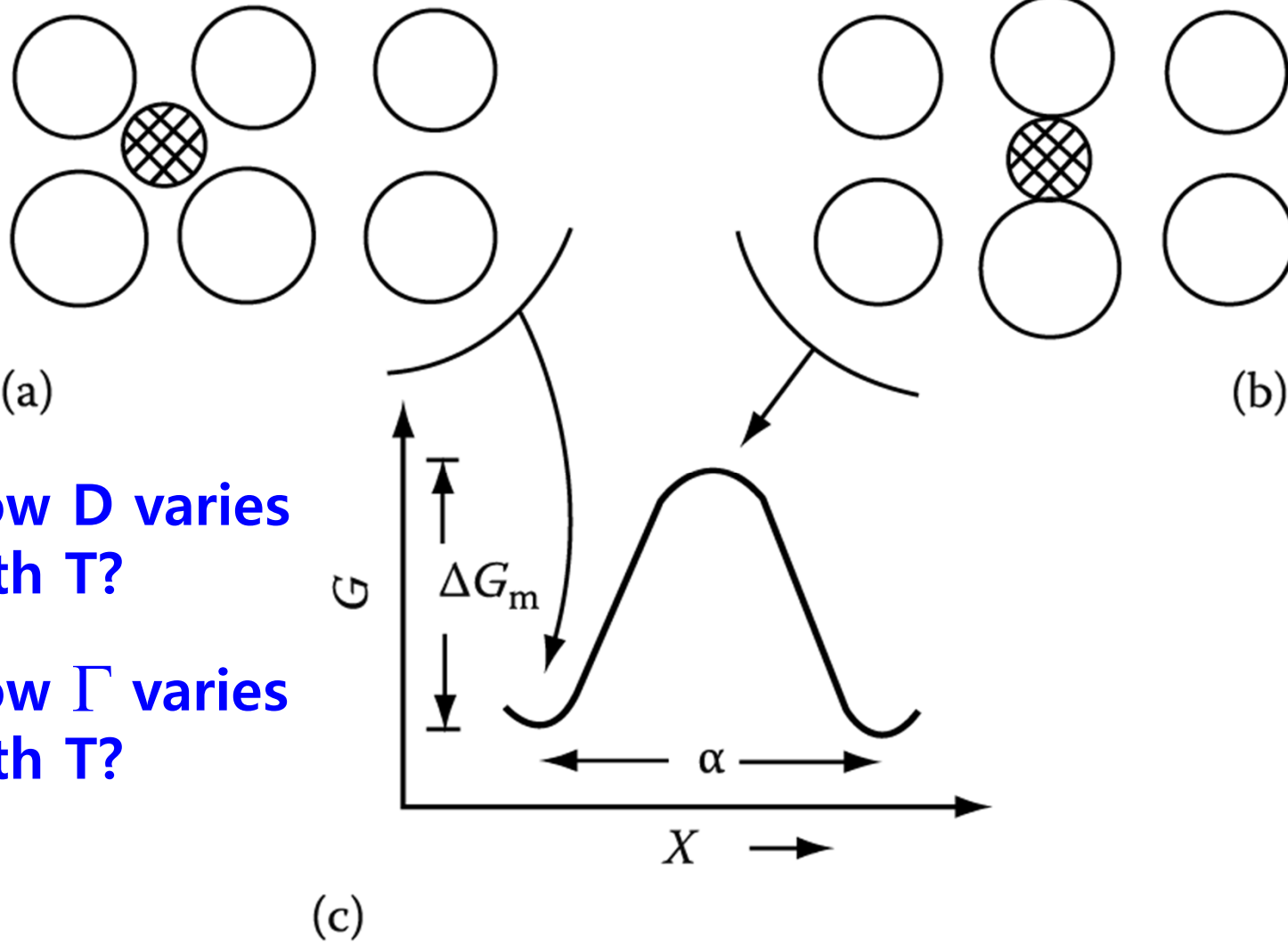


Fig. 2.6 Interstitial atom, (a) in equilibrium position, (b) at the position of maximum lattice distortion. (c) Variation of the free energy of the lattice as a function of the position of interstitial.

Thermally activated process *jump frequency* Γ_B ?

$$\Gamma_B = Z \nu \exp(-\Delta G_m / RT)$$

Z : nearest neighbor sites

ν : vibration frequency

ΔG_m : activation energy for moving

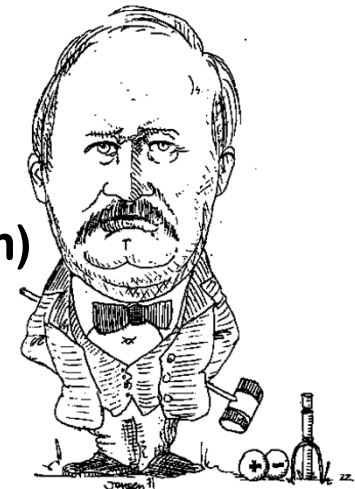
$$\left\{ \Delta G_m = \Delta H_m - T \Delta S_m, D_B = \frac{1}{6} \Gamma_B \alpha^2 \right\}$$

$$D_B = \left[\frac{1}{6} \alpha^2 Z \nu \exp(\Delta S_m / R) \right] \exp(-\Delta H_m / RT)$$

$$\Delta H_m \equiv Q_{ID}$$

$$D_B = D_{B0} \exp \frac{-Q_{ID}}{RT}$$

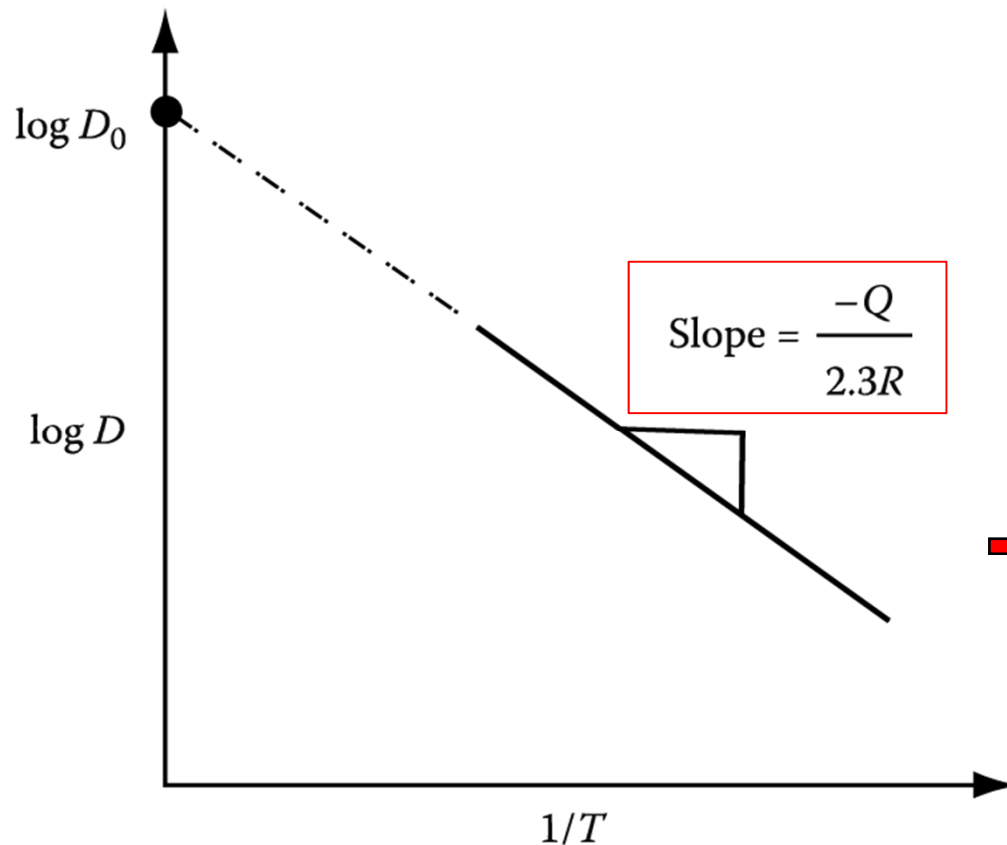
(Arrhenius-type equation)



Temperature Dependence of Diffusion

$$D_B = D_{B0} \exp \frac{-Q_{ID}}{RT}$$

How to determine Q_{ID} experimentally?



$$\log D = \log D_0 - \frac{Q}{2.3R} \left(\frac{1}{T} \right)$$

➔ Therefore, from the slope of the D -curve in an $\log D$ vs $1/T$ coordinate, the activation energy may be found.

Fig. 2.7 The slope of $\log D$ v. $1/T$ gives the activation energy for diffusion Q .

Q. Steady state diffusion vs Non-steady state diffusion?

Fick's first Law

$$J_B = -D_B \frac{\partial C_B}{\partial x}$$

➔ Concentration varies with
“position”

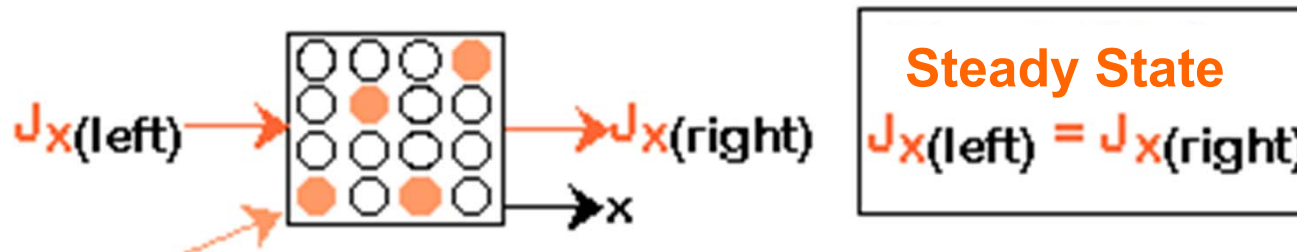
Fick's Second Law

$$\frac{\partial C_B}{\partial t} = D_B \frac{\partial^2 C_B}{\partial x^2}$$

“Both position and time”

Steady-state diffusion

The simplest type of diffusion to deal with is when the concentration at every point does not change with time.



Concentration, C , in the box

- Apply Fick's First Law:

$$J_x = -D \frac{dC}{dx}$$

If $J_x(\text{left}) = J_x(\text{right})$, then

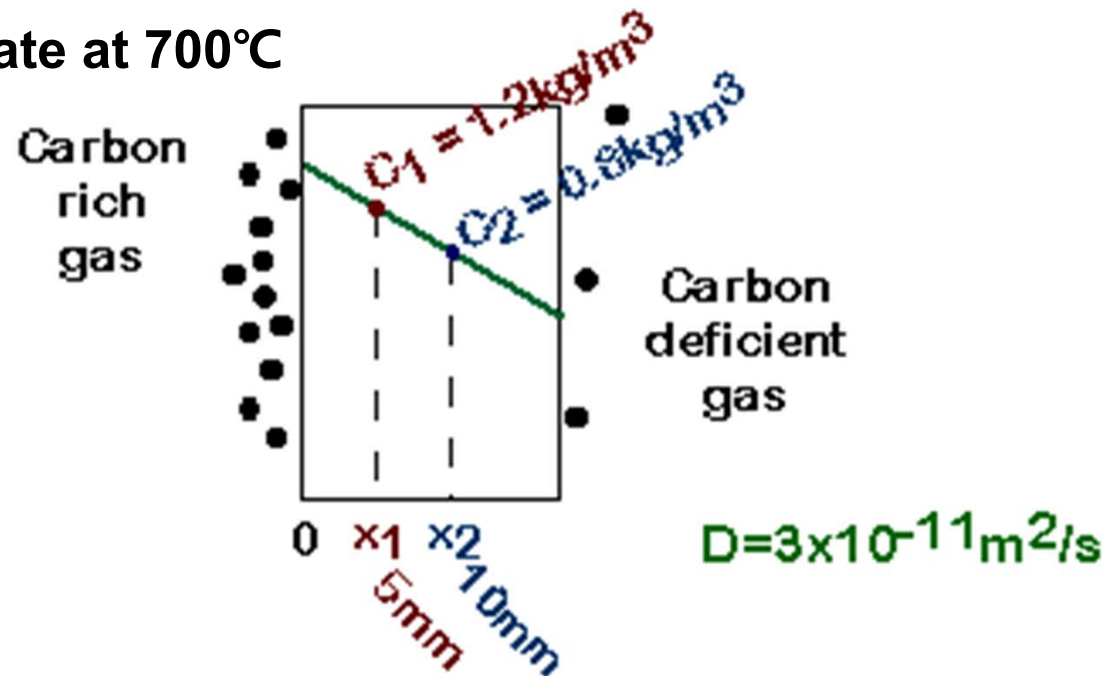
$$\frac{dC}{dx_{(\text{left})}} = \frac{dC}{dx_{(\text{right})}}$$

⇒ the slope, dC/dx , is constant
(does not vary with position)!

Steady-state diffusion

The simplest type of diffusion to deal with is when the concentration at every point does not change with time.

- Steel plate at 700°C



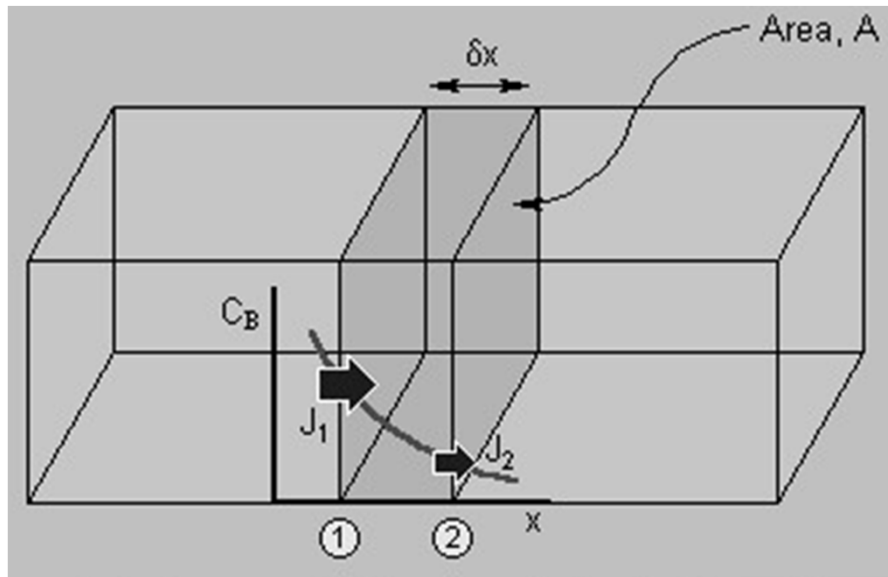
Q: How much carbon is transferring from the rich to deficient side?

$$J = -D \frac{C_2 - C_1}{x_2 - x_1} = 2.4 \times 10^{-9} \frac{\text{kg}}{\text{m}^2 \text{s}}$$

Nonsteady-state diffusion

In most practical situations steady-state conditions are not established, i.e. concentration varies with both distance and time, and Fick's 1st law can no longer be used.

How do we know the variation of C_B with time? → Fick's 2nd law



The number of interstitial B atoms that diffuse into the slice across plane (1) in a small time interval dt :

$$\rightarrow J_1 A dt \quad \text{Likewise : } J_2 A dt$$

Since $J_2 < J_1$, the concentration of B within the slice will have increased by

Due to mass conservation

$$(J_1 - J_2) A \delta t = \delta C_B A \delta x \quad \delta C_B = \frac{(J_1 - J_2) A \delta t}{A \delta x}$$

Nonsteady-state diffusion

$$J_2 = J_1 + \frac{\partial J}{\partial x} \delta x \quad (\delta x \text{ is small})$$

as $\delta t \rightarrow 0$

$$\delta C_B = \frac{(J_1 - J_2) A \delta t}{A \delta x}$$

$$\frac{\partial C_B}{\partial t} = - \frac{\partial J_B}{\partial x}$$

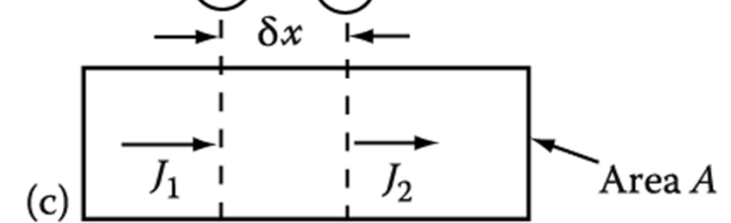
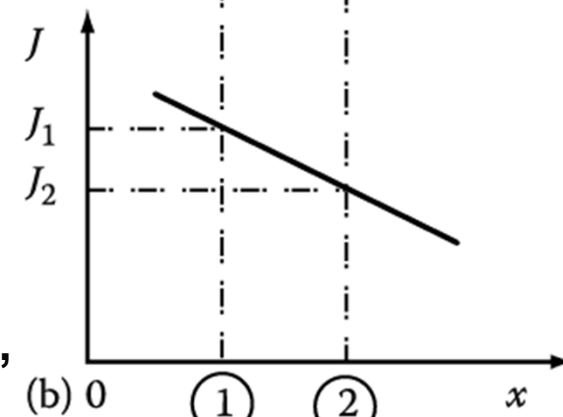
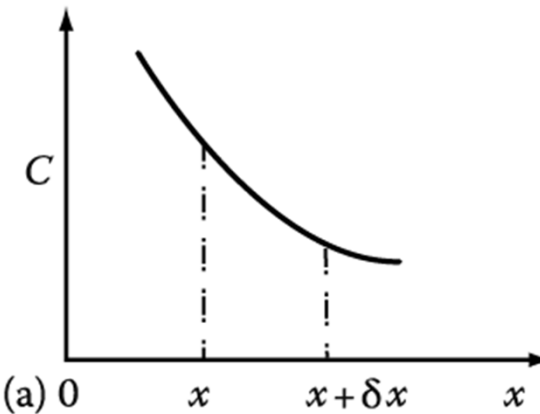
substituting Fick's 1st law gives

$$J_B = -D \frac{\partial C_B}{\partial x} \quad \frac{\partial C_B}{\partial t} = \frac{\partial}{\partial x} \left(D_B \frac{\partial C_B}{\partial x} \right)$$

If variation of D with concentration can be ignored,

$$\frac{\partial C_B}{\partial t} = D_B \frac{\partial^2 C_B}{\partial x^2}$$

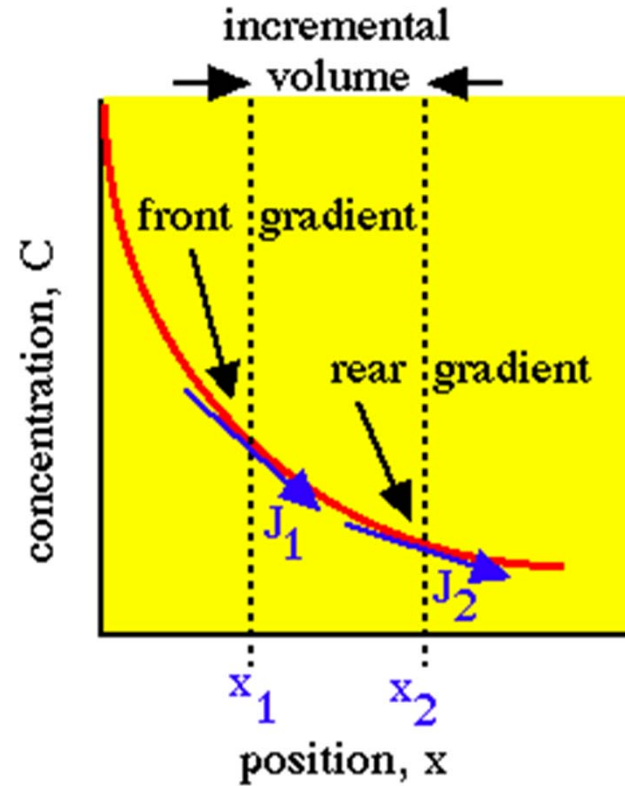
Fick's Second Law



Fick's Second Law

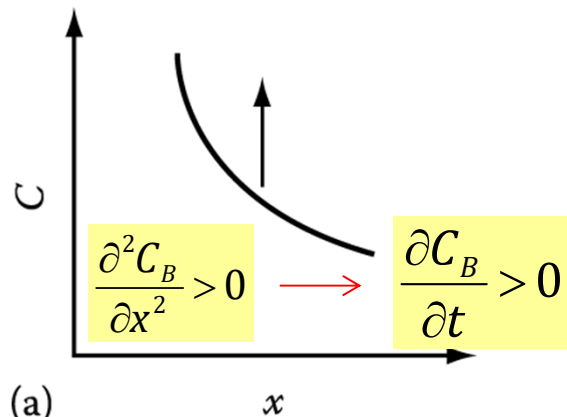
$$\frac{\partial C_B}{\partial t} = D_B \frac{\partial^2 C_B}{\partial x^2}$$

➔ Concentration varies with time and position.

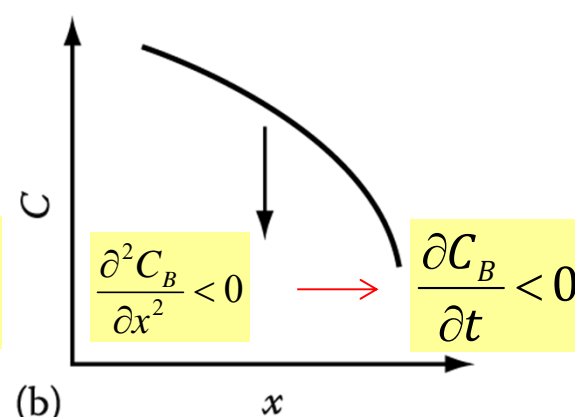


$$\frac{\partial^2 C_B}{\partial x^2} = ?$$

Note that $\frac{\partial^2 C_B}{\partial x^2}$ is the curvature of the C_B versus x curve.



(a) All concentration increase with time



(b) All concentration decrease with time

Q. How to solve the diffusion equations?

: Application of Fick's 2nd law

homogenization, carburization, decarburization, diffusion across a couple

Solutions to the diffusion equations (Application of Fick's 2nd law)

$$\frac{\partial C_B}{\partial t} = D_B \frac{\partial^2 C_B}{\partial x^2}$$

Ex1. Homogenization

of sinusoidal varying composition

in the elimination of segregation in casting

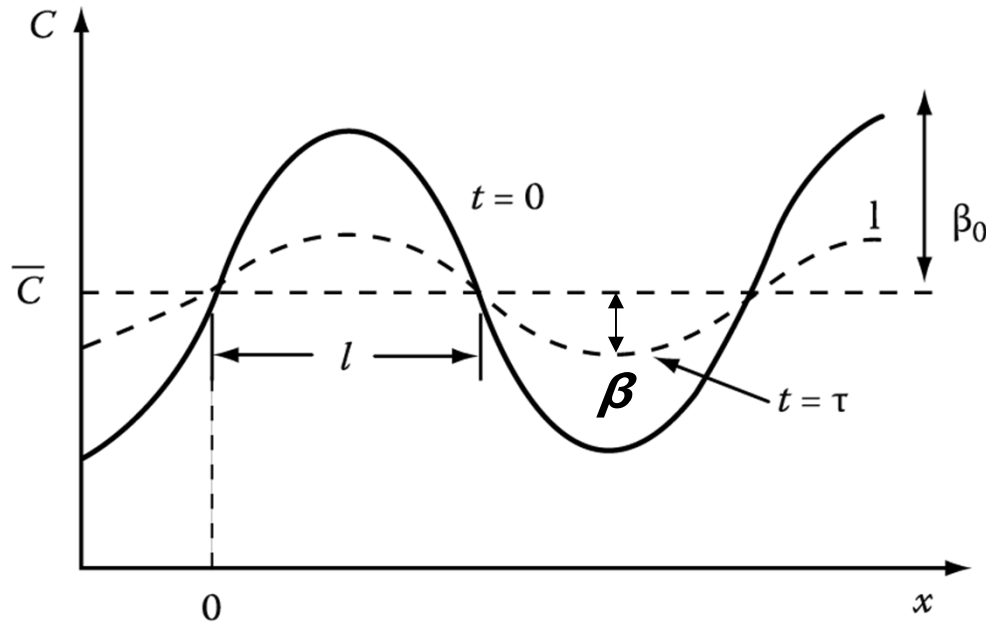


Fig. 2.10 The effect of diffusion on a sinusoidal variation of composition.

\bar{C} : the mean composition l : half wavelength

β_0 : the amplitude of the initial concentration profile

Initial or Boundary Cond.?

$$C = \bar{C} + \beta_0 \sin \frac{\pi x}{l} \quad \text{at } t=0$$

Rigorous solution of $\frac{\partial C}{\partial t} = D \frac{\partial^2 C}{\partial x^2}$

Using a method of variable separation

Let $C = XT \rightarrow \frac{\partial C}{\partial t} = D \frac{\partial^2 C}{\partial x^2}$

$$X \frac{dT}{dt} = DT \frac{d^2 X}{dx^2}$$

$$\frac{1}{DT} \frac{dT}{dt} = \frac{1}{X} \frac{d^2 X}{dx^2} = -\lambda^2$$

$\frac{1}{T} \frac{dT}{dt} = -\lambda^2 D$ $\frac{d \ln T}{dt} = -\lambda^2 D$ $T = T_0 e^{-\lambda^2 D t}$	$\frac{d^2 X}{dx^2} + \lambda^2 X = 0$ $X = A' \cos \lambda x + B' \sin \lambda x$ <div style="border: 1px dashed red; padding: 5px; margin-top: 10px;"> $X(x,0) \equiv \bar{C} + \beta_0 \sin \frac{\pi x}{l}$ </div>
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for $C(x,0) = \bar{C} + \beta_0 \sin \frac{\pi x}{l}$

$$\therefore C = (A \cos \lambda x + B \sin \lambda x) e^{-\lambda^2 D t}$$

$$\therefore C = A_0 + \sum_{n=1}^{\infty} (A_n \sin \lambda_n x + B_n \cos \lambda_n x)$$

$$t = 0 \rightarrow C \equiv \bar{C} + \beta_0 \sin \frac{\pi x}{l} \quad \lambda = \frac{\pi}{l}$$

$$; A_0 = \bar{C}, B_n = 0, A_1 = \beta_0$$

($A_n = 0$ for all others)

$$\therefore C \equiv \bar{C} + \beta_0 \sin \frac{\pi x}{l} e^{-\frac{t}{l^2 / \pi^2 D}}$$

$$C = \bar{C} + \beta_0 \sin \frac{\pi x}{l} \exp\left(\frac{-t}{\tau}\right)$$

$$\tau = \frac{l^2}{\pi^2 D} \quad \tau : \text{relaxation time}$$

$$\beta = \beta_0 \exp(-t / \tau) \quad \text{at } x = \frac{l}{2}$$

Solutions to the diffusion equations

Ex1. Homogenization of sinusoidal varying composition in the **elimination of segregation in casting**

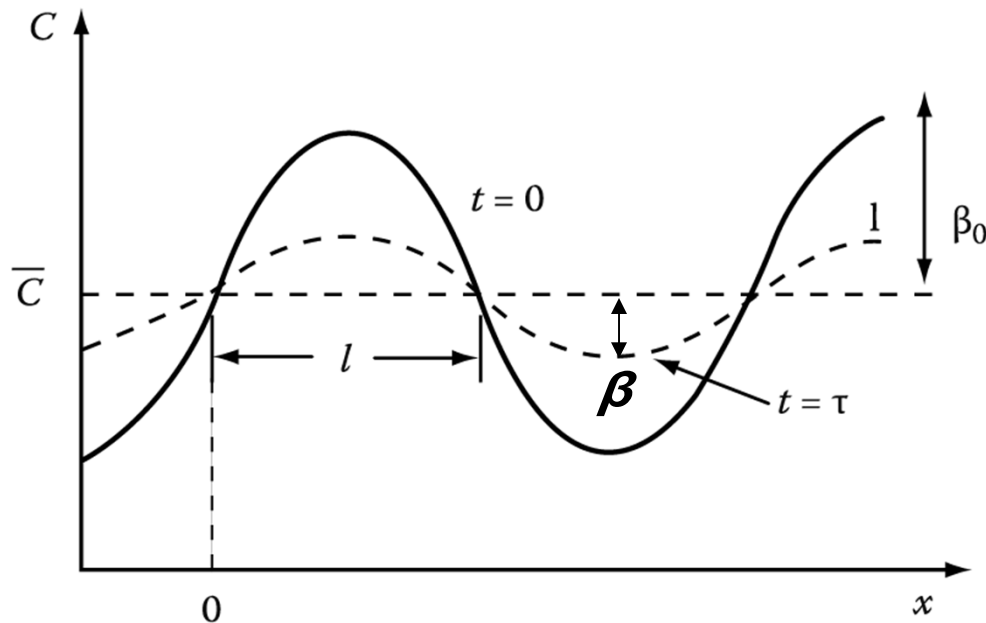


Fig. 2.10 The effect of diffusion on a sinusoidal variation of composition.

$$C = \bar{C} + \beta_0 \sin \frac{\pi x}{l} \quad \text{at } t=0$$

$$C = \bar{C} + \beta_0 \sin \frac{\pi x}{l} \exp\left(\frac{-t}{\tau}\right)$$

$$\beta = \beta_0 \exp(-t / \tau) \quad \text{at } x = \frac{l}{2}$$

Amplitude of the concentration profile (β) decreases exponentially with time, $C \rightarrow \bar{C}$.

$$\tau = \frac{l^2}{\pi^2 D} \quad \tau : \text{relaxation time}$$

decide homogenization rate

The initial concentration profile will not usually be sinusoidal, but in general any concentration profile can be considered as the sum of an infinite series of sine waves of varying wavelength and amplitude, and each wave decays at a rate determined by its own " τ ". Thus, the short wavelength terms die away very rapidly and **the homogenization will ultimately be determined by τ for the longest wavelength component.**

Solutions to the diffusion equations

Ex2. **Carburization** of Steel

The aim of carburization is **to increase the carbon concentration in the surface layers of a steel product in order to achieve a harder wear-resistant surface.**

1. **Holding the steel in CH_4 and/or CO at an austenitic temperature.**
2. **By controlling gases the concentration of C at the surface of the steel can be maintained at a suitable constant value.**
3. **At the same time carbon continually diffuses from the surface into the steel.**



Carburizing of steel

Depth of Carburization?

- 1) Specimen ~ infinitely long.
- 2) Diffusion coefficient of carbon in austenite increases with increasing concentration, but an approximate solution: Taking an average value

The error function solution:

$$C = C_s - (C_s - C_0) \operatorname{erf}\left(\frac{x}{2\sqrt{Dt}}\right) + \frac{1}{2}(C_s + C_0)$$

$$\operatorname{erf}(Z) = \frac{2}{\sqrt{\pi}} \int_0^Z e^{-y^2} dy$$

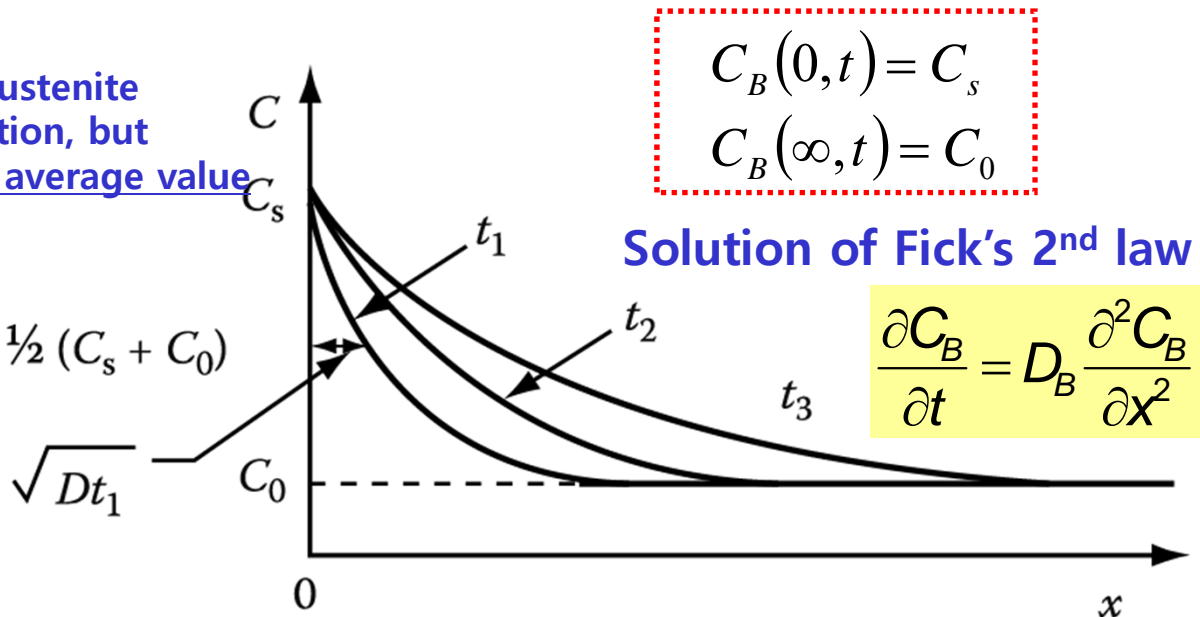


Fig. 2.11 Concentration profiles at successive times ($t_3 > t_2 > t_1$) for diffusion into a semi-infinite bar when the surface concentration C_s is maintained constant.

$$\frac{C_s - C}{C_s - C_0} = \operatorname{erf}\left(\frac{x}{2\sqrt{Dt}}\right) \quad \operatorname{erf}(0.5) \approx 0.5 \quad C = \frac{C_s + C_0}{2}$$

- Since $\operatorname{erf}(0.5) \approx 0.5$, the depth at which the carbon concentration is midway between C_s and C_0 is given $(x / 2\sqrt{Dt}) \cong 0.5$

that is $x \cong \sqrt{Dt}$ → **Depth of Carburization**

Error function

In mathematics, the error function (also called the Gauss error function) is a non-elementary function which occurs in probability, statistics and partial differential equations.

It is defined as:

$$\operatorname{erf}(x) = \frac{2}{\sqrt{\pi}} \int_0^x e^{-t^2} dt.$$

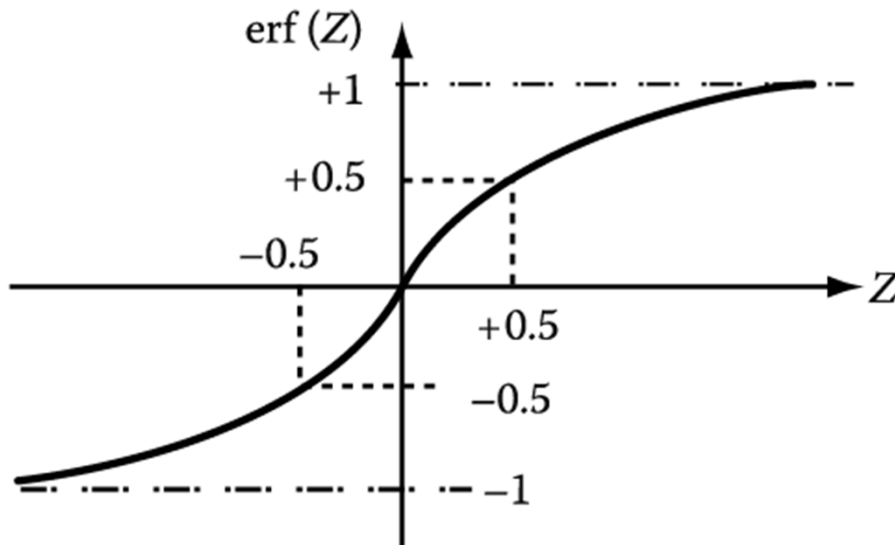
By expanding the right-hand side in a Taylor series and integrating, one can express it in the form

$$\operatorname{erf}(x) = \frac{2}{\sqrt{\pi}} \sum_{n=0}^{\infty} \frac{(-1)^n x^{2n+1}}{(2n+1)n!} = \frac{2}{\sqrt{\pi}} \left(x - \frac{x^3}{3} + \frac{x^5}{10} - \frac{x^7}{42} + \frac{x^9}{216} - \dots \right)$$

for every real number x . (From Wikipedia, the free encyclopedia)

Error function

$$\text{erf}(z) = \frac{2}{\sqrt{\pi}} \int_0^z \exp(-y^2) dy$$



(a) Schematic diagram illustrating the main features of the error function.

Table 1-1. The Error Function

z	$\text{erf}(z)$	z	$\text{erf}(z)$
0	0	0.85	0.7707
0.025	0.0282	0.90	0.7969
0.05	0.0564	0.95	0.8209
0.10	0.1125	1.0	0.8427
0.15	0.1680	1.1	0.8802
0.20	0.2227	1.2	0.9103
0.25	0.2763	1.3	0.9340
0.30	0.3286	1.4	0.9523
0.35	0.3794	1.5	0.9661
0.40	0.4284	1.6	0.9763
0.45	0.4755	1.7	0.9838
0.50	0.5205	1.8	0.9891
0.55	0.5633	1.9	0.9928
0.60	0.6039	2.0	0.9953
0.65	0.6420	2.2	0.9981
0.70	0.6778	2.4	0.9993
0.75	0.7112	2.6	0.9998
0.80	0.7421	2.8	0.9999

Carburizing of steel

Depth of Carburization?

- 1) Specimen ~ infinitely long.
- 2) Diffusion coefficient of carbon in austenite increases with increasing concentration, but an approximate solution: Taking an average value

* Concentration profile : using boundary conditions

$$C_B(0, t) = C_s$$

$$C_B(\infty, t) = C_0$$

→ Solution of Fick's 2nd law

The error function solution:

$$C = C_s - (C_s - C_0) \operatorname{erf}\left(\frac{x}{2\sqrt{Dt}}\right)$$

$$\operatorname{erf}(Z) = \frac{2}{\sqrt{\pi}} \int_0^Z e^{-y^2} dy$$

$$\frac{1}{2} (C_s + C_0)$$

$$\sqrt{Dt_1}$$

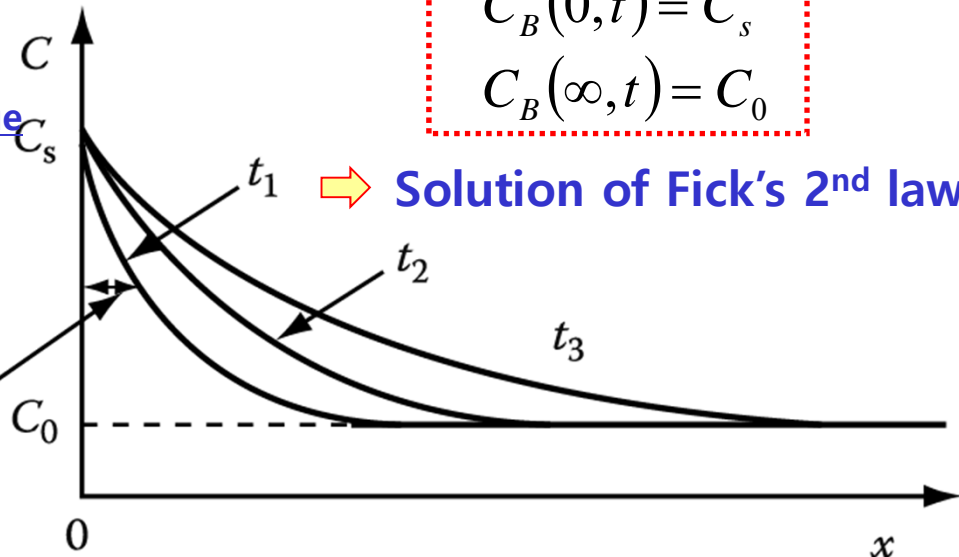


Fig. 2.11 Concentration profiles at successive times ($t_3 > t_2 > t_1$) for diffusion into a semi-infinite bar when the surface concentration C_s is maintained constant.

$$\frac{C_s - C}{C_s - C_0} = \operatorname{erf}\left(\frac{x}{2\sqrt{Dt}}\right)$$

$$\operatorname{erf}(0.5) \approx 0.5$$

$$C = \frac{C_s + C_0}{2}$$

- Since $\operatorname{erf}(0.5) \approx 0.5$, the depth at which the carbon concentration is midway between C_s and C_0 is given $(x / 2\sqrt{Dt}) \cong 0.5$

that is $x \cong \sqrt{Dt}$ → Depth of Carburization

Carburizing of steel

Thus the thickness of the carburized layer is $\cong \sqrt{Dt}$.

Note also that the depth of any is concentration line is directly proportion to \sqrt{Dt} , i.e. *to obtain a twofold increase in penetration requires a fourfold increase in time.*

(2배의 침투 깊이 → 4배의 시간)

Ex.3 Decarburization of Steel?

$$C = C_0 \operatorname{erf}\left(\frac{x}{2\sqrt{Dt}}\right)$$

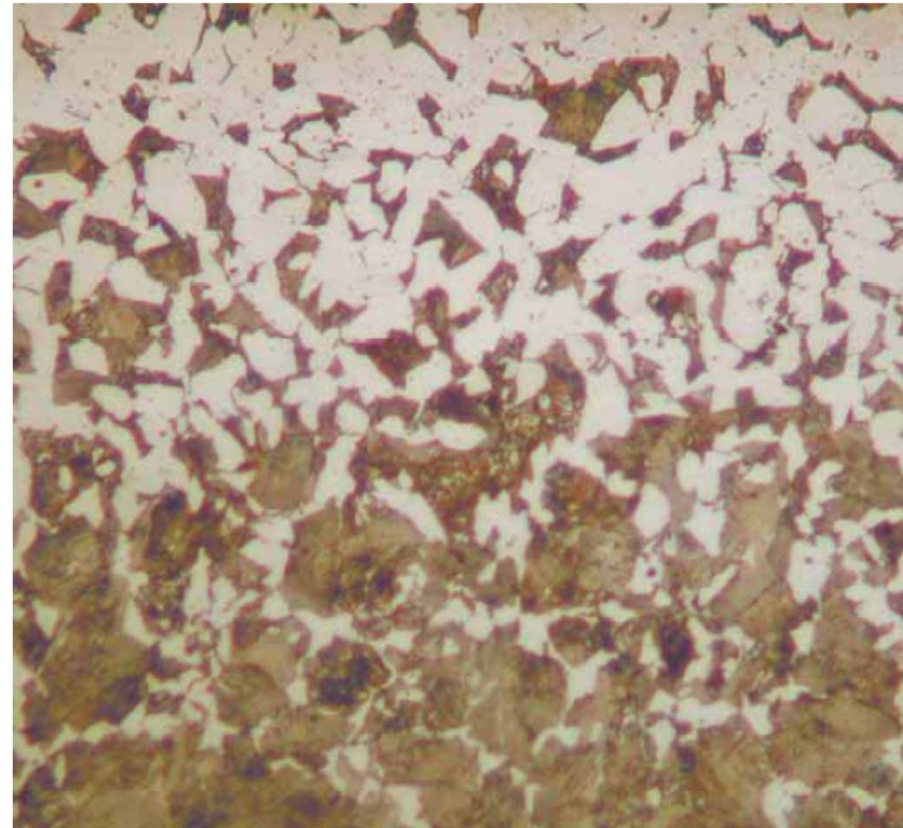
$$C = C_s - (C_s - C_0) \operatorname{erf}\left(\frac{x}{2\sqrt{Dt}}\right)$$

Carburization

$$\operatorname{erf}(Z) = \frac{2}{\sqrt{\pi}} \int_0^Z e^{-y^2} dy$$

C_s = Surface concentration

C_0 = Initial bulk concentration



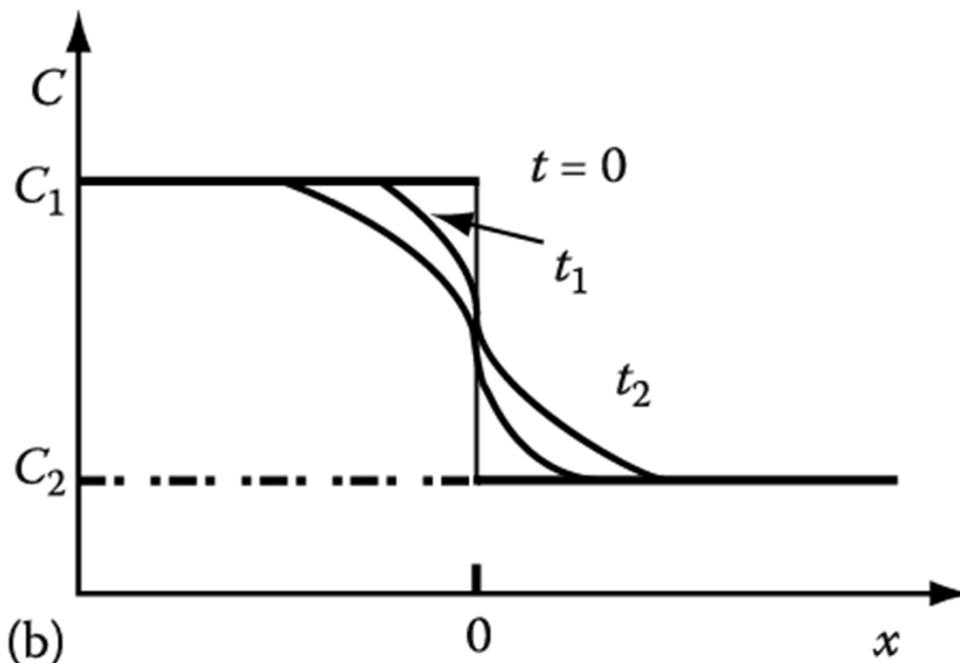
Solutions to the diffusion equations

Ex4. Diffusion Couple

Joining of two semi-infinite specimens of compositions C_1 and C_2 ($C_1 > C_2$)

Draw C vs. x with time $t = 0$ and $t > 0$.

Boundary conditions?



$$C = C_1 \quad t = 0, \quad -\infty < x < 0$$

$$C = C_2 \quad t = 0, \quad 0 < x < \infty$$

$$C = C_1 \quad x = -\infty, \quad 0 < t < \infty$$

$$C = C_2 \quad x = \infty, \quad 0 < t < \infty$$

$$C = \left(\frac{C_1 + C_2}{2} \right) - \left(\frac{C_1 - C_2}{2} \right) \operatorname{erf} \left(\frac{x}{2\sqrt{Dt}} \right)$$

Fig. 2.12 (b) concentration profiles at successive times ($t_2 > t_1 > 0$) when two semi-infinite bars of different composition are annealed after welding.

The section is completed with **4 example solutions to Fick's 2nd law:** *carburisation, decarburisation, diffusion across a couple and homogenisation.*

The solutions given are as follows:

Process	Solution
Homogenization	$C = C_{\text{mean}} + \beta_0 \sin\left(\frac{\pi x}{l}\right) \exp\left(-\frac{t}{\tau}\right)$ <p> C_{mean} = Mean concentration β_0 = Initial concentration amplitude l = half-wavelength of cells t = relaxation time </p>
Carburization	$C = C_s - (C_s - C_0) \operatorname{erf}\left(\frac{x}{2\sqrt{Dt}}\right)$ <p> C_s = Surface concentration C_0 = Initial bulk concentration </p>
Decarburization	$C = C_0 \operatorname{erf}\left(\frac{x}{2\sqrt{Dt}}\right)$ <p> C_0 = Initial bulk concentration </p>
Diffusion Couple	$C = \left(\frac{C_1 + C_2}{2}\right) - \left(\frac{C_1 - C_2}{2}\right) \operatorname{erf}\left(\frac{x}{2\sqrt{Dt}}\right)$ <p> C_1 = Concentration of steel 1 C_2 = Concentration of steel 2 </p>

Contents for today's class

- **Diffusion** Movement of atoms to reduce its chemical potential μ .

→ driving force: Reduction of G

Down-hill diffusion movement of atoms from a high C_B region to low C_B region.

Up-hill diffusion movement of atoms from a low C_B region to high C_B region.

- **Interstitial Diffusion / Substitution Diffusion**

- **Steady-state diffusion– Fick's First Law**

$$J_B = \frac{1}{6} \Gamma_B (n_1 - n_2) = - \left(\frac{1}{6} \Gamma_B \alpha^2 \right) \frac{\partial C_B}{\partial x} = -D_B \frac{\partial C_B}{\partial x}$$

→ Concentration varies with position. (atoms $m^{-2} s^{-1}$)

- **Nonsteady-state diffusion – Fick's Second Law**

$$\frac{\partial C_B}{\partial t} = D_B \frac{\partial^2 C_B}{\partial x^2} \rightarrow \text{Concentration varies with time and position.}$$

- **For random walk in 3 dimensions,**
after n steps of length α

$$r = \alpha \sqrt{(\Gamma t)} = f(D) = 2.4 \sqrt{Dt}$$

- **Effect of Temperature on Diffusivity**

$$\log D = \log D_0 - \frac{Q}{2.3R} \left(\frac{1}{T} \right)$$