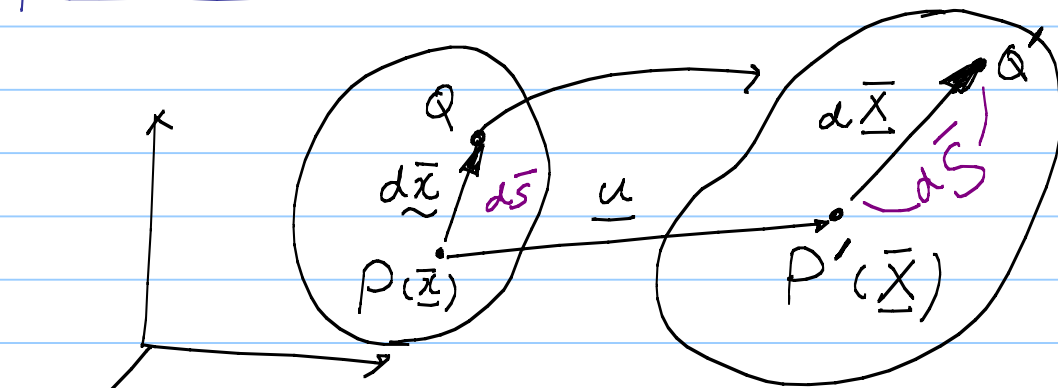


Lecture 1-3 : Deformation



P : point in a space before deformation
 (\bar{x} : coordinate "values" for P)

P' : point in a space after deformation
 (\bar{X} : " " " for P')

< length change & strain >

$$d\bar{S}^2 - d\bar{s}^2 \triangleq 2 E_{jk} d\bar{x}_j d\bar{x}_k = 2 d\bar{x} \cdot \underline{E} \cdot d\bar{x}$$

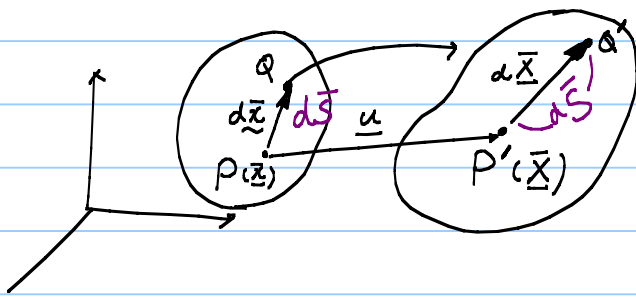
↑ Lagrangian strain
 "relative to undeformed coordinates"

$$\triangleq 2 E_{jk}^* d\bar{X}_j d\bar{X}_k = 2 d\bar{X} \cdot \underline{E}^* \cdot d\bar{X}$$

↑ Eulerian strain
 "relative to deformed coord."

< Analysis of $\underline{\underline{E}}$ >

2



Note

$$\begin{cases} d\bar{s}^2 = d\bar{x} \cdot d\bar{x} \\ d\bar{S}^2 = d\bar{X} \cdot d\bar{X} \end{cases}$$

with $\underline{X} = \underline{x} + \underline{u}$ ← displacement

$$d\bar{S}^2 = d\bar{X} \cdot d\bar{X} = d\bar{X}_i d\bar{X}_i$$

$$= \left(\frac{\partial X_i}{\partial x_j} d\bar{x}_j \right) \left(\frac{\partial X_i}{\partial x_k} d\bar{x}_k \right)$$

$$= \frac{\partial}{\partial x_j} (x_i + u_i) \frac{\partial}{\partial x_k} (x_i + u_i) d\bar{x}_j d\bar{x}_k$$

$$= (\delta_{ij} + u_{i,j}) (\delta_{ik} + u_{i,k}) d\bar{x}_j d\bar{x}_k$$

$$= \left(\underbrace{\delta_{ij} \delta_{ik}}_{\delta_{jk}} + \underbrace{\delta_{ij} u_{i,k}}_{u_{j,k}} + \underbrace{u_{i,j} \delta_{ik}}_{u_{k,j}} + u_{i,j} u_{i,k} \right) d\bar{x}_j d\bar{x}_k^3$$

$$= \underbrace{d\bar{x}_j \cdot d\bar{x}_j}_{d\bar{s}^2} + \underbrace{\left(u_{j,k} + u_{k,j} + u_{i,j} u_{i,k} \right)}_{\triangleq 2 E_{jk}} d\bar{x}_j \cdot d\bar{x}_k$$

From,

$$\begin{aligned} d\bar{S}^2 - d\bar{s}^2 &= 2 E_{jk} d\bar{x}_j \cdot d\bar{x}_k \\ &= 2 \underline{d\bar{x}} \cdot \underline{E} \cdot \underline{d\bar{x}} \end{aligned}$$

We can find

$$\underline{E} = E_{ij} \underline{e}_i \otimes \underline{e}_j$$

$$E_{ij} = (u_{i,j} + u_{j,i} + u_{i,k} u_{j,k}) / 2$$

\Rightarrow Lagrangian strain tensor
Defined in "undeformed"
coordinate system

Remark:

Eulerian strain tensor: E_{ij}^*

$$d\bar{S}^2 - d\bar{s}^2 = 2 d\bar{X} \cdot \underline{E}^* \cdot d\bar{X}$$

→ defined in deformed coord. sys.

$$E_{ij}^* = \frac{1}{2} (u_{i,j} + u_{j,i} - u_{\kappa,i} u_{\kappa,j})$$

$$E_{ij} = \frac{1}{2} (u_{i,j} + u_{j,i} + u_{\kappa,i} u_{\kappa,j})$$

$$(\)_{,i} = \frac{\partial}{\partial x_i}$$

$$(\)_{,i} = \frac{\partial}{\partial x_i}$$

For small deformation;

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$$\frac{\partial}{\partial x_j}(\cdot) \approx \frac{\partial}{\partial X_j}(\cdot) \quad \mathbb{F}_{ij} \approx \mathbb{F}_{ij}^* \Rightarrow \epsilon_{ij}$$

linear strain

$$\epsilon_{ij} = \frac{1}{2} \left(\frac{\partial u_j}{\partial x_i} + \frac{\partial u_i}{\partial x_j} \right)$$

$$\underline{\underline{\epsilon}} = \epsilon_{ij} \underline{e}_i \otimes \underline{e}_j = \frac{1}{2} [\underline{\nabla} \underline{u} + (\underline{\nabla} \underline{u})^T]$$

$$[\underline{\underline{\epsilon}} = \underline{\underline{\epsilon}}^T]$$

Displacement gradient $\underline{\nabla} \underline{u}$

$$\underline{\nabla} \underline{u} = \frac{1}{2} [\underline{\nabla} \underline{u} + (\underline{\nabla} \underline{u})^T] + \frac{1}{2} [\underline{\nabla} \underline{u} - (\underline{\nabla} \underline{u})^T]$$

Symmetric part
($\underline{\nabla} \underline{u}$)_{sym}

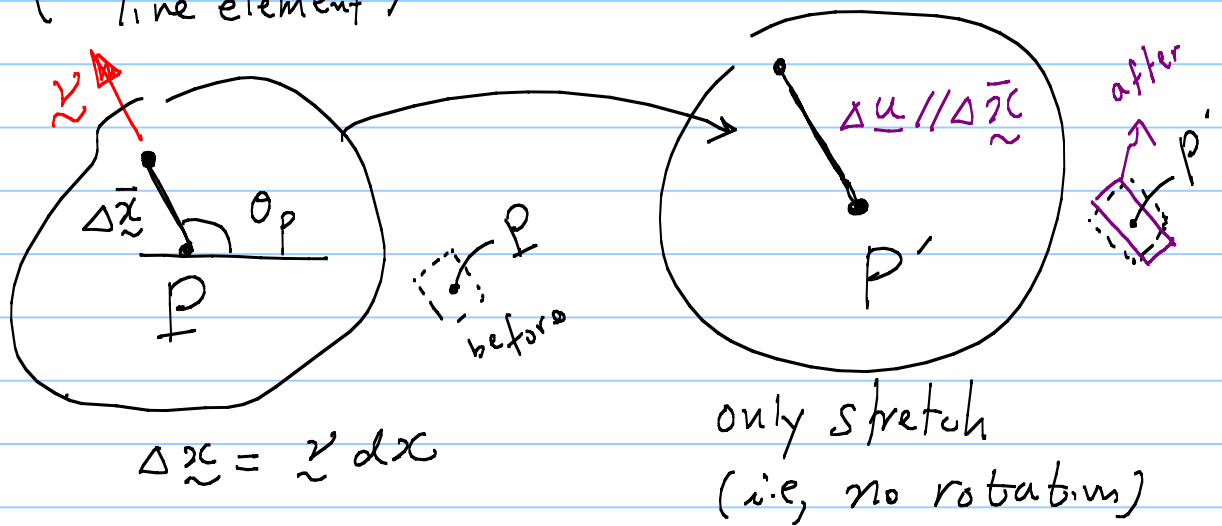
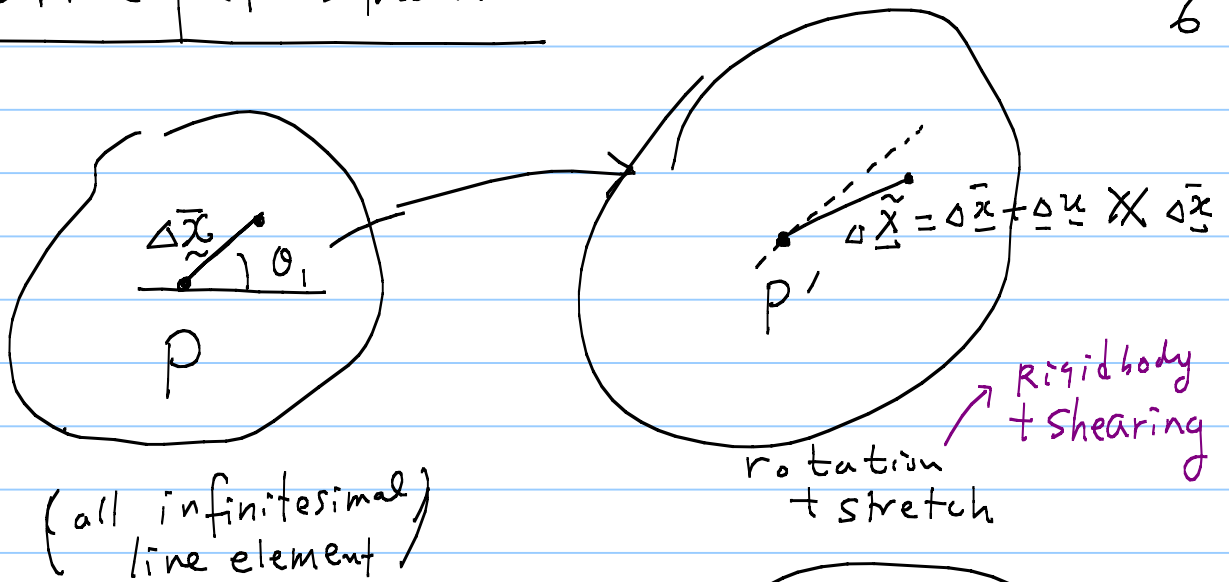
antisymmetric part

Rotation

(see page 49 - 50)

Principal Strain

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For some orientation θ_p , or unit tangent \underline{u} ,

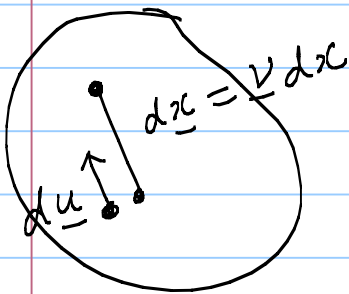
if $d\underline{u} \parallel d\underline{x}$ holds,

$\theta_p \mid \underline{u} \Rightarrow$ principal direction

In this case,
one can set:

$$d\underline{u} \equiv \lambda d\underline{x} \equiv \lambda \underline{v} dx \quad (a)$$

Excluding
Infinitesimal
Rigid-body
Rotation



Along $\underline{v} dx$ direction,

$$d\underline{u} = \underline{v} \cdot \underline{\nabla} \underline{u} dx$$

Recall $\left. \frac{df}{ds} \right|_{\text{along } c} = \hat{\underline{t}} \cdot \underline{\nabla} f, \text{ Eq (1.1.99)}$
 choose: $f \Rightarrow \underline{u}$; $\hat{\underline{t}} \Rightarrow \underline{v}$
 $ds = dx$

Since $\underline{\nabla} \underline{u} = \underline{\underline{\epsilon}} + \underline{\underline{\omega}}$
 sym part \uparrow anti-sym
 Infinitesimal
Rigid-Body Rotation

$$d\underline{u} = \underline{v} \cdot \nabla \underline{u} \, dx$$

$$= (\underline{v} \cdot \underline{\varepsilon} + \underline{v} \cdot \underline{\omega}) \, dx$$

$$\Rightarrow (\underline{v} \cdot \underline{\varepsilon}) \, dx \quad (b)$$

↑ (not concerned with
an infinitesimal rigid-body
rotation)

(a, b) →

$$\underline{v} \cdot \underline{\varepsilon} = \lambda \underline{v}$$

OR

$$\underline{\varepsilon}^T \cdot \underline{v} = \lambda \underline{v} = \lambda \underline{\mathbb{1}} \cdot \underline{v}$$

using

$$\underline{\varepsilon}^T = \underline{\varepsilon}$$

we have

$$\underline{\varepsilon} \cdot \underline{v} = \lambda \underline{\mathbb{1}} \cdot \underline{v} \quad (c1)$$

$$(\underline{\varepsilon} - \lambda \underline{\mathbb{1}}) \cdot \underline{v} = 0 \quad (c2)$$

In component

$$(\varepsilon_{ij} - \lambda \delta_{ij}) v_j = 0 \quad (c3)$$

↑ "Eigenvalue problem"

For non trivial solution of (c);

$$|\epsilon_{ij} - \lambda \delta_{ij}| = 0 \quad (\lambda_1, \lambda_2, \lambda_3) \quad (d)$$

FACT FOR SYMMETRIC TENSORS $\underline{\epsilon}$

i) λ (eigenvalue) = Real

↑ Here, principal strain

ii) $\underline{v} = v_i \underline{e}_i$ (Eigenvector) = Real

↑ Here, principal strain direction

* iii) Three eigenvectors =

Mutually orthonormal

(However, eigenvalues may be not all distinct, i.e., $\lambda_1 \neq \lambda_2 = \lambda_3$)

To expand (d), use Eq. (1.1.13) of book

$$\det A = \begin{vmatrix} A_{11} & A_{12} & A_{13} \\ A_{21} & A_{22} & A_{23} \\ A_{31} & A_{32} & A_{33} \end{vmatrix} = \epsilon_{ijk} A_{1i} A_{2j} A_{3k} \quad (*)$$

Book (1.1.93)

Inserting in (1.1.93):

$$A_{ij} = \epsilon_{ij} - \lambda \delta_{ij}$$

$$0 = |\epsilon_{ij} - \lambda \delta_{ij}| \quad \leftarrow \text{characteristic eq}$$

$$= \epsilon_{ijk} (\epsilon_{1i} - \lambda \delta_{1i}) (\epsilon_{2j} - \lambda \delta_{2j}) (\epsilon_{3k} - \lambda \delta_{3k})$$

$$= -\lambda^3 \epsilon_{ijk} \delta_{1i} \delta_{2j} \delta_{3k} \\ + \lambda^2 \epsilon_{ijk} (\epsilon_{1i} \delta_{2j} \delta_{3k} + \delta_{1i} \epsilon_{2j} \delta_{3k} \\ + \delta_{1i} \delta_{2j} \epsilon_{3k}) \\ - \lambda^1 \epsilon_{ijk} (\epsilon_{1i} \epsilon_{2j} \delta_{3k} + \delta_{1i} \epsilon_{2j} \epsilon_{3k} \\ + \epsilon_{1i} \delta_{2j} \epsilon_{3k})$$

$$+ \lambda^0 (\epsilon_{ijk} \epsilon_{1i} \epsilon_{2j} \epsilon_{3k})$$

//

$$= -\lambda^3 \underbrace{(\epsilon_{123})}_{\rightarrow 1} + \lambda^2 (\epsilon_{i23} \epsilon_{1i} + \epsilon_{ij3} \epsilon_{2j} + \epsilon_{i2k} \epsilon_{3k})$$

$$\underbrace{\epsilon_{11} + \epsilon_{22} + \epsilon_{33}}_{\text{tr } \underline{\epsilon} \triangleq I_1}$$

$$- \lambda (\epsilon_{ij3} \epsilon_{1i} \epsilon_{2j} + \epsilon_{ijk} \epsilon_{2j} \epsilon_{3k} + \epsilon_{i2k} \epsilon_{1i} \epsilon_{3k})$$

$$\Rightarrow [\epsilon_{11} \epsilon_{22} - \epsilon_{12}^2]$$

$$+ [\epsilon_{22} \epsilon_{33} - \epsilon_{23}^2]$$

$$+ [\epsilon_{11} \epsilon_{33} - \epsilon_{13}^2]$$

$$\Rightarrow \frac{1}{2} \left\{ (\epsilon_{11} + \epsilon_{22} + \epsilon_{33})^2 - (\epsilon_{11}^2 + \epsilon_{22}^2 + \epsilon_{33}^2 + 2\epsilon_{12}^2 + 2\epsilon_{23}^2 + 2\epsilon_{31}^2) \right\}$$

$$\Rightarrow \frac{1}{2} \left\{ (\text{tr } \underline{\epsilon})^2 - (\epsilon_{ij} \epsilon_{ij})^2 \right\} \triangleq I_2$$

$$+ \lambda^0 \underbrace{\det \underline{\epsilon}}_{\triangleq I_3}$$

$$= - \left[\lambda^3 - I_1 \lambda^2 + I_2 \lambda - I_3 \right]$$

$$\bullet \begin{cases} I_1 = \text{tr } \underline{\underline{\epsilon}} \quad (= \epsilon_{ii}) \\ I_2 = \frac{1}{2} (I_1^2 - \underline{\underline{\epsilon}} : \underline{\underline{\epsilon}}) = \frac{1}{2} \left[(\epsilon_{ii})^2 - \epsilon_{ij} \epsilon_{ij} \right] \\ I_3 = \det \underline{\underline{\epsilon}} \end{cases}$$

↑↑ three "Invariants" → Indep of coordinates
 same values in any coordinate system

Principal Strain and Direction

$$(\underline{\underline{\epsilon}} - \lambda \underline{\underline{I}}) \cdot \underline{\underline{v}} = 0$$

① need to solve

$$\lambda^3 - I_1 \lambda^2 + I_2 \lambda - I_3 = 0$$

② After finding three λ 's, then find $\underline{\underline{v}}$.

Ex 1.2.2

$$\underline{\underline{\varepsilon}} \Big|_{a+p} = \left[\underline{e}_1 \otimes \underline{e}_1 + \underline{e}_3 \otimes \underline{e}_3 + \sqrt{3} (\underline{e}_1 \otimes \underline{e}_2 + \underline{e}_2 \otimes \underline{e}_1) \right] \times 10^{-3}$$

$$\sim 10^{-3} \begin{bmatrix} 1 & \sqrt{3} & 0 \\ \sqrt{3} & 0 & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

Find the principal strains and directions.

Sol.

$$\underline{\underline{\varepsilon}} \cdot \underline{v} = \lambda \underline{\underline{1}} \cdot \underline{v}$$

① for λ , solve

$$\bullet \quad |\varepsilon_{ij} - \lambda \delta_{ij}| = 0$$

$$\bullet \quad \lambda^3 - I_1 \lambda^2 + I_2 \lambda - I_3 = 0$$

$$I_1 = \text{tr} \underline{\underline{\varepsilon}} = 2 \times 10^{-3}$$

$$I_2 = (I_1^2 - \underline{\underline{\varepsilon}} : \underline{\underline{\varepsilon}}) / 2$$

$$\left(\begin{aligned} \underline{\underline{\varepsilon}} : \underline{\underline{\varepsilon}} &= \text{tr} (\underline{\underline{\varepsilon}}^T \underline{\underline{\varepsilon}}) \\ &= \text{tr} \begin{bmatrix} 4 & \sqrt{3} & 0 \\ \sqrt{3} & 3 & 0 \\ 0 & 0 & 1 \end{bmatrix} \times 10^{-6} \\ &= 8 \times 10^{-6} \end{aligned} \right)$$

$$= -2 \times 10^{-6}$$

$$I_3 = \det \underline{\underline{\varepsilon}} = -3 \times 10^{-9}$$

Solving

$$\lambda^3 - 2 \times 10^{-3} \lambda^2 - 2 \times 10^{-6} \lambda + 3 \times 10^{-9} = 0$$

$$\lambda = \underbrace{1 \times 10^{-3}}_{\parallel \varepsilon^{(1)}}, \quad \underbrace{\frac{1 + \sqrt{13}}{2} \times 10^{-3}}_{\varepsilon^{(2)}}, \quad \underbrace{\frac{1 - \sqrt{13}}{2} \times 10^{-3}}_{\varepsilon^{(3)}}$$

← principal direction

② find $\underline{v}^{(i)}$ for each $\varepsilon^{(i)}$

i) $\varepsilon^{(1)} = 1 \times 10^{-3}$

$$(\underline{\underline{\varepsilon}} - \underline{\underline{\varepsilon}}^{(1)} \underline{\underline{1}}) \cdot \underline{v}^{(1)} = 0$$

$$10^{-3} \begin{bmatrix} 0 & \sqrt{3} & 0 \\ \sqrt{3} & -1 & 0 \\ 0 & 0 & 0 \end{bmatrix} \begin{Bmatrix} v_1^{(1)} \\ v_2^{(1)} \\ v_3^{(1)} \end{Bmatrix} = \begin{Bmatrix} 0 \\ 0 \\ 0 \end{Bmatrix}$$

$$\Rightarrow \begin{cases} \sqrt{3} v_2^{(1)} = 0 & \rightarrow v_2^{(1)} = 0 \\ \sqrt{3} v_1^{(1)} - v_2^{(1)} = 0 & v_1^{(1)} = 0 \\ 0 v_3^{(1)} = 0 & v_3^{(1)} = \text{arbitrary} \neq 0 \\ & \equiv 1 \end{cases}$$

$$\|\underline{v}\| = 1 \quad (\text{normalization})$$

$$\underline{v}^{(1)} \sim \{0, 0, 1\}^T$$

ii) Repeat for $\underline{\epsilon}^{(2)}$ and $\underline{\epsilon}^{(3)}$

$$\underline{v}^{(2)} = \{0.8, 0.6, 0\}^T$$

$$\underline{v}^{(3)} = \{0.6, -0.8, 0\}^T$$

Check:

$\underline{v}^{(i)}$; orthogonal to each other

< Spectral Decomposition >

Base vector at P was $\{\underline{e}_1, \underline{e}_2, \underline{e}_3\}$.

$$\underline{\underline{\epsilon}} = \epsilon_{ij} \underline{e}_i \otimes \underline{e}_j$$

How about using $\underline{v}^{(1)}, \underline{v}^{(2)}, \underline{v}^{(3)}$ as basis vector?

$$\underline{\underline{\epsilon}} \stackrel{\Delta}{=} \sum_{i,j} \hat{\epsilon}_{ij} \underline{v}^{(i)} \otimes \underline{v}^{(j)}$$

i^{th} principal strain

$$\Rightarrow \underline{\underline{\epsilon}} = \sum_{i=1}^3 \epsilon^{(i)} \underline{v}^{(i)} \otimes \underline{v}^{(i)} \quad (*)$$

"spectral decomposition"

proof of (*)

$$\textcircled{1} \quad \underline{\underline{\epsilon}} \cdot \underline{v}^{(i)} = \epsilon^{(i)} \underline{v}^{(i)}, \text{ or } \epsilon_{kl} v_l^{(i)} = \epsilon^{(i)} v_k^{(i)}$$

(no sum on i)

$$\textcircled{2} \text{ orthogonality: } \underline{v}^{(i)} \cdot \underline{v}^{(j)} = \delta^{(ij)}$$

$$v_k^{(i)} v_k^{(j)} = \delta^{(ij)}$$

③ coordinate transform

$$\begin{array}{ccc} \underline{e}_i & \longleftrightarrow & \underline{e}_{i'} \\ \{ \underline{e}_1, \underline{e}_2, \underline{e}_3 \} & & \{ \underline{e}_{1'}, \underline{e}_{2'}, \underline{e}_{3'} \} \\ & & \parallel \parallel \parallel \\ & & \underline{v}^{(1')} \quad \underline{v}^{(2')} \quad \underline{v}^{(3')} \end{array}$$

$$\left. \begin{array}{l} \underline{e}_{i'} \stackrel{\Delta}{=} \beta_{ji'} \underline{e}_j \\ \underline{v}^{(i')} = v_j^{(i')} \underline{e}_j \end{array} \right\} \underline{\beta}_{ji'} = v_j^{(i')} \quad (*)$$

Then

$$\epsilon_{ij} \text{ (in } \underline{e}_i) \rightarrow \epsilon_{i'j'} \text{ (in } \underline{e}_{i'} = \underline{v}^{(i')})$$

$$\epsilon_{i'j'} = \beta_{i'k} \beta_{j'l} \epsilon_{kl}$$

In orthonormal coord (e.g., Cartesian)

$$\beta_{i'k} = \beta_{ki'} \text{ (see probl. 1.1)}$$

Thus

$$\epsilon_{i'j'} = \beta_{ki'} \beta_{lj'} \epsilon_{kl} \stackrel{(*)}{=} v_k^{(i')} v_l^{(j')} \epsilon_{kl}$$

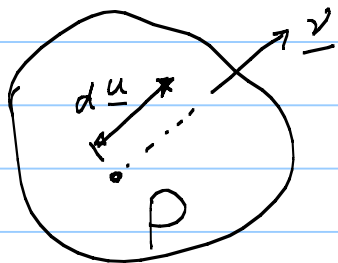
$$\begin{aligned}
 &= v_K^{(i')} \varepsilon_{K\ell} v_\ell^{(j')} \\
 &\quad \searrow \textcircled{\ominus} \rightarrow \varepsilon^{(i'j')} v_K^{(j')} \quad (\text{no sum on } j') \\
 &= \varepsilon^{(i'j')} v_K^{(j')} v_K^{(i')} \quad (\text{no sum on } j') \\
 &\quad \searrow \rightarrow \delta^{(i'j')} \\
 &= \delta^{(i'j')} \varepsilon^{(j')} \quad (\text{no sum on } j')
 \end{aligned}$$

Thus

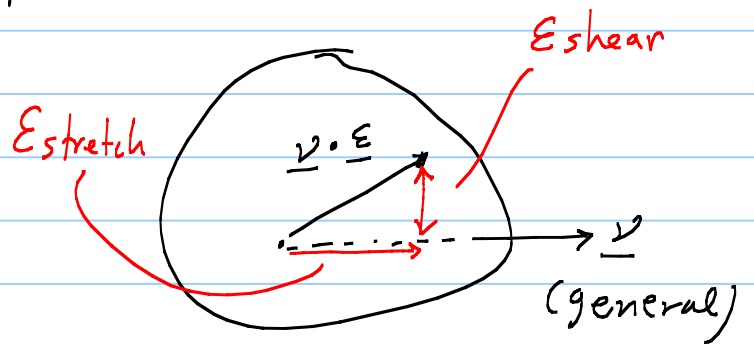
$$\begin{aligned}
 \varepsilon &= \varepsilon_{ij} e_i \otimes e_j \\
 &= \sum_{i,j} \delta^{(i'j')} \varepsilon^{(j')} v^{(i')} \otimes v^{(j')} \\
 &= \sum_{i,j} \varepsilon^{(j')} v^{(j')} \otimes v^{(j')} \\
 &= \sum_{j'} \varepsilon^{(i')} v^{(i')} \otimes v^{(j')}
 \end{aligned}$$



Deformation along general direction \underline{v}
 (not the principal direction)



if \underline{v} : principal

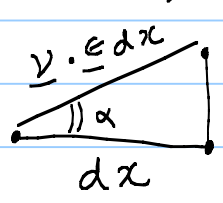


$$\begin{aligned} d\underline{u} &= \underline{v} \cdot \underline{\nabla} \underline{u} \\ &= \underline{v} \cdot (\underline{E} + \underline{\omega}) dx \\ &\quad \text{[Rigid-Body rotation]} \end{aligned}$$

$\underline{v} \cdot \underline{E}$: deformation measure along \underline{v} excluding Rigid-Body Rotation

$$E_{stretch} = \underline{v} \cdot \underline{E} \cdot \underline{v}$$

$$E_{shear}^2 = (\underline{v} \cdot \underline{E})^2 - E_{stretch}^2$$

($\tan \alpha \approx \alpha$ (for small α)  $\leftarrow E_{shear} dx$; $\alpha \approx \frac{E_{shear} \cdot dx}{d\underline{u}}$)

One can show:

$$\begin{cases} \varepsilon_{\text{stretch}}|_{\text{max}} = \text{Largest of } (\varepsilon^{(1)}, \varepsilon^{(2)}, \varepsilon^{(3)}) \\ \varepsilon_{\text{stretch}}|_{\text{min}} = \text{Smallest of } \varepsilon^{(1)}, \varepsilon^{(2)}, \varepsilon^{(3)} \end{cases}$$

$$\varepsilon_{\text{shear}}|_{\text{max}} = \max \frac{1}{2} \left\{ |\varepsilon^{(1)} - \varepsilon^{(2)}|, |\varepsilon^{(2)} - \varepsilon^{(3)}|, |\varepsilon^{(3)} - \varepsilon^{(1)}| \right\}$$

Back to Ex 1.2.2

$$\underline{\varepsilon} = \left[\underline{e}_1 \otimes \underline{e}_1 + \underline{e}_3 \otimes \underline{e}_3 + \sqrt{3} (\underline{e}_1 \otimes \underline{e}_2 + \underline{e}_2 \otimes \underline{e}_3) \right] \times 10^{-3}$$

problem c). $\underline{v} = \frac{1}{\sqrt{2}} (\underline{e}_1 + \underline{e}_2)$

$$\varepsilon_{\text{stretch}} = \underline{v} \cdot \underline{\varepsilon} \cdot \underline{v} = \frac{1}{2} (1 + 2\sqrt{3}) \times 10^{-3}$$

$$\varepsilon_{\text{shear}} = \left((\underline{v} \cdot \underline{\varepsilon})^2 - \varepsilon_{\text{stretch}}^2 \right)^{\frac{1}{2}}$$

$$= \frac{1}{2} \times 10^{-3} \quad (\text{Rotation angle})$$

(problem c) $\varepsilon_{\text{shear}}|_{\text{max}} = \frac{1}{2} (\varepsilon^{(2)} - \varepsilon^{(3)}) = \frac{\sqrt{13}}{2} \times 10^{-3}$

Compatibility of Strain

(Read p 56-61
for derivation) ^{2/}

3 Displacement \rightarrow 6 Strains
(u_1, u_2, u_3) ($\epsilon_{11}, \epsilon_{22}, \epsilon_{33},$
 $\epsilon_{12}, \epsilon_{23}, \epsilon_{31}$)

$$\begin{cases} \epsilon_{33,22} + \epsilon_{22,33} - 2\epsilon_{23,23} = 0 \\ \epsilon_{11,33} + \epsilon_{33,11} - 2\epsilon_{31,31} = 0 \\ \epsilon_{22,11} + \epsilon_{11,22} - 2\epsilon_{12,12} = 0 \end{cases}$$

$$\begin{cases} -\epsilon_{11,23} + (-\epsilon_{23,1} + \epsilon_{31,2} + \epsilon_{12,3}),_1 = 0 \\ -\epsilon_{22,31} + (-\epsilon_{31,2} + \epsilon_{12,3} + \epsilon_{23,1}),_2 = 0 \\ -\epsilon_{33,12} + (-\epsilon_{12,3} + \epsilon_{23,1} + \epsilon_{31,2}),_3 = 0 \end{cases}$$

May be written as

$$\begin{cases} \nabla \times (\nabla \times \underline{\epsilon})^T = 0 \\ \text{or} \\ \epsilon_{ilk} \epsilon_{jmp} \epsilon_{lm, kp} = 0 \end{cases}$$