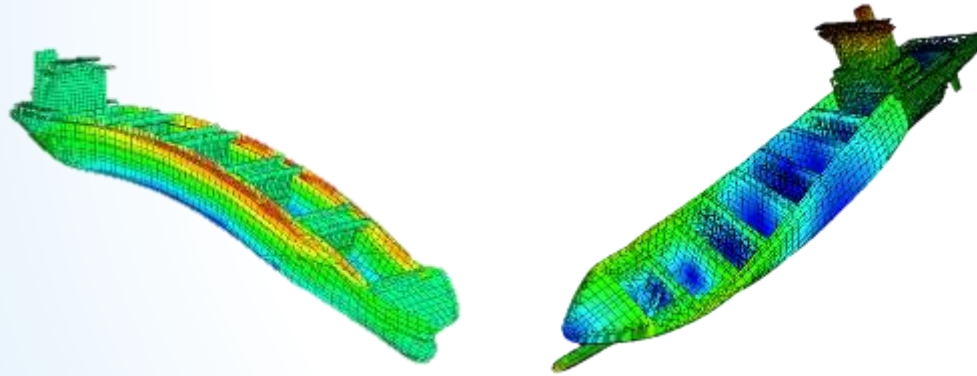


Topics in Ship Structures

(Advanced Local Structural Design & Analysis of Marine Structures)



* Elastic Plate Theory - Basic (Topic 4)

Do Kyun Kim
Seoul National University



[Theory of Plates and Grillages]

[Part I] Plastic Design of Structures

- Plastic theory of bending (Topic 1)
- Ultimate loads on beams (Topic 2)
- Collapse of frames and grillage structures (Topic 3)

[Part II] Elastic Plate Theory under Pressure

- Basic (Topic 4)
- Simply supported plates under Sinusoidal Loading (Topic 5)
- Long clamped plates (Topic 6)
- Short clamped plates (Topic 7)
- Low aspect ratio plates, strength & permanent set (Topic 7A)

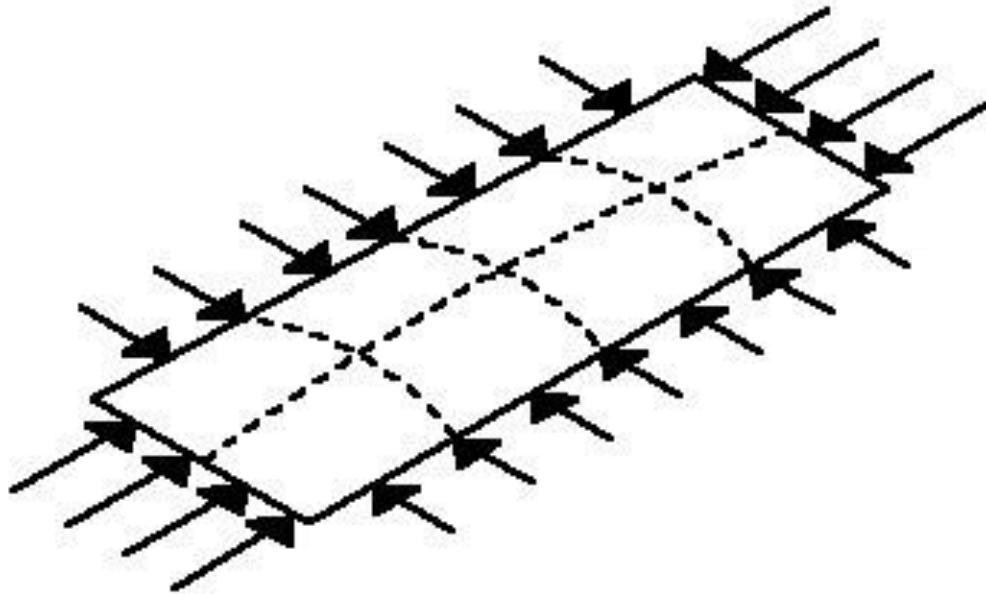
[Part III] Buckling of Stiffened Panels

- Failure modes (Topic 8)
- Tripping (Topic 9) + Post-buckling strength of plate (Topic 9A)
- Post-buckling behaviour (Topic 10)



The aim of this lecture is:

- To **equip** you with the **knowledge and understanding** of **elastic plate theory**.



Picture from:

<http://fgg-web.fgg.uni-lj.si/~pmoze/esdep/master/wg08/10100.htm>

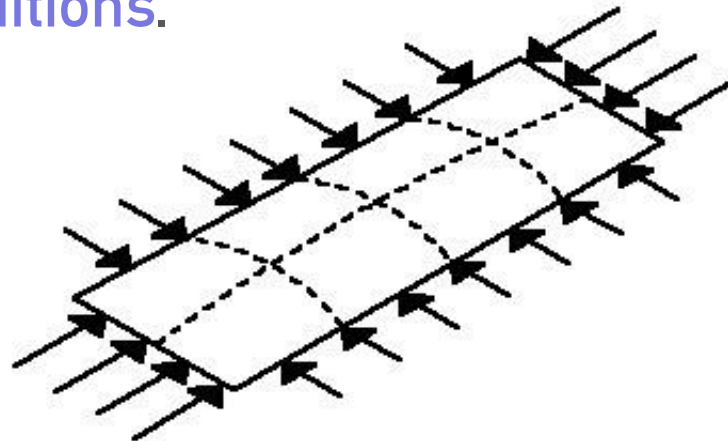
http://homepages.engineering.auckland.ac.nz/~pkel015/SolidMechanicsBooks/Part_I/

http://homepages.engineering.auckland.ac.nz/~pkel015/SolidMechanicsBooks/Part_II/

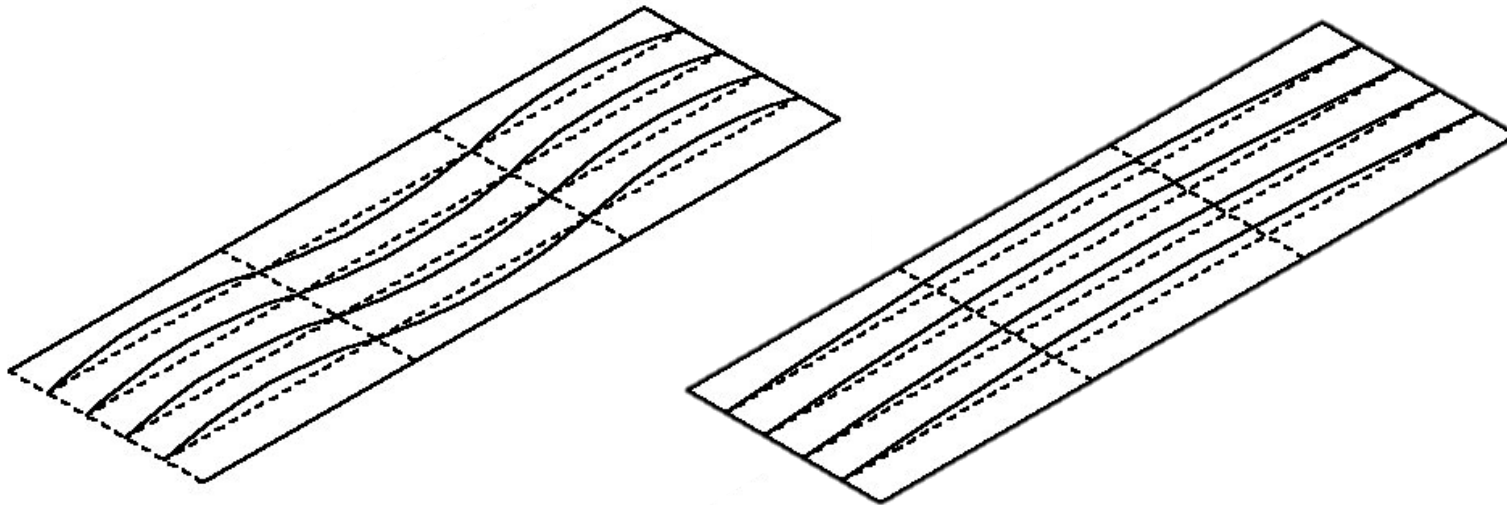


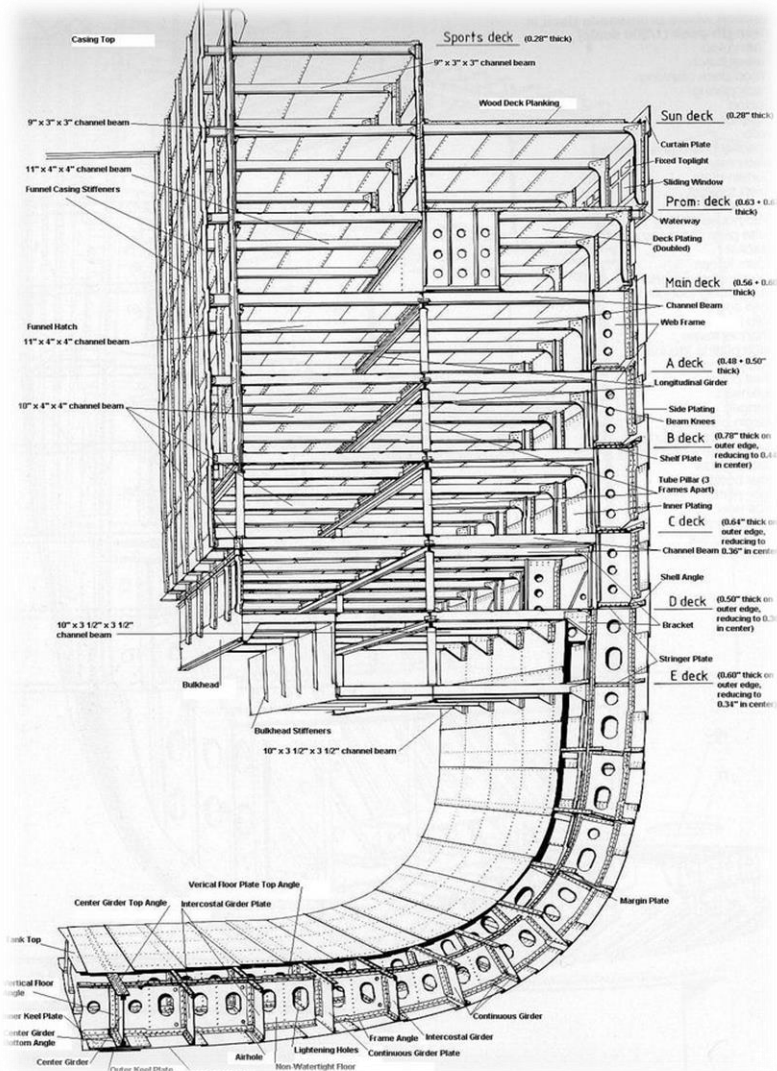
At the end of this lecture, you should be able to:

- Appreciate the assumptions of **elastic plate theory**
- Be familiar with the **bending moment curvature** relationship of **plate** and relationship between **twisting moment** and **twist** in a plate.
- Establish **plate governing equilibrium equation**.
- Be aware of **plate edge conditions**.



- Ships are cut, almost entirely, from **steel plate**. Only a **few** structural elements are **NOT** made from **plate** (e.g. a few castings and some rolled sections).
- An **understanding** of **plate behaviour** is therefore **Crucial**.

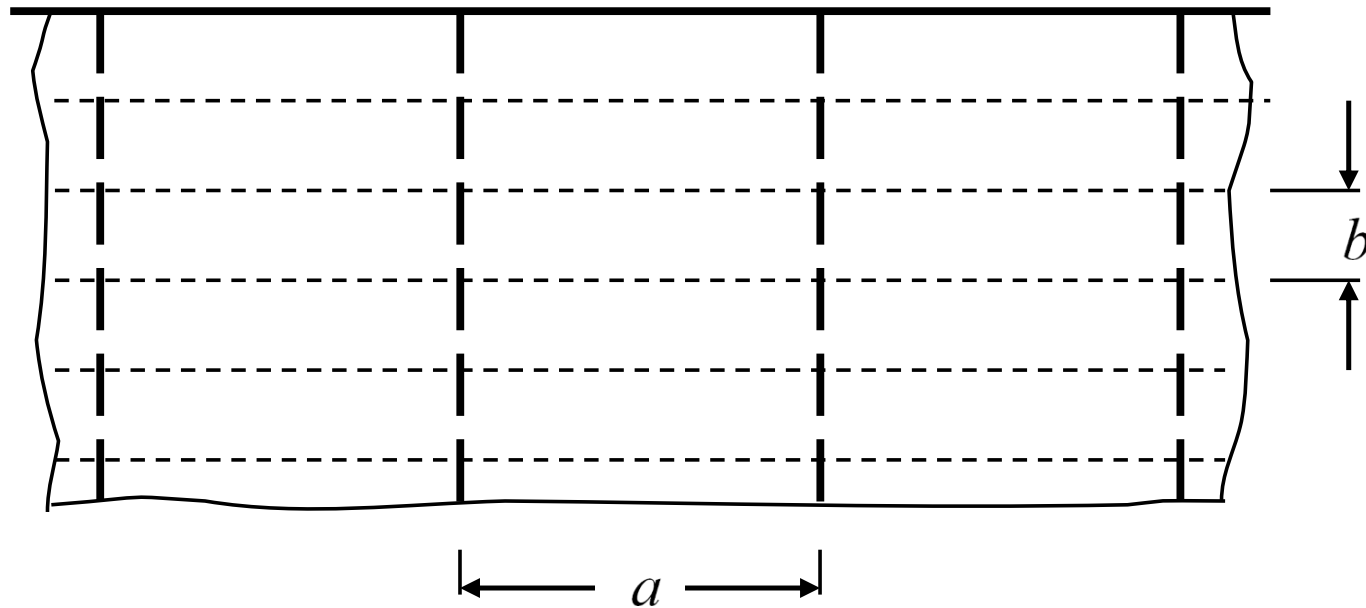




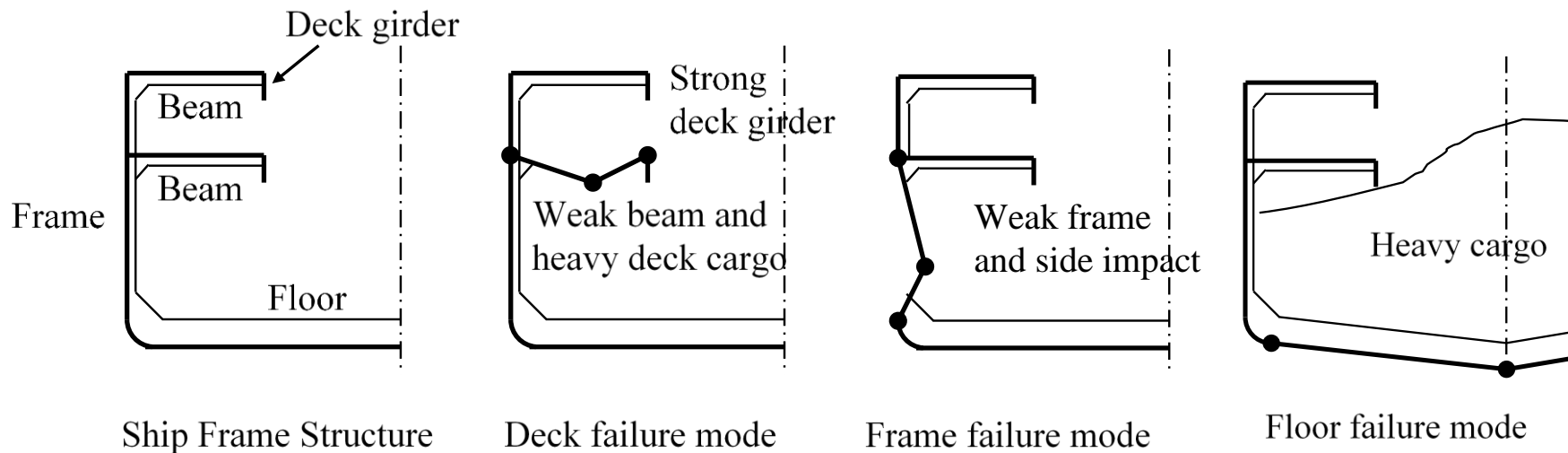
- Some components such as **deck plate** between supports, behave as **ideal plates** under **simple boundary conditions and loads**.
- Other components, such as the **webs** of large frames are subject to **more complex loading**.
- Still, other components such as **deck/stiffener combinations** behave somewhat like **simple plates**, but with **special aspects**.



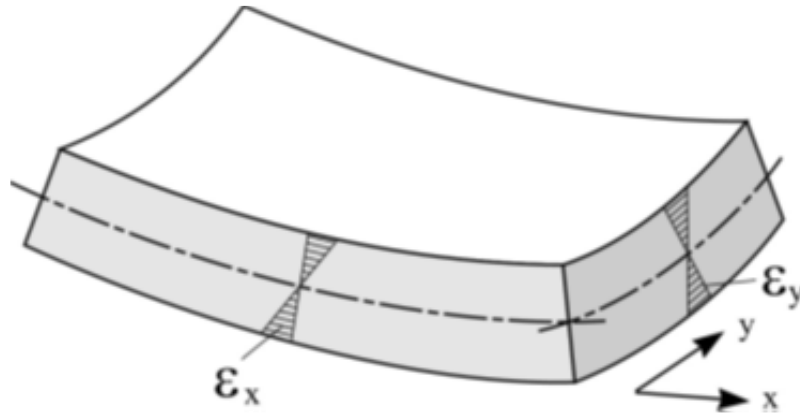
- Main **hull structures** are made up with **plating** supported by **stiffeners**.
- For example, **deck plates** are supported by **beams/longitudinals** and **girders/transverses**, **bottom plates** supported by **longitudinals, floors** and **bottom girders**, etc.



- These **plates** must have **sufficient strengths** to **withstand lateral loads** due to water pressures and/or cargoes, and **stiffness** to **resist deformations**.
- Now you should think how to **relate** the **bending/twisting moment, curvature** and **plate flexural rigidity** together in order to **calculate stresses** and **strains** arisen in the plates.



- We start by sketching a **small portion of plate bending in two directions**, but **with no in-plane stress**.



- In other words, the **stress at the mid-plane** will be **ZERO**, and (equivalently) the **average stress** in the **x** or **y** direction will be **ZERO**.
- There is **bending stress** and **strain** in both **x** and **y** directions. We **assume** that $\sigma_z \approx 0$.

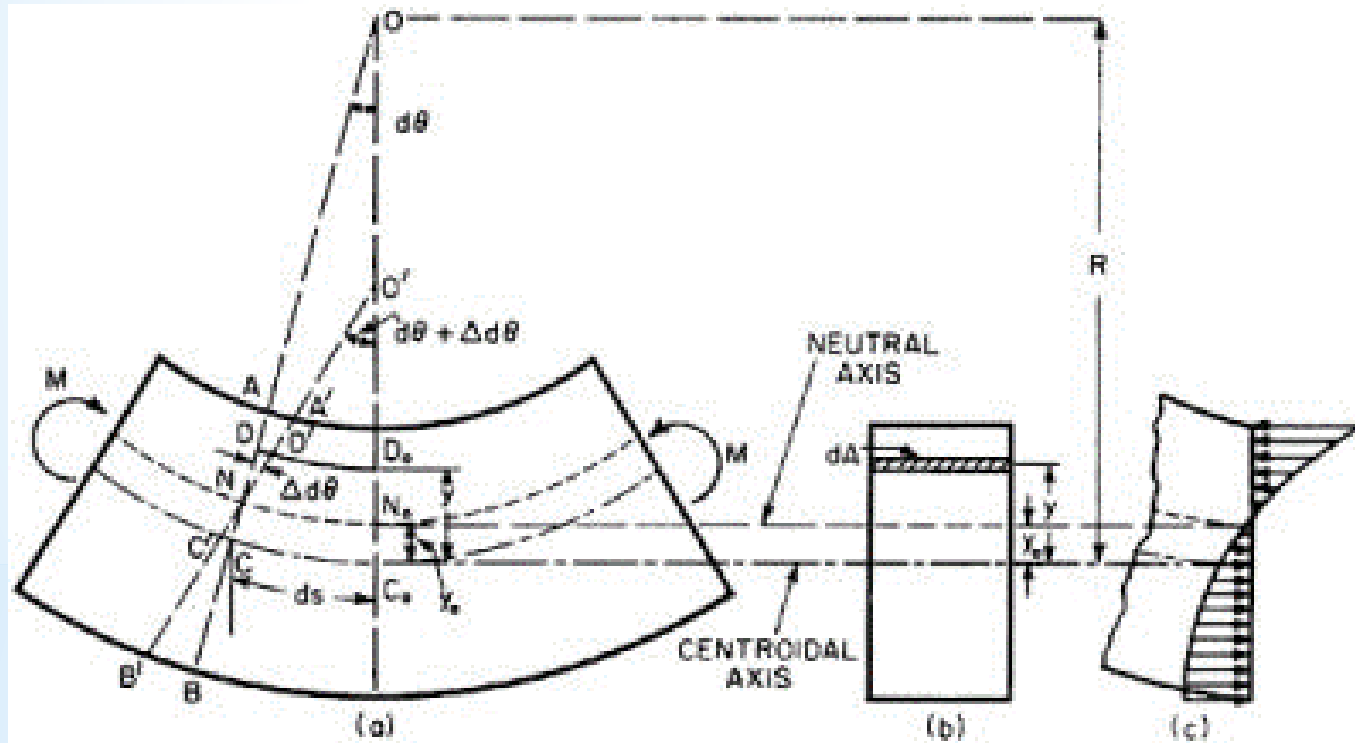


It's usually assumed that in the **bending** of **rectangular plate** elements under normal **lateral load** that:

1. The **boundary supports** (provided by the stiffeners) Do Not Deflect normal to the plate.
2. **Stress** and **strain** in a direction **perpendicular** to the **plate** may be Neglected.
3. Plane sections remain plane, with the middle surface as the neutral surface of the plate in the special case of pure bending.
4. **Deflection** w is Small (less than $t/2$).
5. **Material** remains Elastic.

-
- The resulting theory based on the above assumptions is known as “**small deflection**” or “**linear theory**”
 - The results for a variety of loads are **superposable** always providing of course that the **limit of proportionality of the material** is **not exceeded**.





*Topic 4.1:

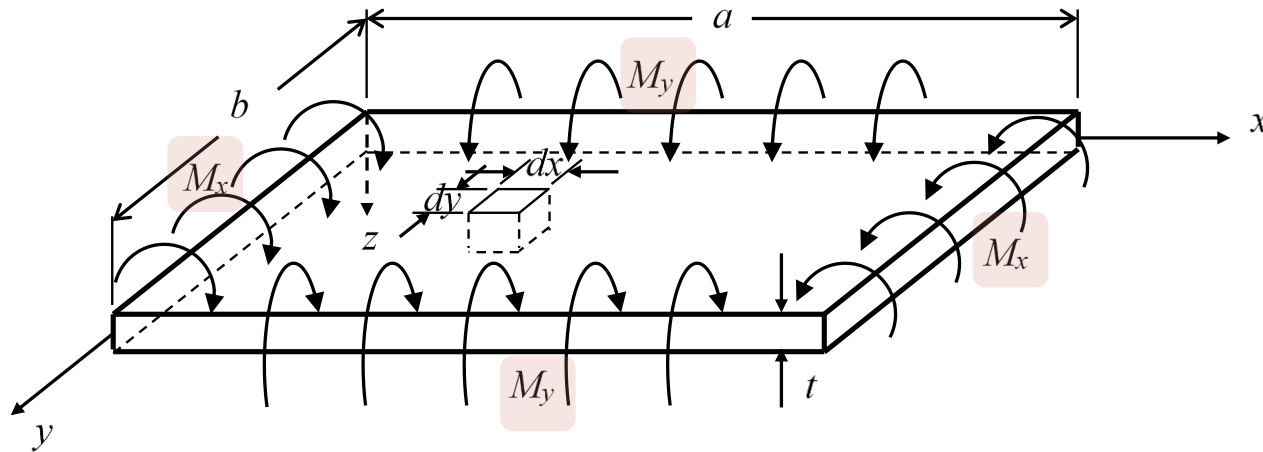
Bending Moment-Curvature Relationships

Picture from:

<https://physics.stackexchange.com/questions/160497/is-the-moment-curvature-relation-for-an-elastic-beam-general>



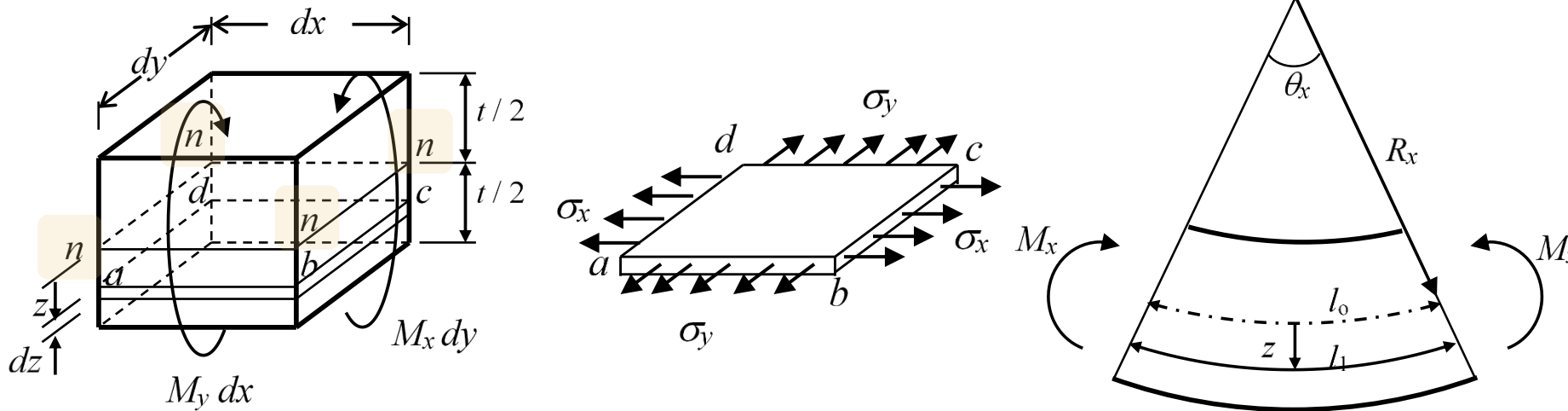
- A plate of thickness t , length a and breadth b is subjected to bending moments on its edges as shown below.



- M_x and M_y are the bending moments per unit length on the edges.
- The subscript indicates the direction of normal of the edge surface on which the bending moment acts.



- Consider a **small element** $dx \times dy \times t$ in the plate.

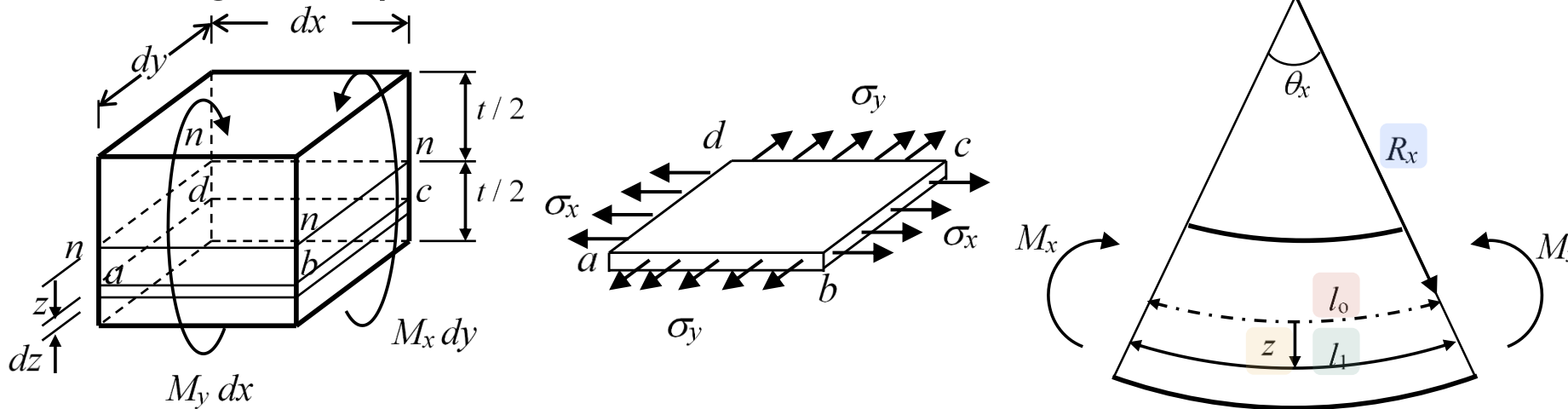


- nn is the **neutral plane**.
- Due to M_x and M_y , **bending stresses** σ_x and σ_y act on lamina $abcd$ of distance z from the **neutral plane**.



[Assumption 3] Plane sections remain plane, with the middle surface as the neutral surface of the plate in the special case of pure bending.

- Using assumption 3, we have



- bending strain** in x direction:

$$\varepsilon_x = (l_1 - l_0) / l_0 = [(R_x + z)\theta_x - R_x\theta_x] / (R_x\theta_x) = z / R_x$$

- bending strain** in y direction: $\varepsilon_y = z / R_y$

where R_x and R_y are radii of curvature of neutral plane in xz and yz plane.



Moment-Curvature Relationships

Required Basics

- $\sigma = \frac{F}{A}$
- $\sigma = E\varepsilon$ (Hooke's Law)
- $l = R\theta$
- $M = F \times y = \sigma A y$

ε (strain) = $\frac{\Delta \text{ (elongation)}}{l_0 \text{ (original length)}}$
 where, $\Delta = l_1 - l_0$

$\varepsilon = \frac{\Delta}{l_0} = \frac{l_1 - l_0}{l_0} = \frac{(R+y)\theta - R\theta}{R\theta} = \frac{y}{R}$
 where, $l_1 = (R+y)\theta$
 $l_0 = R\theta$

$\sigma = E\varepsilon = E\left(\frac{y}{R}\right) = \frac{Ey}{R} \Rightarrow \frac{E}{R} = \frac{\sigma}{y}$

$M = Fy = \sigma Ay = \frac{E Ay^2}{R}$
 if we consider whole sectional area

$M = \frac{E}{R} \int y^2 dA = \frac{EI}{R} \Rightarrow \frac{1}{R} = \frac{M}{EI}$
 I (moment of inertia)

Curvature (c) relationship with "Arc" s_{ds} according to the change of $(d\theta)$

$c = \frac{d\theta}{ds} = \frac{d\theta}{dx} \times \frac{dx}{ds}$
 $\frac{dx}{ds} = \frac{dx}{\sqrt{dx^2 + dy^2}} = \frac{1}{\sqrt{1 + \left(\frac{dy}{dx}\right)^2}}$

$\left(\frac{1}{R}\right) = \frac{d}{dx} \left[\tan^{-1} \left(\frac{dy}{dx} \right) \right] \times \frac{1}{\sqrt{1 + \left(\frac{dy}{dx}\right)^2}}$

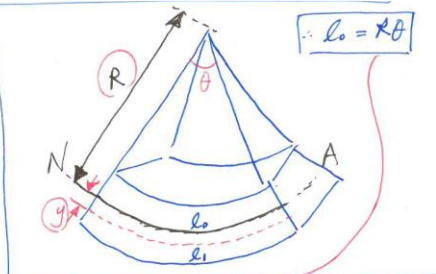
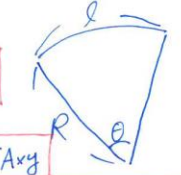
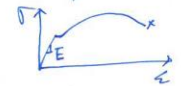
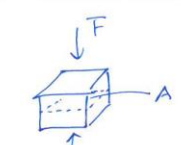
$= \frac{1}{1 + \left(\frac{dy}{dx}\right)^2} \times \frac{d^2y}{dx^2} \times \frac{1}{\sqrt{1 + \left(\frac{dy}{dx}\right)^2}}$

$\therefore C = \frac{d^2y/dx^2}{\left[1 + \left(\frac{dy}{dx}\right)^2\right]^{3/2}}$ if $\frac{dy}{dx} \ll 1$ (very small)
 $C = \frac{d\theta}{ds} = \frac{1}{R} = \frac{d^2y}{dx^2}$

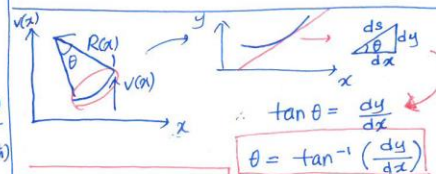
$\therefore \frac{1}{R} = \frac{M}{EI} = \frac{d^2y}{dx^2}$ or $\frac{d^2w}{dx^2}$

$\therefore \frac{M}{EI} = \frac{\sigma}{y}$

Additional tip:
 if $y = \tan^{-1}(x) \mid y' = \frac{dy}{dx} = \frac{1}{x^2+1}(x')$



$$\theta = \frac{l_0}{R} = \left(\frac{1}{R}\right) \times l_0 = \frac{M}{EI} l_0 = \frac{M l_0}{EI}$$



$$ds = \sqrt{(dx)^2 + (dy)^2} \Rightarrow \text{Pythagorean theorem}$$

curvature (c) \Rightarrow relationship with "Arc" s_{ds} according to the change of $(d\theta)$

$$c = \frac{d\theta}{ds} = \frac{d\theta}{dx} \times \frac{dx}{ds}$$

$$\left(\frac{1}{R}\right) = \frac{d}{dx} \left[\tan^{-1} \left(\frac{dy}{dx} \right) \right] \times \frac{1}{\sqrt{1 + \left(\frac{dy}{dx}\right)^2}}$$

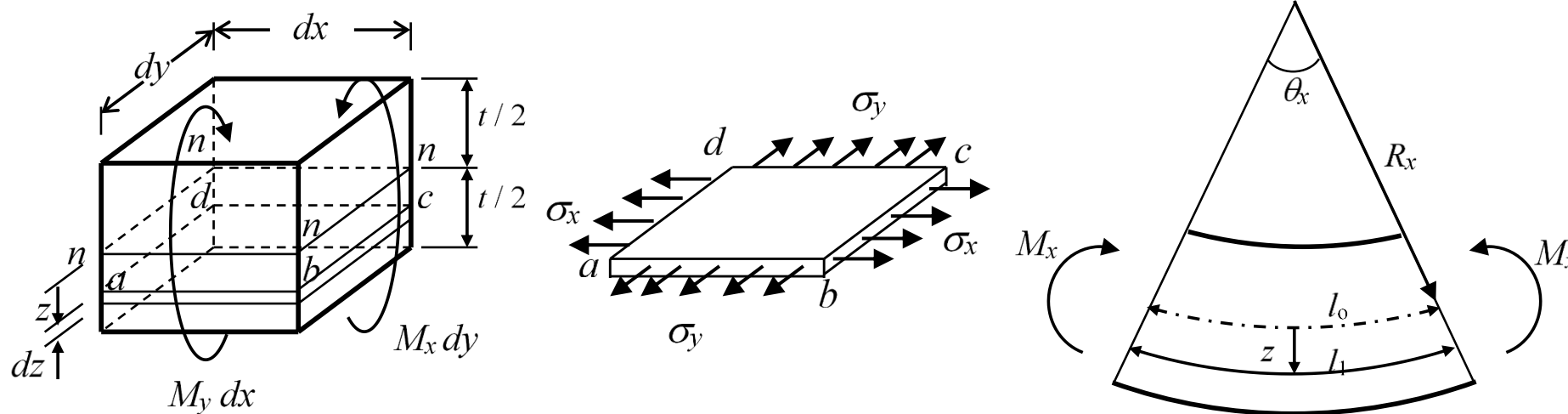
$$C = \frac{d^2y/dx^2}{\left[1 + \left(\frac{dy}{dx}\right)^2\right]^{3/2}}$$

$$\therefore \frac{1}{R} = \frac{M}{EI} = \frac{d^2y}{dx^2} \text{ or } \frac{d^2w}{dx^2}$$



[Assumption 4] Deflection w is **Small** (less than $t/2$).

- Using assumption 4, we have



Curvature (in x direction)

$$\frac{1}{R_x} = -\frac{\partial^2 w}{\partial x^2}$$

and

Curvature (in y direction)

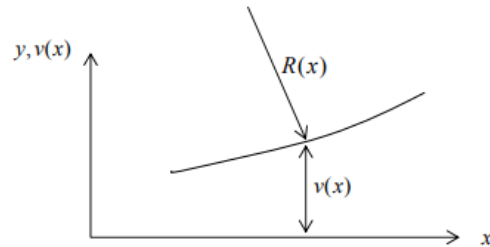
$$\frac{1}{R_y} = -\frac{\partial^2 w}{\partial y^2}$$



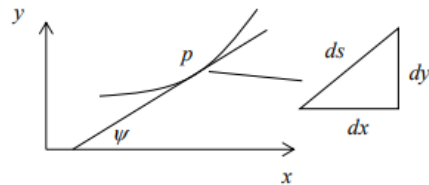
7.4.11 Appendix to §7.4

Curvature of the deflection curve

Consider a deflection curve with deflection $v(x)$ and radius of curvature $R(x)$, as shown in the figure below. Here, *deflection* is the transverse displacement (in the y direction) of the points that lie along the axis of the beam. A relationship between $v(x)$ and $R(x)$ is derived in what follows.



First, consider a curve (arc) s . The tangent to some point p makes an angle ψ with the x -axis, as shown below. As one moves along the arc, ψ changes.



Define the **curvature** κ of the curve to be the rate at which ψ increases relative to s ,

$$\kappa = \frac{d\psi}{ds}$$

Thus if the curve is very “curved”, ψ is changing rapidly as one moves along the curve (as one increase s) and the curvature will be large.

From the above figure,

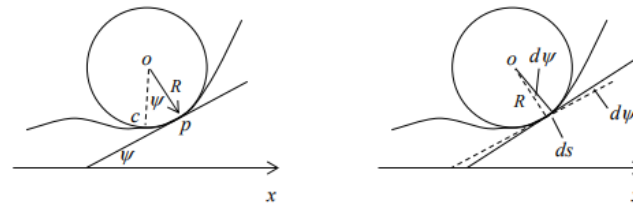
$$\tan \psi = \frac{dy}{dx}, \quad \frac{ds}{dx} = \frac{\sqrt{(dx)^2 + (dy)^2}}{dx} = \sqrt{1 + (dy/dx)^2},$$

so that

$$\begin{aligned} \kappa &= \frac{d\psi}{ds} = \frac{d\psi}{dx} \frac{dx}{ds} = \frac{d(\arctan(dy/dx))}{dx} \frac{dx}{ds} = \frac{1}{1 + (dy/dx)^2} \frac{d^2y}{dx^2} \frac{dx}{ds} \\ &= \frac{d^2y/dx^2}{[1 + (dy/dx)^2]^{3/2}} \end{aligned}$$

Finally, it will be shown that the curvature is simply the reciprocal of the radius of curvature. Draw a circle to the point p with radius R . Arbitrarily measure the arc length s from the point c , which is a point on the circle such that $\angle cop = \psi$. Then arc length $s = R\psi$, so that

$$\kappa = \frac{d\psi}{ds} = \frac{1}{R}$$



Thus

$$\frac{1}{R} = \frac{\frac{d^2v}{dx^2}}{\left[1 + \left(\frac{dv}{dx}\right)^2\right]^{3/2}}$$

If one assumes now that the slopes of the deflection curve are small, then $dv/dx \ll 1$ and

$$\frac{1}{R} \approx \frac{d^2v}{dx^2}$$



[Assumption 5] Material remains Elastic.

$$\varepsilon_x = \frac{z}{R_x} = z \left(-\frac{\partial^2 w}{\partial x^2} \right)$$

&

$$\varepsilon_y = \frac{z}{R_y} = z \left(-\frac{\partial^2 w}{\partial y^2} \right)$$

- Using assumption 5, we have the stress-strain relations for in-plane stresses (i.e. $\sigma_z = 0$) are:

$$\begin{Bmatrix} \varepsilon_x \\ \varepsilon_y \end{Bmatrix} = \frac{1}{E} \begin{bmatrix} 1 & -\nu \\ -\nu & 1 \end{bmatrix} \begin{Bmatrix} \sigma_x \\ \sigma_y \end{Bmatrix}$$

where ν is **Poisson's ratio**.

- Re-arranging the above equations yields

$$\begin{Bmatrix} \sigma_x \\ \sigma_y \end{Bmatrix} = \frac{E}{1-\nu^2} \begin{bmatrix} 1 & \nu \\ \nu & 1 \end{bmatrix} \begin{Bmatrix} \varepsilon_x \\ \varepsilon_y \end{Bmatrix}$$

Hence,

$$\begin{Bmatrix} \sigma_x \\ \sigma_y \end{Bmatrix} = -\frac{Ez}{1-\nu^2} \begin{bmatrix} 1 & \nu \\ \nu & 1 \end{bmatrix} \begin{Bmatrix} \partial^2 w / \partial x^2 \\ \partial^2 w / \partial y^2 \end{Bmatrix} \quad (i)$$

https://en.wikipedia.org/wiki/Sim%C3%A9on_Denis_Poisson

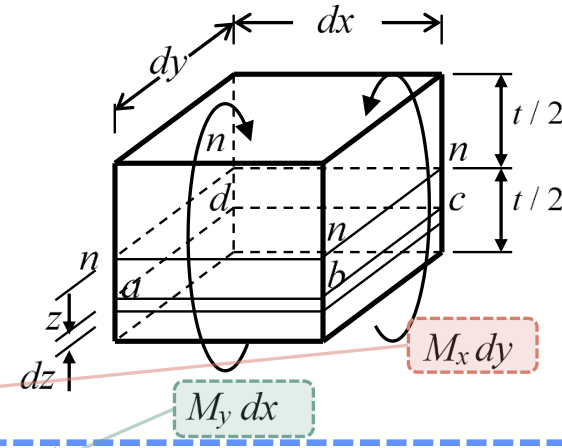
https://en.wikipedia.org/wiki/Poisson%27s_ratio

$$\begin{Bmatrix} \varepsilon_{xx} \\ \varepsilon_{yy} \\ \varepsilon_{zz} \end{Bmatrix} = \begin{Bmatrix} \frac{1}{E} [\sigma_{xx} - \nu(\sigma_{yy} + \cancel{\sigma_{zz}})] \\ \frac{1}{E} [\sigma_{yy} - \nu(\sigma_{xx} + \cancel{\sigma_{zz}})] \\ \frac{1}{E} [\cancel{\sigma_{zz}} - \nu(\sigma_{xx} + \sigma_{yy})] \end{Bmatrix}$$

$$= \frac{1}{E} \begin{bmatrix} \sigma_{xx} & -\nu\sigma_{yy} \\ -\nu\sigma_{xx} & \sigma_{yy} \end{bmatrix} = \frac{1}{E} \begin{bmatrix} 1 & -\nu \\ -\nu & 1 \end{bmatrix} \begin{Bmatrix} \sigma_{xx} \\ \sigma_{yy} \end{Bmatrix}$$



$$\begin{Bmatrix} \sigma_x \\ \sigma_y \end{Bmatrix} = -\frac{Ez}{1-\nu^2} \begin{bmatrix} 1 & \nu \\ \nu & 1 \end{bmatrix} \begin{Bmatrix} \partial^2 w / \partial x^2 \\ \partial^2 w / \partial y^2 \end{Bmatrix} \quad \text{Eqn (i)}$$



- The bending moments resulting from the stresses are

$$M_x dy = \int_{-t/2}^{t/2} z \sigma_x dy dz$$

and

$$M_y dx = \int_{-t/2}^{t/2} z \sigma_y dx dz$$

or

$$M_x = \int_{-t/2}^{t/2} z \sigma_x dz$$

and

$$M_y = \int_{-t/2}^{t/2} z \sigma_y dz$$

- Substituting equations (i) into the above equations and integrating give

$$M_x = -D \left(\frac{\partial^2 w}{\partial x^2} + \nu \frac{\partial^2 w}{\partial y^2} \right)$$

$$M_y = -D \left(\frac{\partial^2 w}{\partial y^2} + \nu \frac{\partial^2 w}{\partial x^2} \right)$$

in which $D = \frac{Et^3}{12(1-\nu^2)}$ = plate flexural rigidity



$$\begin{Bmatrix} \sigma_x \\ \sigma_y \end{Bmatrix} = -\frac{Ez}{1-\nu^2} \begin{bmatrix} 1 & \nu \\ \nu & 1 \end{bmatrix} \begin{Bmatrix} \partial^2 w / \partial x^2 \\ \partial^2 w / \partial y^2 \end{Bmatrix}$$

$$M_x = -D \left(\frac{\partial^2 w}{\partial x^2} + \nu \frac{\partial^2 w}{\partial y^2} \right)$$

$$M_y = -D \left(\frac{\partial^2 w}{\partial y^2} + \nu \frac{\partial^2 w}{\partial x^2} \right)$$

$$D = \frac{Et^3}{12(1-\nu^2)}$$

Notes: • D is sometimes expressed: $D = E'I'$

where, $E' = \frac{E}{1-\nu^2}$ is effective elastic modulus for plates

$I' = t^3 / 12$ is the second moment of area per unit length

- The moment-curvature relation for a beam is given by

$$M_x = -EI \frac{d^2 w}{dx^2} \text{ in which } EI \text{ is beam flexural rigidity.}$$

- Substituting the bending moment-curvature equations into (i) gets

$$\sigma_x = \frac{M_x z}{t^3 / 12}$$

and

$$\sigma_y = \frac{M_y z}{t^3 / 12}$$

- These are analogous to

$$\frac{\sigma}{z} = \frac{M}{I'}$$



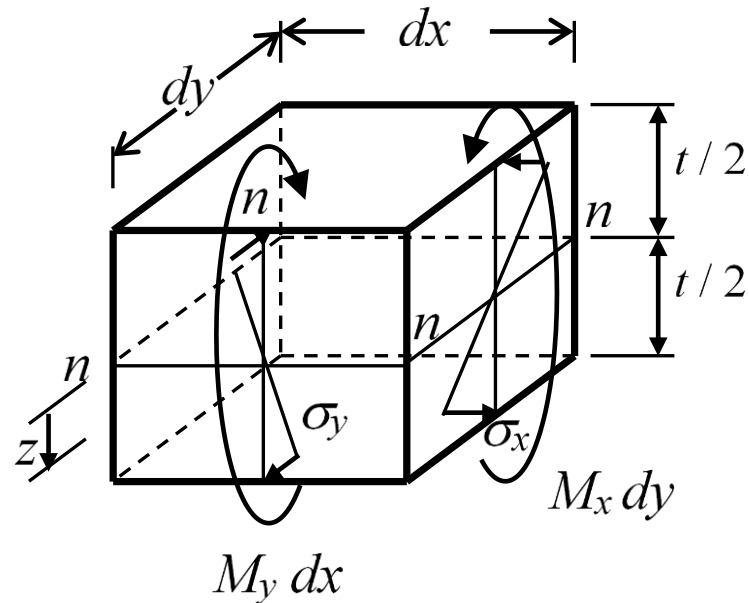
$$\sigma_x = \frac{M_x z}{t^3 / 12}$$

$$\sigma_y = \frac{M_y z}{t^3 / 12}$$

- Maximum stresses occur on **plate surface**:

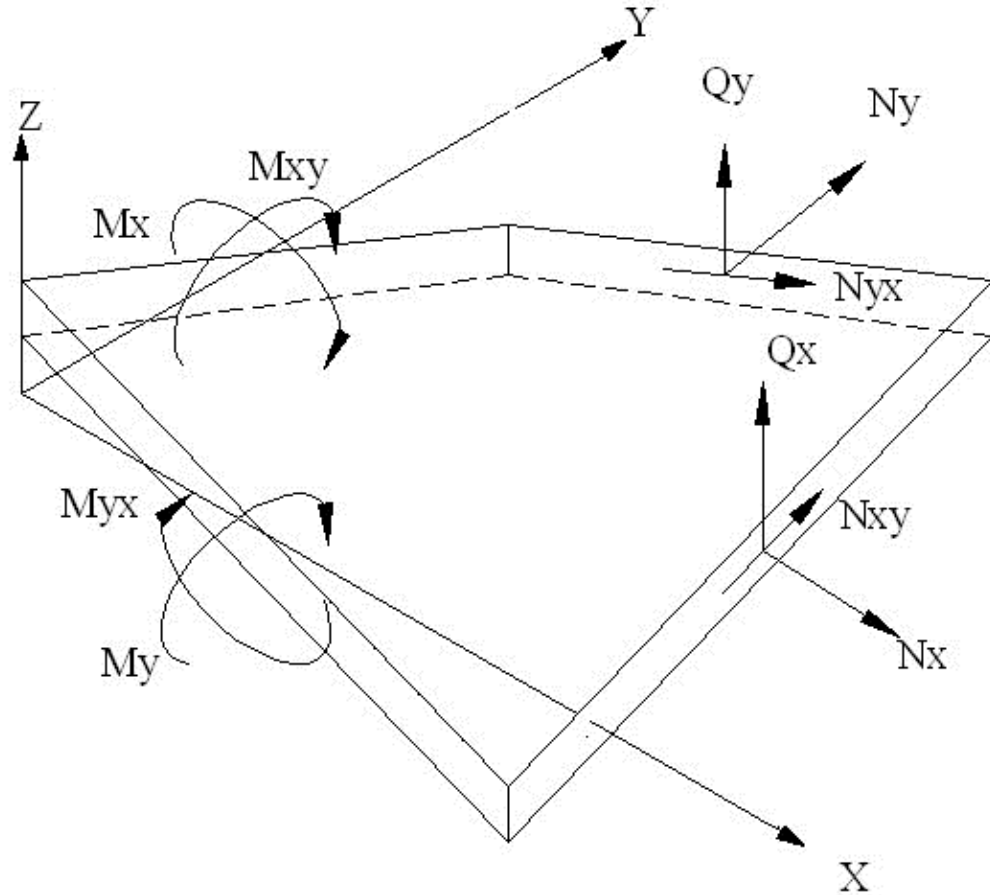
$$(\sigma_x)_{\max} = \pm \frac{6M_x}{t^2}$$

$$(\sigma_y)_{\max} = \pm \frac{6M_y}{t^2}$$



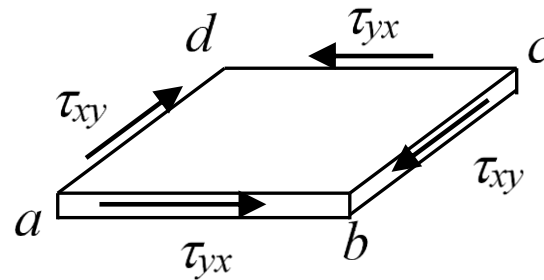
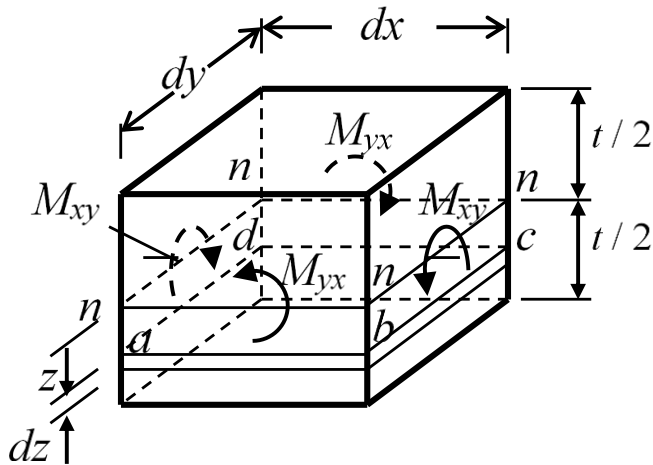
*Topic 4.2:

Twisting Moment-Twist of Plates Relationships



Relations between twisting moment and twist in a plane

- At a point in a plate subjected to **lateral load** there will in general act **both** bending and twisting moments.
- Bending moments are related to **deflection** w through **curvature**. A similar relation is required relating **twisting moments** to **deflection**.



- M_{xy} and M_{yx} are **twisting moments per unit length**.
- The directions for **twisting moments** and **shear stresses** shown above are **Positive**.



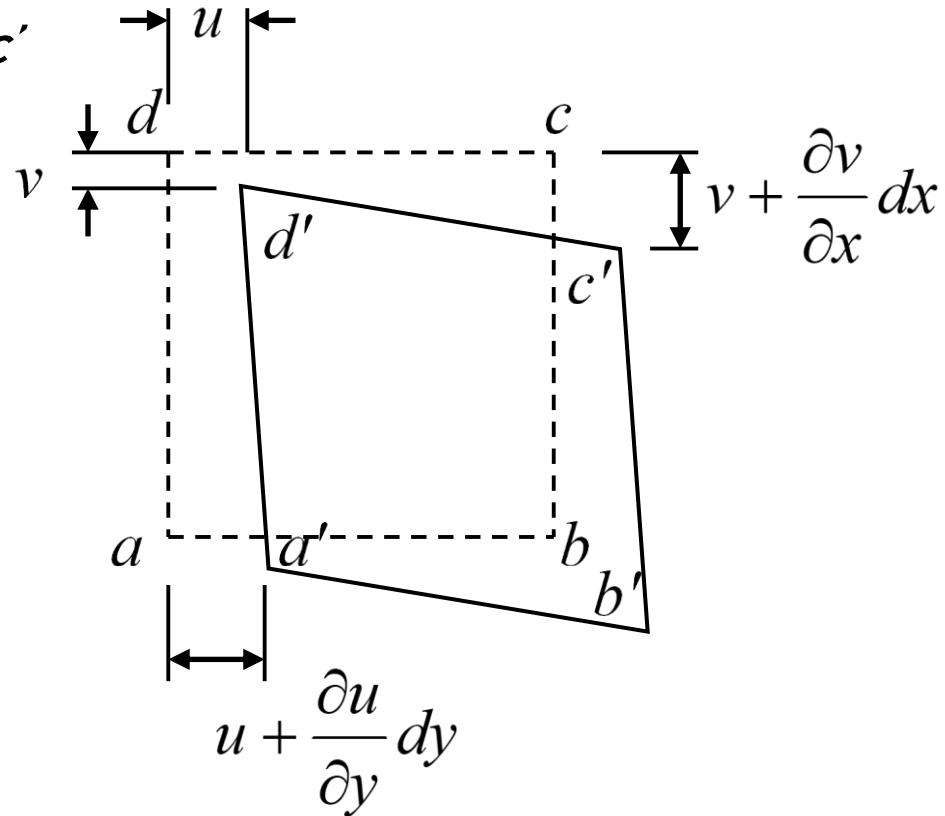
- Shear strain $\gamma_{xy} = \text{angle } adc - \text{angle } a'd'c'$

$$\gamma_{xy} = \frac{\partial v}{\partial x} + \frac{\partial u}{\partial y}$$

u and v are **displacements** in x and y direction.

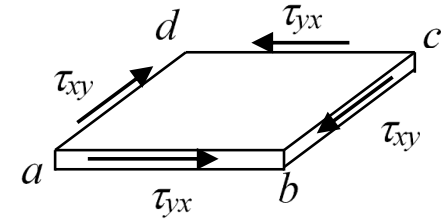
- Hence,

$$\gamma_{xy} = \gamma_{yx} = -2z \frac{\partial^2 w}{\partial x \partial y}$$



$$\gamma_{xy} = \frac{\partial v}{\partial x} + \frac{\partial u}{\partial y}$$

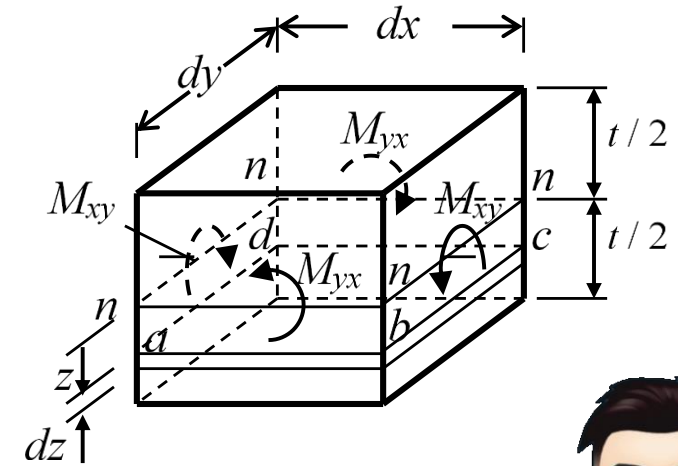
$$\gamma_{xy} = \gamma_{yx} = -2z \frac{\partial^2 w}{\partial x \partial y}$$



- From **stress-strain relations** for elastic material:

Thus,
$$\tau_{xy} = G\gamma_{xy}$$

$$\tau_{xy} = \tau_{yx} = -2zG \frac{\partial^2 w}{\partial x \partial y} \quad (\text{ii})$$



- The applied **twisting moments** must **equal** the **moments** set up by **shear stress**:

$$M_{xy} dy = - \int_{-t/2}^{t/2} z \tau_{xy} dy dz$$

$$\text{and } M_{yx} dx = \int_{-t/2}^{t/2} z \tau_{yx} dx dz$$

or

$$M_{xy} = - \int_{-t/2}^{t/2} z \tau_{xy} dz$$

$$\text{and } M_{yx} = \int_{-t/2}^{t/2} z \tau_{yx} dz$$



$$\tau_{xy} = \tau_{yx} = -2zG \frac{\partial^2 w}{\partial x \partial y} \quad (\text{ii})$$

$$M_{xy} = -\int_{-t/2}^{t/2} z \gamma_{xy} dz$$

$$M_{yx} = \int_{-t/2}^{t/2} z \gamma_{yx} dz$$

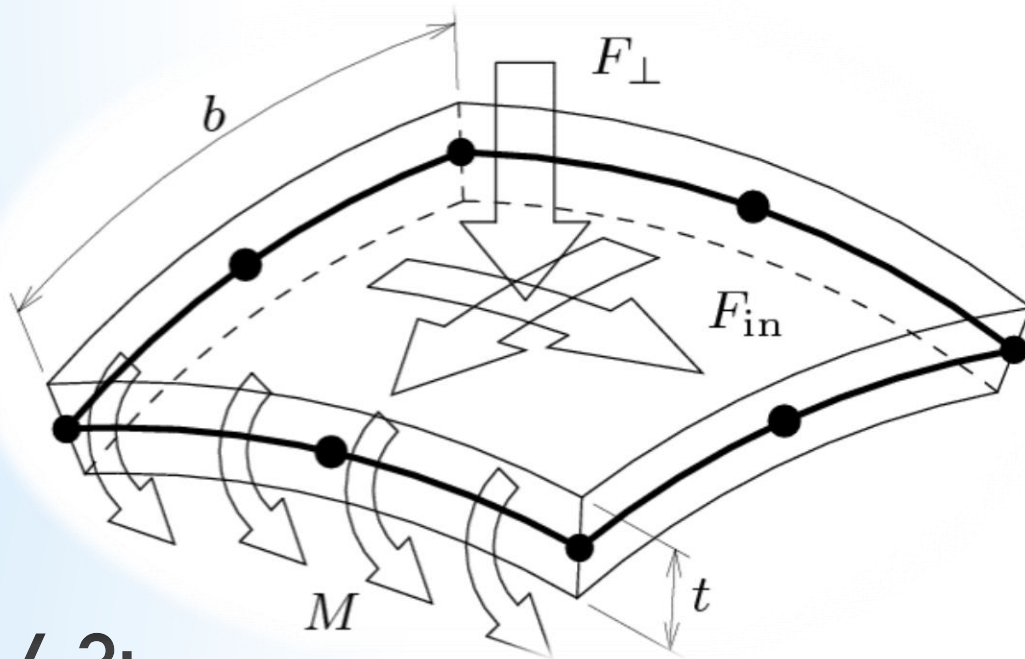
- Substituting **equation (ii)** into the foregoing equations and integrating give

$$M_{xy} = -M_{yx} = \frac{Gt^3}{6} \frac{\partial^2 w}{\partial x \partial y} \quad (\text{iii})$$

- Putting $G = \frac{E}{2(1+\nu)}$ and $D = \frac{Et^3}{12(1-\nu^2)}$ into (iii)

$$M_{xy} = -M_{yx} = D(1-\nu) \frac{\partial^2 w}{\partial x \partial y}$$





*Topic 4.3:

Plate Governing Equilibrium Equation



- In **beams** the **lateral load** q per unit length is related to the resulting deflection w through

$$EI \frac{\partial^4 w}{\partial x^4} = q$$

- A **similar relation** between pressure p on a plate and the **resulting deflection** $w(x,y)$ can be established.

Expected outcome (**Important !!**)

$$v(x) \text{ or } w(x) = \text{deflection}(m)$$

$$\theta(x) = \frac{dv(x)}{dx} = v'(x) = \text{slope}(rad)$$

$$M(x) = EI \frac{d^2v(x)}{dx^2} = EIv''(x) = \text{moment}(Nm)$$

$$V(x) = EI \frac{d^3v(x)}{dx^3} = EIv'''(x) = \text{shear}(N)$$

$$q(x) = EI \frac{d^4v(x)}{dx^4} = EIv''''(x) = \text{Load}(N/m)$$



Plate Equilibrium Governing Equation

- A **small element cut** from the plate is subjected to the actions shown below. These actions (moments, twists, shear forces, etc) vary from point to point, i.e. they are **functions** of **both x and y** .

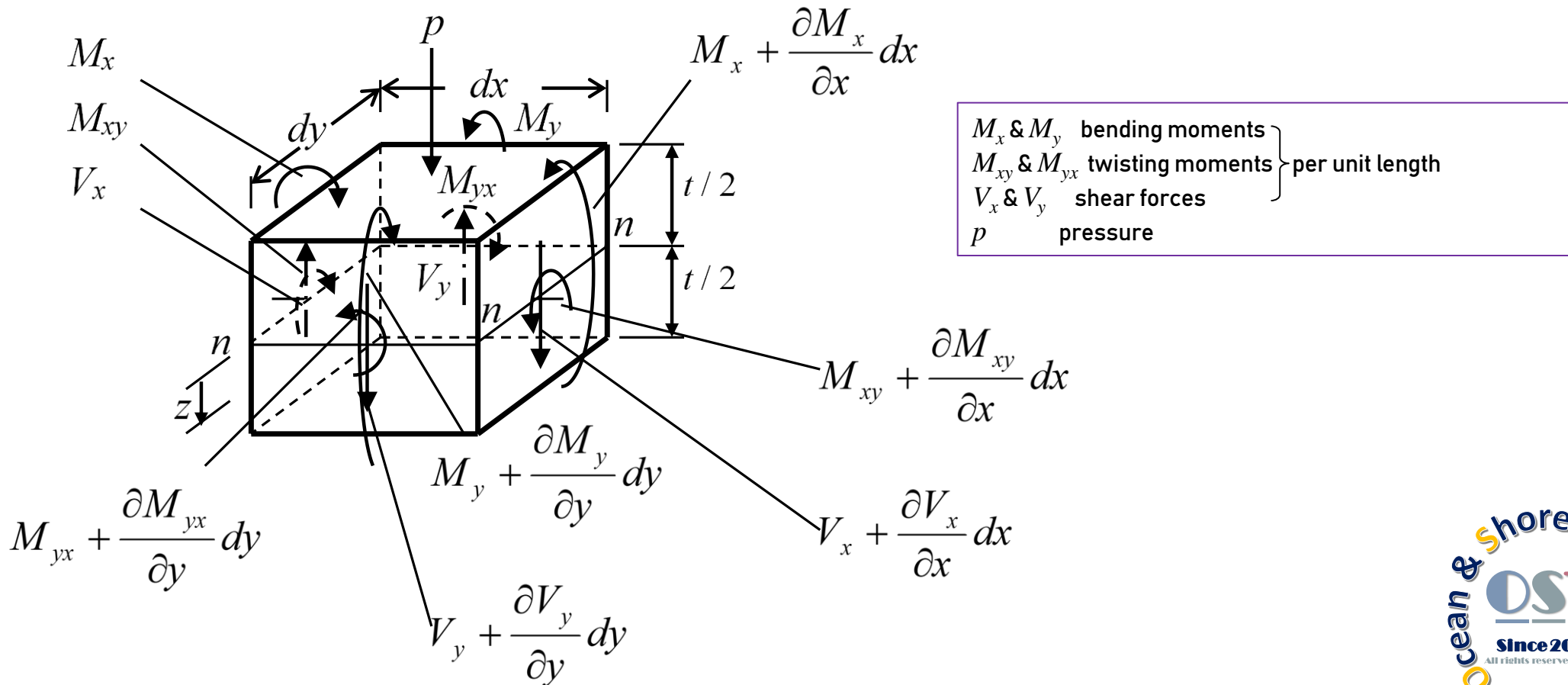


Plate Equilibrium Governing Equation

- For equilibrium

Net **force** in z direction must be zero

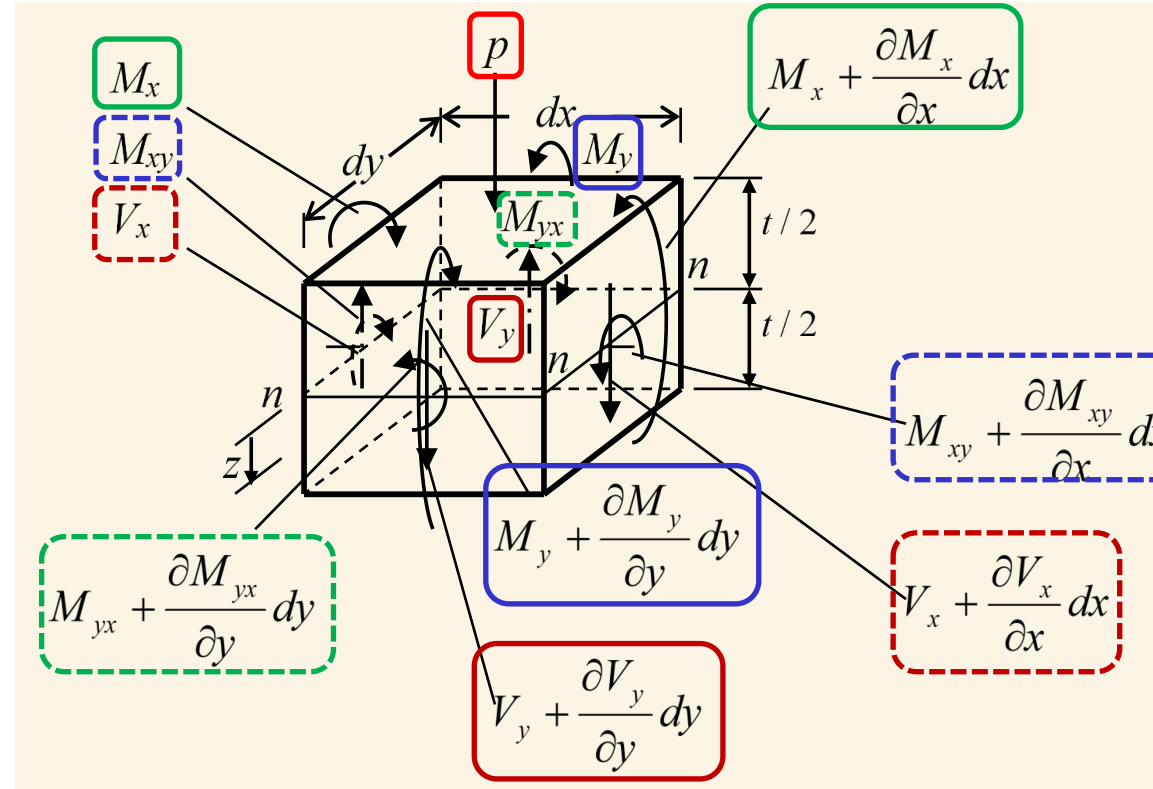
$$\frac{\partial V_x}{\partial x} + \frac{\partial V_y}{\partial y} + p = 0$$

Net **moment** about x axis must be zero

$$\frac{\partial M_{xy}}{\partial x} - \frac{\partial M_y}{\partial y} + V_y = 0$$

Net **moment** about y axis must be zero

$$\frac{\partial M_{yx}}{\partial y} + \frac{\partial M_x}{\partial x} - V_x = 0$$



M_x & M_y bending moments
 M_{xy} & M_{yx} twisting moments per unit length
 V_x & V_y shear forces
 p pressure



Plate Equilibrium Governing Equation

$$\frac{\partial V_x}{\partial x} + \frac{\partial V_y}{\partial y} + p = 0 \quad \frac{\partial M_{xy}}{\partial x} - \frac{\partial M_y}{\partial y} + V_y = 0 \quad \frac{\partial M_{yx}}{\partial y} + \frac{\partial M_x}{\partial x} - V_x = 0$$

Eliminating V_x and V_y gives

$$\frac{\partial^2 M_x}{\partial x^2} + \frac{\partial^2 M_y}{\partial y^2} + \frac{\partial^2 M_{yx}}{\partial x \partial y} - \frac{\partial^2 M_{xy}}{\partial x \partial y} + p = 0$$

Substituting the **moment-curvature relations** into this equation gets

$$\frac{\partial^4 w}{\partial x^4} + 2 \frac{\partial^4 w}{\partial x^2 \partial y^2} + \frac{\partial^4 w}{\partial y^4} = \frac{p}{D}$$

$$\frac{\partial^2 M}{\partial x^2} = \frac{\partial^4 w}{\partial x^4}$$

or

$$D \nabla^4 w = p$$

in which

$$\nabla^4 = \frac{\partial^4}{\partial x^4} + 2 \frac{\partial^4}{\partial x^2 \partial y^2} + \frac{\partial^4}{\partial y^4}$$



Edge conditions

- **Simple support:** $w = 0$ and $M = 0$.
- **Edge clamped:** $w = 0$ and slope = 0.
- **Elastic rotational restraint:** $w = 0$.



- The **major problems** in applying plate theory to design of ship plates, or to predicting their performance under load, is **the uncertainty** about the boundary conditions particularly the **in-plane ones**.
- With **large areas of plate** under pressure load it is **reasonable** by symmetry to assume the edges of plate elements are **clamped against rotation** at the stiffeners, and with single loaded plates the rotational conditions are probably nearer those of simple support.



- We have **investigated** the **Elastic Plate Theory**.
- Now we are able to:
 - **Appreciate** the **assumptions** of elastic plate theory
 - Be **familiar** with the **bending moment curvature relationship** of plate and relationship between twisting moment and twist in a plate.
 - **Establish** plate governing **equilibrium equation**.
 - Be **aware** of plate **edge conditions**.
- Details can be referred to **topics 4** in the lecture notes.



[Theory of Plates and Grillages]

[Part I] Plastic Design of Structures

- Plastic theory of bending (Topic 1)
- Ultimate loads on beams (Topic 2)
- Collapse of frames and grillage structures (Topic 3)

[Part II] Elastic Plate Theory

- Basic (Topic 4)
- Simply supported plates under Sinusoidal Loading (Topic 5)
- Long clamped plates (Topic 6)
- Short Clamped plates (Topic 7)
- Additional (Low aspect ratio plates, strength & permanent set)

[Part III] Buckling of Stiffened Panels

- Failure modes (Topic 8)
- Tripping (Topic 9)
- Post-buckling behaviour (Topic 10)





Kam Sa Hab Ni Da

감사합니다

Thank you!

Questions?

Aerial View of Korean Presidential Archives in Sejong city
(Construction Completed in 2014)

