

**457.643 Structural Random Vibrations**  
**In-Class Material: Class 17**

**III-2. Random Vibration Analysis of Linear Structures (contd.)**

**◎ Response of a linear system to weakly stationary input**

$$\kappa_{FF}(t_1, t_2) = \Gamma_{FF}(\tau) \text{ where } \tau = t_1 - t_2$$

Assuming zero initial conditions,

$$\begin{aligned} \kappa_{XX}(t_1, t_2) &= \int_0^{t_1} \int_0^{t_2} \kappa_{FF}(\tau_1, \tau_2) h(t_1 - \tau_1) h(t_2 - \tau_2) d\tau_2 d\tau_1 \\ &= \int_0^{t_1} \int_0^{t_2} \Gamma_{FF}(\tau) h(t_1 - \tau_1) h(t_2 - \tau_2) d\tau_2 d\tau_1 \end{aligned}$$

where  $\tau = \tau_1 - \tau_2$

Note  $\Gamma_{FF}(\tau) = \int_{-\infty}^{\infty} \Phi_{FF}(\omega) e^{i\omega\tau} d\omega$

Thus,

$$\begin{aligned} \kappa_{XX}(t_1, t_2) &= \int_{-\infty}^{\infty} \int_0^{t_1} \int_0^{t_2} \Phi_{FF}(\omega) h(t_1 - \tau_1) h(t_2 - \tau_2) e^{i\omega\tau} d\tau_2 d\tau_1 d\omega \\ &= \int_{-\infty}^{\infty} \int_0^{t_1} \int_0^{t_2} \Phi_{FF}(\omega) h(t_1 - \tau_1) h(t_2 - \tau_2) e^{-i\omega(t_1 - \tau_1)} e^{i\omega(t_2 - \tau_2)} e^{i\omega(t_1 - t_2)} d\tau_2 d\tau_1 d\omega \end{aligned}$$

By changing variable  $u = t_1 - \tau_1$ , one can show

$$\begin{aligned} \int_0^{t_1} h(t_1 - \tau_1) e^{-i\omega(t_1 - \tau_1)} d\tau_1 &= \int_0^{t_1} h(u) e^{-i\omega u} du \\ &= \int_{-\infty}^{t_1} h(u) e^{-i\omega u} du \\ &= \mathcal{H}(\omega, t_1) \end{aligned}$$

This is so-called “incomplete” Fourier transform of the impulse response function.

cf. “complete” FT of IRF gives the FRF

$$H(\omega) = \int_{-\infty}^{\infty} h(t) e^{-i\omega t} dt$$

Therefore,  $\kappa_{XX}(t_1, t_2)$  for a weakly stationary input  $F(t)$  is expressed as

$$\kappa_{XX}(t_1, t_2) = \int_{-\infty}^{\infty} \Phi_{FF}(\omega) \mathcal{H}(\omega, t_1) \mathcal{H}^*(\omega, t_2) e^{i\omega\tau} d\omega$$

**Note:**

- ♦ The response of a linear system to a stationary input is \_\_\_\_\_ stationary necessarily.
- ♦ However, as  $t_1, t_2 \rightarrow \infty$ , the incomplete FTs becomes independent of  $t_1$  and  $t_2$ , Therefore,  $\kappa_{XX}(t_1, t_2)$  depends only on  $\tau = t_1 - t_2$

**Observations:**

1.  $\lim_{t \rightarrow \infty} \mathcal{H}(\omega, t) = H(\omega)$  for a “stable” system

Therefore, the response of a linear system to a stationary input becomes \_\_\_\_\_ e \_\_\_\_\_

2.  $\kappa_{XX}(0,0)$  must be \_\_\_\_\_ and it means  $\sigma_X^2(0) =$  \_\_\_\_\_. This makes sense because we assumed \_\_\_\_\_ IC's
3. For the stationary response, i.e.  $t_1, t_2 \rightarrow \infty$

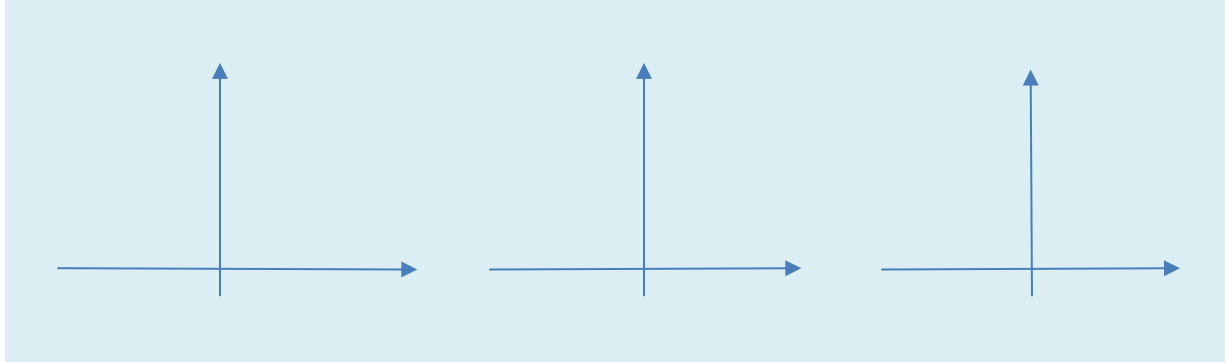
$$\begin{aligned} \kappa_{XX}(t_1, t_2) &= \int_{-\infty}^{\infty} \Phi_{FF}(\omega) \mathcal{H}(\omega, t_1) \mathcal{H}^*(\omega, t_2) e^{i\omega\tau} d\omega \\ &= \int_{-\infty}^{\infty} \Phi_{FF}(\omega) |H(\omega)|^2 e^{i\omega\tau} d\omega \end{aligned}$$

That is,

$$\Gamma_{XX}(\tau) = \int_{-\infty}^{\infty} \Phi_{FF}(\omega) |H(\omega)|^2 e^{i\omega\tau} d\omega$$

4. From this result, for a stationary response, it is found that

$$\Phi_{XX}(\omega) =$$



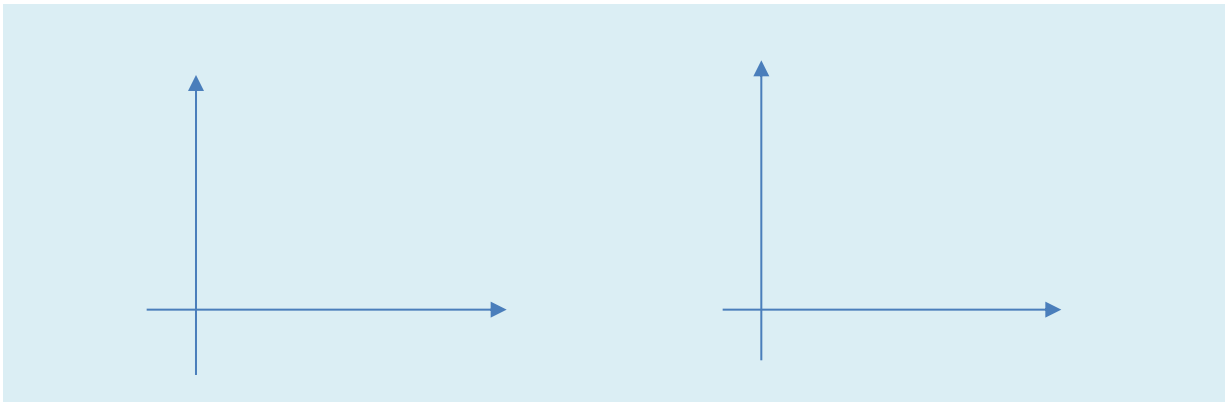
For example, let us consider...

◎  $\mathcal{H}(\omega, t)$  and  $H(\omega)$  of standard SDOF oscillator

Recall

- ♦  $h(t) = \frac{1}{\omega_D} e^{-\xi\omega_0 t} \sin \omega_D t$
- ♦  $H(\omega) = \frac{1}{\omega_0^2 - \omega^2 + 2i\xi\omega_0\omega} = \frac{\omega_0^2 - \omega^2 - 2i\xi\omega_0\omega}{(\omega_0^2 - \omega^2)^2 + (2\xi\omega_0\omega)^2}$

$$|H(\omega)|^2 = \frac{1}{(\omega_0^2 - \omega^2)^2 + 4\xi^2\omega_0^2\omega^2}$$



$$\begin{aligned} \mathcal{H}(\omega, t) &= \int_{-\infty}^t \frac{1}{\omega_D} e^{-\xi\omega_0\tau} \sin \omega_D \tau \cdot U(\tau) \cdot e^{-i\omega\tau} d\tau \\ &= H(\omega) \left[ 1 - \left( \cos \omega_D t + \frac{\xi\omega_0 + i\omega}{\omega_D} \sin \omega_D t \right) \cdot e^{-\xi\omega_0 t} \cdot e^{-i\omega t} \right] \end{aligned}$$

From this result, the terms in ( ) has the same order as 1, and  $e^{-i\omega t}$  oscillates. Therefore, the rate of the convergence of the terms in [ ] to \_\_\_\_\_ is determined by \_\_\_\_\_



In other words, “sufficient” time to achieve stationarity depends on  $\xi \omega_0 t = \xi \frac{2\pi}{T_0} t$

Suppose we set  $\xi \omega_0 t = 2\pi \xi \frac{t}{T_0} = \pi$  (note  $e^{-\pi} = 4\%$ ) and solve it for  $t$ , i.e. time to make the exponentially decaying term as 4%,  $t_{4\%} = \frac{T_0}{2\xi}$

e.g.  $\xi = 0.1 \rightarrow t_{4\%} \cong 5T_0$ ,  $\xi = 0.05 \rightarrow t_{4\%} \cong 10T_0$

※ Alternative (empirical) method:

Wang, Z., and Song, J. (2017) Equivalent linearization method using Gaussian mixture (GM-ELM) for nonlinear random vibration analysis, *Structural Safety*, Vol. 64, 9-19. (<http://dx.doi.org/10.1016/j.strusafe.2016.08.005>)

### 3.1.1. Remark 1: Selecting sample points

One issue in selecting sample points in the aforementioned algorithm is that the nonlinear response takes a certain amount of time to achieve stationarity, thus using the whole time series including a nonstationary part will introduce errors to the estimated PDF. To reduce this error, for each of the  $M$  response histories obtained from the first step of the algorithm, we need to select  $\bar{N}$  stationary response values as the sample points.

Here we provide a method to crudely estimate the time that the system would take to achieve stationarity. To begin with, the standard deviation of the response at a sequence of time points, denoted as  $\text{std}[Z(j\Delta t)]$ , in which  $j = 1, 2, \dots$  and  $\Delta t$  is the time step of the nonlinear analysis, is estimated using the recorded  $M$  response histories, and then a sigmoid function expressed as

$$f_{\text{fit}}(j) = \frac{1}{1 + e^{-aj\Delta t + b}} \quad (12)$$

is employed to fit the  $\text{std}[Z(j\Delta t)]$  curve. Note that  $f_{\text{fit}}(\cdot) \in (0, 1)$ , thus the  $\text{std}[Z(j\Delta t)]$  curve should be scaled by a factor  $J / \sum_{j=1}^J \text{std}[Z(j\Delta t)]$  ( $J\Delta t$  is the duration of the excitation) so that it approximately ranges from 0 to 1. The parameters  $a$  and  $b$  in Eq. (12) can be determined from a least-square regression analysis. A typical scaled  $\text{std}[Z(j\Delta t)]$  curve and its corresponding fitting function  $f_{\text{fit}}(\cdot)$  is illustrated in Fig. 3. With  $f_{\text{fit}}(t)$  available, the time the system takes to achieve stationarity, denoted by  $j_{ns}\Delta t$ , can be estimated via

$$j_{ns} = \text{argmin}\{j | 1 - f(j\Delta t) \leq \text{ToI}, j = 1, 2, \dots\} \quad (13)$$

where  $\text{ToI}$  denotes a specified tolerance. With  $j_{ns}$  determined, for each of the  $M$  response histories,  $\bar{N} = J - j_{ns}$  time points corresponding to the stationary responses are selected to be the sample points, and the total number of sample points is  $N = M \cdot \bar{N} = M \cdot (J - j_{ns})$ .

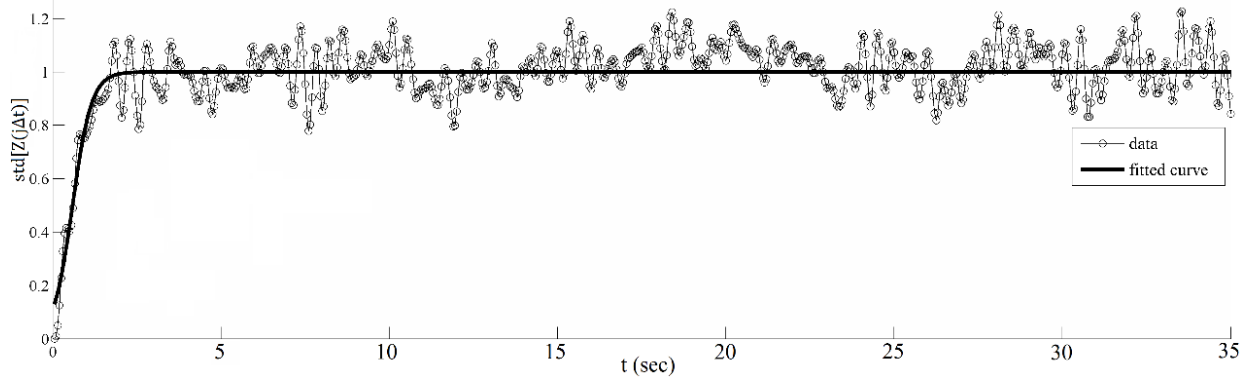


Figure 3. A typical scaled  $\text{std}[Z(j\Delta t)]$  curve and the fitting function

◎ Stationary response of standard SDOF oscillator to “white noise”

Useful for linear random vibration analysis of MDOF systems using modal combination, i.e. each mode is represented by a standard SDOF oscillator (will be shown later)

$$\Phi_{FF}(\omega) = \Phi_0$$

PSD of the stationary response

$$\begin{aligned} \Phi_{XX}(\omega) &= \Phi_0 |H(\omega)|^2 \\ &= \frac{\Phi_0}{(\omega_0^2 - \omega^2)^2 + 4\xi^2 \omega_0^2 \omega^2} \end{aligned}$$

Thus,

$$\begin{aligned} \Gamma_{XX}(\tau) &= \int_{-\infty}^{\infty} \Phi_{XX}(\omega) e^{i\omega\tau} d\omega \\ &= \Phi_0 \int_{-\infty}^{\infty} \frac{e^{i\omega\tau}}{(\omega_0^2 - \omega^2)^2 + 4\xi^2 \omega_0^2 \omega^2} d\omega \end{aligned}$$

How? We can use \_\_\_\_\_ theorem

(to be continued...)

