Chapter 6

The Second Law of Thermodynamics

Advanced Thermodynamics Min Soo Kim



6.1 Introduction to the Second Law of Thermodynamics





6.1 Introduction to the Second Law of Thermodynamics

- Is there any way in which we can write the first law in terms of state variables only? → The Second Law of thermodynamics
- Is there some state variable by which we can distinguish between a reversible and an irreversible process?

 \rightarrow The Second Law of thermodynamics

Most general form (for closed system) is,

 $dU = \delta Q - \delta W \quad (eq. 6.1)$

(Neither δQ or δW is an exact differential)



 $\delta W_r = P dV$ (*V* is a state variable and d*V* is an exact differential)

 $\frac{\delta W_r}{P} = dV \qquad (eq. 6.2) \qquad (\frac{1}{p} \text{ is integrating factor })$ $\frac{\delta Q_r}{T} \equiv dS \qquad (eq. 6.3) \qquad (Clasius definition of the entropy S)$

Subtituting eq(6.2) & eq(6.3) in eq(6.1),

dU = TdS - PdV



 Clausius statement : It is impossible to construct a device that operates in a cycle and whose sole effect is to transfer heat from a cooler body to a hotter body



 $\rightarrow If \ T_2 > T_1$ then $Q_2 = Q_1$, with $W = 0.\,is$ impossible

Figure 6.2 Schematic diagram of a device forbidden by the Clausius statement of the second law.



 Kelvin-Planck statement : It is impossible to construct a device that operates in a cycle and produces no other effect than the performance of work and the exchange of heat with a single reservoir.

 \rightarrow It is impossible to have $W = Q_2$



Figure 6.3 Schematic diagram of a device forbidden by the Kelvin-Planck statement of the second law.



6.4 Carnot's Theorem



T $Q_2 - Q_1$ M W C.V.

Figure 6.4 A composition engine in violation of the Clausius statement

Work generation from?

Figure 6.5 The equivalent engine in violation of the Kelvin-Planck statement

Heat is transported from T to where?

Applying Carnot's theorem to both statement, it is impossible to make engine which goes against the statements.



For Carnot cycle,

$$\frac{Q_2}{T_2} + \frac{Q_1}{T_1} = 0$$

$$\frac{\partial Q_2}{T_2} + \frac{\partial Q_1}{T_1} = 0$$

$$\sum \frac{\delta Q_t}{T_t} \to \oint \frac{\delta Q_r}{T} = 0$$



Figure 6.6 Schematic diagram of Carnot's cycle



For irreversible cycle,

$$\frac{Q_1'}{Q_2'} < \frac{Q_1}{Q_2} = -\frac{T_1}{T_2} \qquad \rightarrow \qquad \frac{Q_2'}{T_2} + \frac{Q_1'}{T_1} < \mathbf{0}$$

$$\oint \frac{\delta Q_r}{T} < \mathbf{0} \rightarrow \frac{\delta Q_2'}{T_2} + \frac{\delta Q_1'}{T_1} < \mathbf{0} \rightarrow \oint \frac{\delta Q}{T} \le \mathbf{0} \rightarrow \oint \frac{\delta Q}{T} = \oint_1^2 \frac{\delta Q}{T} + \oint_2^1 \frac{\delta Q_r}{T} \le \mathbf{0}$$

$$\oint_1^2 \frac{\delta Q}{T} \le \oint_2^1 \frac{\delta Q_r}{T} = S_2 - S_1 \qquad \rightarrow \qquad dS \ge \frac{\delta Q}{T}$$

 $\Delta S \equiv S_2 - S_1 \ge 0$ (isolated system)



$\Delta S \equiv S_2 - S_1 \geq 0 \quad (isolated \ system)$

The entropy of an isolated system increases in any irreversible process and is unaltered in any reversible process. This is the principle of increasing entropy.

