Cryogenic Engineering
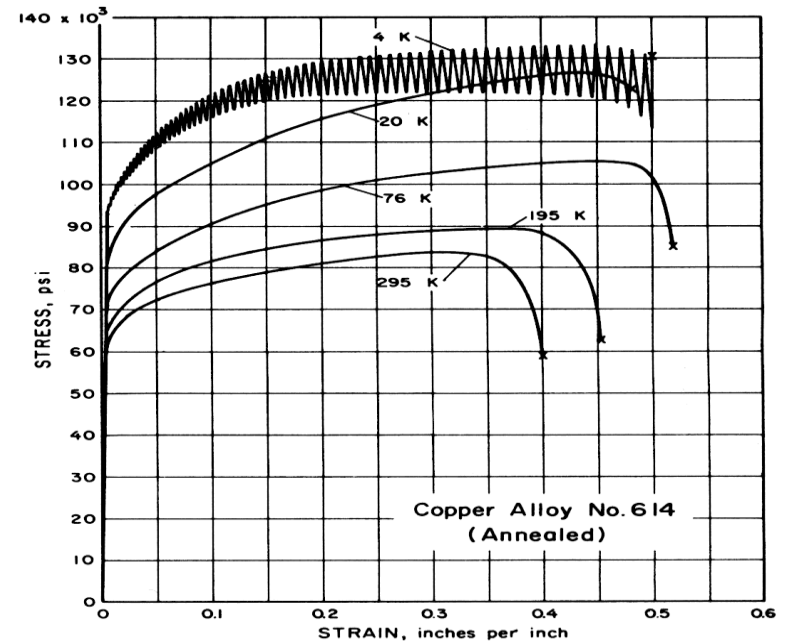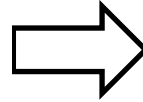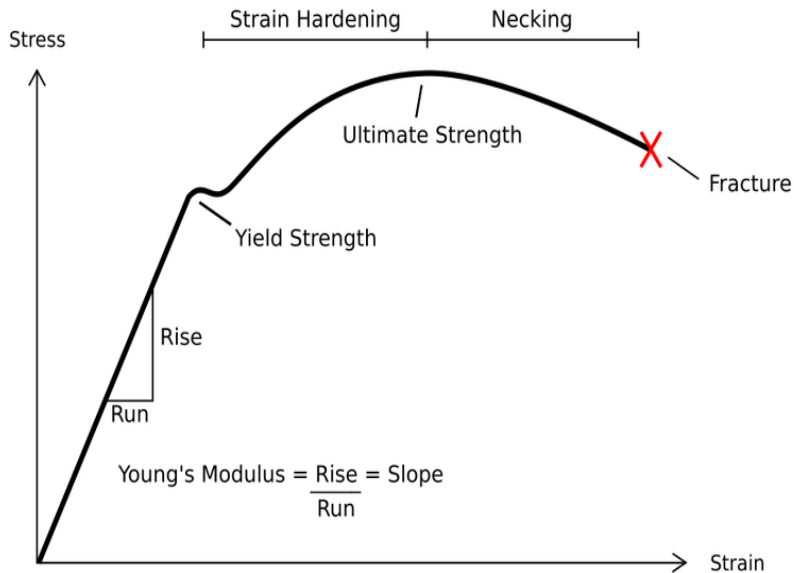
Chapter 2.

LOW TEMPERATURE PROPERTIES OF

ENGINEEING MATERIALS
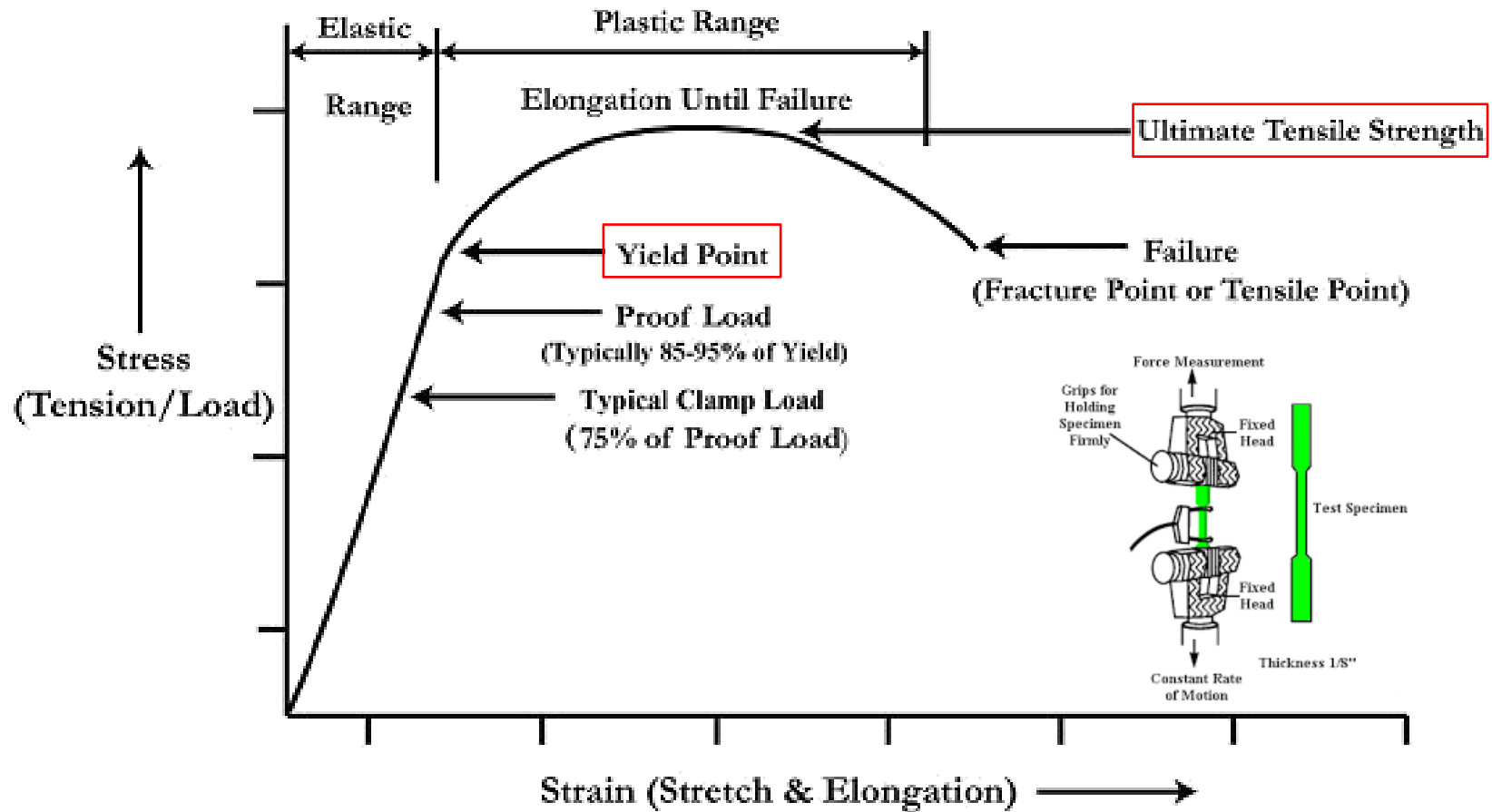
KIM, Min Soo

Stress-strain curve at room temperature

Stress-strain curve for cooper alloy at low temperature

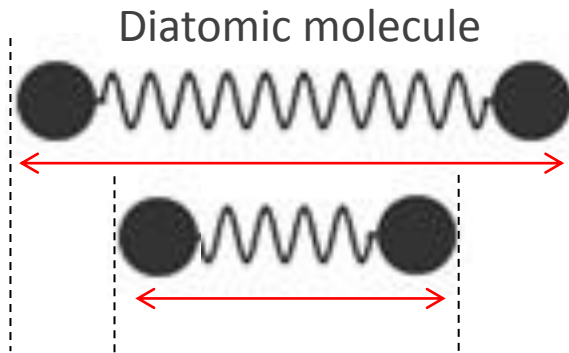Extrapolation for material properties at very low temperature is not exact !!

- Vanishing of specific heat
- Superconductivity phenomenon
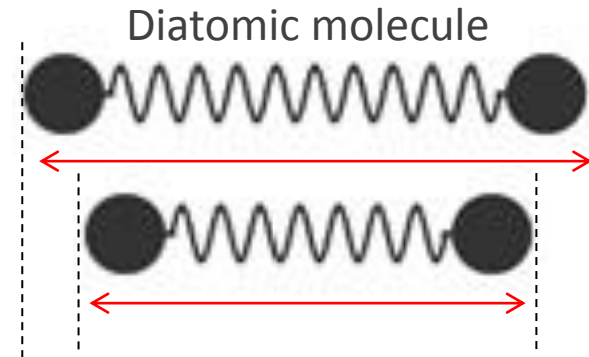- Ductile-brittle transition in carbon steel

Tensile Stress-Strain Diagram
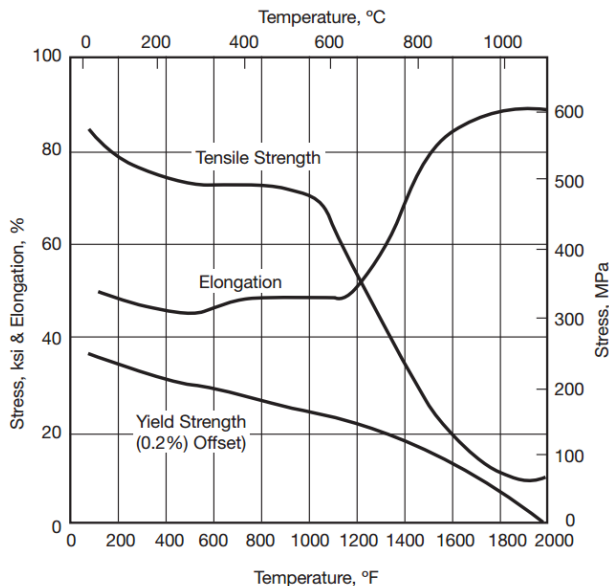
2.2 Ultimate and yield strength

Diatomic molecule

Diatomic molecule

Vigorous vibration at room temperature

Torpid vibration at low temperature

Nickel-Iron-Chromium Alloy

As temperature is lowered, atoms of the material vibrate less rigorously, a larger applied stress is required to tear dislocations from their atmosphere of alloying atoms.

The Charpy and the Izod impact tests

MaterialsScience2000 - Charpy Impact Test
https://www.youtube.com/watch?v=tpGhqQvftAo

2.3.1 Lattice Structure
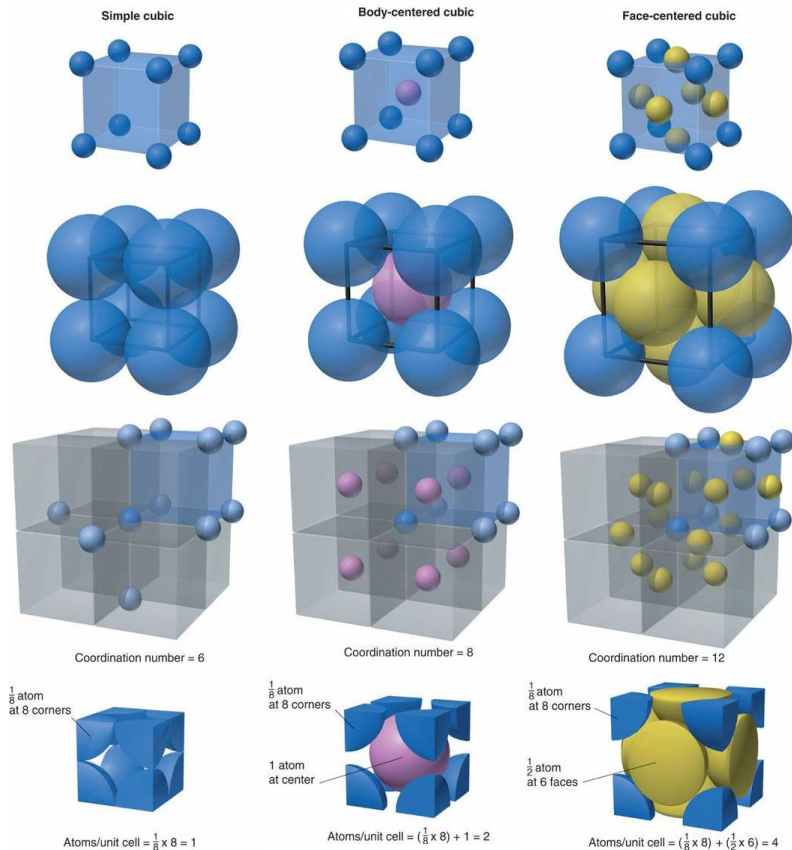
(1) 2024-T4 aluminum; (2) beryllium copper; (3) K Monel; (4) titanium;
(5) 304 stainless steel; (6) C1020 carbon steel; (7) 9 percent Ni steel (Durham et al. 1962)

<Lattice structures>
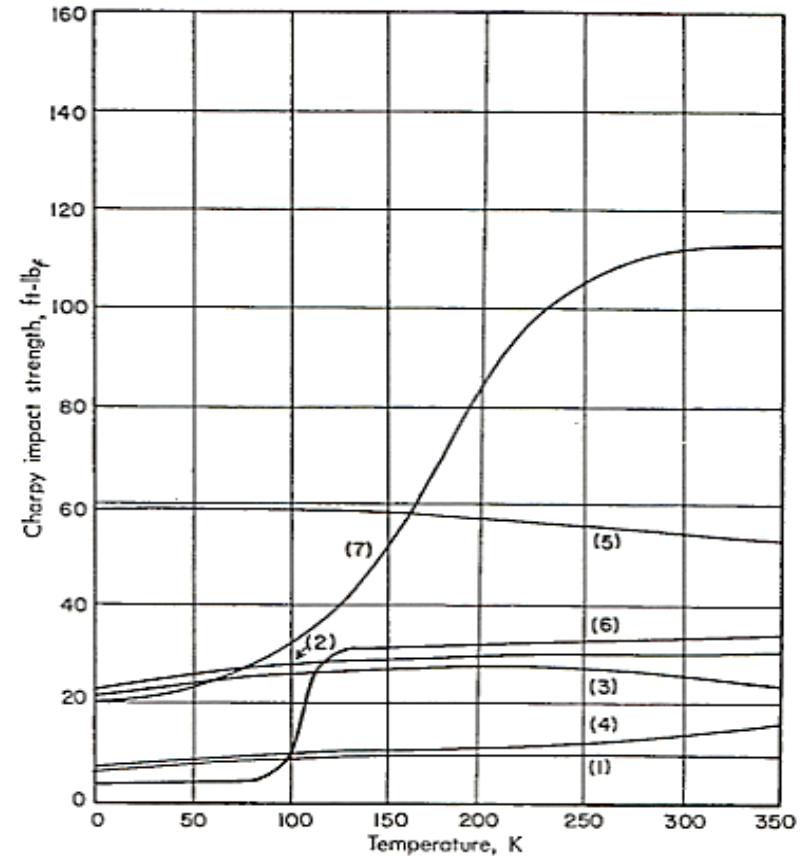
<Charpy impact strength at low temperature>
*Ductile-brittle transition

2.3.1 Lattice Structure

BCC (Body-Centered Cubic) vs FCC (Face-Centered Cubic)
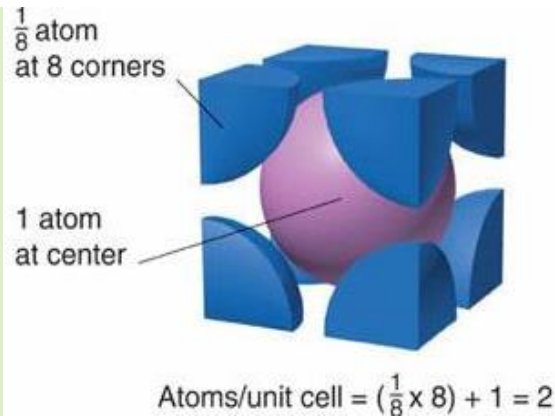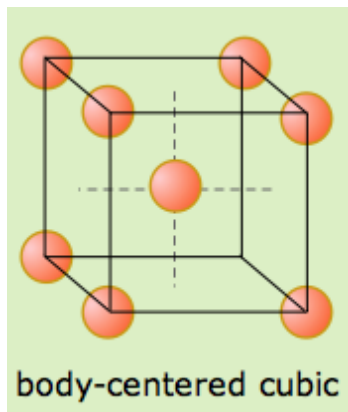
Copper-Nickel alloy
Aluminum alloy
Stainless steel ⟩ Ductile
Zirconium
Titanium

Iron
Carbon alloy
Molybdenum ⟩ Brittle
Zinc
Most plastics

2.4 Hardness and ductility

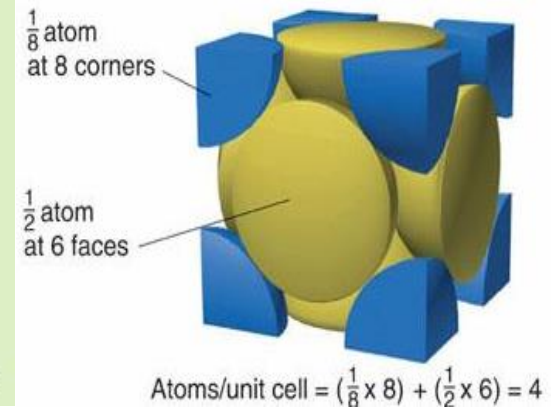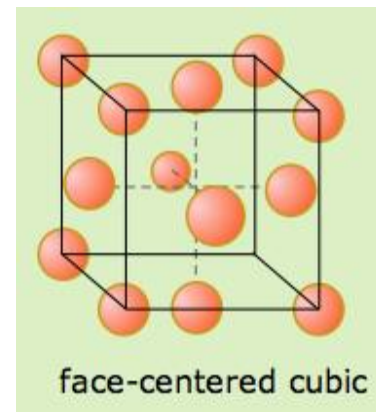**Hardness** is a measure of how resistant solid matter is to various kinds of permanent shape change when a compressive force is applied. Some materials, such as metal, are harder than others. Macroscopic hardness is generally characterized by strong intermolecular bonds, but the behavior of solid materials under force is complex; therefore, there are different measurements of hardness: **scratch hardness, indentation hardness, and rebound hardness.**
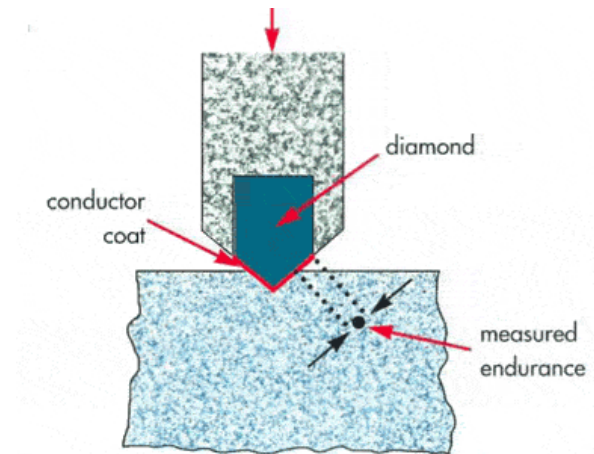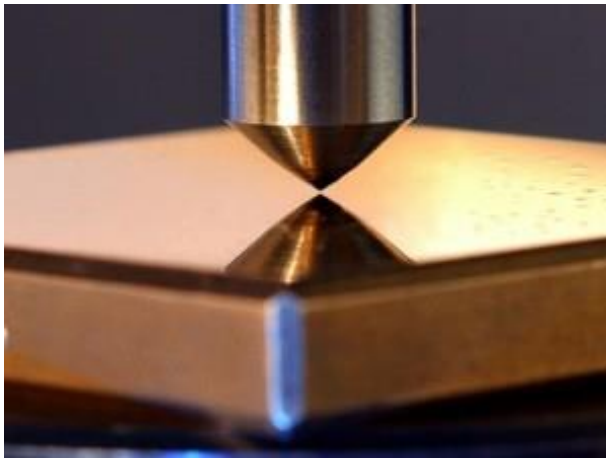
- **Scratch hardness**

  Scratch hardness tests are used to determine the hardness of a material to scratches and abrasion.

Pencil scratch hardness tester

Scratch test

▪ **Indentation hardness**

Indentation hardness tests are used in mechanical engineering to determine the hardness of a material to deformation.

Indentation hardness tester
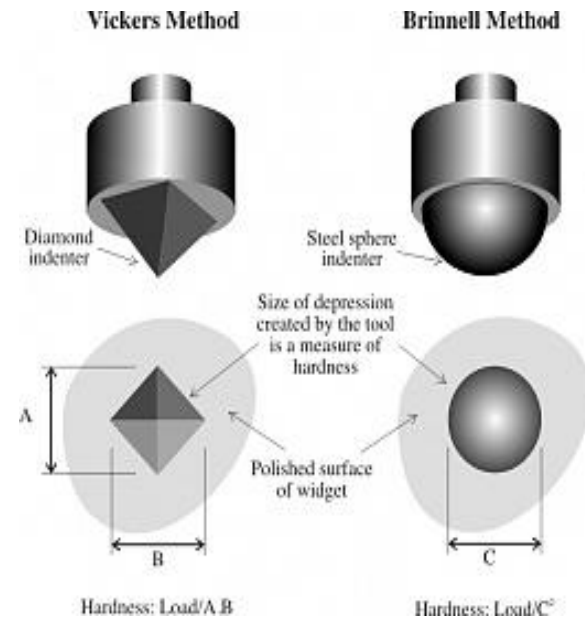
Two types of indentation

▪ **Rebound hardness**

The **Leeb rebound hardness** test is one of the four most used methods for testing metal hardness.
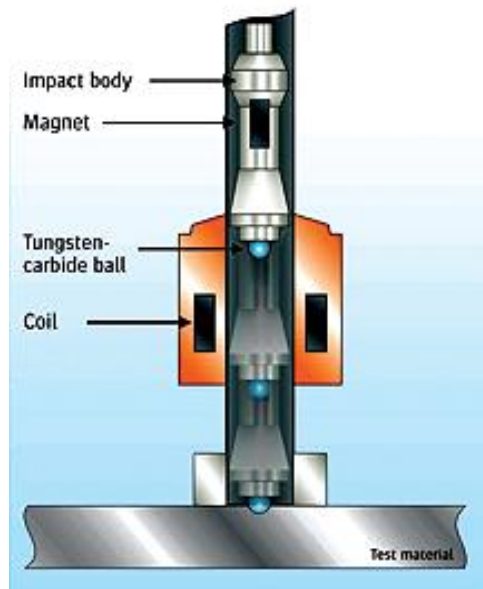
Portable rebound hardness testing machine.
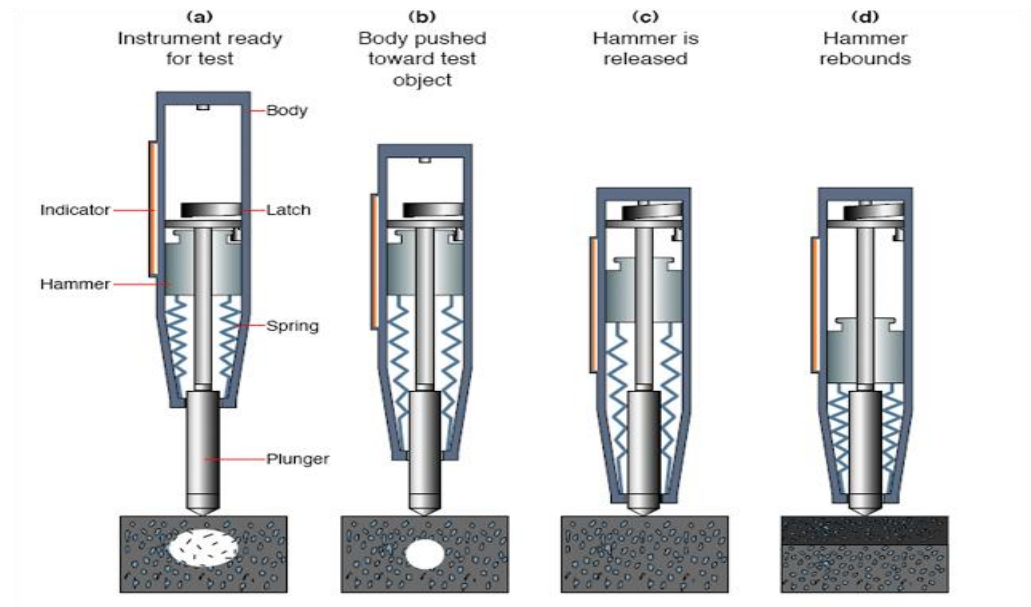
Procedure of rebound hardness test

2.4 Hardness and ductility

In materials science, **ductility** is a solid material's ability to deform under tensile stress; this is often characterized by the material's ability to be stretched into a wire. **Malleability**, a similar property, is a material's ability to deform under compressive stress; this is often characterized by the material's ability to form a thin sheet by hammering or rolling.
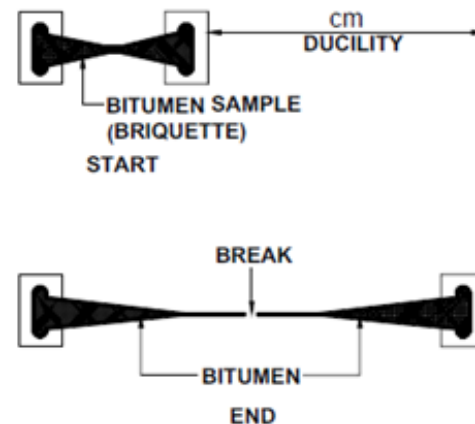
Ductility test and measurement

Tensile test of an AlMgSi alloy. The local necking and the cup and cone fracture surfaces are typical for ductile metals.

This tensile test of a nodular cast iron demonstrates low ductility.

Gold leaf can be produced owing to gold's malleability.
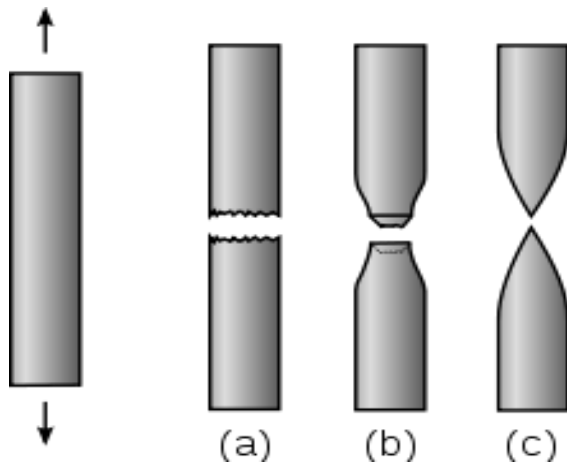
Schematic appearance of round metal bars after tensile testing.

(a) Brittle fracture
(b) Ductile fracture
(c) Completely ductile fracture

(a)   (b)   (c)

2.4 Hardness and ductility

- **Ductile–brittle transition temperature**

At low temperatures some metals that would be ductile at room temperature become brittle. This is known as a ductile to brittle transition.
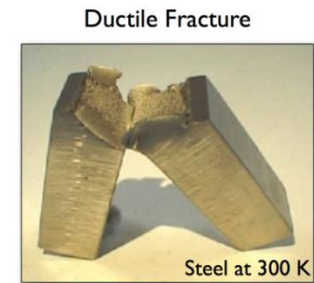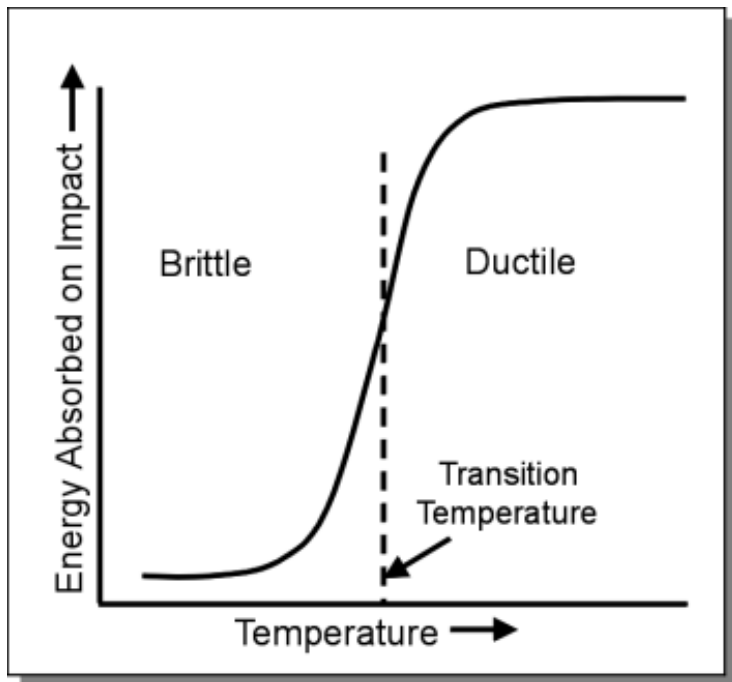
Brittle and ductile failure of steel, at low and high temperature respectively.
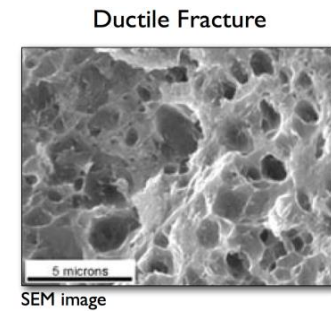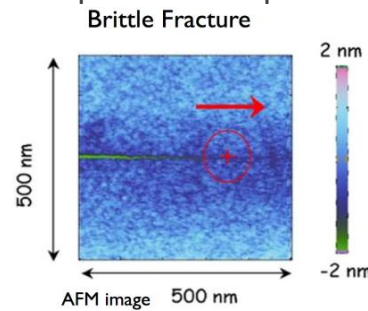
Brittle and ductile failure mechanisms, for silica and aluminium respectively. Image credits

▪ **Stress-strain diagrams for typical brittle and ductile materials**
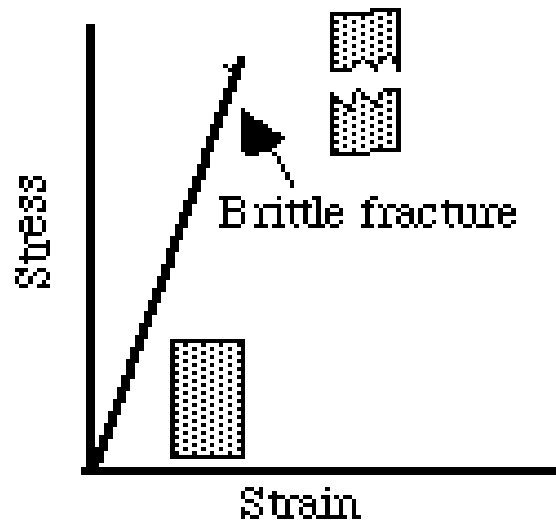
(a) Brittle materials   (b) Ductile materials
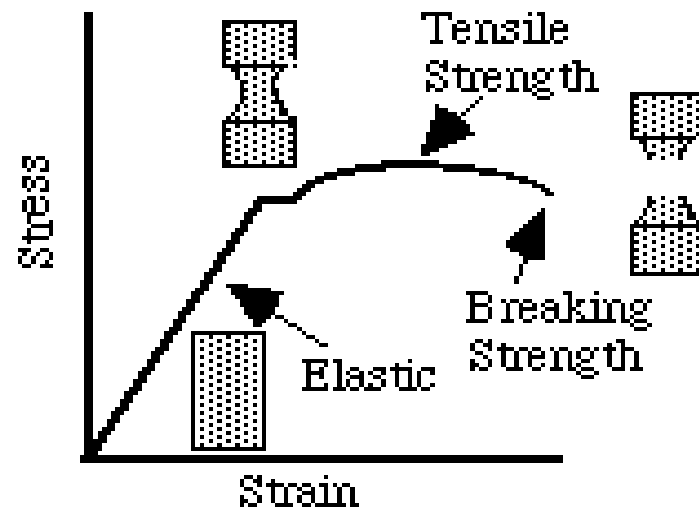
"Young's modulus" or **modulus of elasticity**, is a number that measures an object or substance's resistance to being deformed elastically (i.e., non-permanently) when a force is applied to it.

$$\text{elastic modulus} = \frac{\text{stress}}{\text{strain}}$$

Stress

Yield point
(elastic limit)

Breaking point

Slope = Elastic Modulus

Area = Work / unit volume
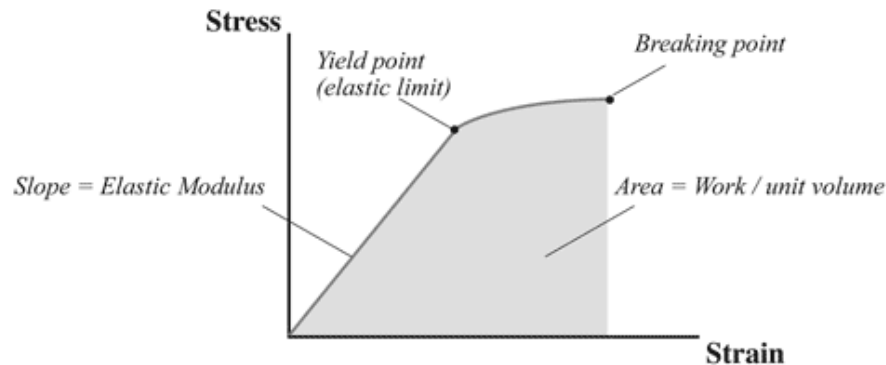
Strain

Curve of **Stress vs. Strain** which describes the mechanical deformation of a solid body.

■ **Young's modulus, E**

Young's Modulus = $\dfrac{\text{Stress}}{\text{Strain}}$ = $\dfrac{P}{\Delta L/L}$
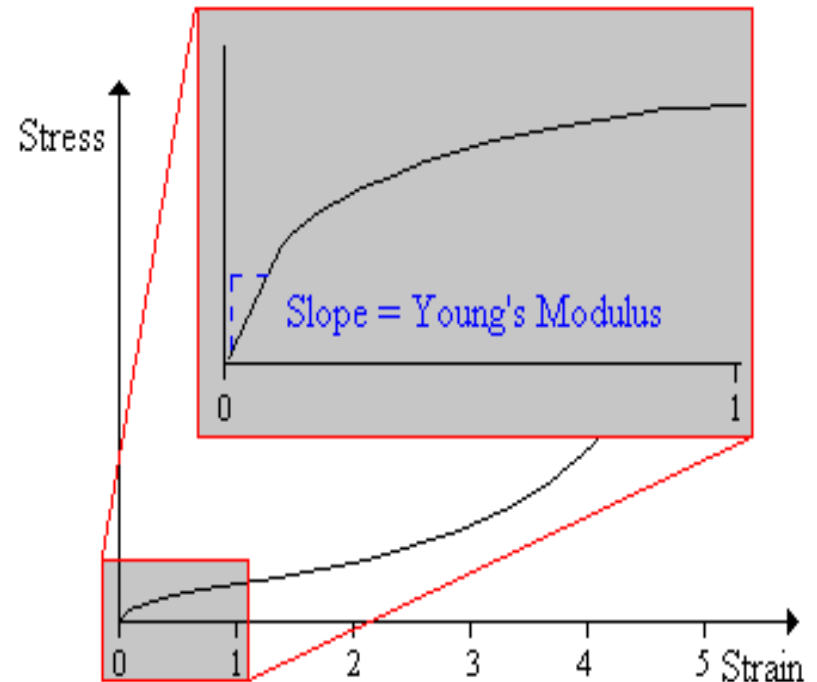
Young's modulus
(slope of a stress-strain curve)

- **Young's modulus, E**

E(0 K) = 1.386 × 10^{11} Pa
E(0°C) = 1.295 × 10^{11} Pa
G(0 K) = 0.518 × 10^{11} Pa
G(0°C) = 0.482 × 10^{11} Pa

2.5 Elastic modulus

- **Young's modulus, E**

Young's Modulus - Density Materials Selection
Chart, showing the classes of materials
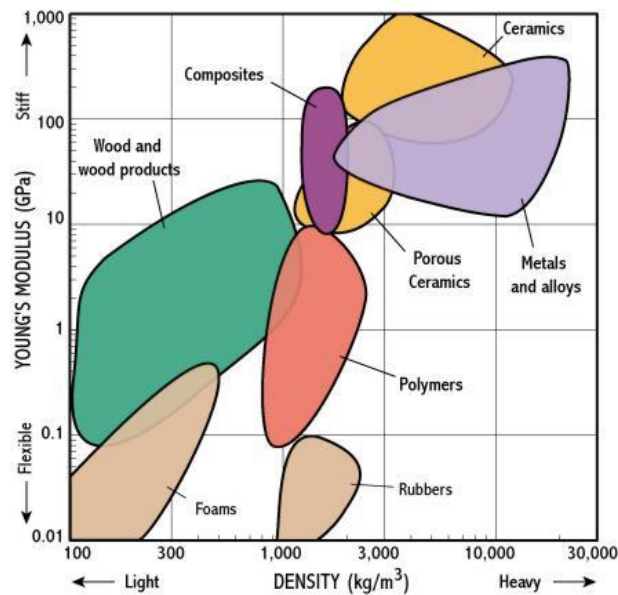
Material	Youngs Modulus /GPa
Mild Steel	210
Copper	120
Bone	18
Plastic	2
Rubber	0.02

2.5 Elastic modulus

■ **Shear modulus, G**

In materials science, shear modulus or modulus of rigidity, denoted by G, or sometimes S or μ, is defined as the ratio of shear stress to the shear strain.

$$G \stackrel{\text{def}}{=} \frac{\tau_{xy}}{\gamma_{xy}} = \frac{F/A}{\Delta x/l} = \frac{Fl}{A\Delta x}$$

where

$\tau_{xy} = F/A$ = shear stress;

F is the force which acts

A is the area on which the force acts

in engineering, $\gamma_{xy} = \Delta x/l = \tan\theta$ = shear strain. Elsewhere, $\gamma_{xy} = \theta$

Δx is the transverse displacement

l is the initial length

■ **Shear modulus, G**

Shear frame for evaluation of the in-plane
shear modulus and strength of FRP

2.5 Elastic modulus

- **Bulk modulus, B**

  The bulk modulus (K or B) of a substance measures the substance's resistance to uniform compression. It is defined as the ratio of the infinitesimal pressure increase to the resulting relative decrease of the volume.

  The bulk modulus $K > 0$ can be formally defined by the equation

  $$K = -V \frac{dP}{dV}$$

  where P is pressure, V is volume, and dP/dV denotes the derivative of pressure with respect to volume. Equivalently

  $$K = \rho \frac{dP}{d\rho}$$

  where ρ is density and dP/dρ denotes the derivative of pressure with respect to density. The inverse of the bulk modulus gives a substance's compressibility.

- **Bulk modulus, B**

$$\text{bulk modulus} = -\frac{F/A}{\Delta V/V}$$

$$\text{volume stress} = \frac{F}{A} \qquad\qquad \text{volume strain} = \frac{\Delta V}{V}$$

F = normal force
A = surface area

ΔV = change in volume
V = original volume

Volume strain produced by a volume stress
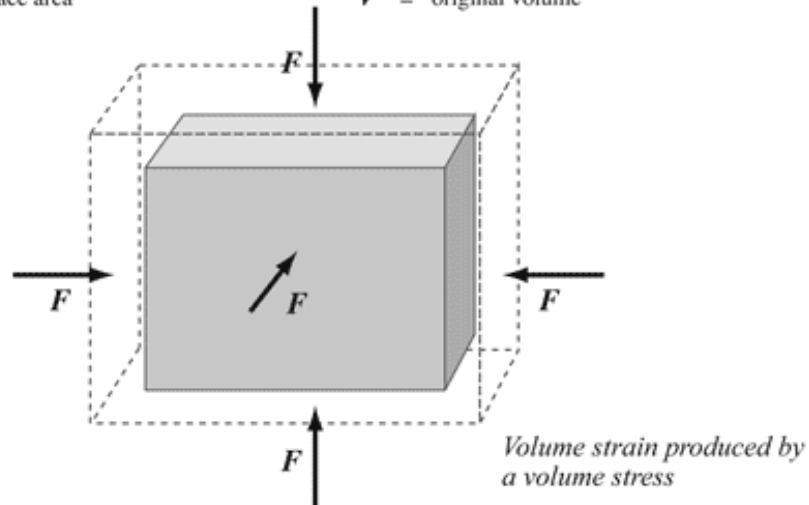
2.5 Elastic modulus

- **Relations between elastic modulus**

For homogeneous isotropic materials simple relations exist between elastic constants (Young's modulus E, shear modulus G, bulk modulus K, and Poisson's ratio ν) that allow calculating them all as long as two are known

$$E = 2G(1 + \nu) = 3K(1 - 2\nu)$$

k_t : Heat-transfer rate per unit area divided by the temperature gradient
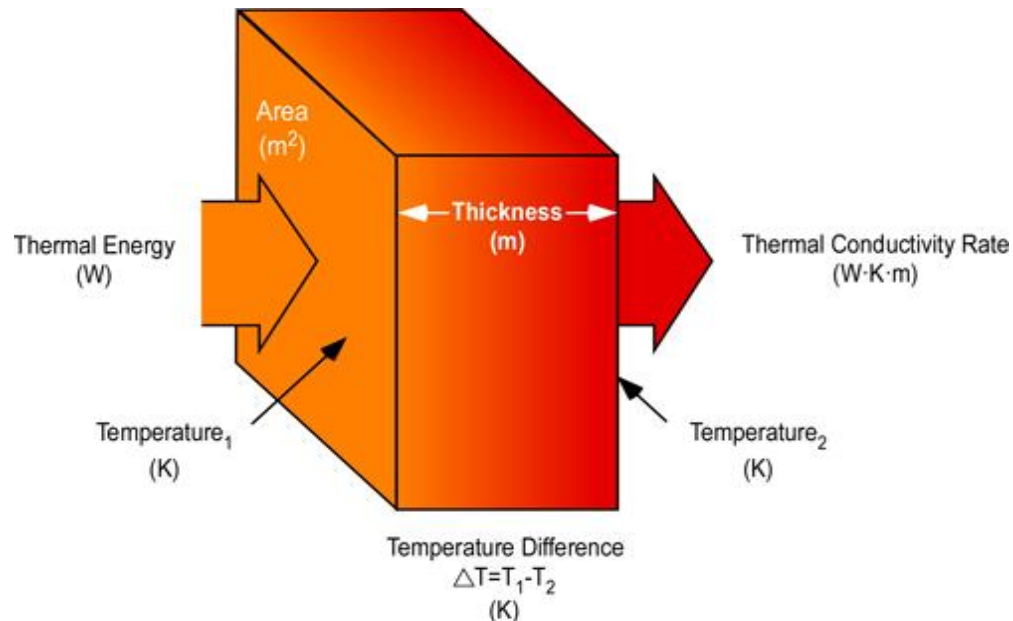
<Heat transfer through conduction>

$$q = -k_t A \frac{dT}{dx}$$

<Fourier's law>

$$k_t = \left[\frac{W}{m \cdot K} \right]$$

$$\delta Q = du + \delta W$$

<1st law of thermodynamics>

Control Volume

<A stainless spoon in hot water>
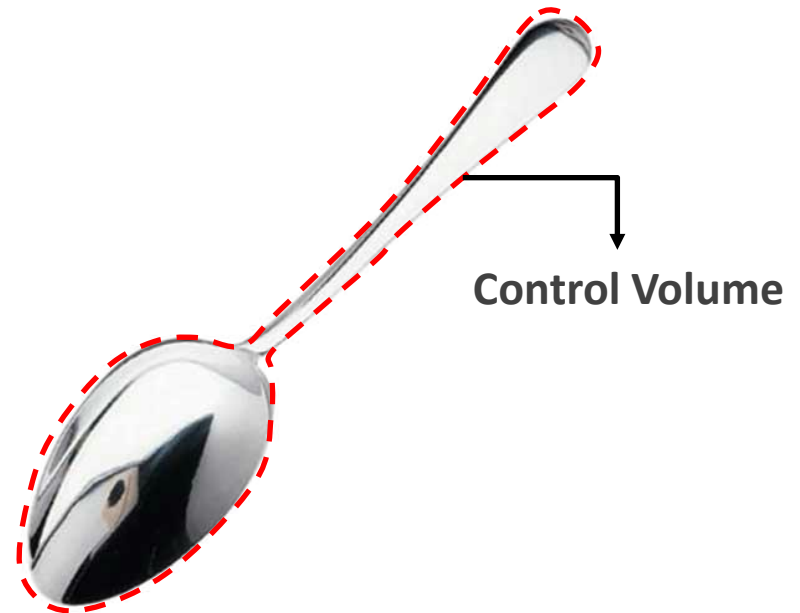
- **Internal energy (u)**

$$\delta Q = du + \cancel{\delta W}$$: 1st law of thermodynamics
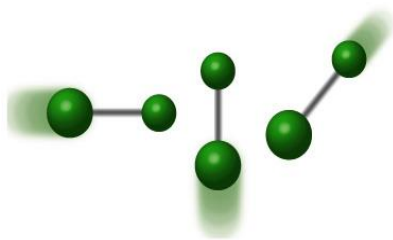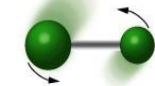
(a) Translational motion   (b) Rotational motion   (c) Vibrational motion
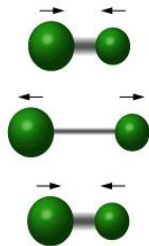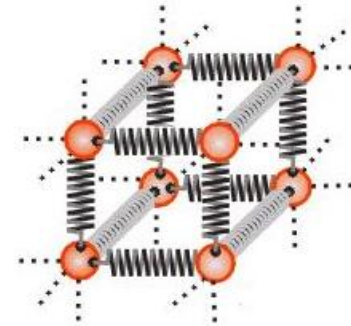
Translational energy	$\frac{1}{2}m(v_x^2 + v_y^2 + v_z^2)$	
Rotational energy	$\frac{1}{2}I(w_x^2 + w_y^2 + w_z^2)$	$\frac{1}{2}(m(\dot{x}^2+\dot{y}^2+\dot{z}^2) + k(x^2+y^2+z^2))$
Vibrational energy	$\frac{1}{2}(m\dot{x}^2 + kx^2)$	\<Internal energy for solid\>

\<Internal energy for dielectric molecule\>

2.6 Thermal conductivity

▪ **Three different mechanisms for conduction in materials**

<A>   <B>   <C>
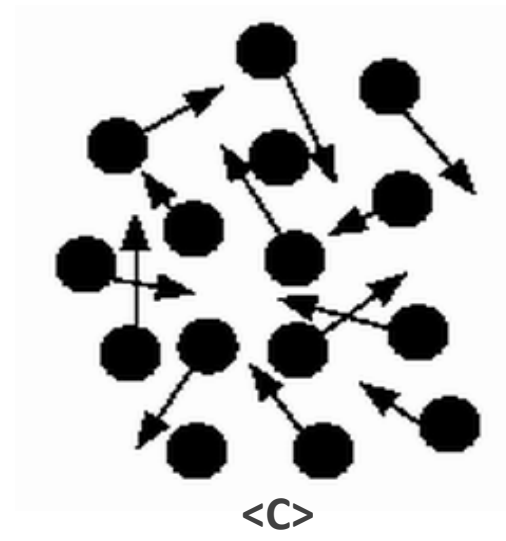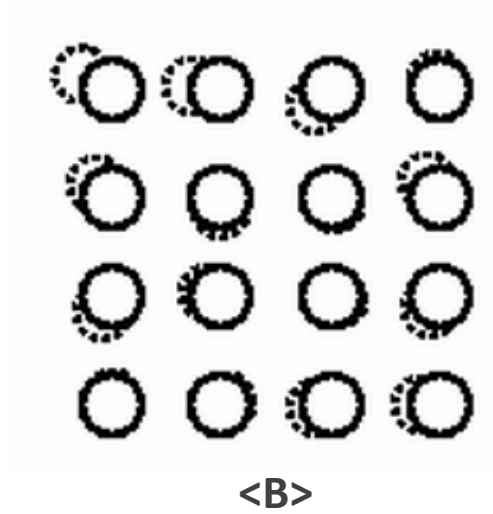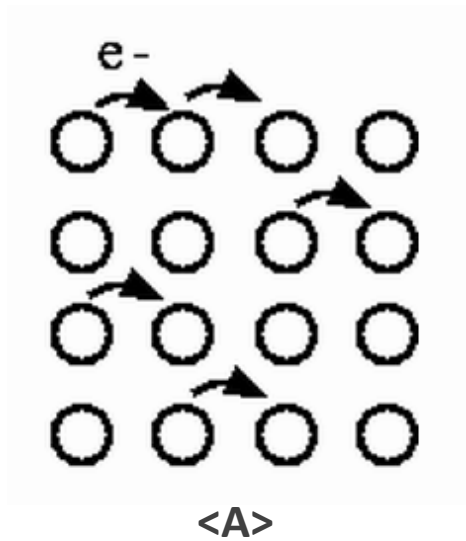
<A> : Electron motion (metallic conductor)

<B> : Lattice vibration-energy transport (phonon motion; only have energy)

<C> : Molecular motion (gases)

Experimental values of thermal conductivity

- **Gases thermal conductivity**

$$k_t = \frac{1}{8}(9\gamma - 5)\rho c_v \bar{v} \lambda \; (Eucken, 1913)$$

γ: specific heat ratio

ρ: density of material

c_v: specific heat at constant volume

$\bar{v}$: average particle velocity

λ: mean free path of particles

▪ **Gases thermal conductivity**

$$\bar{v} = \left(\frac{8 g_c R T}{\pi} \right)^2 \ (Present, 1958)$$

$g_c : 1 \ kg \cdot m/N \cdot s^2$

$R: \ R_u/M, \ R_u = 8.31434 \frac{J}{mol} \cdot K$

$M: \ \text{molecular weight of the gas}$

$T: \text{absolute temperature of the gas}$

Gases thermal conductivity
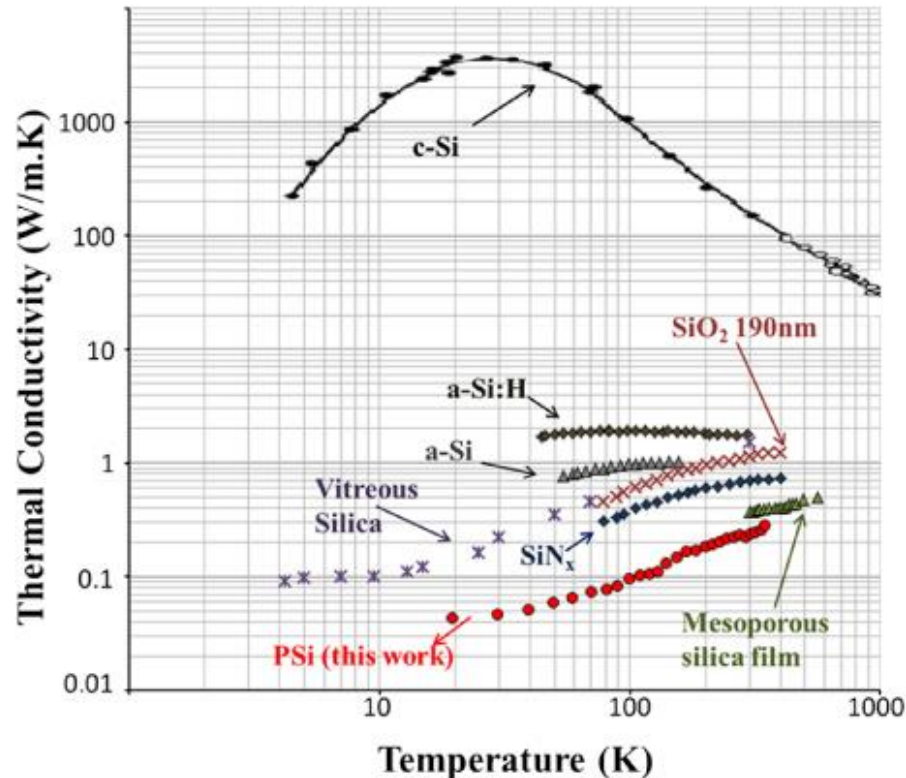
Comparison of the thermal conductivity of porous Si (this work) with that of bulk crystalline silicon (Glassbrenner and Slack 1964), a-Si (Lee et al 1991), a-Si : H (Cahill et al 1989) and different other C-MOS compatible films (mesoporous silica (Shin et al 2008), vitreous silica (Smith et al 1978) and silicon nitride (Lee and Cahill 1997)).

▪ **Solids thermal conductivity**

$$k_t = \frac{1}{3} \rho c_v \bar{v} \lambda$$

ρ: density of material

c_v: specific heat at constant volume

$\bar{v}$: average particle velocity

λ: mean free path of particles

2.6 Thermal conductivity

- **Solids thermal conductivity**
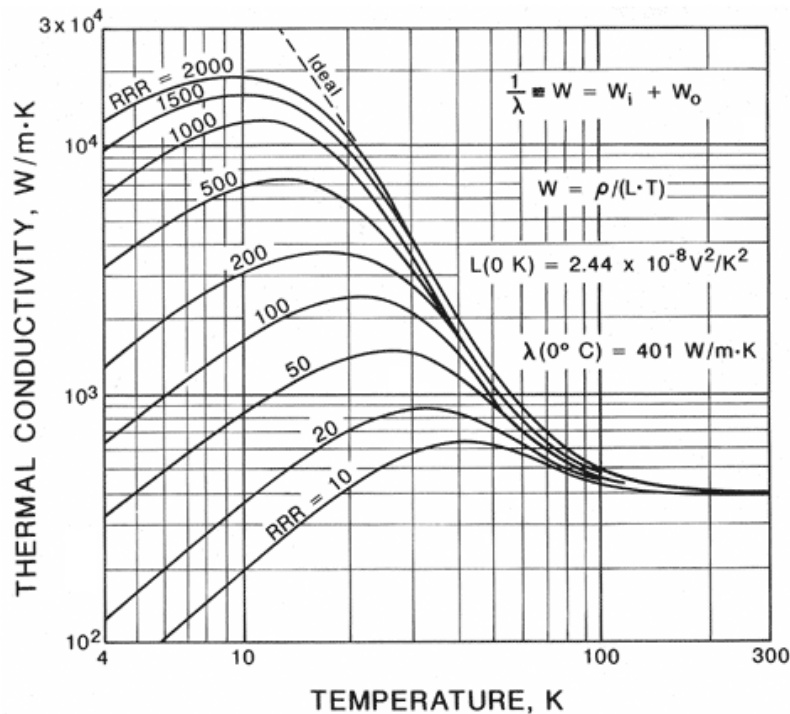
Thermal conductivity of copper alloys

$$\frac{1}{\lambda} \equiv W = W_i + W_o$$

$$W = \rho/(L \cdot T)$$

$$L(0\ K) = 2.44 \times 10^{-8} V^2/K^2$$

$$\lambda(0°\ C) = 401\ W/m \cdot K$$
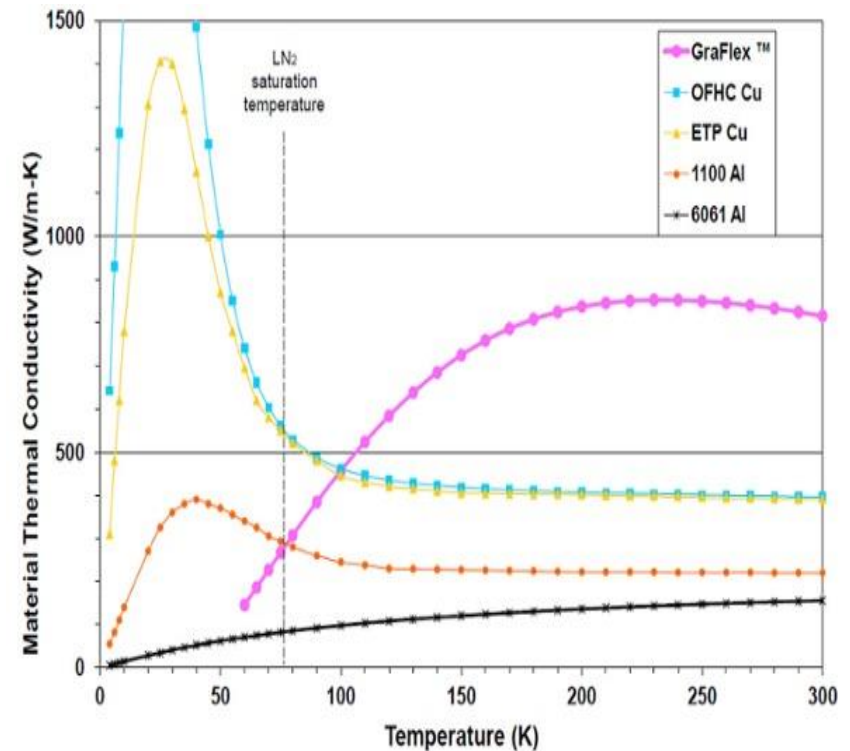
Graphite Thermal Conductivity vs Al and Cu Straps

■ **Specific heat**

The energy required to change the temperature of

the substance by one degree!

- C_p When constant pressure
- C_v When constant volume

Debye specific heat function

■ **Debye model**

For solids, "Debye model" represents how the specific heats change under the temperature variation

$$C_v = \frac{9RT^3}{\theta_D^3} \int_0^{\theta_D/T} \frac{x^4 e^x \, dx}{(e^x - 1)^2} = 3\text{R} \left(\frac{T}{\theta_D}\right)^3 D(\frac{T}{\theta_D})$$

θ_D : Debye Characteristic temperature

$D(\text{T}/\theta_D)$: Debye function

where, $\theta_D = \frac{h v_a}{k} \left(\frac{3N}{4\pi V}\right)^{1/3}$

h : Planck's constant

v_a : Speed of sound in the solid

k : Boltzmann's constant

N/V : Number of atoms per unit volume for the solid

Debye specific heat function
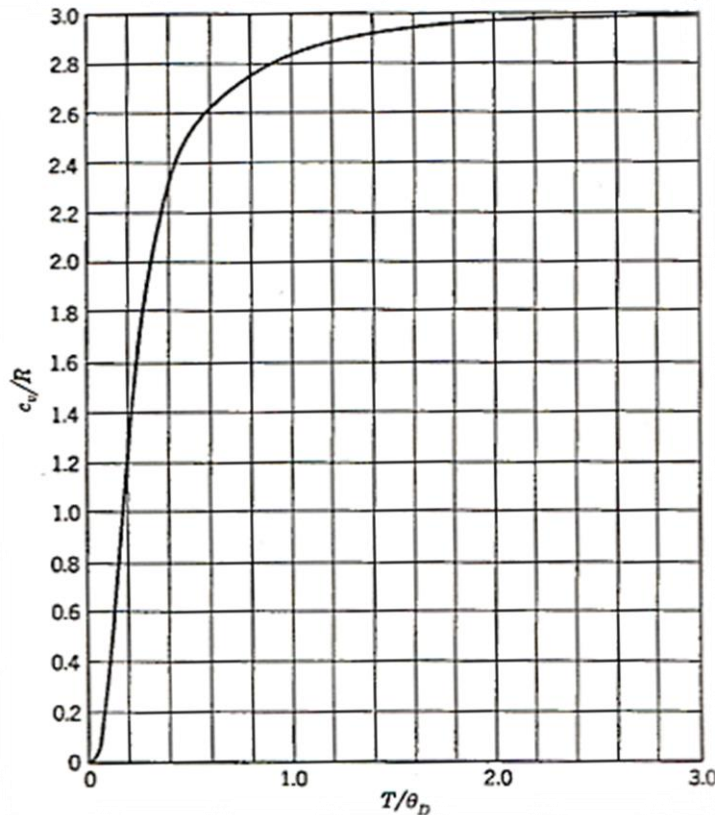
The Debye specific heat function

- $\dfrac{T}{\theta_D} > 3$, $\dfrac{C_v}{R} \approx 3$ (Dulong-Petit value)

- $\dfrac{T}{\theta_D} < \dfrac{1}{12}$, $\dfrac{C_v}{R} \propto T^3$

Specific heat affected by electrons

According to quantum theory,

$$C_{v,e} = \frac{4\pi^4 a m_e M R^2 T}{h^2 N_0 (3\pi^2 \, N/V)^{2/3}} = \gamma_e T$$

$a = Number\ of\ free\ electrons\ per\ atom$
$m_e = Electron\ effective\ mass$
$M = atomic\ weight\ of\ material$
$R = Specific\ gas\ constant\ for\ material$
$T = Absolute\ temperature$
$h = Planck's\ constant$
$N_0 = Avogadro's\ number$
$N/V = Number\ of\ free\ electrons\ per\ unit\ volume$
$\gamma_e = Electronic\ specific-heat\ coefficient$

Specific heat affected by electrons

Table. Electronic specific heat coefficients

Material	γ_e mJ/kg-K^2	γ_e Btu/lb$_m$-°R^2
Aluminum	50.4	6.69×10^{-6}
Beryllium	24.6	3.27×10^{-6}
Chromium	29.8	3.95×10^{-6}
Copper	11.0	1.46×10^{-6}
Gold	3.55	0.471×10^{-6}
Iron	89.9	11.9×10^{-6}
Nickel	124.0	16.5×10^{-6}
Niobium	94.9	12.6×10^{-6}
Platinum	34.0	4.51×10^{-6}
Silver	5.65	0.749×10^{-6}
Tantalum	32.3	4.29×10^{-6}
Titanium	74.1	9.83×10^{-6}
Zirconium	33.2	4.41×10^{-6}

By permission from Gopal (1966).

Ordinary temperature → γ_e is small (ignorable)

Low Temperature → γ_e becomes important!

2.8 Specific heat of liquids and gases

■ **Specific heat of a material (equipartition theorem)**
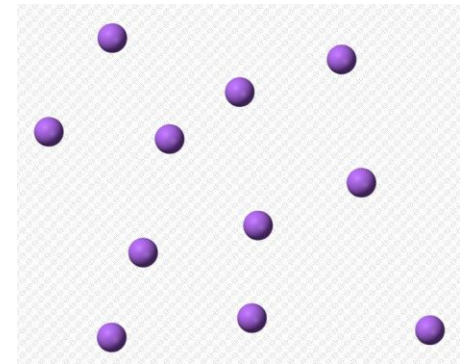
$$C_v = \frac{1}{2}Rf$$

(f : number of degrees of freedom)

① Monatomic gas

• Degree of freedom

- Translational motion : 3

- Rotational motion : 0 $\therefore C_v = \frac{3}{2}R$

- Vibrational motion : 0
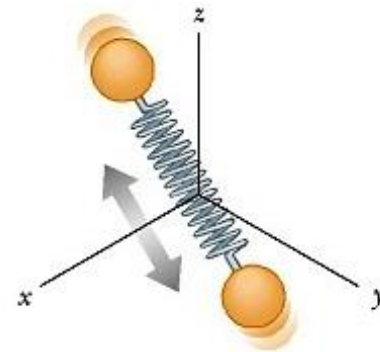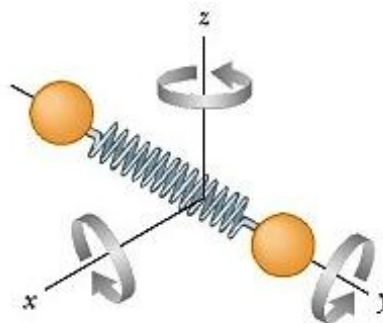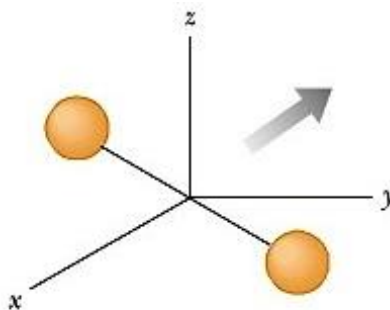
② Diatomic gas

• Degree of freedom

  - Translational motion :   3

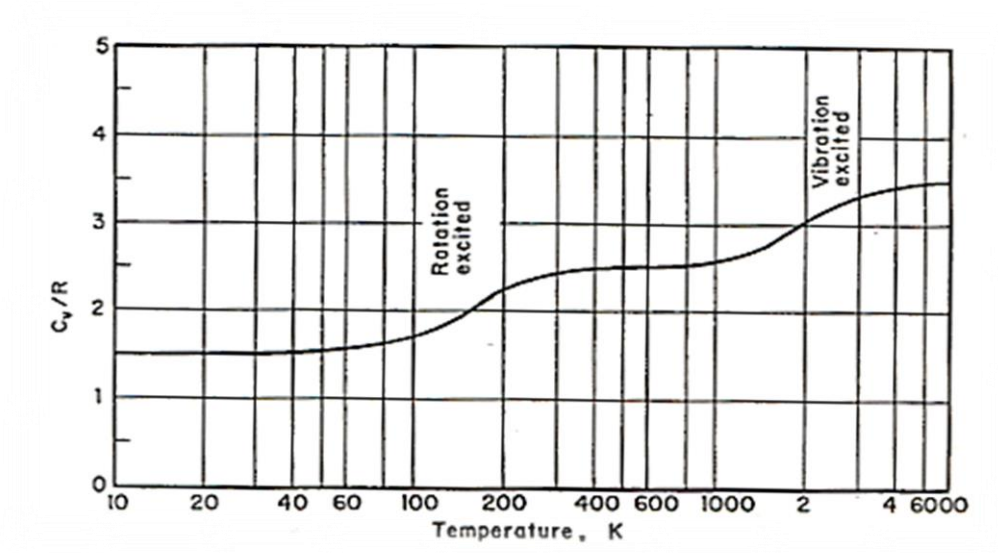 - Rotational motion :   2

  - Vibrational motion :   2

$$\therefore C_v = \frac{7}{2} R$$

(according to the classical theory)

2.8 Specific heat of liquids and gases

In the actual case, rotational, vibrational modes are quantized!

\<Variation of the specific heat C_v for hydrogen gas\>

0~10K	Translational motion only $(\frac{C_v}{R} = \frac{3}{2})$
10~1000K	Translational+Rotational motion $(\frac{C_v}{R} = \frac{5}{2})$
1000K~	Translational+Rotational+Vibrational motion $(\frac{C_v}{R} = \frac{7}{2})$

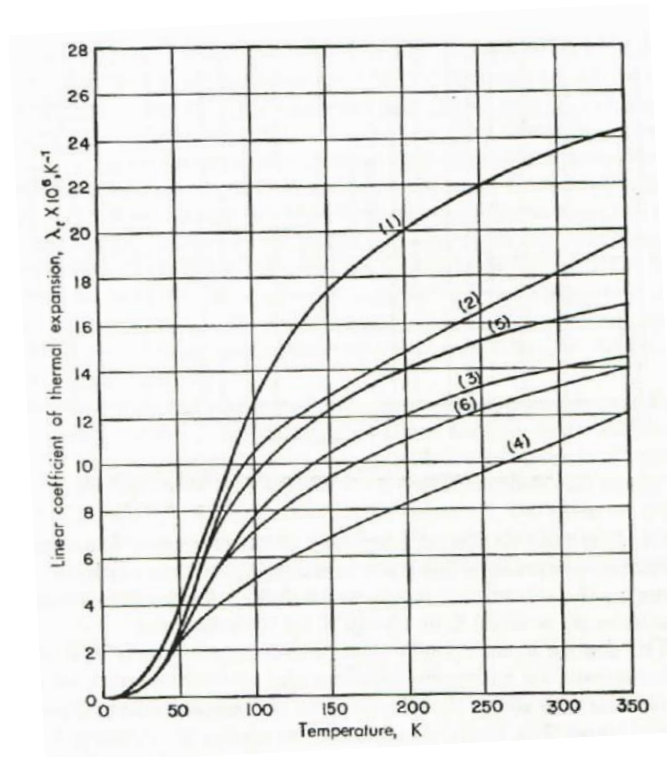■ **Coefficient of thermal expansion**

- For isotropic materials

$$\beta = 3\lambda_t$$

$$\beta = \frac{1}{v}\left(\frac{\partial v}{\partial T}\right)_p \qquad \text{(in the vicinity of the critical point)}$$

β is fractional change in volume per unit change in temperature
λ is the linear coefficient of thermal expansion

2.9 Coefficient of thermal expansion

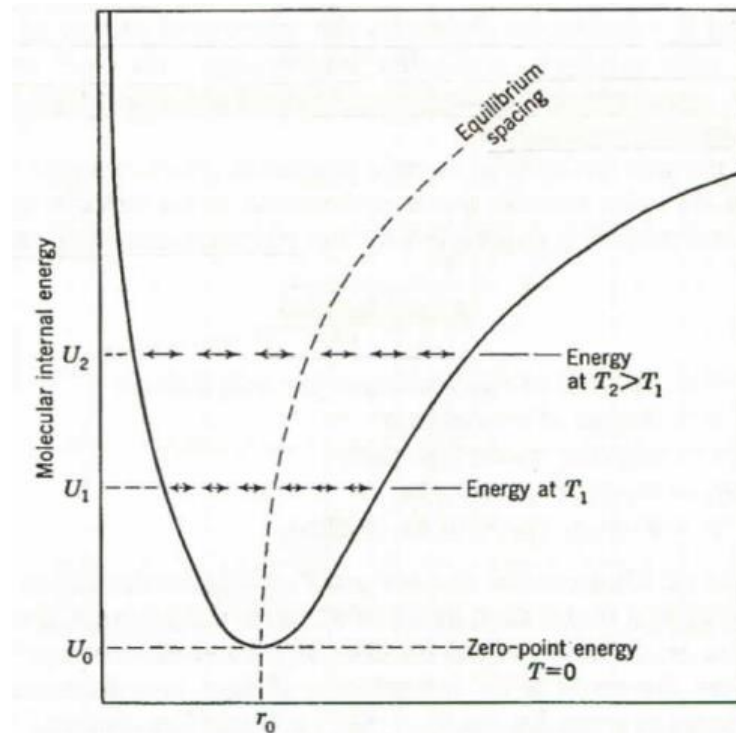- The temperature variation of the linear coefficient for thermal expansion for several materials

Linear coefficient of thermal expansion for several materials at low temperature:
(1) 2024-T4 aluminum (2) beryllium copper (3) K Monel
(2) (4) titanium (5) 304 stainless steel (6) C1020 carbon steel

- **The intermolecular potential-energy curve**

Variation of the intermolecular potential energy for a pair of molecules

2.9 Coefficient of thermal expansion

▪ **The intermolecular forces**

   - The intermolecular potential-energy curve is not symmetrical. Therefore, as the molecule acquires more energy, its mean position relative to its neighbors becomes larger, that is, the material expands.

   - The coefficient of thermal expansion increases as temperature is increased.

2.9 Coefficient of thermal expansion

■ **Coefficient of thermal expansion**

- For crystalline solids, the Gruneisen relation

$$\beta = \frac{\gamma_G c_v \rho}{B}$$

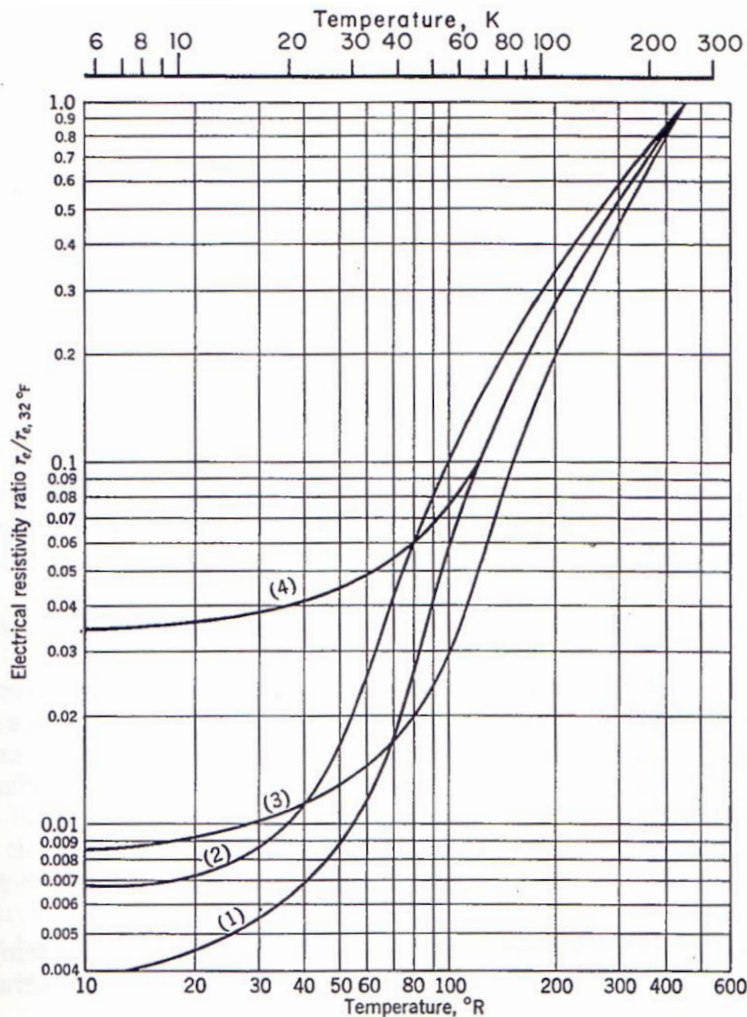ρ is the density of the material
B is the bulk modulus
γ_G is the Gruneisen constant

Material	γ_G
Aluminum	2.17
Copper	1.96
Gold	2.40
Iron	1.60
Lead	2.73
Nickel	1.88
Platinum	2.54
Silver	2.40
Tantalum	1.75
Tungsten	1.62

Values of the Gruneisen constant
for selected solids
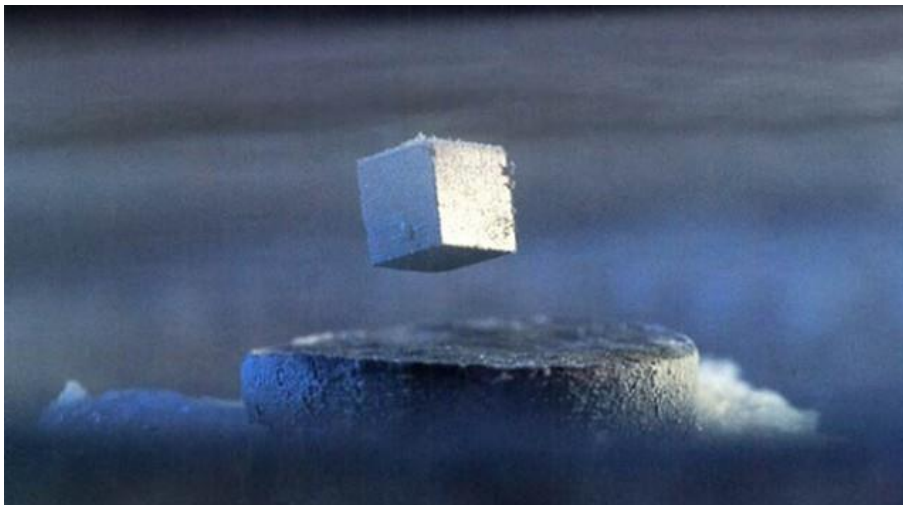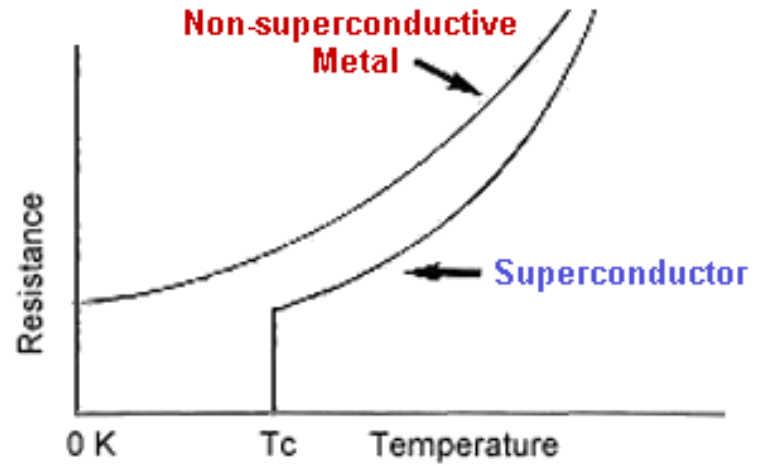
$$k_e = \frac{I/Area}{dV/dx} = \frac{I/A}{V/I} = \frac{L}{RA}$$

Electrical resistivity, $\rho = \dfrac{1}{k_e}$

As T $\downarrow$, Vibrational E $\downarrow$

$$q = kA\frac{dT}{dx} \qquad I = k_eA\frac{dV}{dx}$$

2.11 Superconductivity

- Only at very low T
- Disappearance of all electric resistance
- Appearance of perfect diamagnetism

Paramagnetism ~ magnet
diamagnetism ~ repel
ferromagnetism ~ stick

$T : T < T_o$

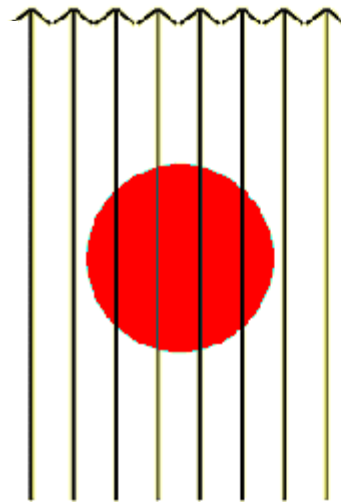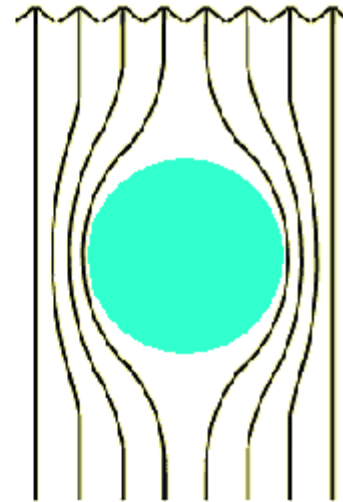$\mathcal{H} : \mathcal{H}_o$ (Critical field)

         └ magnetic field strength required to destroy superconductivity

$I : I_c$ (Critical current)

       └ upper limit to the electric current without destroying superconductivity

Normal   Superconducting

■ **Applications**

MRI (Magnetic Resonance Imaging)
High magnetic field stability

■ **Video: Levitating Superconductor on a Möbius strip**

The Royal Institution - Levitating Superconductor on a Möbius strip
https://www.youtube.com/watch?v=zPqEEZa2Gis

$LN_2 (liquid\ N_2)$

$\rightarrow$ Clear, Colorless
$\rightarrow$ N.B.P. (Normal Boiling Point) : $77K$
$\rightarrow$ Produced by the distiliation of air
$\rightarrow$ Small heat of vaporization

2.12 Cryogenic fluid property

Usage of $LN_2 (liquid\ N_2)$

Quick freezing of food, Drying etc.

$LO_2 (liquid\ O_2)$

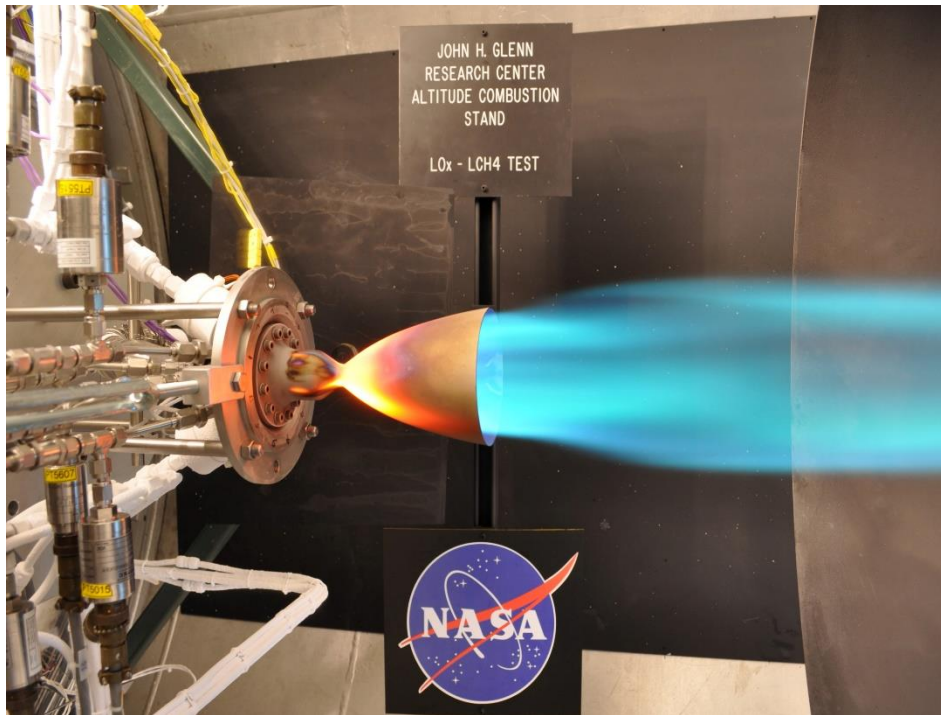$\rightarrow$ Slightly magnetic (paramagnetic)
$\rightarrow$ N.B.P. (Normal Boiling Point) : $90K$
$\rightarrow$ Produced by the distiliation of air
$\rightarrow$ Slightly magnetic (paramagnetic)

Usage of $LO_2 (liquid\ O_2)$

Fuel of rocket, welding etc.

N.B.P. = $20.3K$

Types of hydrogen atom

$$H - H \quad H - D$$
$$D - D \quad T - T$$
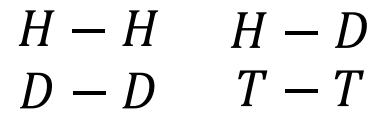
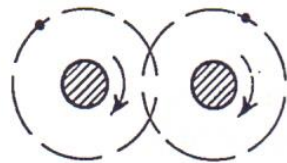|	H hydrogen	D deuterium	T tritium
proton	1	1	1
neutron	0	1	2
electron	1	1	1

2.13 H_2

N.B.P. = $20.3K$

Types of hydrogen molecules

$ortho - H_2$
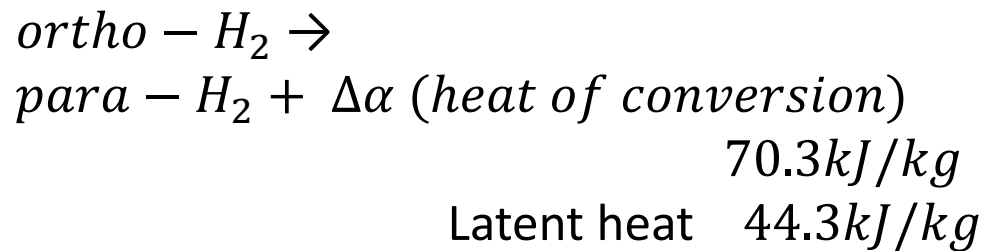(Spins aligned, high energy)

$para - H_2$
(Opposite spins, low energy)

Ortho-hydrogen

Para-hydrogen

$$ortho - H_2 \rightarrow$$
$$para - H_2 + \Delta\alpha \; (heat \; of \; conversion)$$
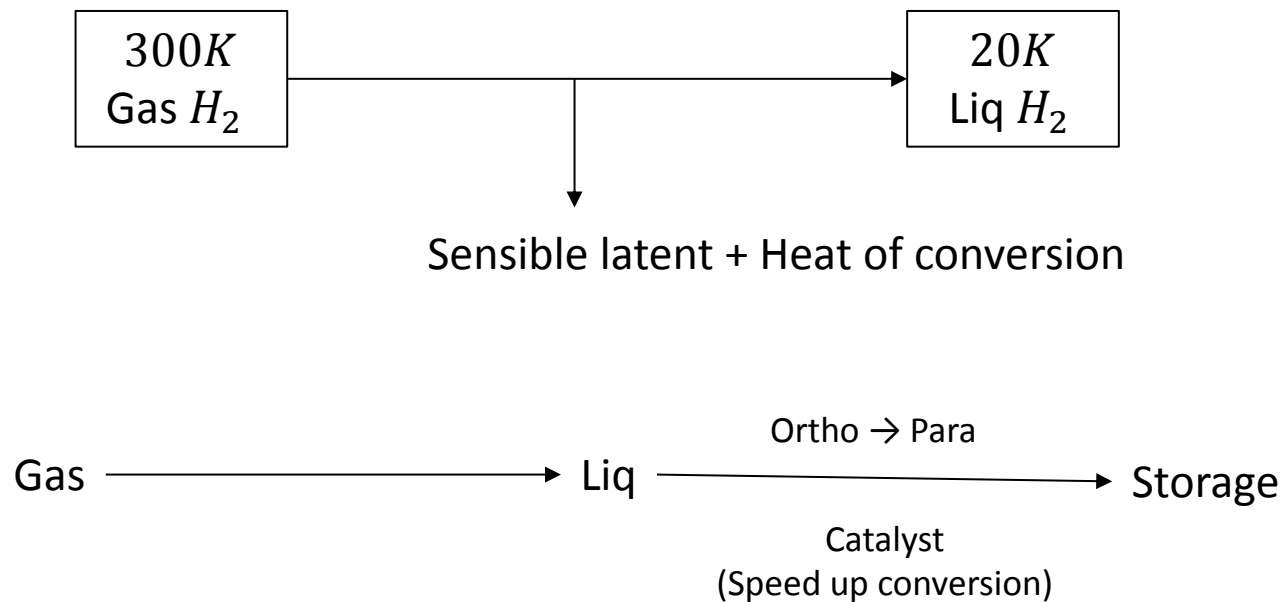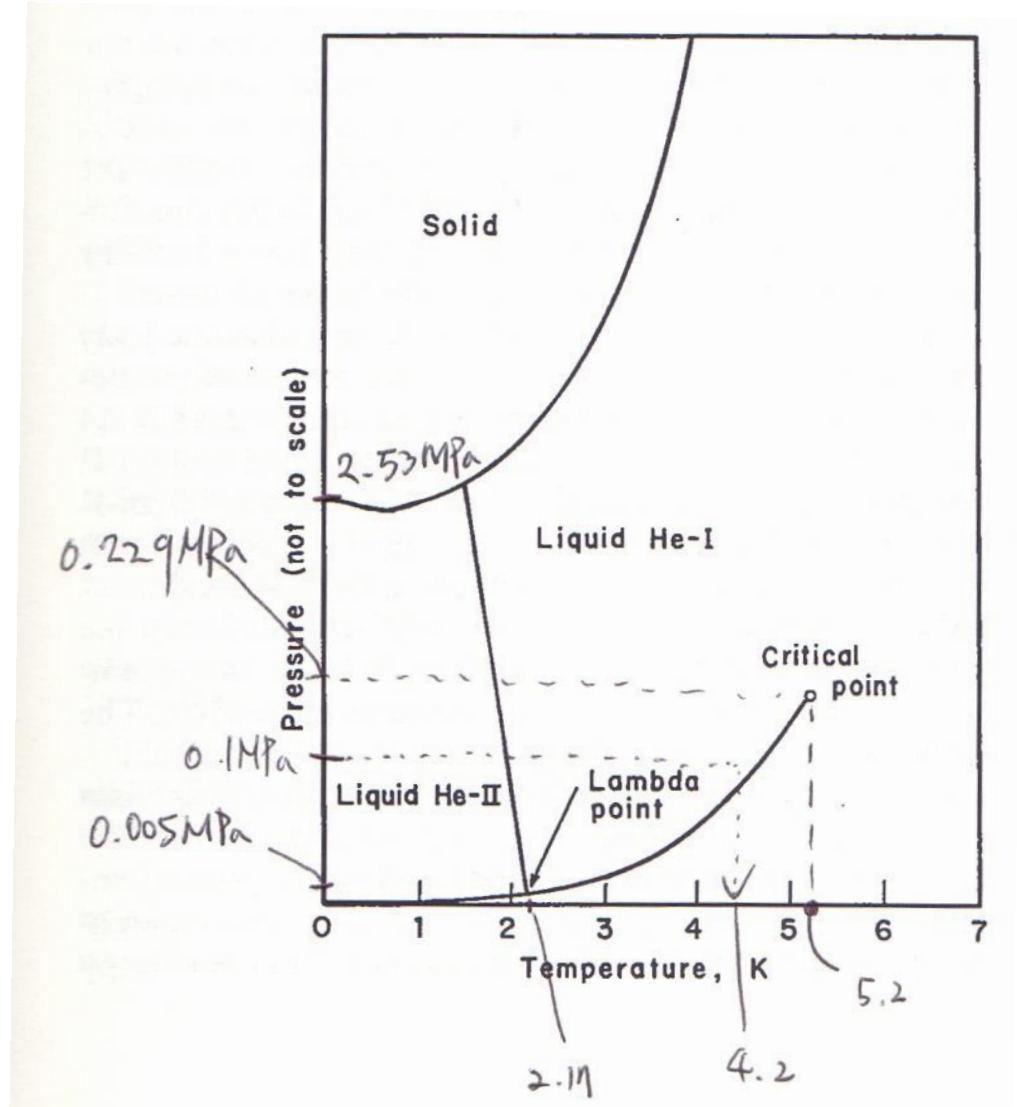$$70.3 kJ/kg$$

Latent heat $44.3 kJ/kg$
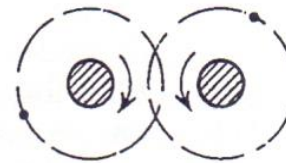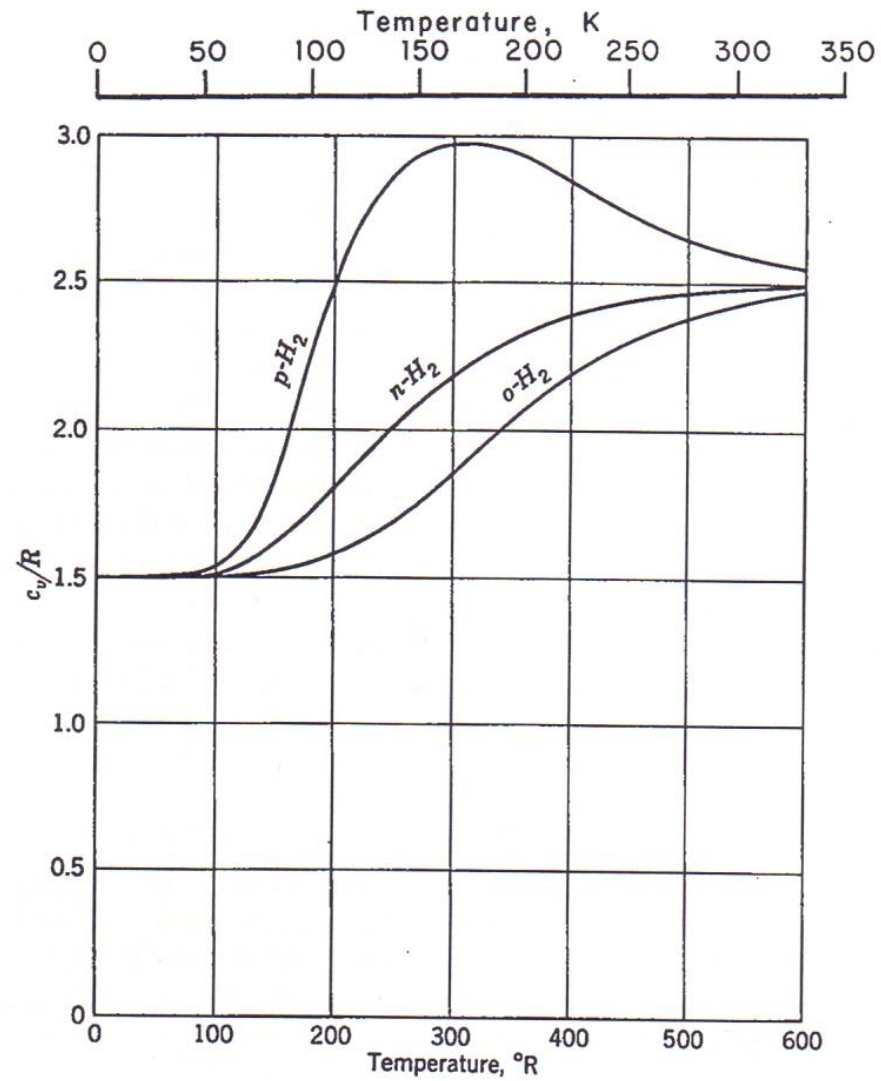
At high temperature is a mixture of 75% $ortho - H_2$ and 25% $para - H_2$
As temperature is cooled to the normal boiling point of hydrogen, the
$ortho - H_2$ concentration decreases from $75 \; to \; 0.2\%$

2.13 H_2

300K
Gas H_2

20K
Liq H_2

Sensible latent + Heat of conversion

Gas ⟶ Liq

Ortho → Para

Catalyst
(Speed up conversion)

Storage

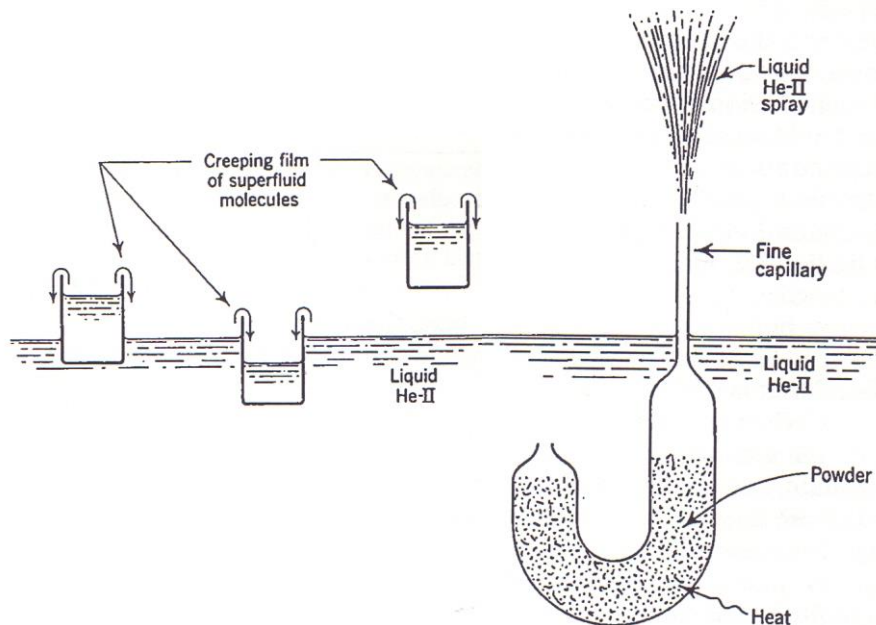# 2.13 He^4

- Difficult to liquefy
- N.B.P. = $4.2K$
- No triple point

$He - $ Ⅰ : Normal fluid

$He - $ Ⅱ : Super fluid

└ act as if it has zero viscosity

$$\left(viscosity, \ \tau = \mu \frac{dV}{dr}\right)$$

Behavior of superfluid

- **Video: Superfluid helium**

BBC - Ben Miller experiments with superfluid helium - Horizon: What is One Degree? - BBC Two
https://www.youtube.com/watch?v=9FudzqfpLLs

2.14 He^3

- N.B.P. = $3.19K$
- Super fluid transition = $3.5mK$