

Electromagnetics:

Circular Waveguides

(10-5)

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Circular Waveguides

Just recall:

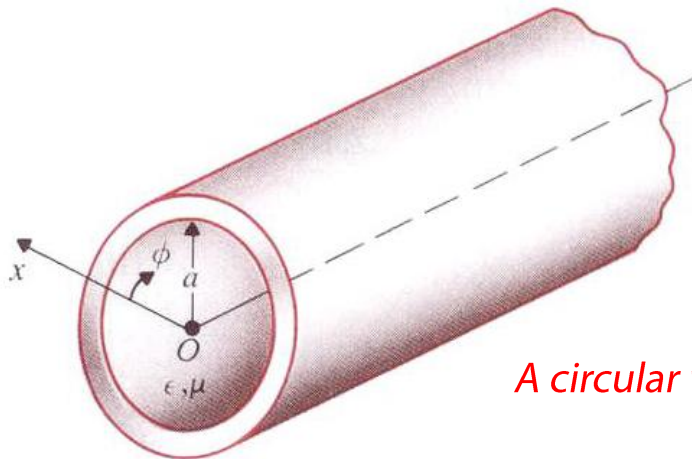
$$\gamma = \alpha + j\beta \quad \rightarrow \text{Propagation constant}$$

$$\mathbf{E}(x, y, z; t) = \text{Re}[\mathbf{E}^0(x, y)e^{(j\omega t - \gamma z)}]$$

$$\text{TM: } H_z = 0 \quad \rightarrow \nabla_{xy}^2 E_z^0 + (\gamma^2 + k^2)E_z^0 = 0 \quad \rightarrow \nabla_{xy}^2 E_z^0 + h^2 E_z^0 = 0$$

$$\text{TE: } E_z = 0 \quad \rightarrow \nabla_{xy}^2 H_z^0 + (\gamma^2 + k^2)H_z^0 = 0 \quad \rightarrow \nabla_{xy}^2 H_z^0 + h^2 H_z^0 = 0$$

$\leftarrow h^2 = \gamma^2 + k^2$



A circular waveguide

→ Choice of the coordinate system:

$$\rightarrow \nabla_{r\phi}^2 E_z^0 + h^2 E_z^0 = 0$$

$$\rightarrow \nabla_{r\phi}^2 H_z^0 + h^2 H_z^0 = 0$$

Bessel's Differential Equation

Consider:

$$\nabla_{r\phi}^2 E_z^0 + h^2 E_z^0 = 0$$

$$\rightarrow \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial E_z^0}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 E_z^0}{\partial \phi^2} + h^2 E_z^0 = 0$$

Suppose: $E_z^0 = R(r)\Phi(\phi)$

$$\rightarrow \frac{d^2 \Phi(\phi)}{d\phi^2} + n^2 \Phi(\phi) = 0$$

$$\rightarrow \frac{d^2 R(r)}{dr^2} + \frac{1}{r} \frac{dR(r)}{dr} + \left(h^2 - \frac{n^2}{r^2} \right) R(r) = 0$$

\rightarrow Bessel's differential equation

cf.

$$H_n^{(1)}(hr) = J_n(hr) + jN_n(hr)$$

$$H_n^{(2)}(hr) = J_n(hr) - jN_n(hr)$$

$$R(r) = \begin{cases} CJ_n(hr) + DN_n(hr) & \text{for } h^2 > 0 \\ CI_n(hr) + DK_n(hr) & \text{for } h^2 < 0 \end{cases}$$

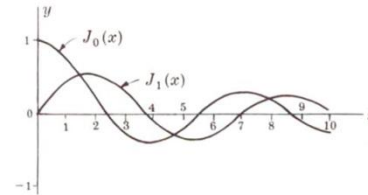


Fig. 24-1

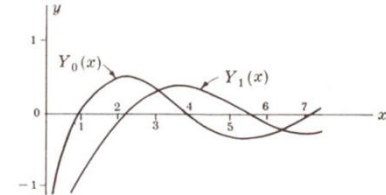


Fig. 24-2

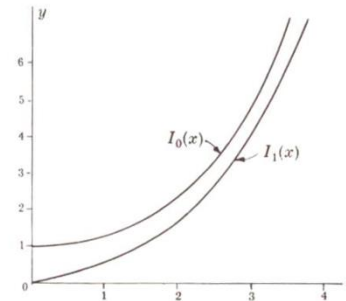


Fig. 24-3

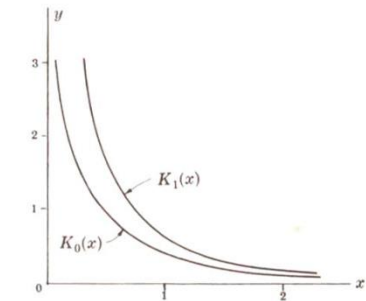


Fig. 24-4

M. R. Spiegel, Mathematical Handbook of Formulas and Tables, McGraw-Hill, 1990.

TM Waves in Circular Waveguides (1)

Let: $E_z(r, \phi, z) = E_z^0(r, \phi)e^{-\gamma z}$

Recall:

$$R(r) = \begin{cases} CJ_n(hr) + DN_n(hr) & \text{for } h^2 > 0 \\ CI_n(hr) + DK_n(hr) & \text{for } h^2 < 0 \end{cases}$$

$$\rightarrow E_z^0 = C_n J_n(hr) \cos n\phi$$

Also recall: $\mathbf{E}_T^0 = -\frac{\gamma}{h^2} \nabla_T E_z^0$

$$\rightarrow \nabla_T E_z^0 = \left(\mathbf{a}_r \frac{\partial}{\partial r} + \mathbf{a}_\phi \frac{\partial}{r \partial \phi} \right) E_z^0$$

Solution:

$$E_r^0 = -\frac{j\beta}{h} C_n J'_n(hr) \cos n\phi$$

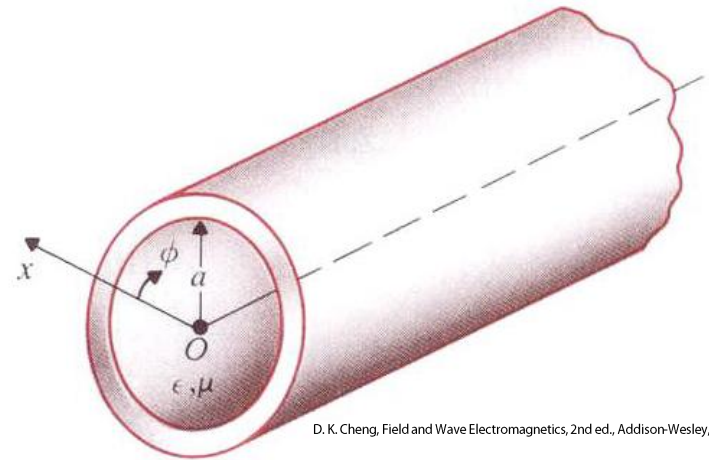
$$E_\phi^0 = \frac{j\beta n}{h^2 r} C_n J_n(hr) \sin n\phi$$

$$H_r^0 = -\frac{j\omega\epsilon n}{h^2 r} C_n J_n(hr) \sin n\phi$$

$$H_\phi^0 = -\frac{j\omega\epsilon}{h} C_n J'_n(hr) \cos n\phi$$

$$\rightarrow \mathbf{H} = \frac{1}{Z_{TM}} \mathbf{a}_z \times \mathbf{E}$$

$$\leftarrow \gamma = j\beta$$



D. K. Cheng, Field and Wave Electromagnetics, 2nd ed., Addison-Wesley, 1989.

A circular waveguide with a conducting wall

TM Waves in Circular Waveguides (2)

Solution:

$$E_r^0 = -\frac{j\beta}{h} C_n J'_n(hr) \cos n\phi$$

$$E_\phi^0 = \frac{j\beta n}{h^2 r} C_n J_n(hr) \sin n\phi$$

$$E_z^0 = C_n J_n(hr) \cos n\phi$$

$$H_r^0 = -\frac{j\omega\varepsilon n}{h^2 r} C_n J_n(hr) \sin n\phi$$

$$H_\phi^0 = -\frac{j\omega\varepsilon}{h} C_n J'_n(hr) \cos n\phi$$

$$H_z^0 = 0$$

Boundary condition:

$$E_z^0 = E_\phi^0 = 0 \quad \text{at } r = a$$

$$\rightarrow J_n(ha) = 0 \quad \text{at } r = a$$

$$\rightarrow ha = x_{nm}$$

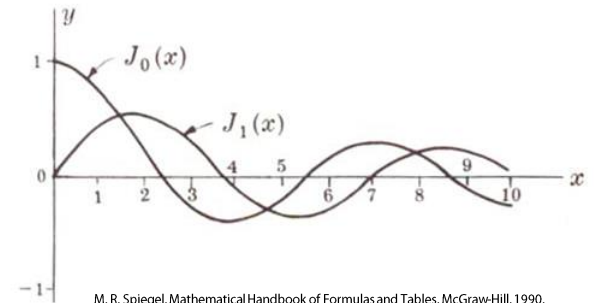
where x_{nm} represents the m^{th} zero of $J_n(x)$.

TABLE 10-2
Zeros of $J_n(x)$, x_{np}

| $p \backslash n$ | $n = 0$ | $n = 1$ | $n = 2$ |
|------------------|---------|---------|---------|
| 1 | 2.405 | 3.832 | 5.136 |
| 2 | 5.520 | 7.016 | 8.417 |

D. K. Cheng, Field and Wave Electromagnetics, 2nd ed., Addison-Wesley, 1989.

Recall:



M. R. Spiegel, Mathematical Handbook of Formulas and Tables, McGraw-Hill, 1990.

For TM_{01} mode:

$$h_{TM_{01}} = \frac{2.405}{a} \quad \rightarrow \quad f_{c, TM_{01}} = \frac{h_{TM_{01}}}{2\pi\sqrt{\mu\varepsilon}} = \frac{0.383}{a\sqrt{\mu\varepsilon}}$$

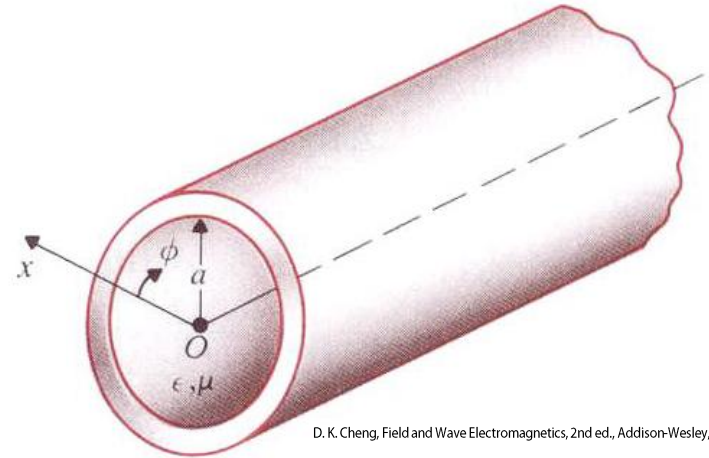
TE Waves in Circular Waveguides (1)

Let: $H_z(r, \phi, z) = H_z^0(r, \phi)e^{-\gamma z}$

Recall again:

$$R(r) = \begin{cases} C J_n(hr) + D N_n(hr) & \text{for } h^2 > 0 \\ C I_n(hr) + D K_n(hr) & \text{for } h^2 < 0 \end{cases}$$

$$\rightarrow H_z^0 = C'_n J_n(hr) \cos n\phi$$



D. K. Cheng, Field and Wave Electromagnetics, 2nd ed., Addison-Wesley, 1989.

A circular waveguide with a conducting wall

Also recall: $\mathbf{H}_T^0 = -\frac{\gamma}{h^2} \nabla_T H_z^0$ & $\mathbf{E} = -Z_{TE} (\mathbf{a}_z \times \mathbf{H})$

Solution:

$$H_r^0 = -\frac{j\beta}{h} C'_n J'_n(hr) \cos n\phi$$

$$H_\phi^0 = \frac{j\beta n}{h^2 r} C'_n J_n(hr) \sin n\phi$$

$$E_r^0 = \frac{j\omega\mu n}{h^2 r} C'_n J_n(hr) \sin n\phi$$

$$E_\phi^0 = \frac{j\omega\mu}{h} C'_n J'_n(hr) \cos n\phi$$

$$\leftarrow \gamma = j\beta$$

TE Waves in Circular Waveguides (2)

Solution:

$$H_r^0 = -\frac{j\beta}{h} C'_n J'_n(hr) \cos n\phi$$

$$H_\phi^0 = \frac{j\beta n}{h^2 r} C'_n J_n(hr) \sin n\phi$$

$$H_z^0 = C'_n J_n(hr) \cos n\phi$$

$$E_r^0 = \frac{j\omega\mu n}{h^2 r} C'_n J_n(hr) \sin n\phi$$

$$E_\phi^0 = \frac{j\omega\mu}{h} C'_n J'_n(hr) \cos n\phi$$

$$E_z^0 = 0$$

Boundary condition:

$$E_z^0 = E_\phi^0 = 0 \quad \text{at } r = a$$

$$\rightarrow J'_n(ha) = 0 \quad \text{at } r = a \quad \rightarrow J'_n(x) = \frac{1}{2} [J_{n-1}(x) - J_{n+1}(x)]$$

$$\rightarrow ha = x'_{nm}$$

where x'_{nm} represents the m^{th} zero of $J'_n(x)$.

For TE_{11} mode:

$$h_{TE_{11}} = \frac{1.841}{a} \quad \rightarrow f_{c,TE_{11}} = \frac{h_{TE_{11}}}{2\pi\sqrt{\mu\epsilon}} = \frac{0.293}{a\sqrt{\mu\epsilon}}$$

→ Dominant mode!

TABLE 10-3
Zeros of $J'_n(x)$, x'_{np}

| $p \backslash n$ | $n = 0$ | $n = 1$ | $n = 2$ |
|------------------|---------|---------|---------|
| 1 | 3.832 | 1.841 | 3.054 |
| 2 | 7.016 | 5.331 | 6.706 |

D. K. Cheng, Field and Wave Electromagnetics, 2nd ed., Addison-Wesley, 1989.

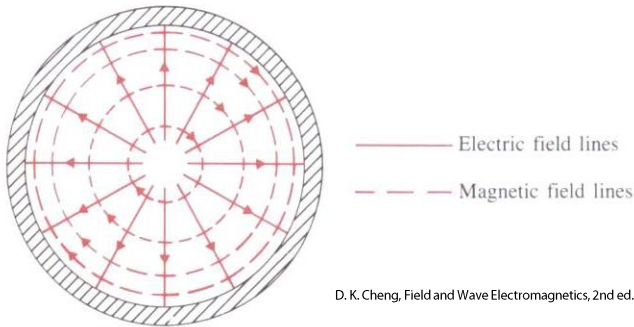
TM/TE Waves in Circular Waveguides

TM₀₁ mode:

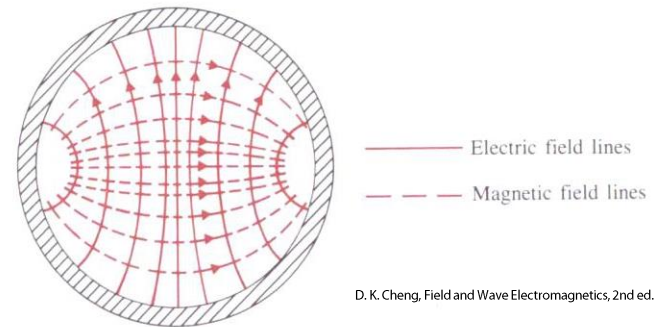
| | |
|--|--|
| $E_r^0 = -\frac{j\beta}{h} C_0 J'_0(hr)$ | $H_r^0 = 0$ |
| $E_\phi^0 = 0$ | $H_\phi^0 = -\frac{j\omega\epsilon}{h} C_0 J'_0(hr)$ |
| $E_z^0 = C_0 J_0(hr)$ | $H_z^0 = 0$ |

TE₁₁ mode:

| | |
|---|--|
| $E_r^0 = \frac{j\omega\mu}{h^2 r} C'_1 J_1(hr) \sin \phi$ | $H_r^0 = -\frac{j\beta}{h} C'_1 J'_1(hr) \cos \phi$ |
| $E_\phi^0 = \frac{j\omega\mu}{h} C'_1 J'_1(hr) \cos \phi$ | $H_\phi^0 = \frac{j\beta h}{h^2 r} C'_1 J_1(hr) \sin \phi$ |
| $E_z^0 = 0$ | $H_z^0 = C'_1 J_1(hr) \cos \phi$ |



D. K. Cheng, Field and Wave Electromagnetics, 2nd ed., Addison-Wesley, 1989.

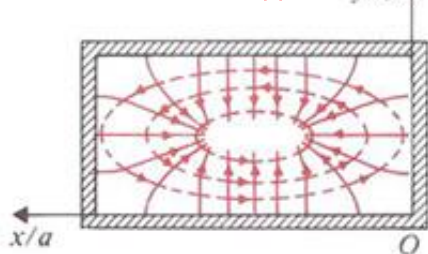


D. K. Cheng, Field and Wave Electromagnetics, 2nd ed., Addison-Wesley, 1989.

$$h_{TM_{01}} = \frac{2.405}{a} \rightarrow f_{c, TM_{01}} = \frac{0.383}{a\sqrt{\mu\epsilon}}$$

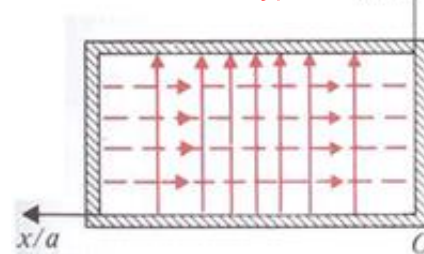
$$h_{TE_{11}} = \frac{1.841}{a} \rightarrow f_{c, TE_{11}} = \frac{0.293}{a\sqrt{\mu\epsilon}}$$

cf. TM₁₁



D. K. Cheng, Field and Wave Electromagnetics, 2nd ed., Addison-Wesley, 1989.

cf. TE₀₁



D. K. Cheng, Field and Wave Electromagnetics, 2nd ed., Addison-Wesley, 1989.

→ Dominant mode!