

Electro-Optics

Theory of Lasers (3)

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Multimode Laser Oscillation

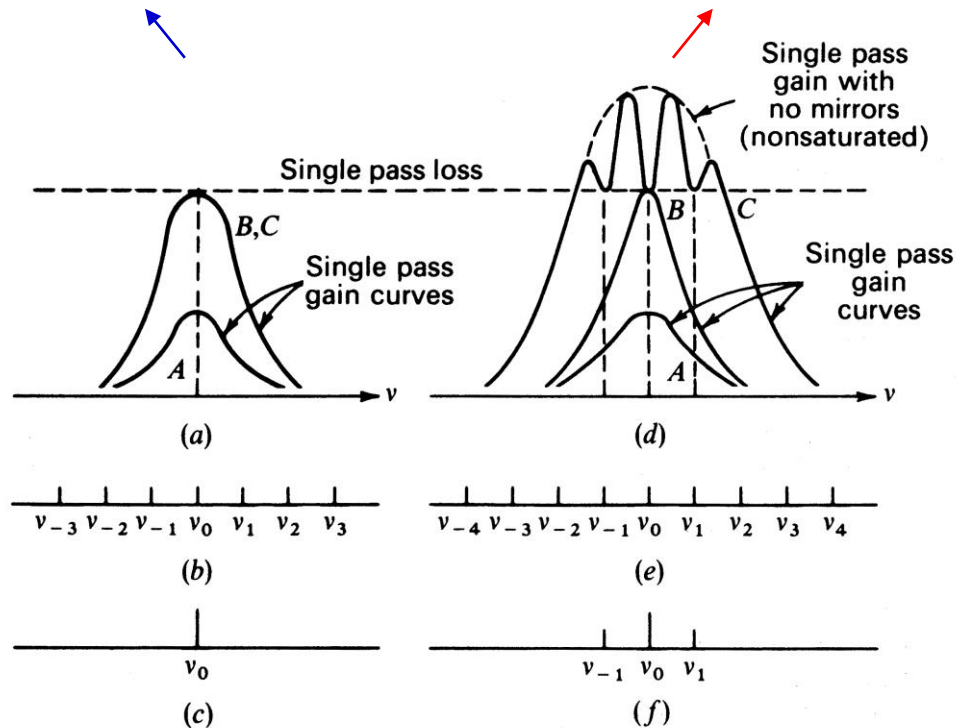
Laser oscillation condition:

$$\gamma_{th}(\nu) = \alpha - \frac{1}{l} \ln r_1 r_2 \rightarrow \gamma(\nu) = (N_2 - N_1) \frac{c^2}{8\pi n^2 \nu^2 \tau_{sp}} g(\nu)$$

$$\nu_{q+1} - \nu_q = \frac{c}{2nl}$$

Homogeneous atomic system:

Inhomogeneous atomic system:



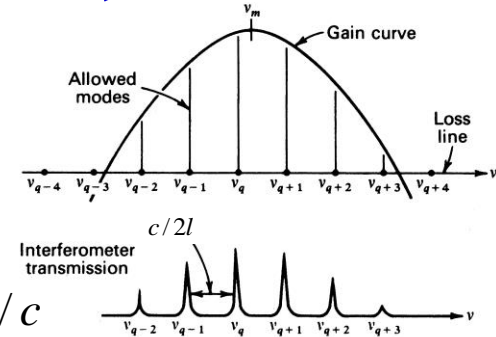
Mode Locking (1)

Mode separation in an **inhomogeneously** broadened system:

$$\omega_{q+1} - \omega_q = \frac{\pi c}{l} \equiv \omega \quad \leftarrow n = 1$$

Total electric field:

$$e(t) = \sum_n E_n e^{i[(\omega_0 + n\omega)t + \phi_n]} \quad \leftarrow \omega_0 = m\pi c / l, \quad T = 2\pi / \omega = 2l / c$$



A. Yariv, *Quantum Electronics*, 3rd ed., Wiley, 1989.

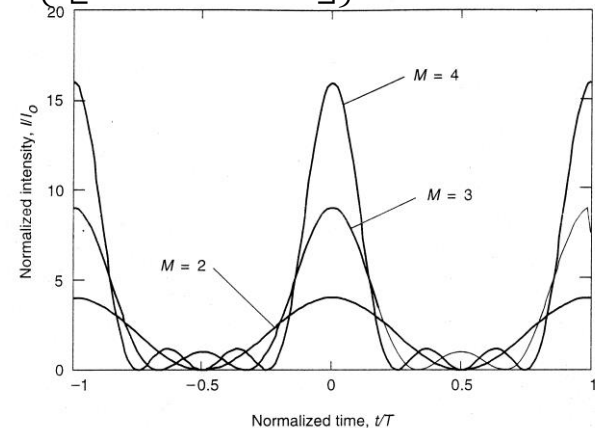
In case E_n and ϕ_n are fixed (locked):

$$\begin{aligned} e(t+T) &= \sum_n E_n \exp\left\{i\left[(\omega_0 + n\omega)\left(t + \frac{2\pi}{\omega}\right) + \phi_n\right]\right\} \\ &= \sum_n E_n \exp\{i[(\omega_0 + n\omega)t + \phi_n]\} \exp\left\{i\left[2\pi\left(\frac{\omega_0}{\omega} + n\right)\right]\right\} = e(t) \end{aligned}$$

“Periodic”

Taking $E_n = 1$ and $\phi_n = 0$:

$$\begin{aligned} e(t) &= \sum_{-(M-1)/2}^{(M-1)/2} \exp[i(\omega_0 + n\omega)t] \\ &= \exp(i\omega_0 t) \frac{\sin(M\omega t / 2)}{\sin(\omega t / 2)} \end{aligned}$$



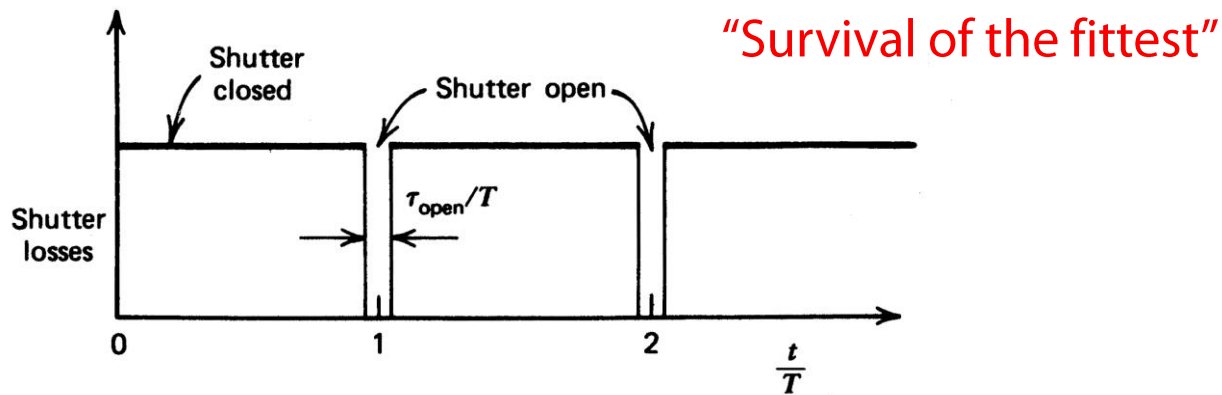
A. Yariv, *Quantum Electronics*, 3rd ed., Wiley, 1989.

Mode-Locked Lasers

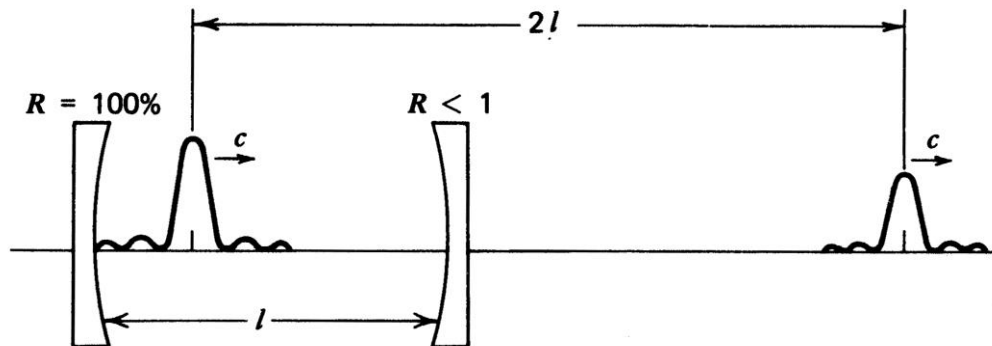
Methods of mode locking: Loss (gain) modulation

Active mode locking: EOM, AOM, etc.

Passive mode locking: Saturable absorber, e.g. dye, SESAM, etc.



A. Yariv, *Quantum Electronics*, 3rd ed., Wiley, 1989.



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Q Switching (1)

Initially lasing is prohibited: $Q \ll 1$

Total number of photons in the cavity when Q is restored:

$$\frac{d\phi}{dt} = \phi \left(\underbrace{\frac{\gamma c L}{n l}}_{\text{Gain}} - \underbrace{\frac{1}{t_c}}_{\text{Loss (Cavity decay time)}} \right) \rightarrow \tau = \frac{t}{t_c}$$

$$\rightarrow \frac{d\phi}{d\tau} = \phi \left(\frac{\gamma}{n l / c L t_c} - 1 \right) = \phi \left(\frac{\gamma}{\gamma_{th}} - 1 \right) \rightarrow \gamma \propto \eta$$

$\leftarrow \eta \equiv (N_2 - N_1)V$

$$= \phi \left(\frac{\eta}{\eta_{th}} - 1 \right)$$

Number of photons generated by induced emission

$$\rightarrow \frac{d\eta}{d\tau} = -2\phi \frac{\eta}{\eta_{th}}$$

$$\rightarrow \frac{d\phi}{d\eta} = \frac{\eta_{th}}{2\eta} - \frac{1}{2} \rightarrow \phi - \phi_i = \frac{1}{2} \left[\eta_{th} \ln \frac{\eta}{\eta_i} - (\eta - \eta_i) \right]$$

$$\rightarrow \phi \cong \frac{1}{2} \left[\eta_{th} \ln \frac{\eta}{\eta_i} - (\eta - \eta_i) \right]$$

Q Switching (2)

Final inversion: $t \gg t_c$

$$\phi \cong \frac{1}{2} \left[\eta_{th} \ln \frac{\eta}{\eta_i} - (\eta - \eta_i) \right] \rightarrow 0$$

$$\rightarrow \frac{\eta_f}{\eta_i} = \exp\left(\frac{\eta_f - \eta_i}{\eta_{th}}\right) \rightarrow \text{To be solved graphically or numerically}$$

Instantaneous power output:

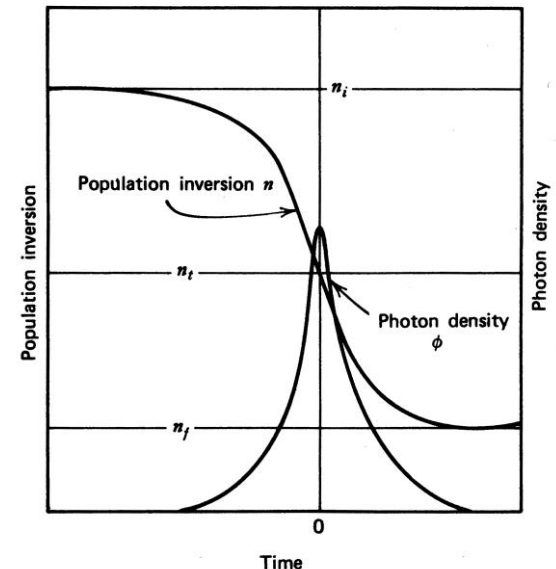
$$P = \phi h \nu / t_c$$

$$= \frac{h \nu}{2 t_c} \left[\eta_{th} \ln \frac{\eta}{\eta_i} - (\eta - \eta_i) \right]$$

Peak power output:

$$\rightarrow \partial P / \partial \eta = 0 \rightarrow \eta = \eta_{th}$$

$$\rightarrow P_{peak} = \frac{h \nu}{2 t_c} \left[\eta_{th} \ln \frac{\eta_{th}}{\eta_i} - (\eta_{th} - \eta_i) \right]$$



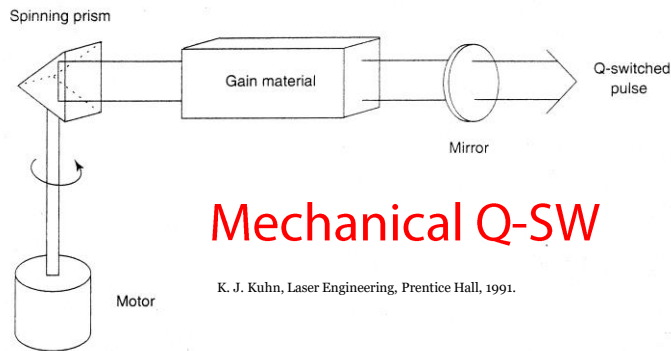
A. Yariv, *Quantum Electronics*, 3rd ed., Wiley, 1989.

Q-Switched (Giant Pulse) Lasers (1)

Methods of Q switching:

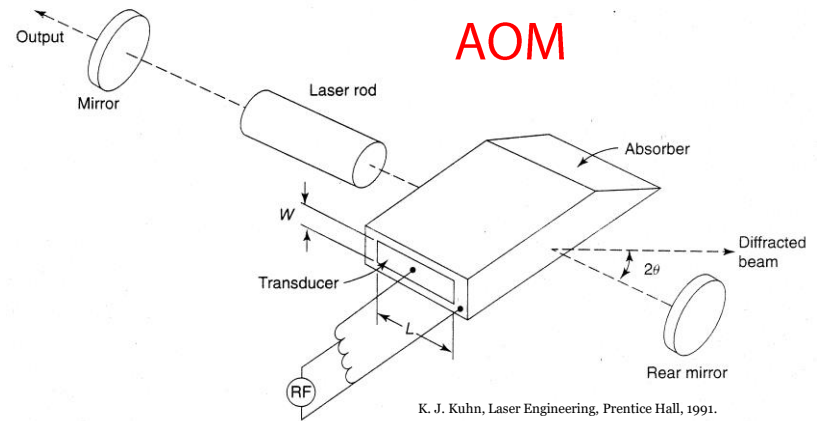
Active Q switching → Mechanical Q-SW, EOM, AOM, etc.

Passive Q switching → Saturable absorber, e.g. dye, SESAM, etc.



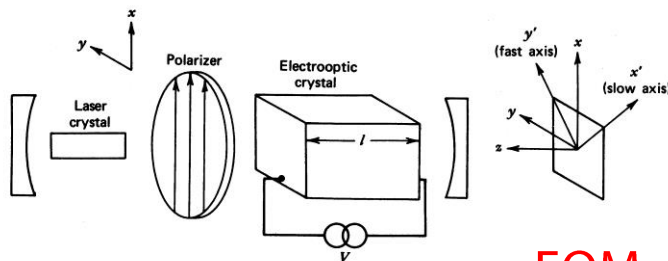
Mechanical Q-SW

K. J. Kuhn, Laser Engineering, Prentice Hall, 1991.

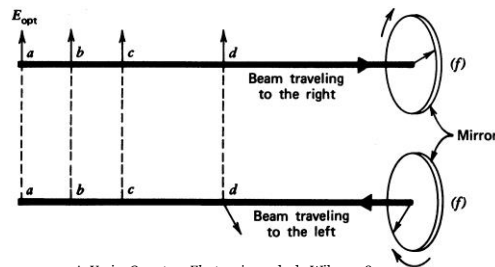


AOM

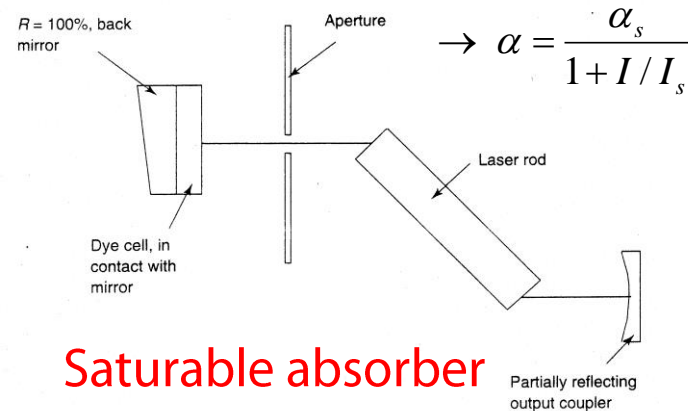
K. J. Kuhn, Laser Engineering, Prentice Hall, 1991.



EOM



A. Yariv, Quantum Electronics, 3rd ed., Wiley, 1989.



Saturable absorber

K. J. Kuhn, Laser Engineering, Prentice Hall, 1991.