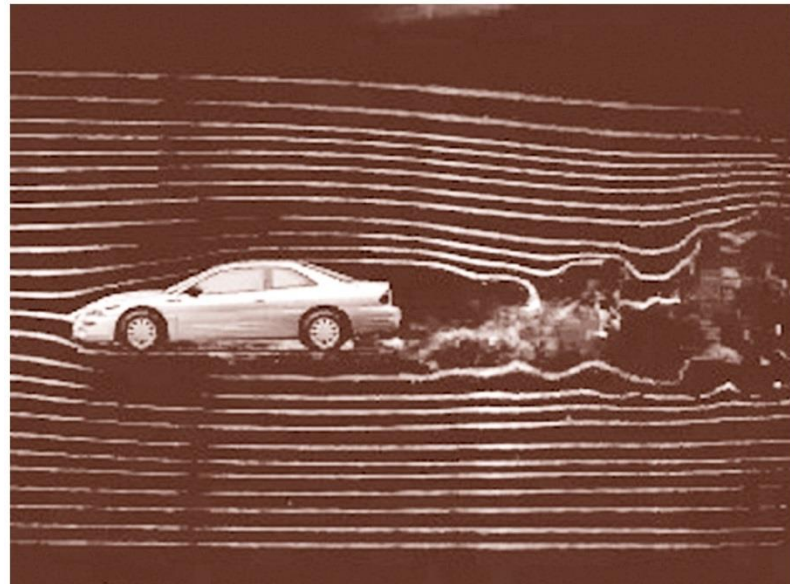


# Lecture 20

## Origin of Turbulence (3)



# Lecture 20 Origin of Turbulence (3)

## Contents

20.1 Velocities and Energies in Turbulence

20.2 Continuity for Turbulent Motion

## Objectives

- Learn fundamental concept of turbulence
- Study Reynolds decomposition
- Derive Reynolds equation from Navier-Stokes equation
- Study eddy viscosity model and mixing length model

# 20.1 Velocities and Energies in Turbulence

## 20.1.1 Reynolds decomposition

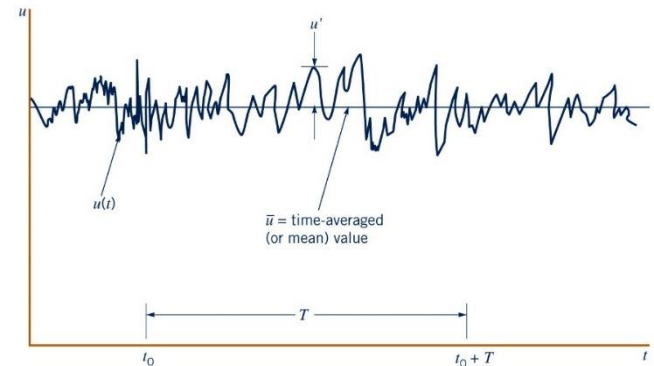
### (1) Velocity decomposition

- The variables in a turbulent flow are described using the theory of stochastic process and random variables even though fluid-dynamic turbulence is not entirely random.

- instantaneous velocity = time-averaged value + fluctuating components

$$u_i(x, y, z, t) = \bar{u}_i(x, y, z) + u_i'(x, y, z, t) \quad (20.1)$$

$$\bar{u}_i = \frac{1}{T} \int_0^T u_i dt \quad (20.2)$$



# 20.1 Velocities and Energies in Turbulence

$T$  = long time compared to the time scale of the turbulence

- $T$  should include (smooth out) all the effects of turbulence fluctuation
- pipe flow:  $10^{-1} \sim 10^0$  s; channel/river flow:  $10^1 \sim 10^2$  s

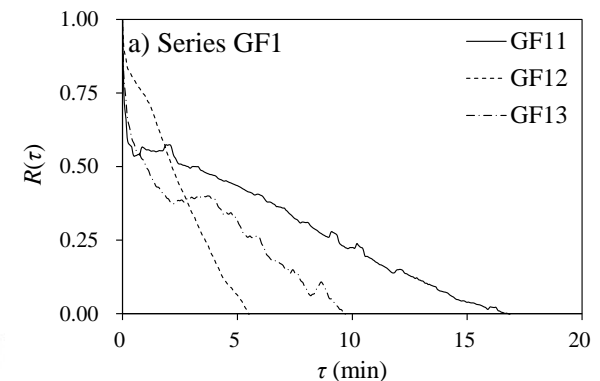
[Re] Integral time scale,  $T_I$

- Time over which a turbulent fluctuation is correlated with itself
- A measure of the memory of the turbulence

$$T_I = \int_0^{\infty} R_I(\tau) d\tau$$

$$R_I(\tau) = \frac{\langle U(t)U(t+\tau) \rangle}{\langle U^2 \rangle}$$

(20.3)



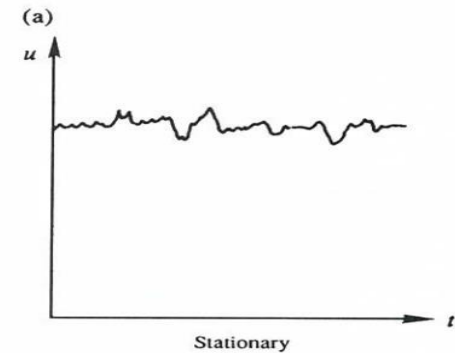
# 20.1 Velocities and Energies in Turbulence

[Re] Averaging of stationary vs nonstationary turbulence

## 1) Stationary turbulence

- Use time average of a single time series data

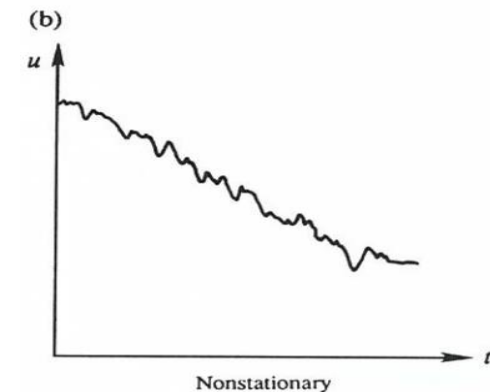
$$\bar{u} = \frac{1}{T} \int_0^T u dt \quad (20.4)$$



## 2) Nonstationary turbulence

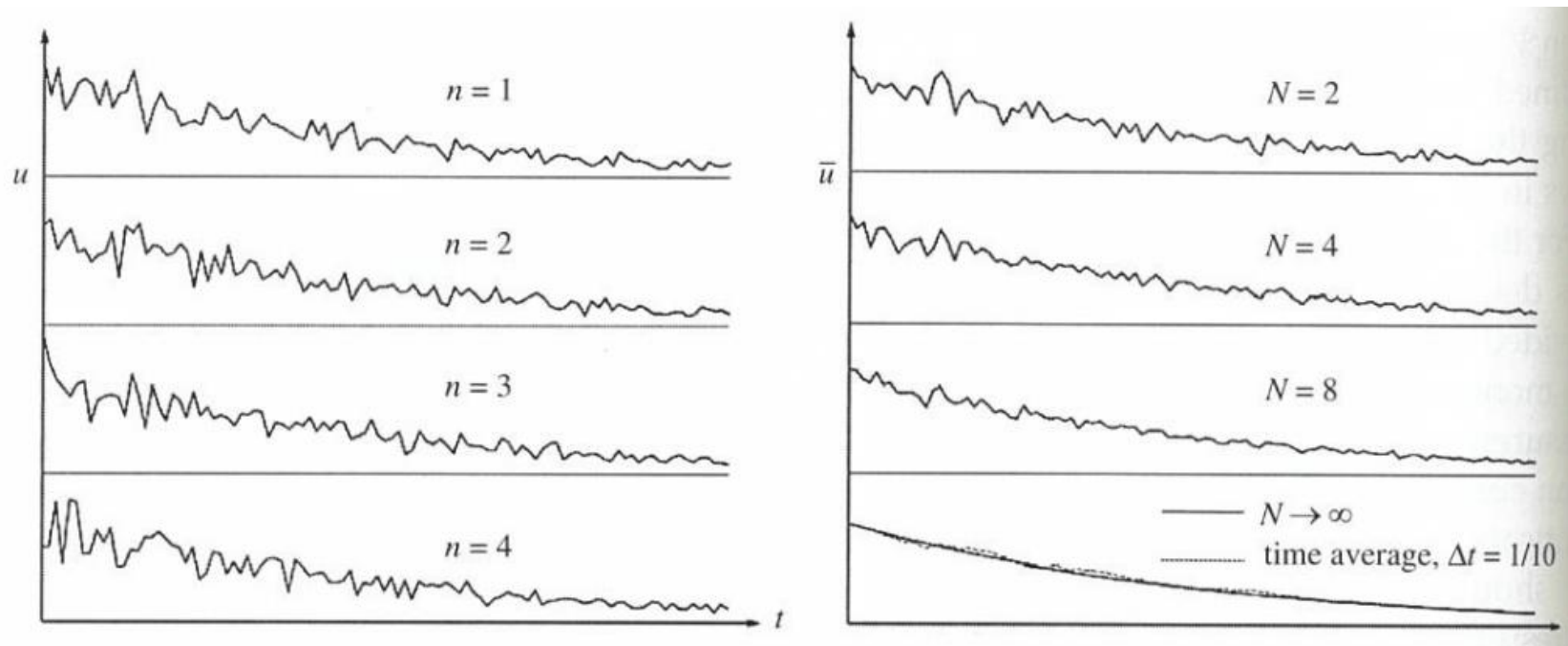
- Use ensemble average of  $N$  independent realization of time series data

$$\langle u \rangle = \lim_{N \rightarrow \infty} \frac{1}{N} \sum_{n=1}^N u \quad (20.5)$$



# 20.1 Velocities and Energies in Turbulence

- Ensemble average of nonstationary turbulence



# 20.1 Velocities and Energies in Turbulence

$$\overline{u'} = \frac{1}{T} \int_0^T u' dt \equiv 0 \quad (\because \text{fluctuations are both plus and minus}) \quad (20.6)$$

$$\left( \frac{1}{T} \int_0^T (u - \bar{u}) dt = \bar{u} - \bar{u} = 0 \right)$$

(2) Pressure and stress decomposition

$$p = \bar{p} + p'$$

$$\sigma_{ij} = \bar{\sigma}_{ij} + \sigma'_{ij} \quad (20.7)$$

# 20.1 Velocities and Energies in Turbulence

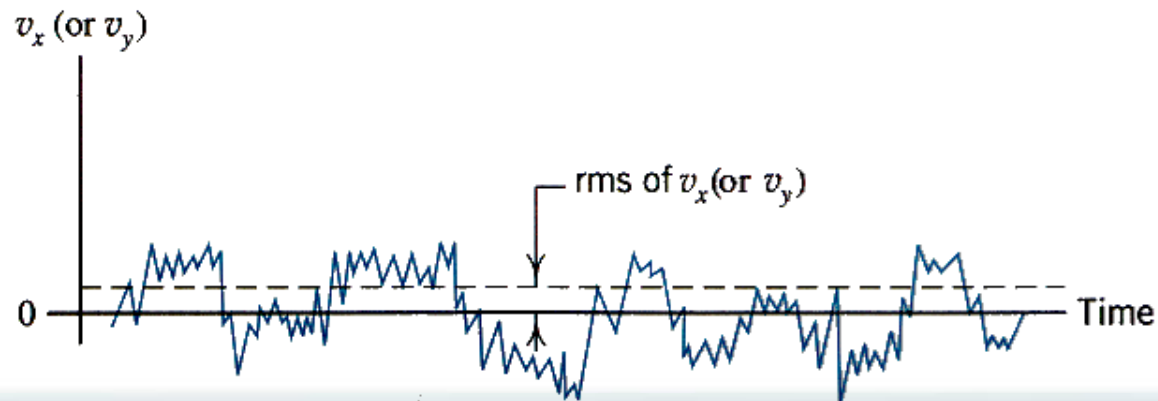
(3) Turbulence intensity - show turbulence effects

→ root-mean-square (rms) = square root of variance = standard deviation

- average intensity of the turbulence = *rms* of  $u'$

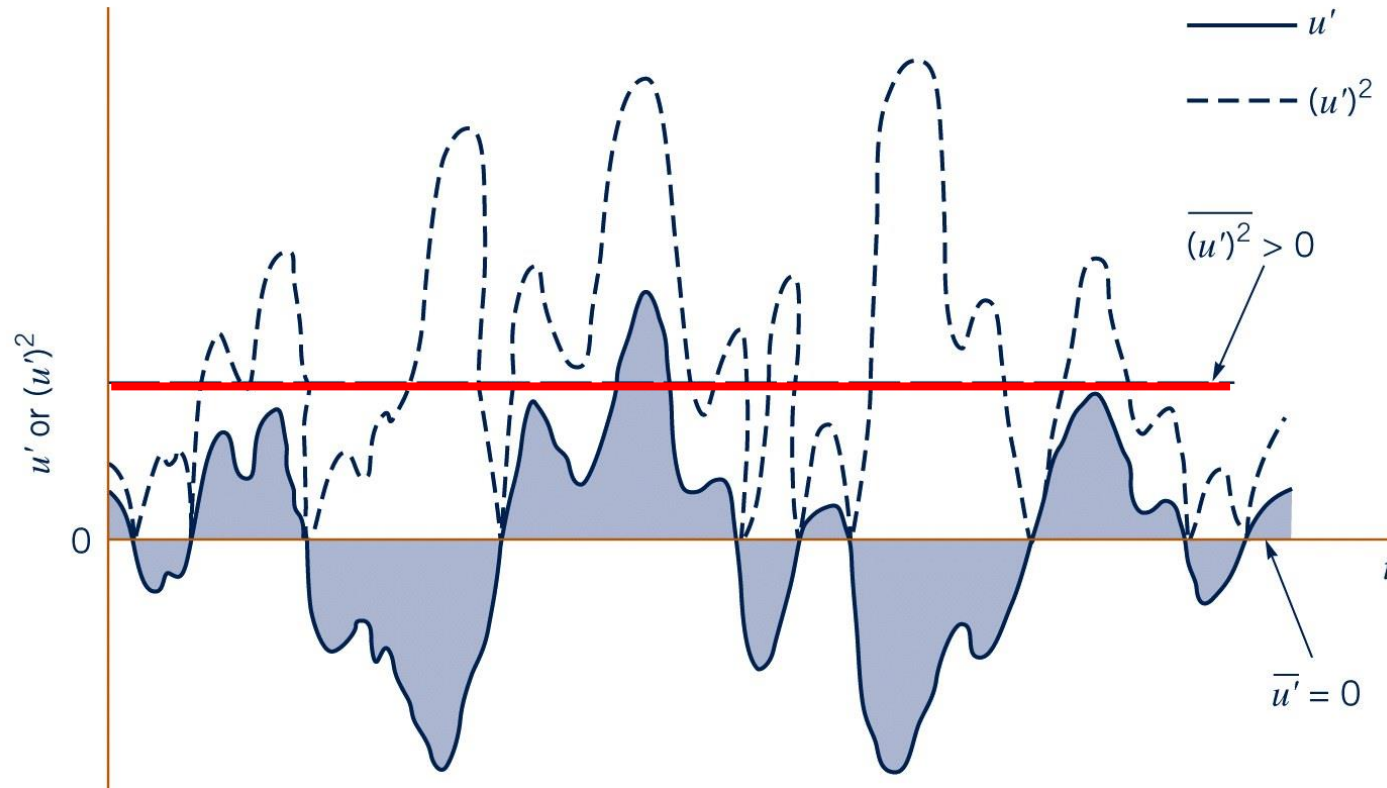
$$TI = \tilde{u} = \sqrt{\overline{u'^2}} = \left\{ \frac{1}{T} \int_0^T u'^2 dt \right\}^{\frac{1}{2}} \quad (20.8)$$

• Relative Turbulence Intensity (RTI) =  $\frac{\sqrt{\overline{u'^2}}}{\bar{u}}$





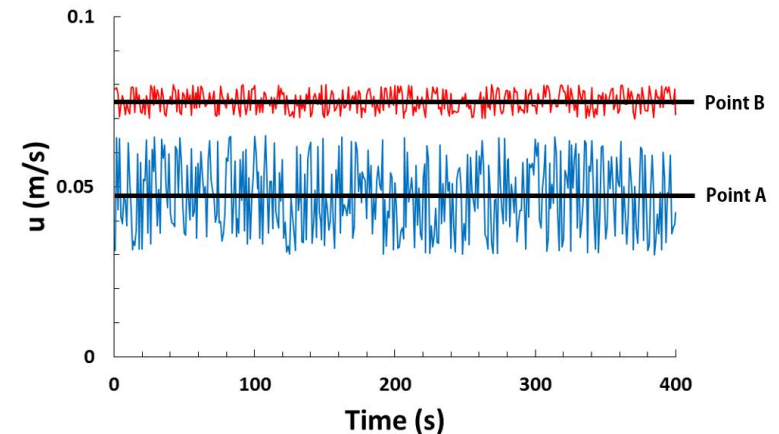
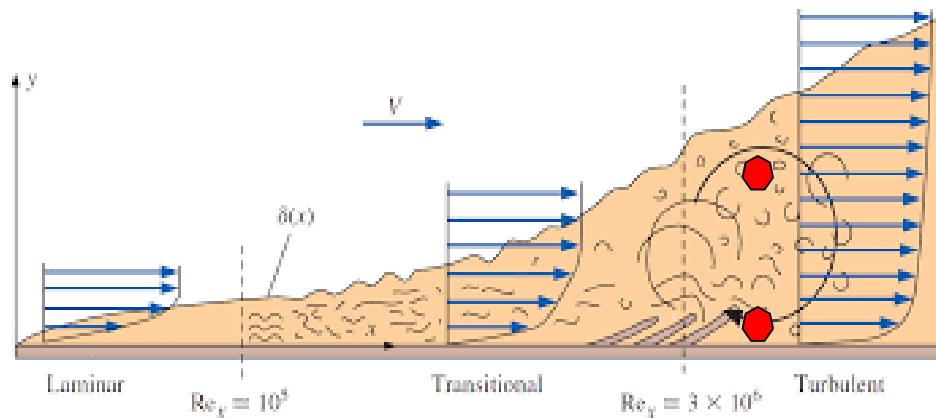
# 20.1 Velocities and Energies in Turbulence



# 20.1 Velocities and Energies in Turbulence

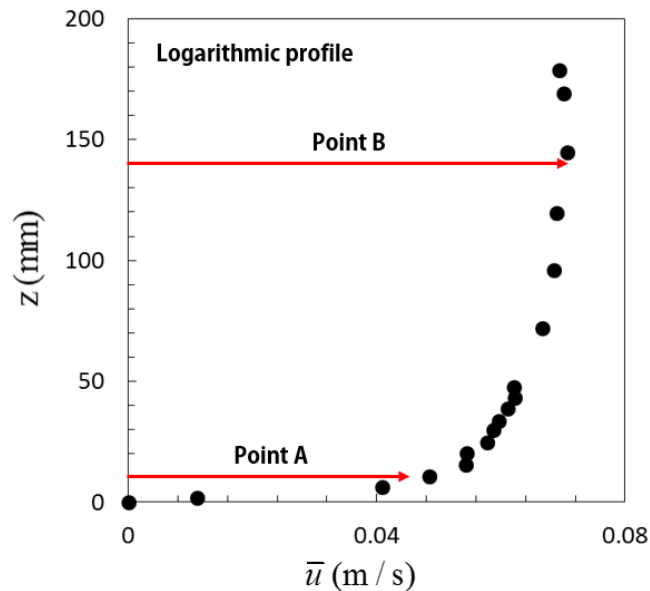
[Ex] Actual velocity records obtained at two depths in the open channel flow (stationary turbulence)

- The time-averaged velocity is greater farther from the wall, but the turbulence intensity is significantly larger near the wall.

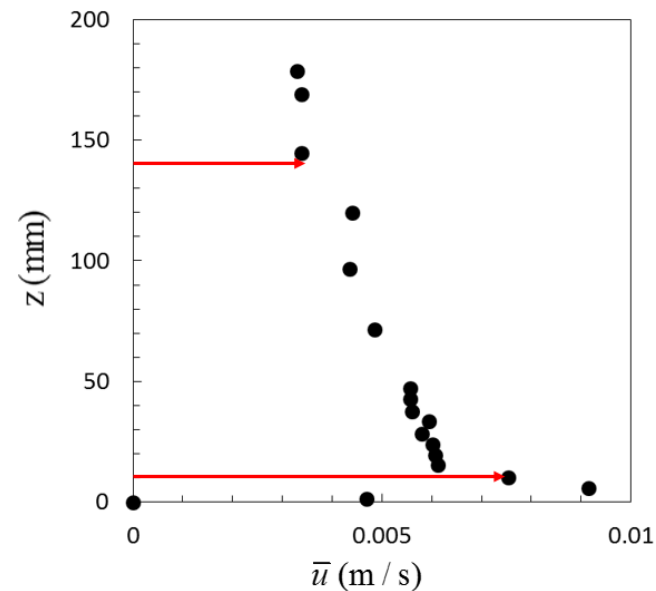


# 20.1 Velocities and Energies in Turbulence

- The time-averaged velocity increases monotonically from zero at the wall to be approximated into the logarithmic profile.
- The turbulence intensity increases rapidly from zero at the wall to a local maximum near the wall and then monotonically decreases.



a) Mean velocity



b) Turbulence intensity

## 20.1 Velocities and Energies in Turbulence

(4) Average kinetic energy of turbulence per unit mass

~ average KE of turbulence / mass

$$TKE = \frac{1}{2} (\overline{u'^2} + \overline{v'^2} + \overline{w'^2}) = \frac{1}{2} \sum (\textit{intensity})^2 \quad (20.9)$$

(5) Energy density,  $\phi(f)$

The kinetic energy is decomposed into an energy spectrum (density) vs. frequency.

≡ limit of average kinetic energy per unit mass divided by the bandwidth  $\Delta f$

$$\phi(f) = \lim_{\Delta f \rightarrow 0} \frac{\textit{average KE / mass contained in } \Delta f}{\Delta f} = \frac{\partial KE}{\partial f} \quad (20.10)$$

## 20.1 Velocities and Energies in Turbulence

where  $f$  = ordinary frequency in cycles per second

$$\therefore \text{average KE of turbulence / mass} = \int_0^{\infty} \phi(f) df = \frac{1}{2} (\overline{u'^2} + \overline{v'^2} + \overline{w'^2})$$

(6) Correlation between  $u'$ ,  $v'$ , and  $w'$

exact correlation = one-to-one correlation

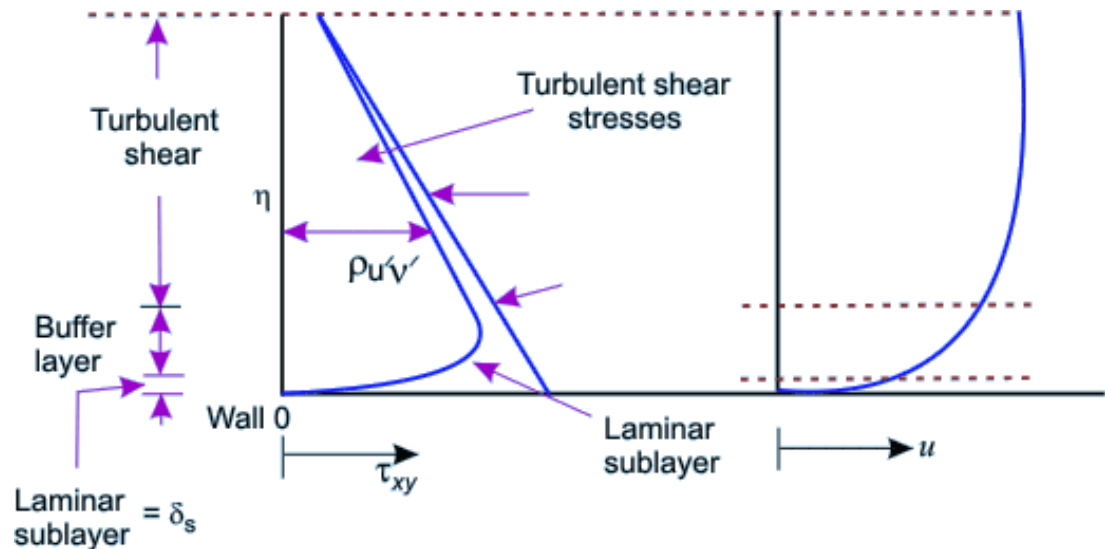
zero correlation = completely independent

$$\overline{u'v'} = \frac{1}{T} \int_0^T u'v' dt \quad \left[ \begin{array}{l} \neq 0 \text{ correlated} \\ = 0 \text{ uncorrelated} \end{array} \right. \quad (20.11)$$

# 20.1 Velocities and Energies in Turbulence

- In a shear flow in an  $xy$ -plane,  $\overline{u'v'}$  is finite, and it is related to the magnitude of the turbulent shear stress.

$$\tau_{tur} = -\rho \overline{u'v'}$$



# 20.1 Velocities and Energies in Turbulence

[Re] Correlated variables

1) Averages of products  $u$

$$\begin{aligned}
 \overline{u_i u_j} &= \overline{(\bar{u}_i + u_i')( \bar{u}_j + u_j' )} \\
 &= \overline{\bar{u}_i \bar{u}_j} + \overline{u_i' u_j'} + \cancel{\overline{\bar{u}_i u_j'}} + \cancel{\overline{\bar{u}_j u_i'}} \\
 &= \overline{\bar{u}_i \bar{u}_j} + \overline{u_i' u_j'} \qquad (20.12)
 \end{aligned}$$

If  $\overline{u_i' u_j'} \neq 0 \rightarrow u_i'$  and  $u_j'$  are said to be correlated.

If  $\overline{u_i' u_j'} = 0 \rightarrow u_i'$  and  $u_j'$  are uncorrelated.

# 20.1 Velocities and Energies in Turbulence

## 2) Correlation coefficient

$$c_{ij} = \frac{\overline{u_i' u_j'}}{\left(\overline{u_i'^2} \cdot \overline{u_j'^2}\right)^{1/2}} \quad (20.13)$$

in which  $\overline{u_i'^2}$ ,  $\overline{u_j'^2}$  = variances

$c_{ij} = \pm 1 \rightarrow$  perfect correlation (같이 움직임)

$c_{ij} = 0 \rightarrow$  independent (독립적 운동)  $c_{ij} = 0.45 \sim 0.55$  for general turbulence

## [Re] Classification of turbulence

### 1) General turbulence

$$\bar{u} \neq \bar{v} \neq \bar{w}$$

$$\overline{u'^2} \neq \overline{v'^2} \neq \overline{w'^2} \quad \overline{u'v'} \neq \overline{v'w'} \neq \overline{w'u'}$$



# 20.1 Velocities and Energies in Turbulence

## 2) Homogeneous turbulence

~ statistically independent of the location

$$\overline{(u_i' u_j')} _a = \overline{(u_i' u_j')} _b$$

## 3) Isotropic turbulence

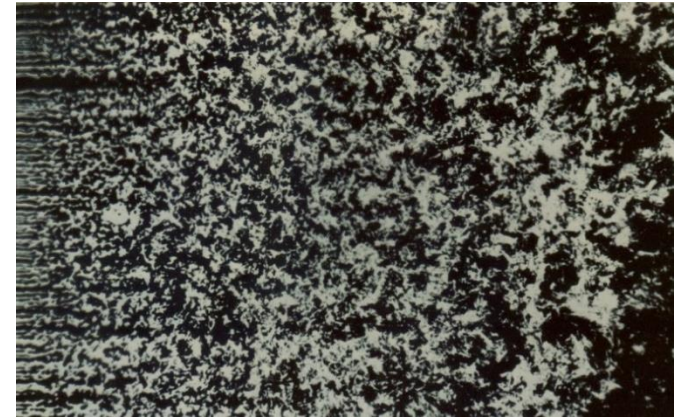
~ statistically independent of the orientation and location of the coordinate axes

$$\overline{u'^2} = \overline{v'^2} = \overline{w'^2} = \text{constant}$$

$$\overline{u'v'} = \overline{v'w'} = \overline{w'u'} = 0$$

~ uncorrelated

~ not coherent structures - small scale eddies



# 20.1 Velocities and Energies in Turbulence

## 20.1.2 Measurement of turbulence

~ measure turbulent fluctuations

(1) Hot-wire (hot-film) anemometer

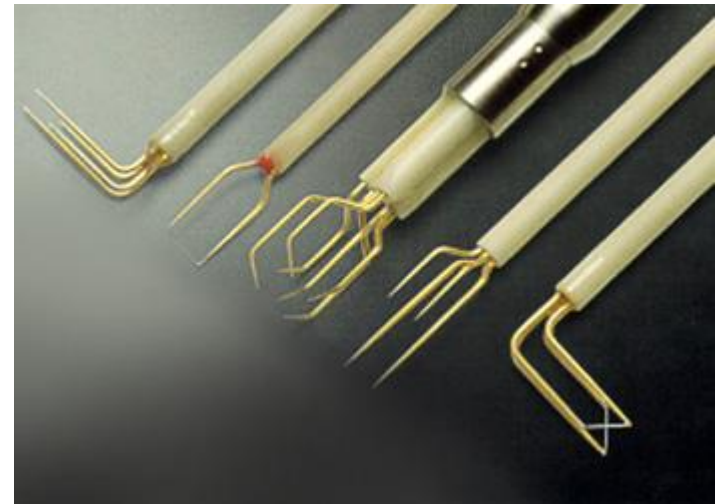
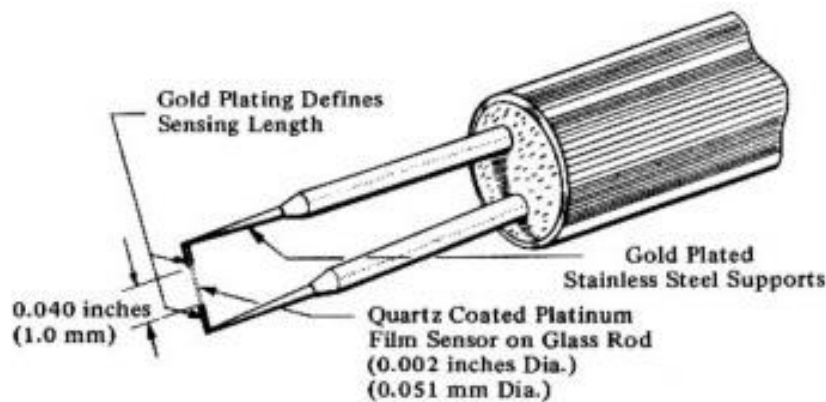
~ Hot-film is usable in contaminated water.

~ Change of temperature affects the electric current flow or voltage drop through wire. Fine platinum wire (film) is heated electrically by a circuit that maintains voltage drop constant.

~ When inserted into the stream, the cooling, which is a function of the velocity, can be detected as variations in voltage.

# 20.1 Velocities and Energies in Turbulence

- ~ Use two or more wires at one point in the flow to make simultaneous measurements of different velocity components.
- After subtracting mean value, rms-values, correlations, and energy spectra can be computed using fluctuation.
- These operations can be performed electronically.



# 20.1 Velocities and Energies in Turbulence

## (2) Laser Doppler Velocimeter (LDV)

~ use Doppler effect

~ A laser (ultrasonic) beam transmitted into the fluid will be reflected by impurities or bubbles in the fluid to a receiving sensor at a different frequency.

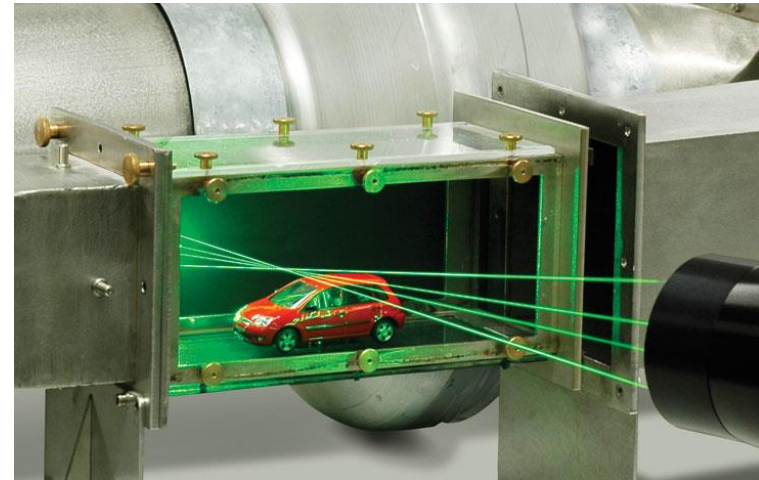
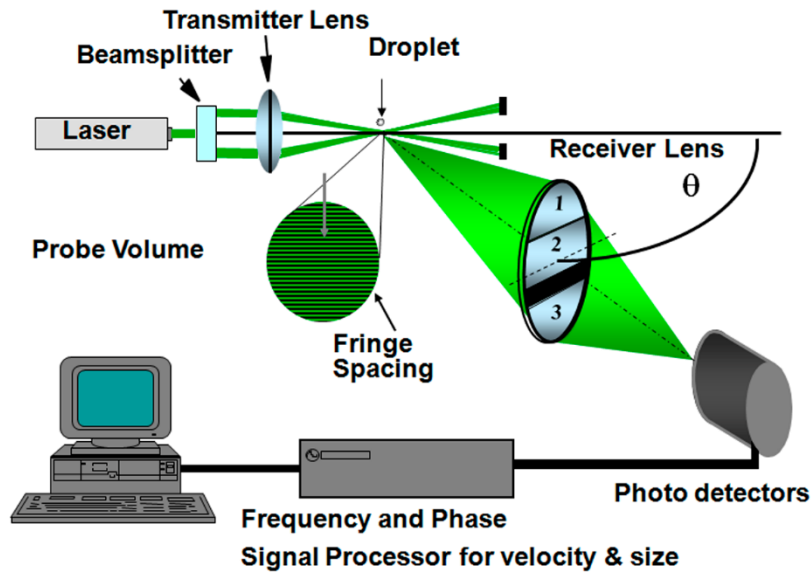
→ The transmitted and reflected signals are then compared by electronic means to calculate the Doppler shift which is proportional to the velocity.

~ non-intrusive sensing (immersion LDA)

~ sampling frequency is up to 20,000 Hz

$$F_{doppler} = -F_{source} \frac{V}{C}$$

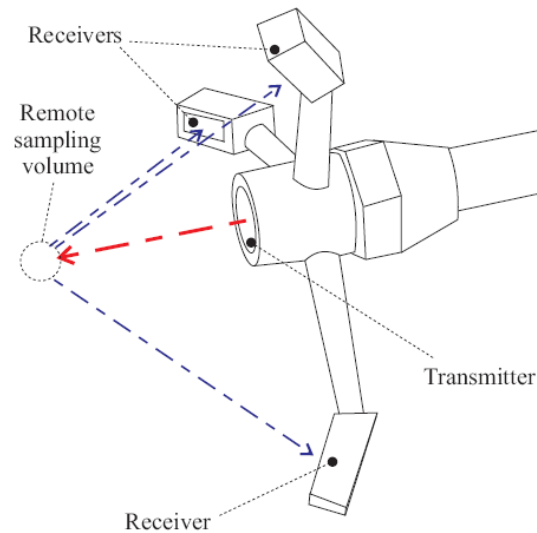
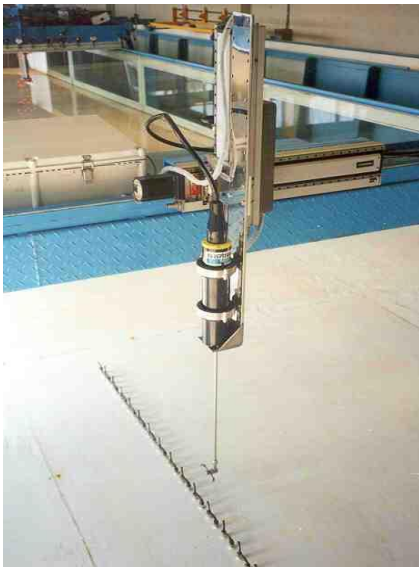
# 20.1 Velocities and Energies in Turbulence



# 20.1 Velocities and Energies in Turbulence

## (3) Acoustic Doppler Velocimeter (ADV)

- ~ use Doppler effect of sonic wave
- ~ intrusive sensing
- ~ sampling frequency = 25-50 Hz

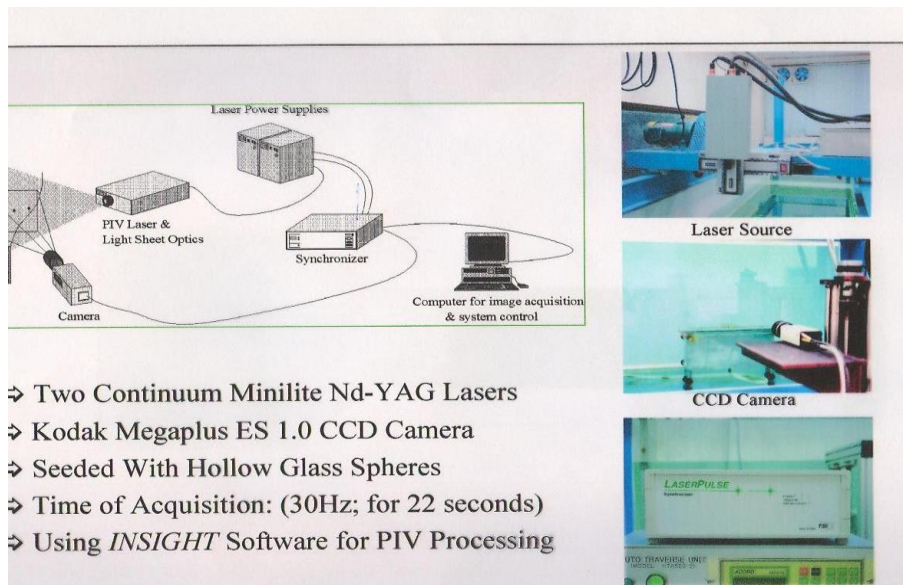




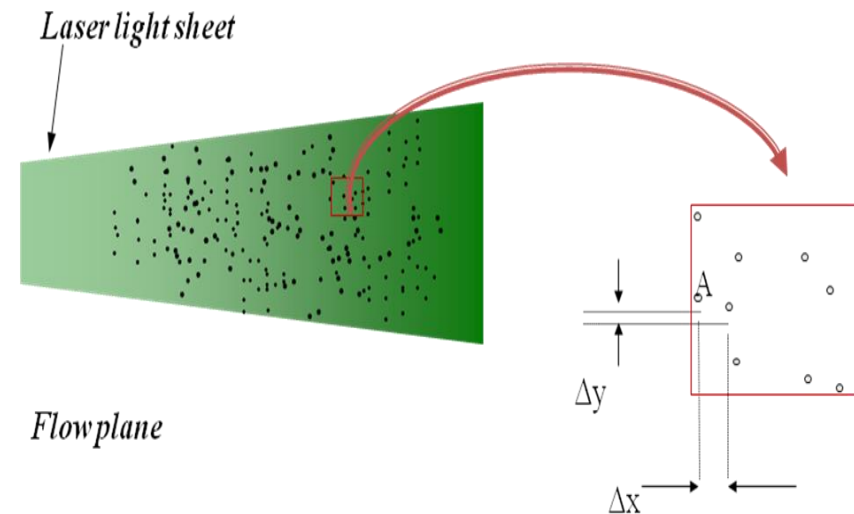
# 20.1 Velocities and Energies in Turbulence

## (4) Particle Image Velocimetry (PIV)

- ~ use Laser and CCD camera
- ~ measure flow field at once
- ~ sampling frequency = 30 Hz

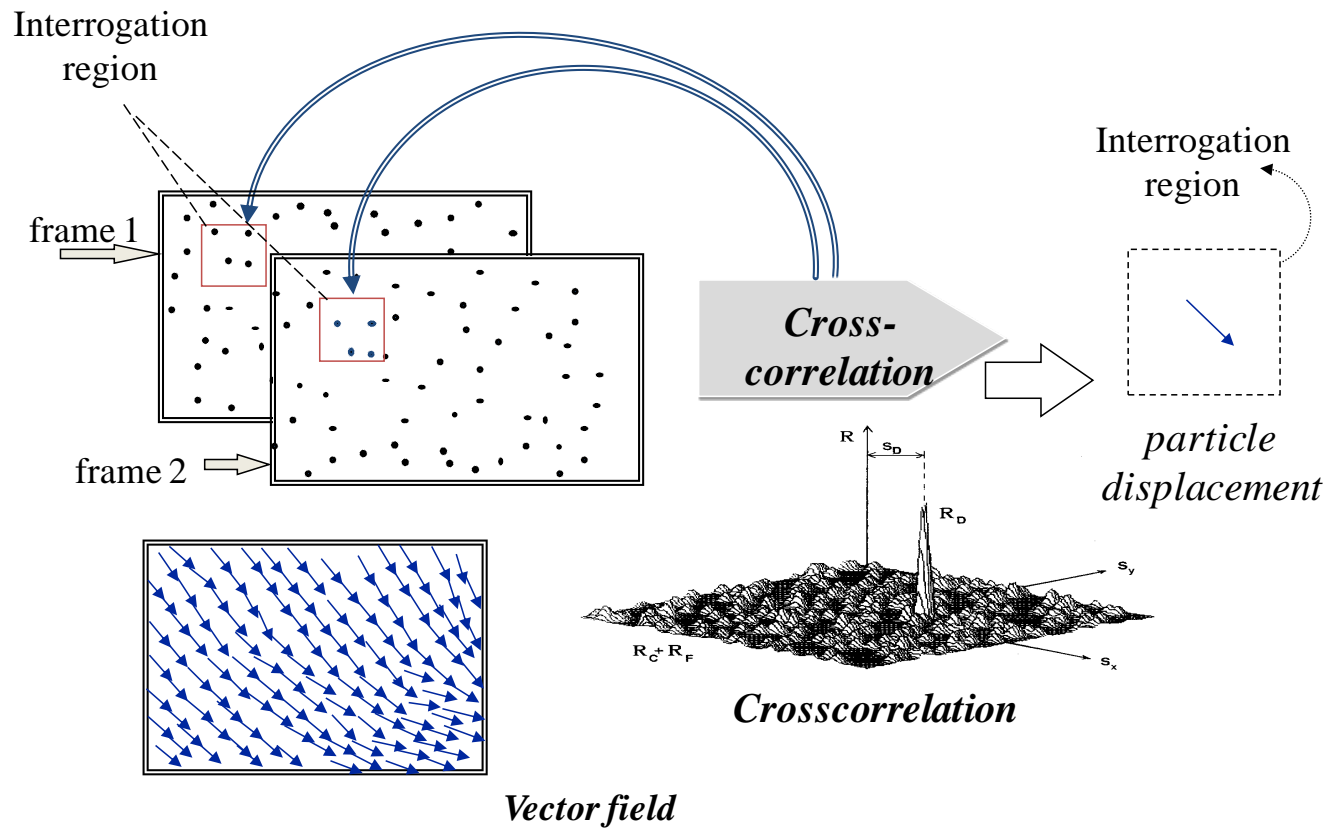


- Two Continuum Minilite Nd-YAG Lasers
- Kodak Megaplus ES 1.0 CCD Camera
- Seeded With Hollow Glass Spheres
- Time of Acquisition: (30Hz; for 22 seconds)
- Using *INSIGHT* Software for PIV Processing



# 20.1 Velocities and Energies in Turbulence

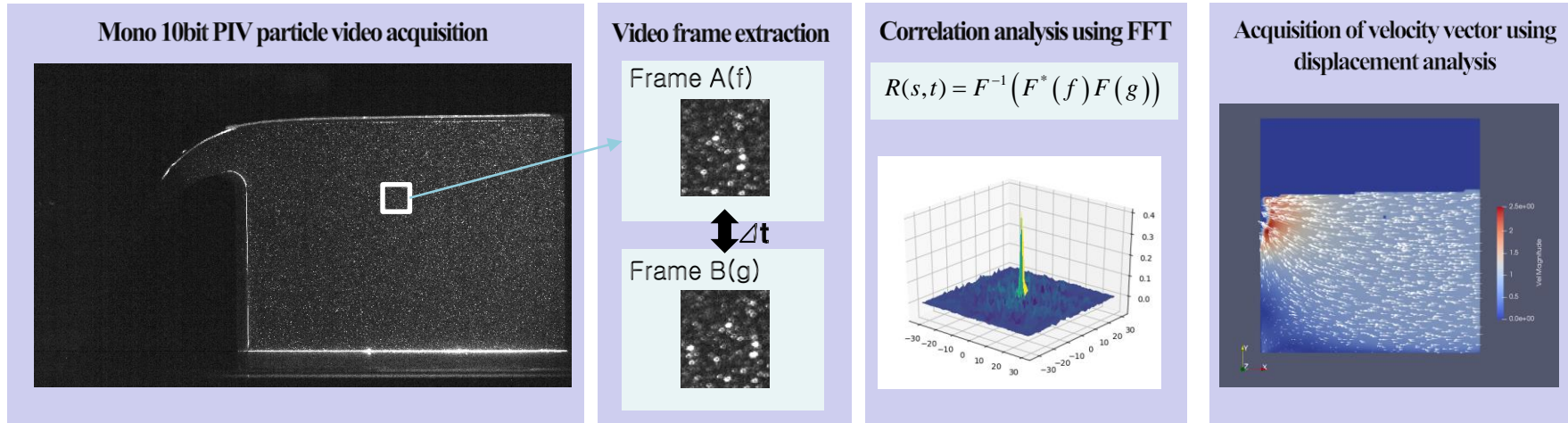
## PIV system



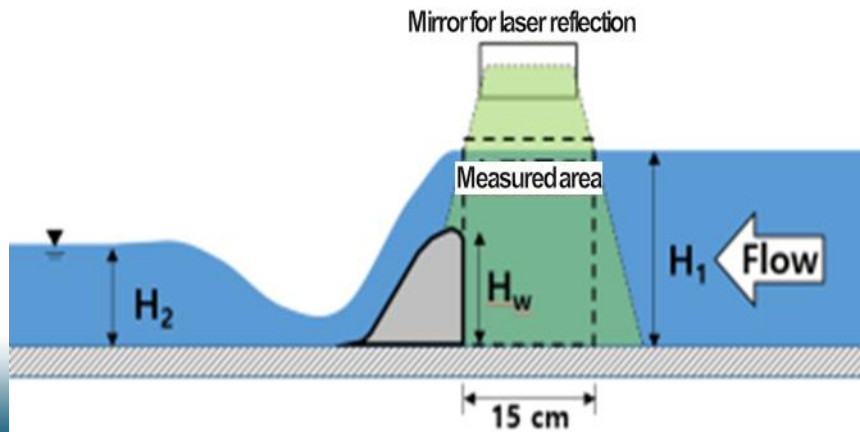


# 20.1 Velocities and Energies in Turbulence

## Velocity measurement using PIV system



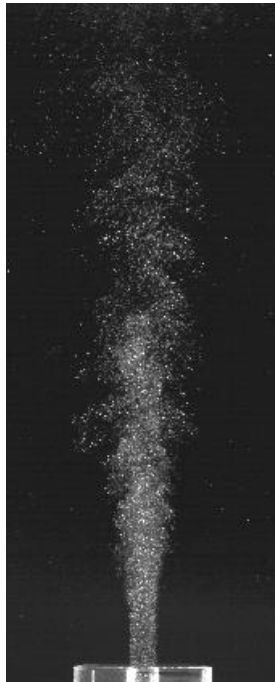
## Velocity fields near ogee weir



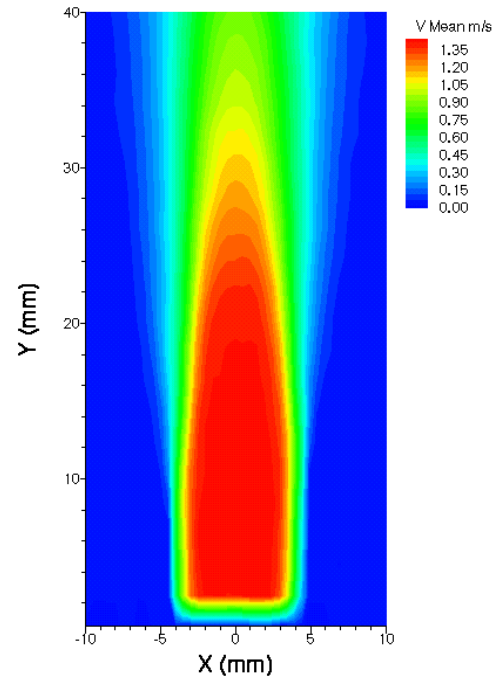
## PIV measurement



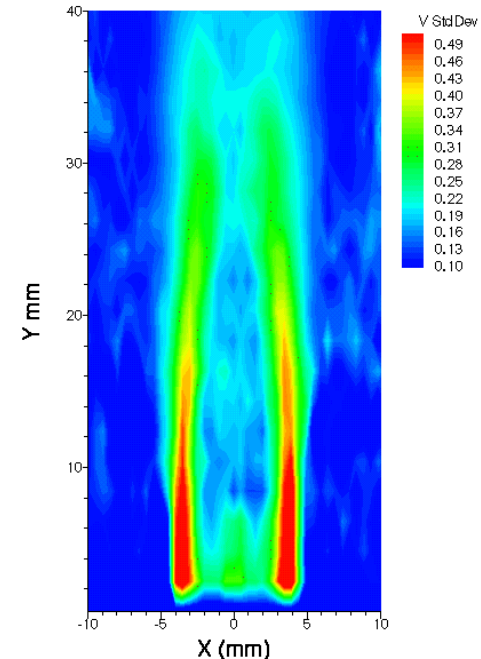
# 20.1 Velocities and Energies in Turbulence



a) Image



b) Velocity



c) Turbulence Intensity

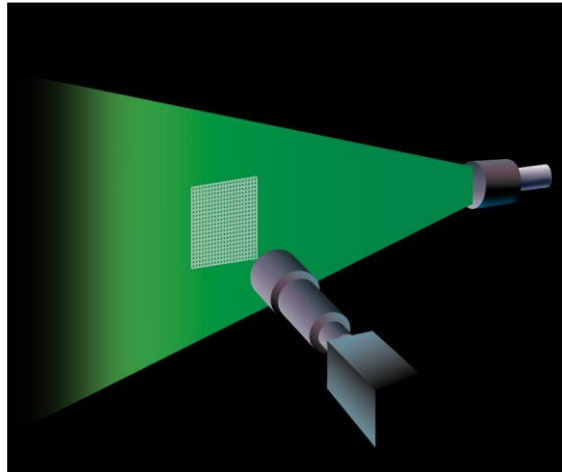
b) Fig. 1 Jet Characteristics Measured by PIV (Seo et al., 2002)

# 20.1 Velocities and Energies in Turbulence

LDV: single point measurement



PIV: field measurement



# 20.1 Velocities and Energies in Turbulence

[Re] Reynolds rules of averages: Schlichting (1979) Boundary-Layer Theory

Let  $f$  and  $g$  are two dependent variables whose time mean values are to be found.  $s$  is any one of the independent variables  $x, y, z, t$ .

$$\overline{\overline{f}} = \overline{f}$$

$$\overline{f + g} = \overline{f} + \overline{g}$$

$$\overline{f \cdot g} = \overline{f} \cdot \overline{g}$$

$$\frac{\partial \overline{f}}{\partial s} = \overline{\frac{\partial f}{\partial s}} \rightarrow \left\{ \begin{array}{l} \text{since time averaging is carried out by integrating over a long} \\ \text{period of time, which commutes with differentiation with respect} \\ \text{to another independent variable} \end{array} \right.$$

$$\overline{\int f ds} = \int \overline{f} ds$$

## 20.2 Continuity for Turbulent Motion

### 20.2.1 Continuity equation for turbulent motion

Continuity equation for incompressible fluid

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} = 0 \quad (\text{A})$$

Substitute velocity decomposition into (A)

$$\frac{\partial(\bar{u} + u')}{\partial x} + \frac{\partial(\bar{v} + v')}{\partial y} + \frac{\partial(\bar{w} + w')}{\partial z} = 0 \quad (20.14)$$

$$\frac{\partial \bar{u}}{\partial x} + \frac{\partial \bar{v}}{\partial y} + \frac{\partial \bar{w}}{\partial z} + \frac{\partial u'}{\partial x} + \frac{\partial v'}{\partial y} + \frac{\partial w'}{\partial z} = 0 \quad (\text{B})$$

## 20.2 Continuity for Turbulent Motion

Take time-averages of each term of (B)

$$\frac{\overline{\partial u}}{\partial x} + \frac{\overline{\partial v}}{\partial y} + \frac{\overline{\partial w}}{\partial z} + \frac{\overline{\cancel{\partial u'}}}{\cancel{\partial x}} + \frac{\overline{\cancel{\partial v'}}}{\cancel{\partial y}} + \frac{\overline{\cancel{\partial w'}}}{\cancel{\partial z}} = 0$$

$$\left( \because \frac{\overline{\partial u'}}{\partial x} = \frac{\partial(\overline{u'})}{\partial x} = 0 \right)$$

$$\frac{\partial \bar{u}}{\partial x} + \frac{\partial \bar{v}}{\partial y} + \frac{\partial \bar{w}}{\partial z} = 0$$

(20.15)

Substitute (20.15) into (B)

$$\frac{\partial u'}{\partial x} + \frac{\partial v'}{\partial y} + \frac{\partial w'}{\partial z} = 0$$

(20.16)

→ Both mean-motion components and the superposed turbulent-motion components must satisfy the continuity equation.

→ Continuity must be satisfied for both turbulent and laminar motions.

## 20.2 Continuity for Turbulent Motion

[Re] Continuity Eq. for compressible fluid

$$\frac{\partial \rho}{\partial t} + \frac{\partial \rho u_i}{\partial x_i} = 0$$

$$\frac{\partial(\bar{\rho} + \rho')}{\partial t} + \frac{\partial\{(\bar{\rho} + \rho')(\bar{u}_i + u_i')\}}{\partial x_i} = 0$$

Time averaging yields

$$\overline{\frac{\partial(\bar{\rho} + \rho')}{\partial t}} + \overline{\frac{\partial\{(\bar{\rho} + \rho')(\bar{u}_i + u_i')\}}{\partial x_i}} = 0$$

$$\overline{\frac{\partial \bar{\rho}}{\partial t}} + \cancel{\overline{\frac{\partial \rho'}{\partial t}}} + \frac{\partial}{\partial x_i} \overline{(\bar{\rho} \bar{u}_i + \cancel{\rho' \bar{u}_i} + \cancel{\bar{\rho} u_i'} + \rho' u_i')} = 0$$

$$\overline{\frac{\partial \bar{\rho}}{\partial t}} + \frac{\partial}{\partial x_i} \overline{(\bar{\rho} \bar{u}_i + \rho' u_i')} = 0$$

$$\overline{\frac{\partial \bar{\rho}}{\partial t}} + \frac{\partial \bar{\rho} \bar{u}}{\partial x} + \frac{\partial \bar{\rho} \bar{v}}{\partial y} + \frac{\partial \bar{\rho} \bar{w}}{\partial z} + \frac{\partial}{\partial x} \overline{(\rho' u')} + \frac{\partial}{\partial y} \overline{(\rho' v')} + \frac{\partial}{\partial z} \overline{(\rho' w')} = 0 \quad (20.17)$$