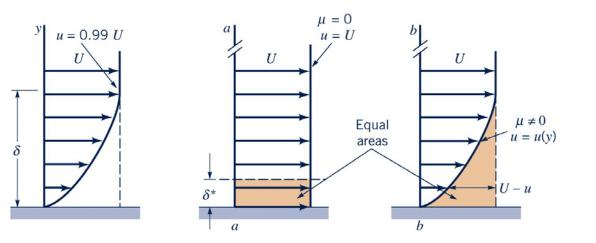


# **Turbulent Boundary-Layer Flows**







# **Chapter 9 Turbulent Boundary-Layer Flows**

#### Contents

- 9.1 Introduction
- 9.2 Structure of a turbulent boundary layer
- 9.3 Mean-flow characteristics for turbulent boundary layer

#### Objectives

- Study wall turbulence
- Derive equations of velocity distribution and friction coefficient for both smooth and rough walls





- Turbulence occurs most commonly in shear flows.
- Shear flow: spatial variation of the mean velocity
- wall turbulence: along solid surface → no-slip condition at surface
   free turbulence: at the interface between fluid zones having
   different velocities, and at boundaries of a jet → jet, wakes
- Turbulent motion in shear flows
- self-sustaining
- Turbulence arises as a consequence of the shear.
- Shear persists as a consequence of the turbulent fluctuations.
- $\rightarrow$  Turbulence can neither arise nor persist without shear.





#### 9.2.1 Boundary layer flows

- (i) Smooth boundary
- Consider a fluid stream flowing past a smooth boundary.
- $\rightarrow$  A boundary-layer zone of viscous influence is developed near the boundary.

1)  $\text{Re} < \text{Re}_{crit}$ 

→ The boundary-layer is initially laminar.

 $\rightarrow$  u = u(y)

2)  $\text{Re} > \text{Re}_{crit}$ 

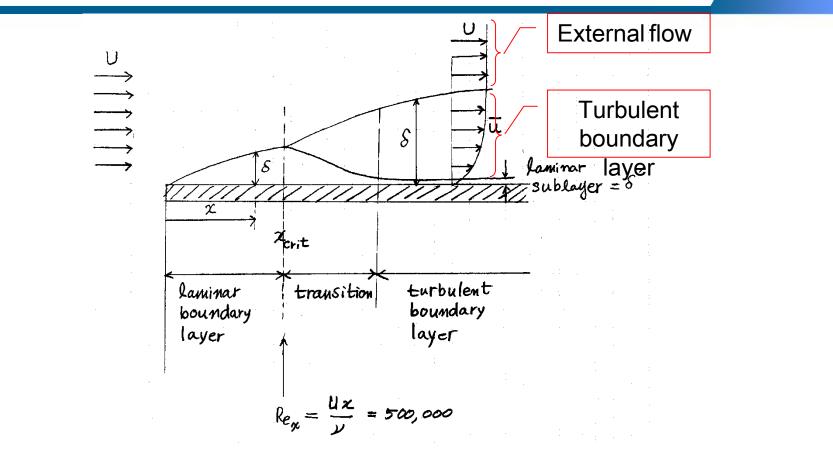
 $\rightarrow$  The boundary-layer is turbulent.

 $\rightarrow \overline{u} = \overline{u}(y)$ 

- $\rightarrow$  Turbulence reaches out into the free stream to entrain and mix more fluid.
- → thicker boundary layer:







- $x < x_{crit}$ , total friction = laminar
- $x > x_{crit}$ , total friction = laminar + turbulent

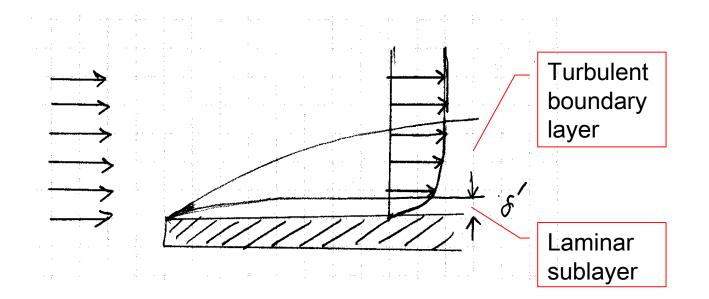




(ii) Rough boundary

 $\rightarrow$  Turbulent boundary layer is established near the <u>leading edge</u> of

the boundary without a preceding stretch of laminar flow.







9.2.2 Comparison of laminar and turbulent boundary-layer profiles

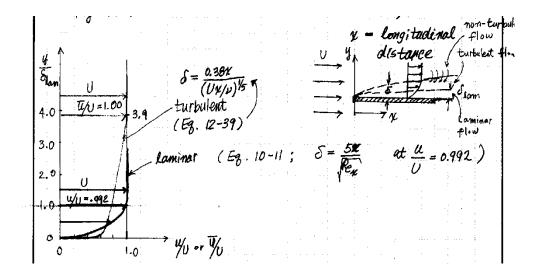
Laminar boundary layer:

study Ch. 10 (D&H)

For the flows of the same Reynolds number

 $\text{Re}_{x} = 500,000$ 

$$\operatorname{Re}_{x} = \frac{Ux\rho}{\mu}$$







1) Boundary layer thickness

$$\frac{\delta_{turb}}{\delta_{lam}} = 3.9$$

2) Mass displacement thickness,

Eq. (8.9): 
$$\delta^* = \int_0^h (1 - \frac{u}{U}) dy$$
$$\frac{\delta^*_{turb}}{\delta^*_{lam}} = 1.41$$

3) Momentum thickness,  $\theta$ 

Eq. (8.10): 
$$\theta = \int_0^h \frac{u}{U} (1 - \frac{u}{U}) dy$$
$$\frac{\theta_{turb}}{\theta_{lam}} = 2.84$$

→ Because of the higher flux of mass and momentum through the zone nearest the

<u>wall</u> for turbulent flow, increases of and  $\theta$  rate are not as large as  $\delta$ .





#### 9.2.3 Intermittent nature of the turbulent layer

- Outside a boundary layer
- $\rightarrow$  free-stream shearless flow (*U*)  $\rightarrow$  potential flow (inviscid)
- $\rightarrow$  slightly turbulent flow
- → considered to be <u>non-turbulent flow</u> relative to higher turbulence inside a turbulent boundary layer
- Interior of the turbulent boundary layer ( $\delta$ )
- ~ consist of regions of different types of flow (laminar, buffer, turbulent)
- <u>Instantaneous border</u> between turbulent and non-turbulent fluid is <u>irregular and</u> <u>changing</u>.
- ~ Border consists of fingers of turbulence extending into the non-turbulent fluid and fingers of non-turbulent fluid extending deep into the turbulent region.
- ~ intermittent nature of the turbulent layer

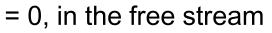


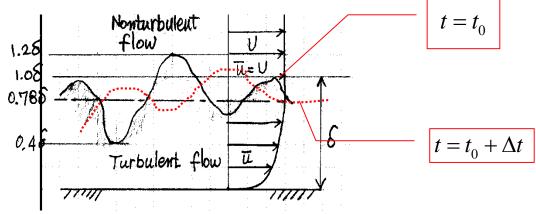


• Intermittency factor,  $\Omega$ 

 $\Omega$  = fraction of time during which the flow is turbulent

 $\Omega$  = 1.0, deep in the boundary layer



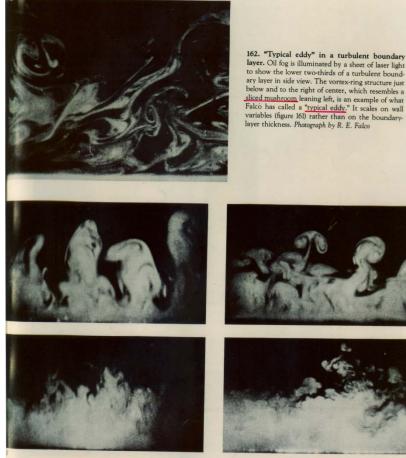


(1) Average position of the turbulent-nonturbulent interface = 0.78  $\delta$ 

②Maximum stretch of interface =  $1.2 \delta$ 

③Minimum stretch of interface = 0.4  $\delta$ 





163. Oblique transverse sections of a turbulent boundary layer. The flow is viewed head-on, with smoke illuminated by a sheet of light that is inclined 45° downstream from the wall on the left and 45° upstream on the right.

The Reynolds number based on momentum thickness is 600 in the upper pair of photographs and 9400 below. Head & Bandyopadhyay 1981





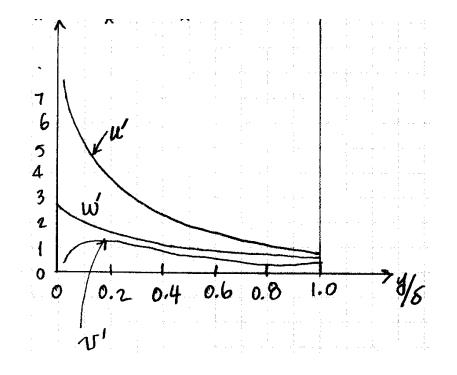
• Turbulent energy in a boundary layer,  $\delta$ 

- Dimensionless energy = 
$$\frac{\overline{u^{12}} + \overline{v^{12}} + \overline{w^{12}}}{u_*^2}$$
(9.1)  
where  $u_* = \sqrt{\frac{\tau_0}{\rho}}$  = shear velocity  
 $\operatorname{Re}_{\delta} = \frac{U\delta}{v} = 73,000 \Leftrightarrow \operatorname{Re}_{x} = 4 \times 10^{6}$  for turbulent layer  
 $u_*^{0}$  for turbulent energy average for both turbulent energy average for turbulent  
 $u_*^{0}$  regions only  
 $u_*^{0}$  for turbulent energy average for turbulent  
 $u_*^{0}$  for turbulent  





$$\frac{\overline{u'^{2}}}{\Omega u'^{2}_{*}}, \frac{\overline{v'^{2}}}{\Omega u'^{2}_{*}}, \frac{\overline{w'^{2}}}{\Omega u'^{2}_{*}}$$



- smooth wall  $\rightarrow v' = 0$  at wall
- rough wall  $\rightarrow v' \neq 0$  at wall
- smooth & rough wall
- $\rightarrow$  turbulent energy  $\neq$  0 at y =  $\delta$





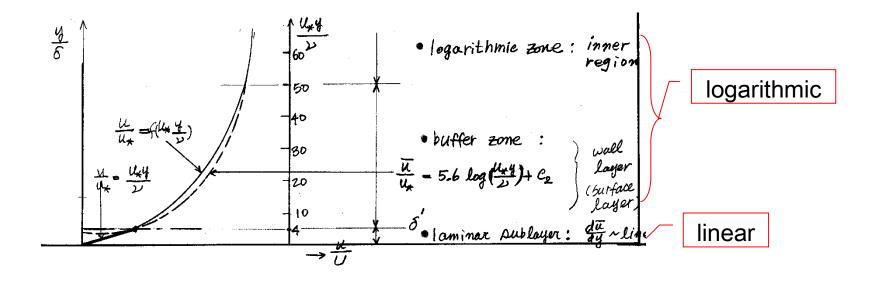
- Relations describing the mean-flow characteristics
- → predict velocity magnitude and relation between velocity and wall
- shear or pressure gradient forces
- $\rightarrow$  It is desirable that these relations should <u>not require knowledge of</u> <u>the turbulence details</u>.
- Turbulent boundary layer
- $\rightarrow$  is composed of zones of different types of flow
- $\rightarrow$  Effective viscosity (  $\mu + \eta\,$  ) varies from wall out through the layer.
- → Theoretical solution is not practical for the general nonuniform boundary layer
- $\rightarrow$  use semiempirical procedure





9.3.1 Universal velocity and friction laws: smooth walls

(1) Velocity-profile regions







1) laminar sublayer: 
$$0 < \frac{u_* y}{v} \le 4$$
  
 $\frac{d\overline{u}}{dy} \sim \text{linear}$   
 $\rightarrow \frac{u}{u_*} = \frac{u_* y}{v}$ 

- $\sim$  Mean shear stress is controlled by the dynamic molecular viscosity .
- $\rightarrow$  Reynolds stress is negligible.  $\rightarrow$  Mean flow is laminar.
- ~ energy of velocity fluctuation  $\approx 0$
- 2) buffer zone:  $4 < \frac{u_*y}{v} < 30 \sim 70$
- ~ Viscous and Reynolds stress are of the same order.
- $\rightarrow$  Both laminar flow and turbulence flow exist.
- ~ Sharp peak in the turbulent energy occurs (Fig. 9.4).





- 3) turbulent zone inner region:  $\frac{u_*y}{y} > 30 \sim 70$ , and  $y < 0.15\delta$
- ~ fully turbulent flow
- ~ inner law zone/logarithmic law
- ~ Intensity of turbulence decreases.
- ~ velocity equation: logarithmic function

 $\frac{\overline{u}}{u_*} = 5.6 \log \frac{u_* y}{v} + C_2$ 

- 4) turbulent zone-<u>outer region</u>:  $0.15\delta < y < 0.4\delta$
- ~ outer law, velocity-defect law
- 5) intermittent zone:  $0.4\delta < y < 1.2\delta$

~ Flow is intermittently turbulent and non-turbulent.

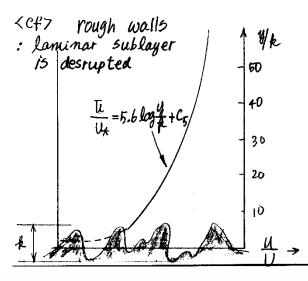




- 6) non-turbulent zone
- ~ external flow zone
- ~ potential flow

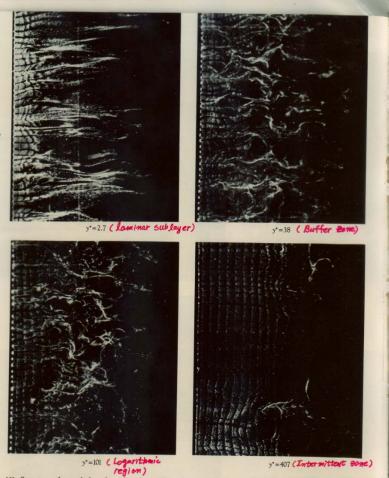
[Cf] Rough wall:

 $\rightarrow$  Laminar sublayer is destroyed by the <u>roughness elements.</u>









161. Structure of a turbulent boundary layer. Successive layers of the flow near a flat plate in a water channel are shown by tim hydrogen bubbles released periodically from a thin platinum wire seen at the left. The height  $\gamma^{a} = \gamma u \cdot \gamma v$  of the wire above the plate is shown in wall variables, where  $u_{-} = (\tau_{a'}/\varrho)^{1/2}$  is the friction velocity. The

characteristic low- and high-speed streaks shown in the viscous sublayer at  $y^+=2.7$  become less noticeable farther away, and have disappeared in the logarithmic region at  $y^+=101$ . In the wake region at  $y^+=407$  the turbulence is seen to be intermittent and of larger scale. Kline, Reynolds, Schraub & Restadler 1967.



(2) Wall law versus velocity-defect law

wall law  $\rightarrow$  inner region

velocity-defect law  $\rightarrow$  outer region

1) Law of wall = inner law;  $\frac{u_*y}{v} > 30 \sim 70$ ;  $y < 0.15\delta$ 

~close to smooth boundaries (molecular viscosity dominant)

~ Law of wall assumes that the relation between wall shear stress and velocity at distance *y* from the wall depends only on fluid density and viscosity;

 $f(\overline{u}\,,\,u_*\,,\,y\,,\,\rho\,,\,\mu)=0$ 

Dimensional analysis yields

$$\frac{\overline{u}}{u_*} = f\left(\frac{u_*y}{v}\right)$$





(9.3)

i) Laminar sublayer

- mean velocity,  $\overline{u} \equiv u$ 

- velcocity gradient, 
$$\frac{\partial u}{\partial y} \sim \text{constant} \equiv \frac{u}{y}$$

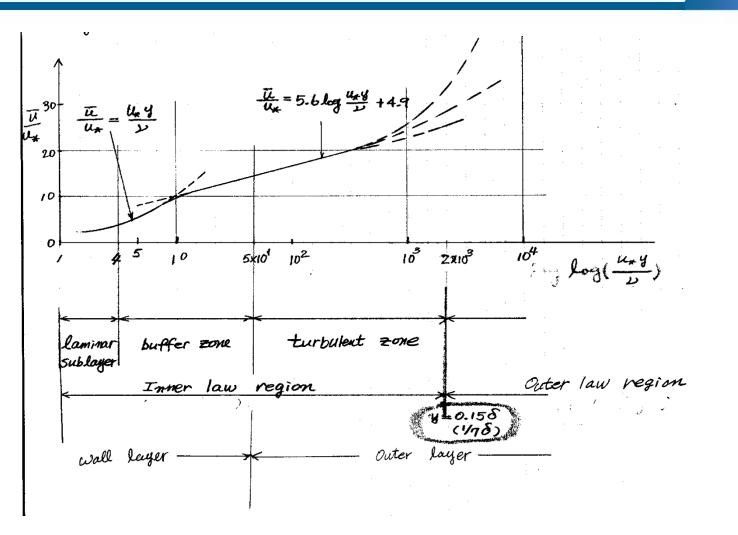
- shear stress, 
$$\tau \approx \tau_0 = \mu \frac{\partial u}{\partial y} \bigg|_{y=0} \equiv \mu \frac{u}{y}$$
 (9.5)

- shear velocity, 
$$u_* = \sqrt{\frac{\tau}{\rho}} \rightarrow u_*^2 = \frac{\tau}{\rho} = \frac{\mu}{\rho} \frac{u}{y}$$

$$\frac{u}{u_*} = \frac{u_* y}{v} \tag{9.6}$$











Thickness of laminar sublayer

define thickness of laminar sublayer ( $\delta'$ ) as the value of y which makes

$$\frac{u_*y}{v} = 4$$

$$\delta' = \frac{4\nu}{u_*} = \frac{4\nu}{\sqrt{\frac{\tau_0}{\rho}}} = \frac{4\nu}{\left(\frac{c_f \rho U^2 / 2}{\rho}\right)^{\frac{1}{2}}} = \frac{4\nu}{U\sqrt{c_f / 2}}$$
(9.7)

where 
$$\tau_0 = c_f \rho \frac{u^2}{2}$$
;  $c_f = local \ shear \ stress \ coeff$ .





Integrate (4)  

$$\overline{u} = \frac{u_*}{\kappa} \ln y + C_1$$

$$\overline{u} = \frac{1}{\kappa}$$

$$\frac{\overline{u}}{u_*} = \frac{1}{\kappa} \ln y + C_1 \tag{9.10}$$

Substitute BC [  $\overline{u} = 0$  at y = y'] into (9.10)

$$0 = \frac{1}{\kappa} \ln y' + C_1$$

$$\therefore C_1 = -\frac{1}{\kappa} \ln y'$$
(5)

Assume  $y' \propto \frac{v(m^2/s)}{u_*(m/s)} \rightarrow y' = C \frac{v}{u_*}$ 

Then (5) becomes

$$C_1 = -\frac{1}{\kappa} \ln y' = -\frac{1}{\kappa} \ln \left( C \frac{\nu}{u_*} \right) = C_2 - \frac{1}{\kappa} \ln \frac{\nu}{u_*}$$





(9.11)

Substitute (9.11) into (9.10)

$$\frac{\overline{u}}{u_*} = \frac{1}{\kappa} \ln y + C_2 - \frac{1}{\kappa} \ln \frac{v}{u_*} = \frac{1}{\kappa} \ln \left(\frac{u_* y}{v}\right) + C_2$$

Changing the logarithm to base 10 yields

$$\frac{\overline{u}}{u_*} = \frac{2.3}{\kappa} \log_{10} \left( \frac{u_* y}{v} \right) + C_2$$
(9.12)

Empirical values of  $\kappa$  and  $C_2$  for inner region of the boundary layer

$$\kappa = 0.41; C_2 = 4.9$$

$$\frac{\overline{u}}{u_*} = 5.6 \log\left(\frac{u_*y}{v}\right) + 4.9, \ 30 \sim 70 < \frac{u_*y}{v}, \ and \ \frac{y}{\delta} < 0.15$$
(9.13)

→ Prandtl's velocity distribution law; inner law; wall law





- 2) Velocity-defect law
- ~outer law

~outer reaches of the turbulent boundary layer for both smooth and rough walls

 $\rightarrow$  Reynolds stresses dominate the viscous stresses.

Assume velocity defect (reduction) at y wall shear stress

$$\frac{U-\overline{u}}{u_*} = g\left(\frac{y}{\delta}\right) \tag{9.14}$$

Substituting BC [] into Eq. (9.12) leads to

$$\frac{U}{u_*} = \frac{2.3}{\kappa} \log\left(\frac{u_*\delta}{v}\right) + C_2^{'}$$
(9.15)





Subtract (9.12) from (9.15)

$$\frac{U}{u_*} = \frac{2.3}{\kappa} \log\left(\frac{u_*\delta}{v}\right) + C_2'$$

$$- \left| \frac{\overline{u}}{u_*} = \frac{2.3}{\kappa} \log\left(\frac{u_*y}{v}\right) + C_2$$

$$\frac{U - \overline{u}}{u_*} = \frac{2.3}{\kappa} \left\{ \log\left(\frac{u_*\delta}{v}\right) - \log\left(\frac{u_*y}{v}\right) \right\} + C_2' - C_2$$

$$= \frac{2.3}{\kappa} \log\left(\frac{u_*\delta}{v}\right) + C_3$$

$$= -\frac{2.3}{\kappa} \log\left(\frac{y}{\delta}\right) + C_3$$

where  $\kappa$  and  $C_3$  are empirical constants





i) Inner region; 
$$\frac{y}{\delta} \le 0.15$$
  
 $\rightarrow \kappa = 0.41, C_3 = 2.5$   
 $\frac{U - \overline{u}}{u_*} = -5.6 \log\left(\frac{y}{\delta}\right) + 2.5$   
ii) Outer region;  $\frac{y}{\delta} > 0.15$ 

$$\rightarrow \kappa = 0.267, \ C_3 = 0$$

$$\frac{U - \overline{u}}{u_*} = -8.6 \log\left(\frac{y}{\delta}\right)$$

(9.17)

(9.16)

 $\rightarrow$  Eqs. (9.16) & (9.17 ) apply to both smooth and rough surfaces.

→ Eq. (9.16) = Eq. (9.13)





Fig. 9.8  $\rightarrow$  velocity-defect law is applicable for both smooth and rough walls

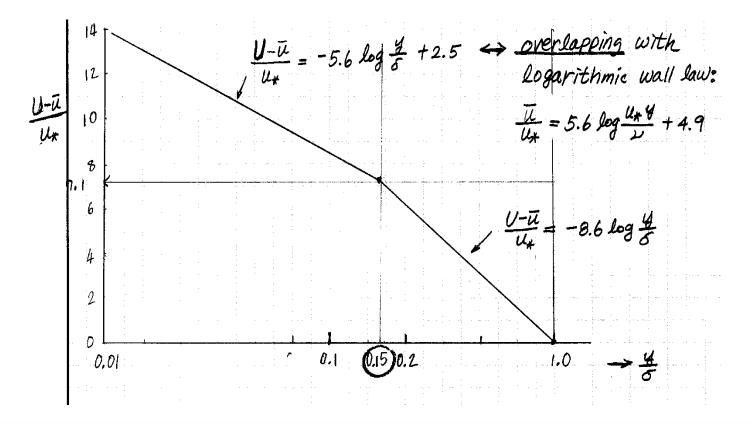
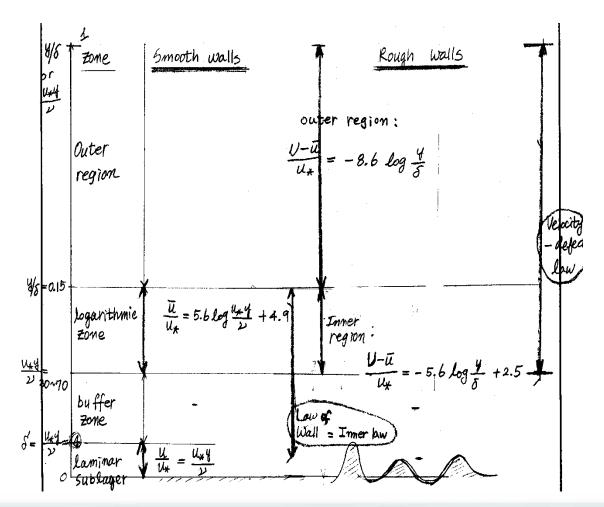






Table 9-1



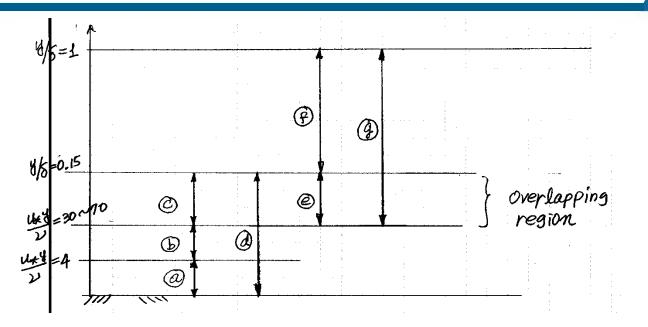




- For smooth walls
- a) laminar sublayer:  $\frac{u}{u_*} = \frac{u_*y}{v}$
- $\rightarrow$  viscous effect dominates.
- b) buffer zone
- c) logarithmic zone:  $\frac{\overline{u}}{u_*} = 5.6 \log \frac{u_* y}{v} + 4.9$
- → turbulence effect dominates.
- d) inner law (law of wall) region







For both smooth and rough walls

e) inner region: 
$$\frac{\overline{U} - \overline{u}}{u_*} = -5.6 \log \frac{y}{\delta} + 2.5$$

f) outer region:  $\frac{\overline{U}-\overline{u}}{u_*} = -8.6 \log \frac{y}{\delta}$ 

g) outer law (velocity - defect law) region





(3) Surface-resistance formulas

1) local shear-stress coefficient on smooth walls

velocity profile ↔ shear-stress equations

$$u_* = \sqrt{\tau / \rho} = \sqrt{\frac{c_f}{2}} U \rightarrow \frac{U}{u_*} = \sqrt{\frac{2}{c_f}}$$

(9.18)

where  $c_f = \underline{\text{local shear - stress coeff.}}$ [Re]  $\tau_0 = \frac{1}{2}\rho c_f U^2 \rightarrow \frac{\tau_0}{\rho} = \frac{c_f}{2} U^2$ 

$$u_* = \sqrt{\frac{\tau_0}{\rho}} = \sqrt{\frac{c_f}{2}} U$$





#### i) Apply logarithmic law

Substituting into @ into Eq. (9.12) yields

$$\frac{U}{u_*} = \frac{2.3}{\kappa} \log\left(\frac{u_*\delta}{v}\right) + C_4 \tag{A}$$

#### Substitute (9.18) into (A)

$$\sqrt{\frac{2}{c_f}} = \frac{2.3}{\kappa} \log\left(\frac{U\delta}{v}\sqrt{\frac{c_f}{2}}\right) + C_4$$
(9.19)

~  $c_f$  is not given explicitly.





ii) For explicit expression, use displacement thickness and momentum thickness  $\theta$  instead of  $\delta$ 

Clauser: 
$$\frac{1}{\sqrt{c_f}} = 3.96 \log \operatorname{Re}_{\delta^*} + 3.04$$
 (9.20)

Squire and Young: 
$$\frac{1}{\sqrt{c_f}} = 4.17 \log \text{Re}_{\theta} + 2.54$$
 (9.21)

V

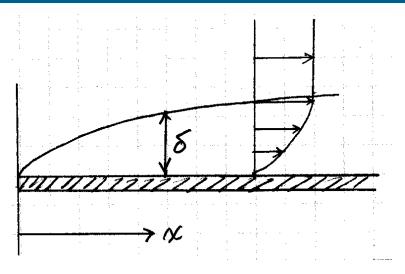
where

**e** 
$$\operatorname{Re}_{\delta^*} = \frac{U\delta^*}{v}$$
;  $\operatorname{Re}_{\theta} = \frac{U\theta}{v}$ 

$$\operatorname{Re}_{\delta^*}, \operatorname{Re}_{\theta} = f(\operatorname{Re}_x), \operatorname{Re}_x = \frac{Ux}{v}$$







iii) Karman's relation

~ assume turbulence boundary layer all the way from the leading

#### edge

(i.e., no preceding stretch of laminar boundary layer)

$$\frac{1}{\sqrt{c_f}} = 4.15\log(\text{Re}_x c_f) + 1.7$$





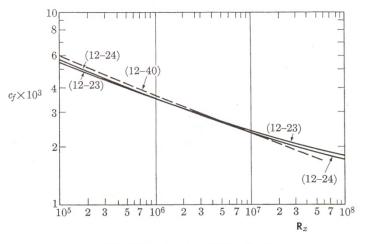


## iv) Schultz-Grunow (1940)

$$c_f = \frac{0.370}{(\log \text{Re}_x)^{2.58}}$$

(9.24)

## Comparison of (9.23) and (9.24)







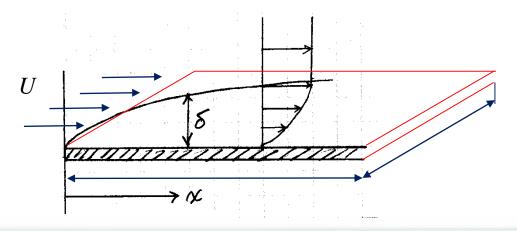


2) Average shear-stress coefficient on smooth walls

Consider average shear-stress coefficient over a distance I along a flat plate of a width b

total drag (D) = 
$$\tau \times bl = \frac{1}{2}C_f \rho U^2 bl$$

$$C_f \equiv \frac{D}{bl\rho U^2/2}$$







i) Schoenherr (1932)

$$\frac{1}{\sqrt{C_f}} = 4.13\log(\operatorname{Re}_l C_f)$$

 $\rightarrow$  implicit

where  $\operatorname{Re}_{l} = \frac{Ul}{v}$ 

ii) Schultz-Grunow

$$C_f = \frac{0.427}{\left(\log \operatorname{Re}_l - 0.407\right)^{2.64}}$$
,  $10^2 < \operatorname{Re}_l < 10^9$  (9.27)





39/76

(9.26)

Comparison of (9.26) and (9.27)

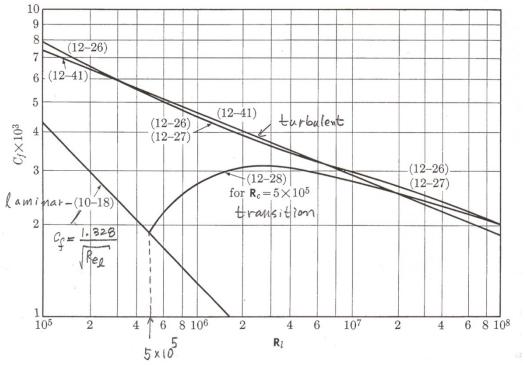
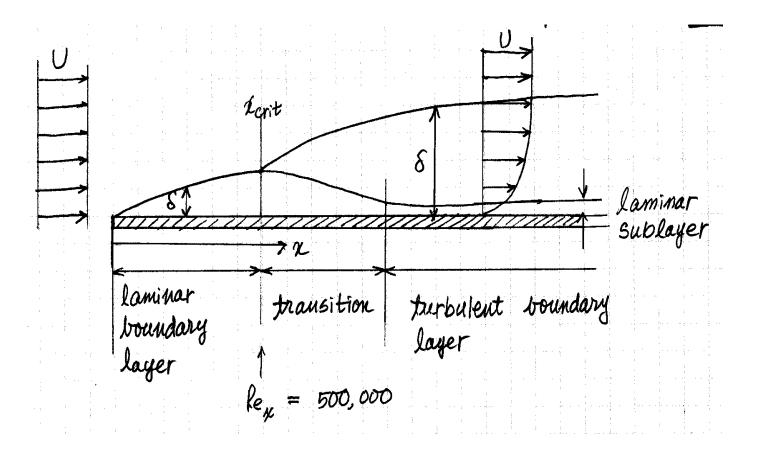


FIG. 12-10. Average coefficient of resistance for flat plates.











- 3) Transition formula
- Boundary layer developing on a smooth flat plate
- ~ At upstream end (leading edge), laminar boundary layer develops because viscous effects dominate due to inhibiting effect of the smooth wall on the development of turbulence.
- ~ Thus, when there is a significant stretch of laminar boundary layer preceding the turbulent layer, total friction is the laminar portion up to  $x_{crit}$  plus the turbulent portion from  $x_{crit}$  to *l*.
- ~ Therefore, average shear-stress coefficient is lower than the prediction by Eqs. (9.26) or (9.27).
- → Use transition formula

$$C_{f} = \frac{0.427}{(\log \operatorname{Re}_{l} - 0.407)^{2.64}} - \frac{A}{\operatorname{Re}_{l}}$$





(9.28)

where  $A/\text{Re}_{l}$  = correction term =  $f(\text{Re}_{crit})$ ,  $\text{Re}_{crit} = \frac{Ux_{crit}}{v}$  $\rightarrow$  A = 1,060~3,340 (Table 9.2, p. 240)

Eq. (9.28) falls between the laminar and turbulent curves. Laminar flow:  $C_f = \frac{1.328}{\text{Re}_l^{1/2}}$ 





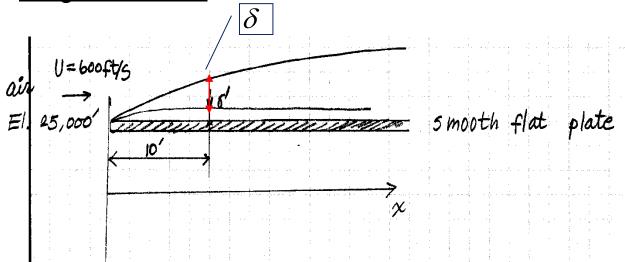
	Smooth walls	Rough walls
LOCAL SHEAR Universal equations		1 8
Clauser (12–20)	$1/\sqrt{c_f} = 3.96 \log \mathbf{R}_{\delta^*} + 3.04$	(12-46) $\frac{1}{\sqrt{c_f}} = 3.96 \log \frac{\delta}{k} + C_8$
		$C_8 = f$ (size, shape, and distribution of roughness)
Squire and Young (12-21)	$1/\sqrt{c_f} = 4.17 \log \mathbf{R}_{\theta} + 2.54$	
von Kármán (12–23)	$1/\sqrt{c_f} = 4.15 \log (\mathbf{R}_x c_f) + 1.7$	
Schultz-Grunow (12–24)	$c_f = \frac{0.370}{(\log \mathbf{R}_x)^{2.58}}$	
Power law (12–40)	$c_f = \frac{0.0466}{R_{\delta}^{1/4}} = \frac{0.059}{R_{x}^{1/5}}$	
Average Shear		
Universal equations		
Schoenherr (12–26)	$1/\sqrt{C_f} = 4.13 \log \left(\mathbf{R}_l C_f\right)$	
Schultz-Grunow (12–27)	$C_f = \frac{0.427}{(\log \mathbf{R}_l - 0.407)^{2.64}}$	
Power law (12–41)	$C_f = \frac{0.074}{\mathbf{R}_l^{1/5}}$	1 A
Transition formula		
Schultz-Grunow-Prandtl (12–28)	$C_f = \frac{0.427}{(\log \mathbf{R}_l - 0.407)^{2.64}} - $	$\frac{A}{\mathbf{R}_{l}}$
	$A = f(\mathbf{R}_{crit}) \text{ as given}$ in Table 12–2	





[Ex. 9.1] Turbulent boundary-layer velocity and thickness An aircraft flies at 25,000 ft with a speed of 410 mph (600 ft/s). Compute the following items for the boundary layer at a distance 10 ft from the leading

edge of the wing of the craft.







(a) Thickness  $\delta'$  (laminar sublayer) at x = 10 ftAir at El. 25,000 ft:  $\nu = 3 \times 10^{-4} ft^2 / s$ 

$$\rho = 1.07 \times 10^{-3} \ slug \ / \ ft^3$$

Select  $\operatorname{Re}_{crit} = 5 \times 10^5 = \frac{600(x)}{3 \times 10^{-4}}$ 

 $\therefore x_{crit} = 0.25 ft \sim \text{negligible comparaed to } l = 10 ft$ 

Therefore, assume that <u>turbulent boundary layer</u> develops all the way from the leading edge. Use Schultz-Grunow Eq., (9.24) to compute  $C_f$  $\operatorname{Re}_x = \frac{Ux}{V} = \frac{600(10)}{(3 \times 10^{-4})} = 2 \times 10^7$ 





(b) Velocity 
$$\bar{u}$$
 at  $y = \delta'$   
Eq. (9.6):  $\frac{u}{u_*} = \frac{u_* y}{v}$   
 $\therefore u = \frac{u_*^2 \delta'}{v} = \frac{(19.8)^2 (0.61 \times 10^{-4})}{(3 \times 10^{-4})} = 79.7 \ ft / s \rightarrow 13\% \ of \ U$   
 $\uparrow \qquad [Cf] \ U = 600 \ ft / s$   
 $y = \delta'$ 

(c) Velocity  $\overline{u}$  at  $y / \delta = 0.15$ 

Use Eq. (9.16) – outer law

$$\frac{U - \overline{u}}{u_*} = -5.6 \log\left(\frac{y}{\delta}\right) + 2.5$$
  
$$\frac{600 - \overline{u}}{19.8} = -5.6 \log(0.15) + 2.5$$
  
$$\overline{u} = 600 - 91.35 - 49.5 = 459.2 \text{ ft / s} \rightarrow 76\% \text{ of } U$$



$$\begin{bmatrix} Cf \end{bmatrix} \quad \overline{u} = U + 5.6 \ u_* \ \log\left(\frac{y}{\delta}\right) - 2.5u_*$$

(d) Distance y at  $y / \delta = 0.15$  and thickness  $\delta$ 

Use Eq. (9.13) – inner law

$$\frac{\overline{u}}{u_*} = 5.6 \log\left(\frac{u_* y}{v}\right) + 4.9$$

$$At \frac{y}{\delta} = 0.15: \frac{459}{19.8} = 5.6 \log\left(\frac{19.8y}{3 \times 10^{-4}}\right) + 4.9$$
$$\log\left(\frac{19.8y}{3 \times 10^{-4}}\right) = 3.26; \quad \frac{19.8y}{3 \times 10^{-4}} = 1839$$
$$y = 0.028' = 0.33in \approx 0.8cm$$





*48/76* 

(B)

## Substitute (B) into

$$\delta = \frac{y}{0.15} = 0.186' = 2.24in \approx 5.7cm$$

At 
$$x = 10'$$
:  $\frac{\delta}{\delta'} = \frac{0.186}{0.61 \times 10^{-4}} = 3049 \approx 3 \times 10^{3}$ 

$$\frac{\delta'}{\delta} = 0.0003$$





[Ex. 9.2] Surface resistance on a smooth boundary given as Ex.9.1 (a) Displacement thickness  $\delta^*$ 

$$\delta^* = \int_0^h \left(1 - \frac{u}{U}\right) dy$$
$$\frac{\delta^*}{\delta} = \int_0^{h/\delta} \left(1 - \frac{\overline{u}}{U}\right) d\left(\frac{y}{\delta}\right), \quad h/\delta \ge 1$$
(A)

Neglect laminar sublayer and approximate buffer zone with Eq. (9.16)

(i) 
$$y/\delta < 0.15$$
,  $\frac{U-\overline{u}}{u_*} = -5.6\log\left(\frac{y}{\delta}\right) + 2.5 \leftarrow (9.16)$   
$$\therefore 1 - \frac{\overline{u}}{U} = -5.6\frac{u_*}{U}\log\frac{y}{\delta} + 2.5\frac{u_*}{U}$$
(B)





(*ii*) 
$$y/\delta > 0.15$$
,  $\frac{U-\overline{u}}{u_*} = -8.6\log\left(\frac{y}{\delta}\right) \quad \leftarrow (9.17)$   
 $\therefore 1 - \frac{\overline{u}}{U} = -8.6\frac{u_*}{U}\log\left(\frac{y}{\delta}\right)$ 





51/76

(C)

$$\delta^* = 0.1184\delta = 0.1184 \ (0.186) = 0.022 \ ft$$
$$\frac{\delta^*}{\delta} = 0.1184 \ \rightarrow 11.8\%$$

(b) Local surface-resistance coeff.  $C_f$ Use Eq. (9.20) by Clauser

$$\frac{1}{\sqrt{c_f}} = 3.96 \log \operatorname{Re}_{\delta^*} + 3.04 \quad \leftarrow \quad \operatorname{Re}_{\delta^*} = \frac{600 \times 0.022}{3 \times 10^{-4}} = 44,000$$
$$= 3.96 \log (44,000) + 3.04$$

 $\therefore c_f = 2.18 \times 10^{-3} = 0.00218$ 

[Cf]  $c_f = 0.00218$  by Schultz-Grunow Eq.





9.3.2. Power-law formulas: Smooth walls

- Logarithmic equations for velocity profile and shear-stress coeff .
- ~ universal
- ~ applicable over almost entire range of Reynolds numbers
- Power-law equations
- ~ applicable over only limited range of Reynolds numbers
- ~ simpler
- ~ explicit relations for  $\overline{u}/U$  and  $c_f$
- ~ explicit relations for  $\delta$  in terms of *Re* and distance *x*





- (1) Assumptions of power-law formulas
- 1) Except very near the wall, mean velocity is closely proportional to a

root of the distance y from the wall.

$$\overline{u} \propto y^{\frac{1}{n}}$$
 (A

2) Shear stress coeff.  $C_f$  is inversely proportional to a root of  $\text{Re}_{\delta}$ 

$$c_f \propto \frac{1}{\operatorname{Re}^m_{\delta}}$$
,  $\operatorname{Re}_{\delta} = \frac{U}{v}$   
 $c_f = \frac{A}{\left(\frac{U\delta}{v}\right)^m}$ 

(9.29)

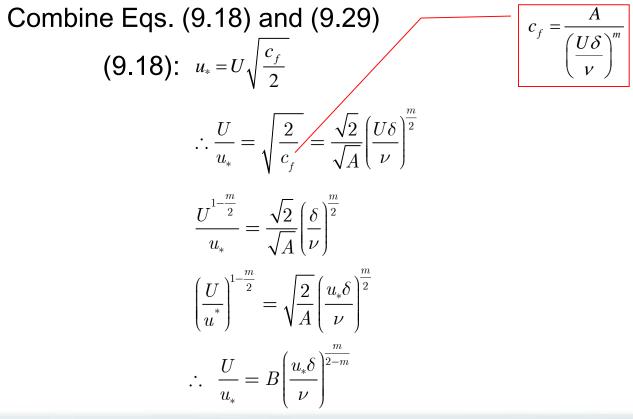
## where A, m = constants





[Cf] Eq. (9.29) is similar to equation for laminar boundary layer,  $c_f = \frac{3.32}{\text{Re}_s}$ 

(2) Derivation of power equation







(9.30)

## 56/76 9.3 Mean-flow characteristics for turbulent boundary layer Substitute Assumption (1) into Eq. (9.30), replace $\delta$ with y $\frac{\overline{u}}{u_*} = B \left( \frac{u_* y}{\nu} \right)^{\frac{m}{2-m}}$ (9.31)Divide (9.31) by (9.30) $\frac{\overline{u}}{U} = \left(\frac{y}{\delta}\right)^{\frac{m}{2-m}}$ (9.32)For 3,000 < $\text{Re}_{\delta}$ < 70,000; $m = \frac{1}{4}$ , A = 0.0466 , B = 8.74 $\frac{\overline{u}}{U} = \left(\frac{y}{\delta}\right)^{\overline{7}}$ (9.33) $\frac{U}{u_*} = 8.74 \left(\frac{u_*\delta}{\nu}\right)^{\frac{1}{7}}$ (9.34) $\frac{\overline{u}}{u_*} = 8.74 \left(\frac{u_* y}{\nu}\right)^{\frac{1}{7}}$ (9.35) $c_f = \frac{0.0466}{\left(Re_{\delta}\right)^{\frac{1}{4}}}$ (9.36)

(3) Relation for  $\delta$ 

Adopt integral-momentum eq. for steady motion with  $\frac{\partial p}{\partial x} = 0$ 

~ integration of Prandtl's 2-D boundary-layer equations

$$\frac{\tau_0}{\rho} = U^2 \frac{\partial \theta}{\partial x} = c_f \frac{U^2}{2} \leftarrow \tau_0 = \frac{\rho}{2} c_f U^2$$
(9.37)

where  $\theta$  = momentum thickness

$$\boldsymbol{\Theta} = \int_{0}^{h} \frac{\overline{u}}{U} \left( 1 - \frac{\overline{u}}{U} \right) dy$$
(A)

Substitute Eq. (9.33) into (A) and integrate





$$\theta = \int_{0}^{h} \left( \frac{y}{\delta} \right)^{\frac{1}{7}} \left\{ 1 - \left( \frac{y}{\delta} \right)^{\frac{1}{7}} \right\} dy$$

$$= \frac{7}{72} \delta \approx 0.1\delta$$
(9.38)

Substitute Eqs. (9.36) and (9.38) into (9.37) and integrate w.r.t. x

$$\delta = \frac{0.318x}{\left(\frac{Ux}{\nu}\right)^{\frac{1}{5}}} = \frac{0.318x}{\left(\operatorname{Re}_{x}\right)^{\frac{1}{5}}} , \operatorname{Re}_{x} < 10^{7}$$

$$c_{f} = \frac{0.059}{\left(\operatorname{Re}_{x}\right)^{\frac{1}{5}}} , \operatorname{Re}_{x} < 10^{7}$$
(9.39)
(9.40)

## Integrate (9.40) over I to get average coefficient

$$c_f = \frac{0.074}{\left(\text{Re}_l\right)^{\frac{1}{5}}}$$
,  $\text{Re}_l < 10^7$  (9.41)





[Re] Derivation of (9.39) and (9.40)

$$U^{2} \frac{\partial \theta}{\partial x} = c_{f} \frac{U^{2}}{2}$$
$$\frac{\partial \theta}{\partial x} = \frac{c_{f}}{2}$$

Substitute (9.38) and (9.36) into (B)

$$\frac{\partial}{\partial x} \left( \frac{7}{72} \delta \right) = \frac{1}{2} \left( 0.0466 / (\operatorname{Re}_{\delta})^{\frac{1}{4}} \right)$$
$$\frac{7}{72} \quad \frac{\partial \delta}{\partial x} = \frac{0.0233}{\left(\operatorname{Re}_{\delta}\right)^{\frac{1}{4}}} = \frac{0.0233}{\left(\frac{U\delta}{\nu}\right)^{\frac{1}{4}}}$$
$$\frac{\partial \delta}{\partial x} = \frac{0.2397}{\left(\frac{U\delta}{\nu}\right)^{\frac{1}{4}}}$$



59/76

(B)

### Integrate once w.r.t. x

$$\delta = \frac{0.2397}{\left(\frac{U\delta}{\nu}\right)^{\frac{1}{4}}}x + C$$

**B.C.:** 
$$\delta \cong 0$$
 at  $\mathbf{x} = \mathbf{0}$ 

$$0 = \frac{0.2397}{\left(\frac{U\delta}{\nu}\right)^{\frac{1}{4}}} 0 + C \quad \longrightarrow \quad C = 0$$

$$\delta = \frac{0.2397}{\left(\frac{U\delta}{\nu}\right)^{\frac{1}{4}}}x$$

$$\delta^{\frac{5}{4}} = \frac{0.2397}{\left(\frac{Ux}{\nu}\right)^{\frac{1}{4}}} x^{\frac{5}{4}}$$





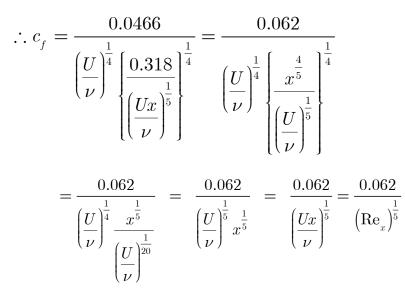
$$\therefore \delta = \left\{ \frac{0.2397}{\left(\frac{Ux}{\nu}\right)^{\frac{1}{4}}} x^{\frac{5}{4}} \right\}^{\frac{4}{5}} = \frac{(0.2397)^{\frac{4}{5}}}{\left(\frac{Ux}{\nu}\right)^{\frac{1}{5}}} x$$
$$\therefore \delta = \frac{0.318}{\left(\frac{Ux}{\nu}\right)^{\frac{1}{5}}} x = \frac{0.318}{\left(\operatorname{Re}_{x}\right)^{\frac{1}{5}}} x \to \operatorname{Eq.}(C)(9.39)$$

(9-36): 
$$c_f = \frac{0.0466}{\left(\operatorname{Re}_{\delta}\right)^{\frac{1}{4}}} = \frac{0.0466}{\left(\frac{U\delta}{\nu}\right)^{\frac{1}{4}}}$$



(C)

Substitute (9.39) into (C)



→ (9.40)

## Integrate (9.40) over *l*

$$\overline{C}_{f} = \frac{1}{l} \int_{0}^{l} \frac{0.062}{\left(\operatorname{Re}_{x}\right)^{\frac{1}{5}}} dx = \frac{0.062}{l} \int_{0}^{l} \frac{1}{\left(\frac{Ux}{\nu}\right)^{\frac{1}{5}}} dx \frac{0.062}{l\left(\frac{U}{\nu}\right)^{\frac{1}{5}}} \int_{0}^{l} \frac{1}{x^{\frac{1}{5}}} dx = \frac{0.076}{\left(\operatorname{Re}_{l}\right)^{\frac{1}{5}}}$$





9.3.3. Laws for rough walls

(1) Effects of roughness

rough walls: velocity distribution and resistance = *f*(Reynolds number, roughness)

smooth walls: velocity distribution and resistance = *f*(Reynolds number)

- For natural roughness, is random, and statistical quantity
- $\rightarrow k = k_s =$  uniform sand grain





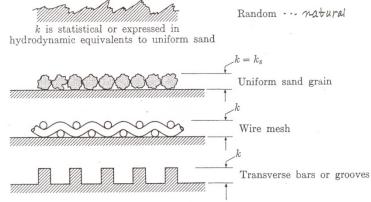
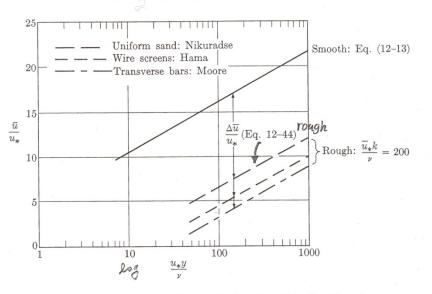


FIG. 12-11. Example of roughness types and definitions of roughness magnitude k.









- Measurement of roughness effects
  - a) experiments with sand grains cemented to smooth surfaces
  - b) evaluate roughness value = height  $k_s$
  - c) compare hydrodynamic behavior with other types and magnitude of roughness
- Effects of roughness
- i)  $\frac{k_s}{\delta'} < 1$
- ~ roughness has negligible effect on the wall shear
- → <u>hydrodynamically smooth</u>

 $\delta' = \frac{4\nu}{u_*}$  = laminar sublayer thickness





ii)  $\frac{k_s}{\delta} > 1$ 

- ~ roughness effects appear
- ~ roughness disrupts the laminar sublayer
- ~ smooth-wall relations for velocity and  $C_f$  no longer hold
- → <u>hydrodynamically rough</u>

iii) 
$$\frac{k_s}{\delta'} > 15 \sim 25$$

~ friction and velocity distribution <u>depend only on roughness</u>rather than Reynolds number

→ <u>fully rough flow</u> condition





Critical roughness, k<sub>crit</sub>

$$\begin{split} k_{crit} &= \delta' \\ &= \frac{4\nu}{u_*} = \frac{4\nu}{U\sqrt{c_f \ / \ 2}} \quad \propto \quad \mathrm{Re}_x \quad \propto \quad x \\ c_f &\propto \ \frac{1}{\mathrm{Re}_x} \end{split}$$

If x increases, then  $C_f$  decreases, and  $\delta'$  increases.

Therefore, for a surface of uniform roughness, it is possible to be <u>hydrodynamically rough upstream</u>, and <u>hydrodynamically smooth</u> <u>downstream</u>.





# 9.3 Mean-flow characteristics for turbulent boundary layer $k < \delta'$ $k > \delta'$ , downstream , upstream k Z **≯**

(2) Rough-wall velocity profiles

Assume roughness height k accounts for magnitude, form, and

distribution of the roughness. Then

$$\frac{\overline{u}}{u_*} = f\left(\frac{y}{k}\right)$$







Make f in Eq. (9.42) be a <u>logarithmic function</u> to overlap the <u>velocity-defect law</u>, Eq. (9.16), which is applicable for both rough and smooth boundaries.

(9.16): 
$$\frac{U - \overline{u}}{u_*} = 5.6 \log\left(\frac{y}{\delta}\right) + 2.5, \ \frac{y}{\delta} < 0.15$$

i) For rough walls, in the wall region

$$\frac{\overline{u}_{rough}}{u_{*}} = -5.6 \log\left(\frac{k}{y}\right) + C_{5}, \quad \frac{u_{*}y}{\nu} > 50 \sim 100, \quad \frac{y}{\delta} < 0.15$$
(9.43)

where  $C_5 = \text{const} = f$  (size, shape, distribution of the roughness)





ii) For smooth walls, in the wall region

(9.13): 
$$\frac{\overline{u}_{smooth}}{u_{*}} = 5.6 \log \left( \frac{u_{*}y}{\nu} \right) + C_{2}, \quad \frac{u_{*}y}{\nu} > 30 \sim 70, \quad \frac{y}{\delta} < 0.15$$

where  $C_2 = 4.9$ 

Subtract Eq. (9.43) from Eq. (9.13)

$$\frac{\Delta \overline{u}}{u_*} = \frac{\overline{u}_{smooth} - \overline{u}_{rough}}{u_*} = 5.6 \log \left(\frac{u_* k}{\nu}\right) + C_6 \tag{9.44}$$

→ Roughness reduces the local mean velocity  $\overline{u}$  in the wall region where  $C_5$  and  $C_6$  → Table 9-4





#### TABLE 12-4

VALUES OF CONSTANTS IN ROUGH-WALL EQUATIONS FOR THE WALL REGION

 $(y/\delta < 0.15; u_*k/\nu > 50 \text{ to } 100)$ 

$egin{array}{c} { m Roughness} \\ { m type} \end{array}$	Source of data	$C_5,$ Eq. (12–43)	C <sub>6</sub> , Eq. (12–44)	C <sub>8</sub> , Eq. (12–46)
Uniform sand grains	Nikuradse [11] (pipes)	8.2	-3.3	7.55
Wire screens	Hama [12] (plates)	. 6.1	-1.2	6.1
Transverse bars	Moore [7] (plates)	4.9	0	5.25
Eq. (1	2–44): $\Delta \overline{u}/u_*$	$= -5.6 \log (k/k) = 5.6 \log (u_*/k) = 3.96 \log (\delta/k)$	$k/\nu)+C_6,$	I

(Constants in this table were evaluated graphically from Fig. 12-12.)





(6) Surface-resistance formulas: rough walls

Combine Eqs. (9.43) and (9.16)

$$\frac{U-\overline{u}}{u_{*}} = -5.6 \log\left(\frac{y}{\delta}\right) + 2.5 , \quad \frac{y}{\delta} < 0.15$$

$$+ \left| \quad \frac{\overline{u}}{u_{*}} = -5.6 \log\left(\frac{k}{y}\right) + C_{5} \right|$$

$$\rightarrow \quad \frac{U}{u_{*}} = -5.6 \log\left(\frac{\delta}{k}\right) + C_{7}$$

$$\frac{U}{u_{*}} = \sqrt{\frac{2}{c_{f}}} = 5.6 \log\left(\frac{\delta}{k}\right) + C_{7} \qquad (9.45)$$

$$\therefore \frac{1}{\sqrt{c_{f}}} = 3.96 \log\left(\frac{\delta}{k}\right) + C_{8} \qquad (9.46)$$



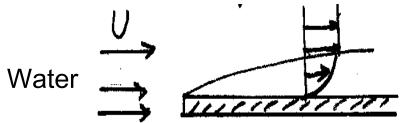


[Ex. 9.3]

Rough wall velocity distribution and local skin friction coefficient

- Comparison of the boundary layers on a smooth plate and a plate

roughened by sand grains



• Given:  $\tau_0 = 0.485 \ lb \ /ft^2$  on both plates

U = 10 ft / sec past the rough plate

 $k_s = 0.001 \text{ ft}$ 

Water temp. = 58 °F on both plates





(a) Velocity reduction  $\Delta u$  due to roughness From Table 1-3:

 $ho = 1.938 \ slug \ / \ ft^3$ ;  $u = 1.25 \times 10^{-5} ft^2 \ / \ {
m sec}$ 

Eq. (9.18)  $\therefore u_* = \sqrt{\frac{\tau_0}{\rho}} = \sqrt{\frac{0.485}{1.938}} = 0.5 ft / sec$  $c_f = 2 \left(\frac{u_*}{U}\right)^2 = 2 \left(\frac{0.5}{10}\right)^2 = 0.005$  $\frac{u_*k_s}{\nu} = \frac{0.5(0.001)}{1.25 \times 10^{-5}} = 40$ Eq. (9.44):  $\frac{\Delta u}{u} = 5.6 \log \left( \frac{u_* k_s}{u} \right) - 3.3$  $\therefore \Delta u = 0.5 \{ 5.6 \log 40 - 3.3 \} = 2.83 ft / sec$ 





(b) Velocity  $\overline{u}$  on each plate at y = 0.007 ft

i) For rough plate

Eq. (9.43): 
$$\frac{\overline{u}}{u_*} = 5.6 \log \frac{y}{k} + 8.2$$
  
 $\therefore \overline{u} = 0.5(5.6 \log \frac{0.007}{0.001} + 8.2) = \underline{6.47 ft / sec}$ 

ii) For smooth plate,

Eq. (9.13): 
$$\frac{\overline{u}}{u_*} = 5.6 \log \frac{u_* y}{\nu} + 4.9$$
  
 $\frac{u_* y}{\nu} = \frac{0.5(0.007)}{1.25 \times 10^{-5}} = 280$   
 $\therefore \overline{u} = 0.5\{5.6 \log(280) + 4.9\} = 9.3 ft / sec$ 

Check  $\Delta \overline{u} = 9.3 - 6.47 = 2.83 \rightarrow$  same result as (a)





(c) Boundary layer thickness  $\delta$  on the rough plate Eq. (9.46):

$$\frac{1}{\sqrt{c_f}} = 3.96 \log \frac{\delta}{k_s} + 7.55$$
$$\therefore \log \frac{\delta}{k_s} = \frac{1}{3.96} \left\{ \frac{1}{\sqrt{0.005}} - 7.55 \right\} = 1.66$$
$$\therefore \frac{\delta}{k_s} = 46 \to \delta = 0.046 ft = 0.52 in = 1.4 cm$$



