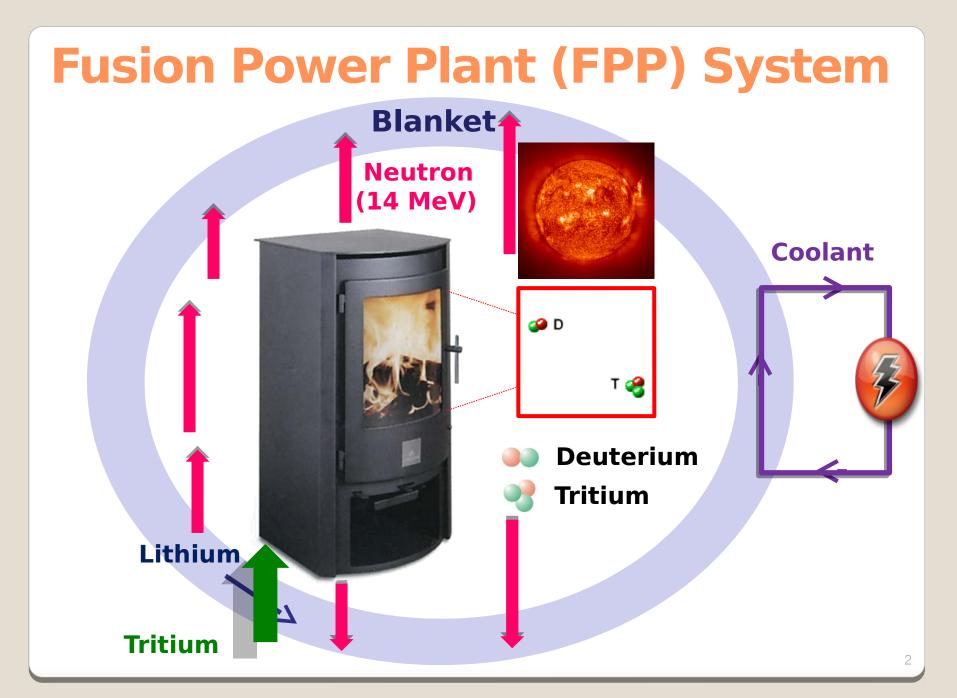
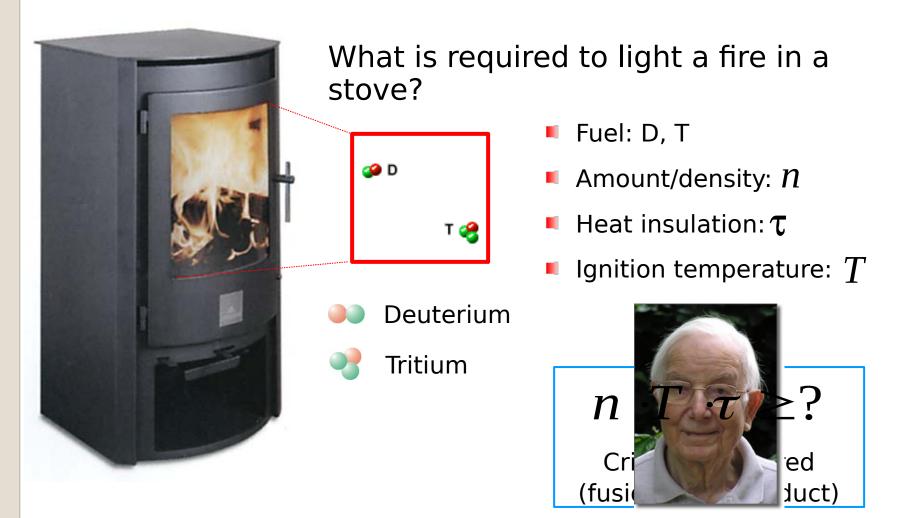
#### Fusion Reactor Technology 2 (459.760, 3 Credits)

# **Prof. Dr. Yong-Su Na** (32-206, Tel. 880-7204)





J. D. Lawson

#### Fundamental requirement of a fusion reactor system

The overall net energy should be larger than the total energy externally supplied to sustain fusion reactions and associated processes subtracted from the total recovered energy

$$E_{net}^* = E_{out}^* - E_{in}^* > 0$$

\*: referring to the entire

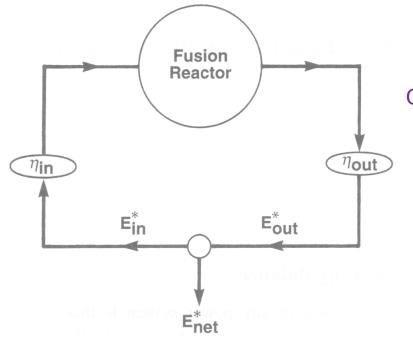
reaction volume

$$\int_{V} E_{out}(r,t) d^{3}r = E_{out}^{*}(t)$$

Considering the time variations of power (Particularly for pulsed systems)

$$\int_{0}^{\tau_{b}} \left(\frac{dE^{*}}{dt}\right)_{net} dt = \int_{0}^{\tau_{b}} \left(\frac{dE^{*}}{dt}\right)_{out} dt - \int_{0}^{\tau_{b}} \left(\frac{dE^{*}}{dt}\right)_{in} dt > 0$$

 $\tau_b$ : burning time



#### • Fusion Plasma Energy Balance

$$\int_{0}^{\tau_{b}} \left(\frac{dE^{*}}{dt}\right)_{net} dt = \int_{0}^{\tau_{b}} \left(\frac{dE^{*}}{dt}\right)_{out} dt - \int_{0}^{\tau_{b}} \left(\frac{dE^{*}}{dt}\right)_{in} dt > 0$$

Thermal energy content in the total plasma volume

$$\int_{0}^{\tau_{b}} \frac{dE_{th}^{*}}{dt} dt = E_{aux}^{*} + E_{fu}^{*} - E_{n}^{*} - E_{rad}^{*} - \int_{0}^{\tau_{b}} \frac{E_{th}^{*}}{\tau_{E^{*}}} dt \qquad E_{out}^{*} = E_{aux}^{*} + E_{alpha}^{*} = E_{fu}^{*} - E_{n}^{*}$$

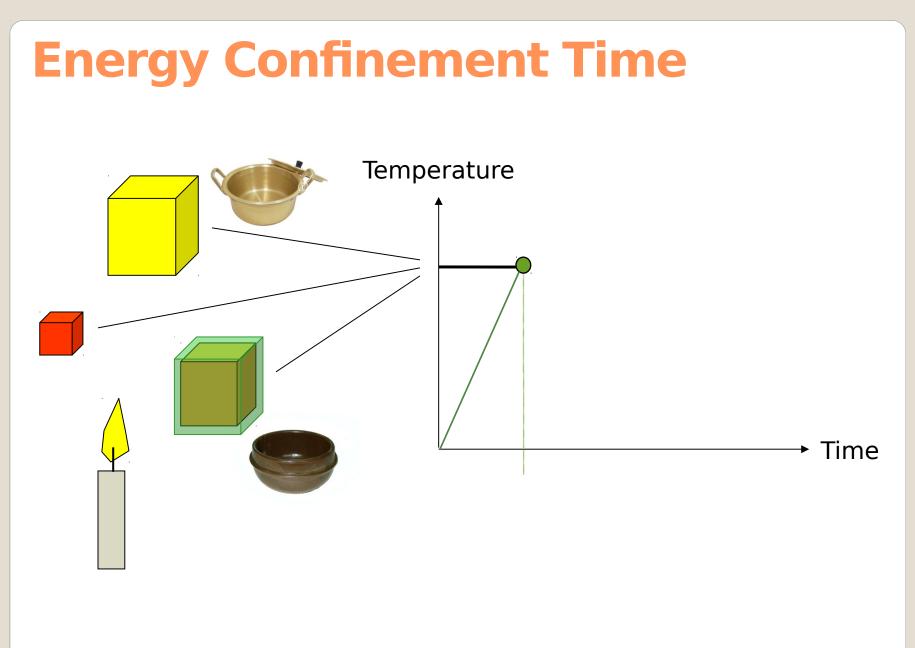
$$E_{out}^{*} = E_{aux}^{*} + E_{alpha}^{*} = E_{fu}^{*} - E_{n}^{*}$$

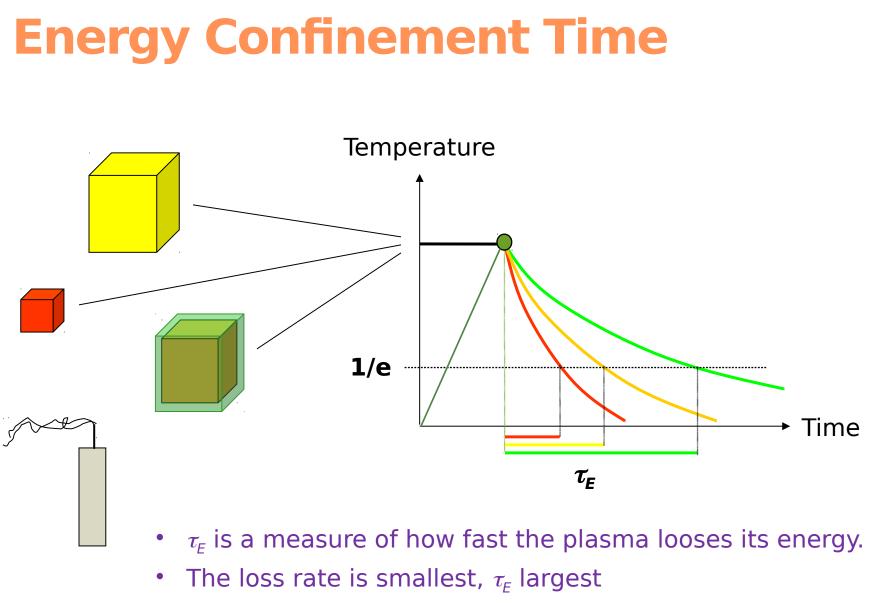
$$E_{in}^{*} = E_{rad}^{*} + \int_{0}^{\tau_{b}} \frac{E_{th}^{*}}{\tau_{E^{*}}} dt$$

$$E_{aux}^{*} = \eta_{in} E_{in}^{*}$$

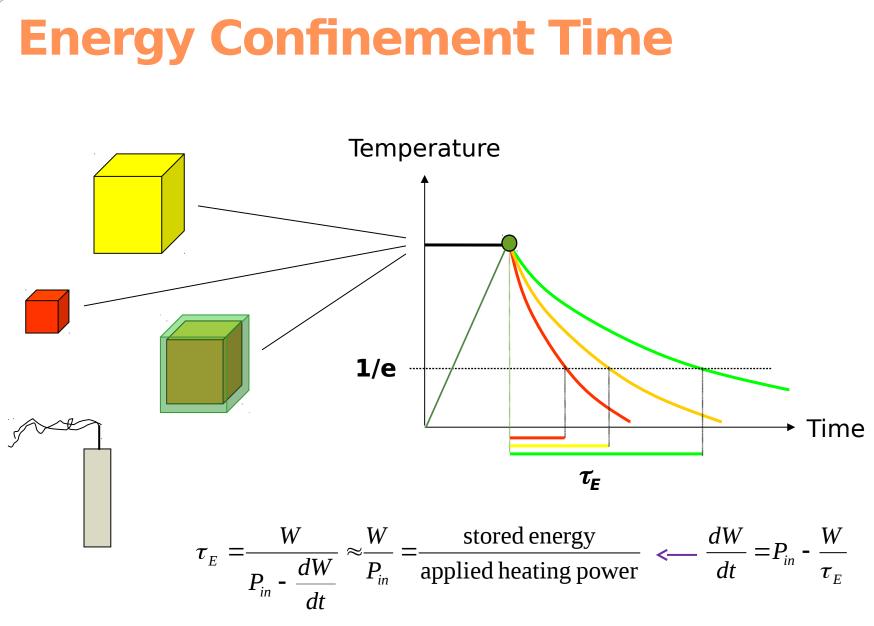
$$\frac{E_{alpha}^{*}}{E_{fu}^{*}} = f_{c} \qquad E_{n}^{*} = (1 - f_{c}) E_{fu}^{*}$$

 $f_c$ : alpha particle fraction of fusion product energy





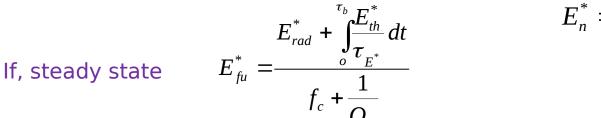
if the fusion plasma is big and well insulated.



$$Q_{p} = \frac{E_{fu}^{*}}{E_{aux}^{*}} = \frac{E_{fu}^{*}}{\eta_{in}E_{in}^{*}}$$

Plasma *Q*-value (fusion multiplication factor): measure for how efficiently an energy input to the plasma is converted into fusion energy

$$\int_{0}^{\tau_{b}} \frac{dE_{th}^{*}}{dt} dt = E_{aux}^{*} + E_{fu}^{*} - E_{n}^{*} - E_{rad}^{*} - \int_{0}^{\tau_{b}} \frac{E_{th}^{*}}{\tau_{E}^{*}} dt \longrightarrow \int_{0}^{\tau_{b}} \frac{dE_{th}^{*}}{dt} dt = \left(\frac{1}{Q_{p}} + f_{c}\right) E_{fu}^{*} - E_{rad}^{*} - \int_{0}^{\tau_{b}} \frac{E_{th}^{*}}{\tau_{E}^{*}} dt$$



 $E_n^* = (1 - f_c) E_{f_u}^*$ 

 if, Q<sub>p</sub> → ∞, the fusion energy delivered to the plasma via the charged reaction products is seen to balance the total energy loss from the plasma.

→ What are requirements of a fusion reactor?

#### • Ignition

Energy viability of the fusion *plasma*:

actual self-sustaining engineering reactor condition with no heating power

$$\frac{E_{fu}^{*}}{\eta_{in}E_{in}^{*}} = Q_{p} \to \infty$$

Considering a D-T plasma with  $Q_p \rightarrow \infty$ ,

$$\int_{0}^{\tau_{b}} \frac{dE_{th}^{*}}{dt} dt = \left(\frac{1}{Q_{p}} + f_{c}\right) E_{fu}^{*} - E_{rad}^{*} - \int_{0}^{\tau_{b}} \frac{E_{th}^{*}}{\tau_{E^{*}}} dt > 0$$

$$f_{c}E_{fu}^{*} > E_{rad}^{*} + \int_{o}^{\tau_{b}} \frac{E_{th}^{*}}{\tau_{E^{*}}} dt$$

$$f_{c}E_{fu}^{*} > E_{rad}^{*} + \int_{o}^{\tau_{b}} \frac{E_{th}^{*}}{\tau_{E^{*}}} dt$$

$$f_{c,dt} \int_{V} d^{3}r \int_{o}^{\tau_{b}} R_{dt}(\vec{r},t) Q_{dt} dt > \int_{V} d^{3}r \left[ \int_{o}^{\tau_{b}} (P_{br} + P_{cyc}^{net}) dt + \int_{o}^{\tau_{b}} \frac{E_{th}(\vec{r},t)}{\tau_{E}(r,t)} dt \right]$$

$$P_{br} = A_{br} n_{i} n_{e} Z^{2} \sqrt{kT_{e}} \qquad A_{br} \approx 1.6 \times 10^{-38} \left[ \frac{m^{3} J}{\sqrt{eVs}} \right]$$
$$P_{cvc}^{net} = A_{cvc} n_{e} B^{2} kT_{e} \psi \qquad A_{cvc} \approx 6.3 \times 10^{-20} \left| JeV^{-1}T^{-2}s^{-1} \right|$$

$$\int_{V} d^{3}r \frac{E_{th}(\underline{r},t)}{\tau_{E}(r,t)} = \frac{E_{th}^{*}(t)}{\tau_{E^{*}}(t)} \quad \text{volume integrated}$$

 In a homogeneous plasma, local D-T fusion ignition condition: Charged particle self-heating power > loss powers (radiation + plasma transport)

$$f_{c,dt} \int_{V} d^{3}r \int_{o}^{\tau_{b}} R_{dt}(\vec{r},t) Q_{dt} dt > \int_{V} d^{3}r \left[ \int_{o}^{\tau_{b}} (P_{br} + P_{cyc}^{net}) dt + \int_{o}^{\tau_{b}} \frac{E_{th}(\vec{r},t)}{\tau_{E}(r,t)} dt \right]$$

$$f_{c,dt}P_{dt}(n_i, T_i) > P_{br}(n_i, n_e, T_e) + P_{cyc}^{net}(n_e, T_e) + \frac{3}{2} \frac{(n_i T_i + n_e T_e)}{\tau_{E^*}} \quad E_{th.j} = \frac{3}{2} n_j T_j , \quad j = i, e$$

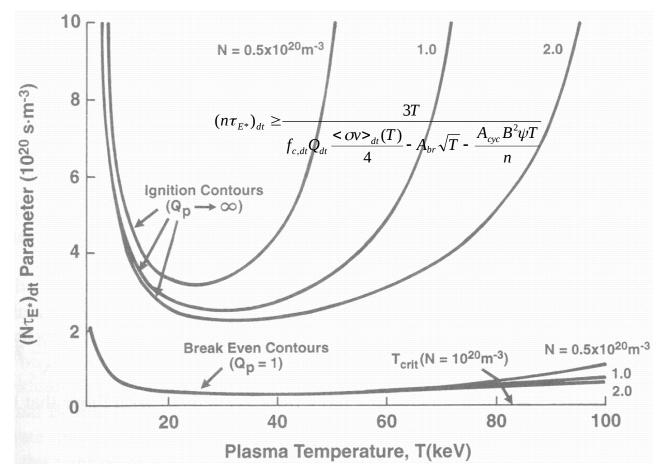
$$f_{c,dt}P_{dt}\tau_{E^*} \ge (P_{br} + P_{cyc}^{net})\tau_{E^*} + 3nT \quad \lt n_i = n_e = n, \ T_i = T_e = T$$

$$(n\tau_{E^*})_{dt} \ge \frac{3T}{\int_{c,dt} Q_{dt} \frac{\langle OV \rangle_{dt}(T)}{4} - A_{br}\sqrt{T} - \frac{A_{cyc}B^2\psi T}{n}}$$

no energy conversion efficiency contained

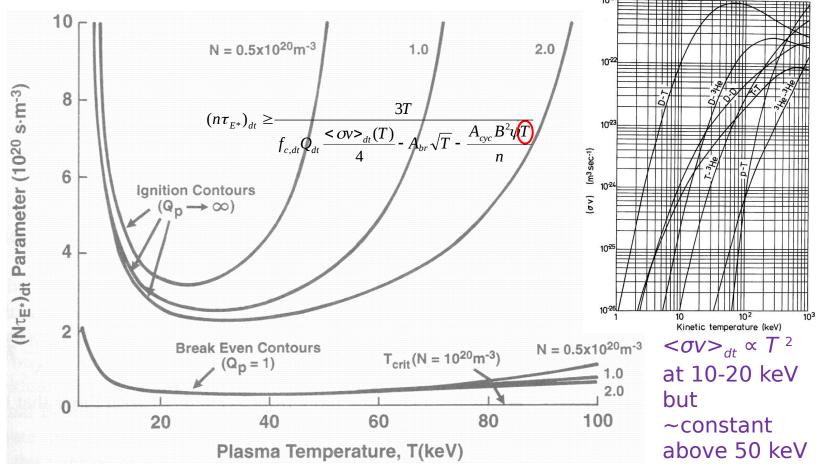
**Plot?** 

- complex interrelation between the plasma density and its temperature as required for ignition



- $n = 10^{20} \text{ m}^{-3}$ :  $T \sim 30 \text{ keV}$ ,  $n\tau_{E^*} \sim 2.7 \times 10^{20} \text{ m}^{-3}$ s,  $\tau_{E^*} \sim 2.7 \text{ s}$
- Ignition contours tend towards infinity as T approaches  $T_{crit}$ .

Why?



- $n = 10^{20} \text{ m}^{-3}$ :  $T \sim 30 \text{ keV}$ ,  $n\tau_{E^*} \sim 2.7 \times 10^{20} \text{ m}^{-3}$ s,  $\tau_{E^*} \sim 2.7 \text{ s}$
- Ignition contours tend towards infinity as T approaches  $T_{crit}$  due to the high Bremmstrahlung and the cyclotron radiation.

#### Break-even (scientific)

The total fusion energy production amounts to a magnitude equal to the

effective plasma energy input.

$$\frac{E_{fu}^*}{\eta_{in}E_{in}^*} = Q_p = 1$$

$$\int_{0}^{\tau_{b}} \frac{dE_{th}^{*}}{dt} dt = \left(\frac{1}{Q_{p}} + f_{c}\right) E_{fu}^{*} - E_{rad}^{*} - \int_{0}^{\tau_{b}} \frac{E_{th}^{*}}{\tau_{E^{*}}} dt > 0$$

 $(n\tau_{E^*})_{dt} \geq ?$  Plot?

#### Lawson criterion from the original paper

- reactor criterion: energy viability of the entire plant

Some Criteria for a Power Producing Thermonuclear Reactor

#### By J. D. LAWSON

Atomic Energy Research Establishment, Harwell, Berks.

Communicated by D. W. Fry; MS. received 2nd November 1956

Abstract. Calculations of the power balance in thermonuclear reactors operating under various idealized conditions are given. Two classes of reactor are considered: first, self-sustaining systems in which the charged reaction products are trapped and, secondly, pulsed systems in which all the reaction products escape so that energy must be supplied continuously during the pulse. It is found that not only must the temperature be sufficiently high, but also the reaction must be sustained long enough for a definite fraction of the fuel to be burnt.

Proceedings of the Physical Society (London), **B70** 6 (1957)

#### Lawson criterion from the original paper

- Deriving some criteria which have to be satisfied in a power producing system by considering power balance for systems in which the reaction products escape, defined as "pulsed systems" by Lawson.
- The gas is heated instantaneously to a temperature T, this temperature is

maintained for a time t, after which the gas is allowed to cool.

- The energy released by the reaction appears as heat generated in the walls of the apparatus (*blanket*), and thus has to be converted to electrical,
- mechanical or chemical energy before it can be fed back into the gas with efficiency  $\eta$ .
- Assumptions:  $P_B = 1.4 \times 10^{-34} n^2 T^{1/2}$  watts cm<sup>-3</sup> considering Bremsstrahlung radiation only (*Spitzer 1956*) (cyclotron radiation neglected) neglecting conduction set  $E_{rad} = e_{rad} e_{th}$ energy used to heat the gas and supply the radiation loss regained as useful heat  $R = E_{fu}^* / E_{aux}^*$ introducing R: ratio of the energy released in the hot gas to the energy

sunnlied

- Lawson criterion from the original paper.

06<sub>10</sub> R

- Condition for a system with net power gain  $E_{out}^* > E_{in}^* = \frac{E_{aux}^*}{\eta_{in}} = \frac{E_{aux}^*}{E_{aux}^*} = \frac{E$  $\eta_{in} \frac{(E_{aux}^* + E_{fu}^*)}{E_{aux}^*} > 1 \qquad \eta(R+1) > 1$  $\eta_{in}$  $R = \frac{E_{fu}^{*}}{E_{aux}^{*}} = \frac{tP_{R}}{tP_{B} + \frac{3}{2}(n_{i}T_{i} + n_{e}T_{i})k} \approx \frac{tP_{R}}{tP_{B} + 3nkT} = \frac{P_{R}/3n^{2}kT}{P_{B}/3n^{2}kT + 1/nt}$ Energy required to heat the gas to a temperature T  $\frac{tP_R + tP_B + 3nkT}{2} > \frac{1}{2}$  $tP_{R} + 3nkT \qquad \eta$  $tP_R + tP_B + 3nkT > \frac{tP_B + 3nkT}{n}$ Total output energy after a pulse ariation of R with T for various values > input energy for heating and of *nt* for T-D reaction compensating loss

#### Lawson criterion from the original paper

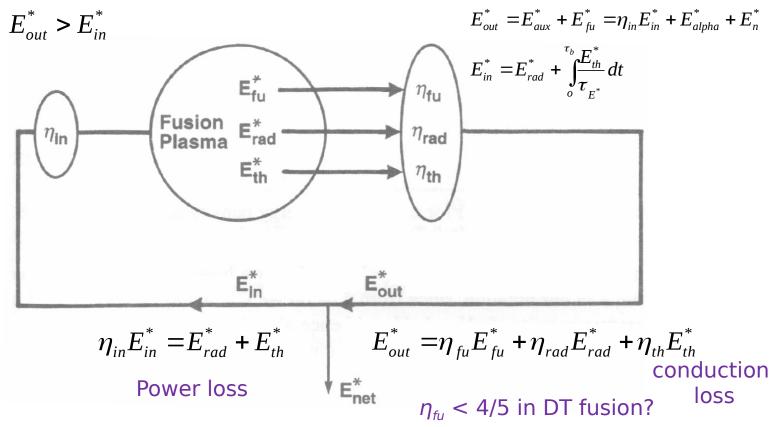
- Conclusion:

For a successful thermonuclear reactor not only has the temperature to be

sufficiently high, but also the reaction has to be sustained for a sufficient time. The reason for this is that the organized energy used to heat the gas is ultimately degraded to the temperature of the walls of the apparatus and, consequently, sufficient thermonuclear energy must be released during each heating cycle to compensate for this degradation.

#### Lawson criterion

- The recoverable energy from a fusion reactor must exceed the energy which is supplied to sustain the fusion reaction.



#### Lawson criterion

- output electric energy (recoverable energy) > required input energy

$$E_{out}^{*} > E_{in}^{*} \qquad E_{out}^{*} = \eta_{fu}E_{fu}^{*} + \eta_{rad}E_{rad}^{*} + \eta_{th}E_{th}^{*} \\ \eta_{in}E_{in}^{*} = E_{rad}^{*} + E_{th}^{*} \\ \eta_{fu}E_{fu}^{*} + \eta_{rad}E_{rad}^{*} + \eta_{th}E_{th}^{*} > \frac{E_{rad}^{*} + E_{th}^{*}}{\eta_{in}} \\ \eta_{in}\eta_{out}(E_{fu}^{*} + E_{rad}^{*} + E_{th}^{*}) > E_{rad}^{*} + E_{th}^{*} \\ \eta_{out} = \frac{\sum_{l} \eta_{l}E_{l}}{\sum_{l} E_{l}}, \ l = fu, rad, th \quad \text{average conversion efficiency} \\ E_{l}^{*} = \tau_{E^{*}}\int_{V}P_{l}(\vec{r})d^{3}r \quad \text{global energy terms} \end{cases}$$

Assuming, Bremsstrahlung only

$$\eta_{in}\eta_{out}\int_{V} d^{3}r(\tau_{E^{*}}P_{fu} + \tau_{E^{*}}P_{br} + 3nT) > \int_{V} d^{3}r(\tau_{E^{*}}P_{br} + 3nT)$$
  
$$\vec{F}_{th}(r) = \frac{3}{2}(n_{i}T_{i} + n_{e}T_{e}) = 3nT$$

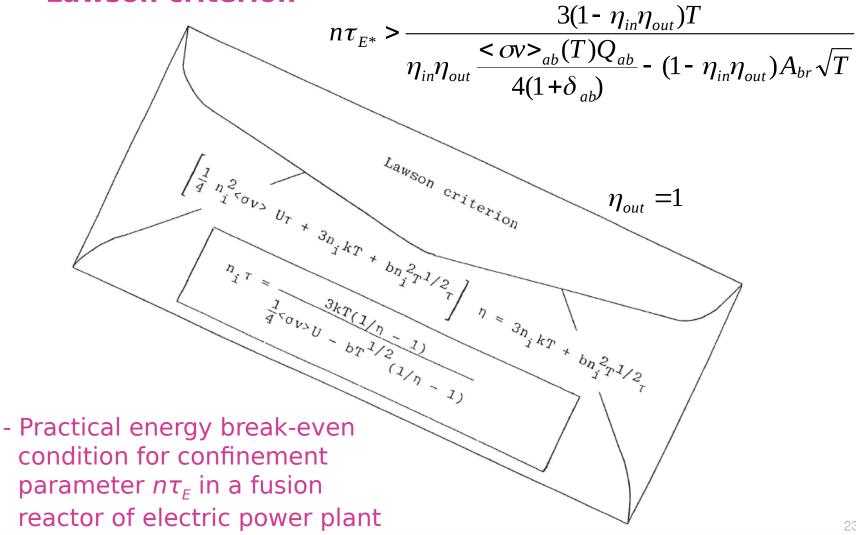
Assuming, homogeneity throughout the plasma volume V

$$\eta_{in}\eta_{out}\left(\frac{n_{a}n_{b}}{1+\delta_{ab}} < \sigma v >_{ab}Q_{ab}\tau_{E^{*}} + A_{br}n^{2}\sqrt{T}\tau_{E^{*}} + 3nT\right) > A_{br}n^{2}\sqrt{T}\tau_{E^{*}} + 3nT$$

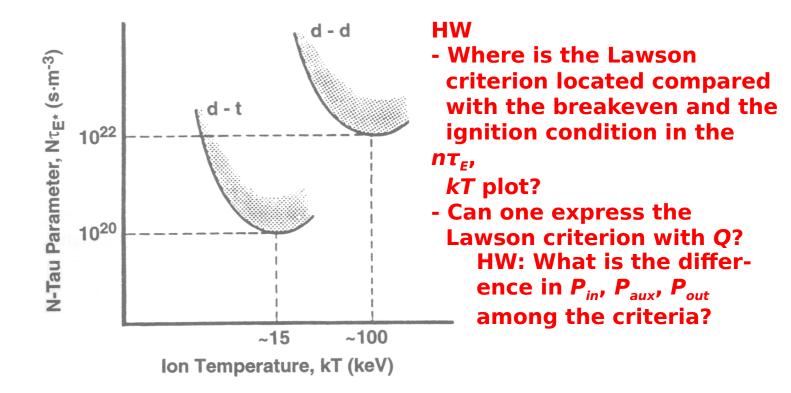
Kronecker- $\delta$  introduced to account for the case of indistinguishable reactants

$$n\tau_{E^*} > \frac{3(1 - \eta_{in}\eta_{out})T}{\eta_{in}\eta_{out} \frac{\langle OV \rangle_{ab}(T)Q_{ab}}{4(1 + \delta_{ab})} - (1 - \eta_{in}\eta_{out})A_{br}\sqrt{T}} \qquad \eta_{in}\eta_{out} \approx 1/3$$

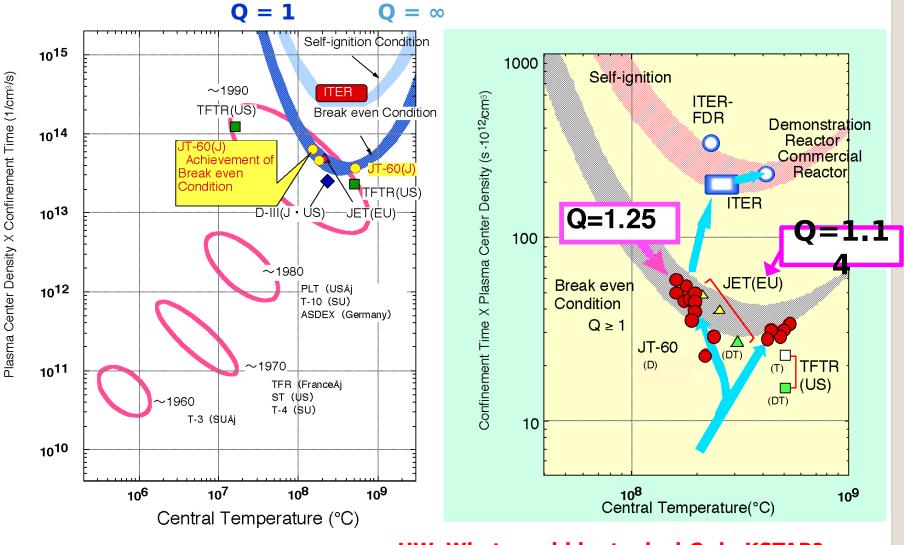
#### **Lawson criterion**



#### Lawson criterion

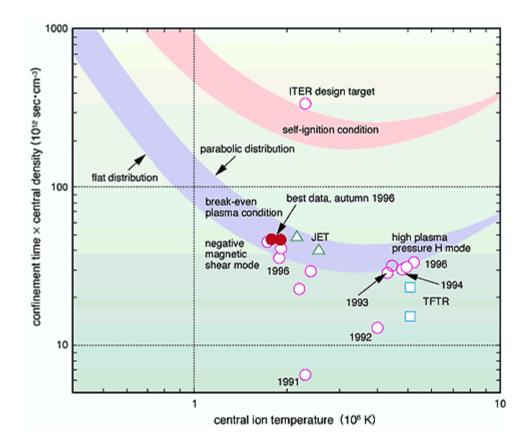


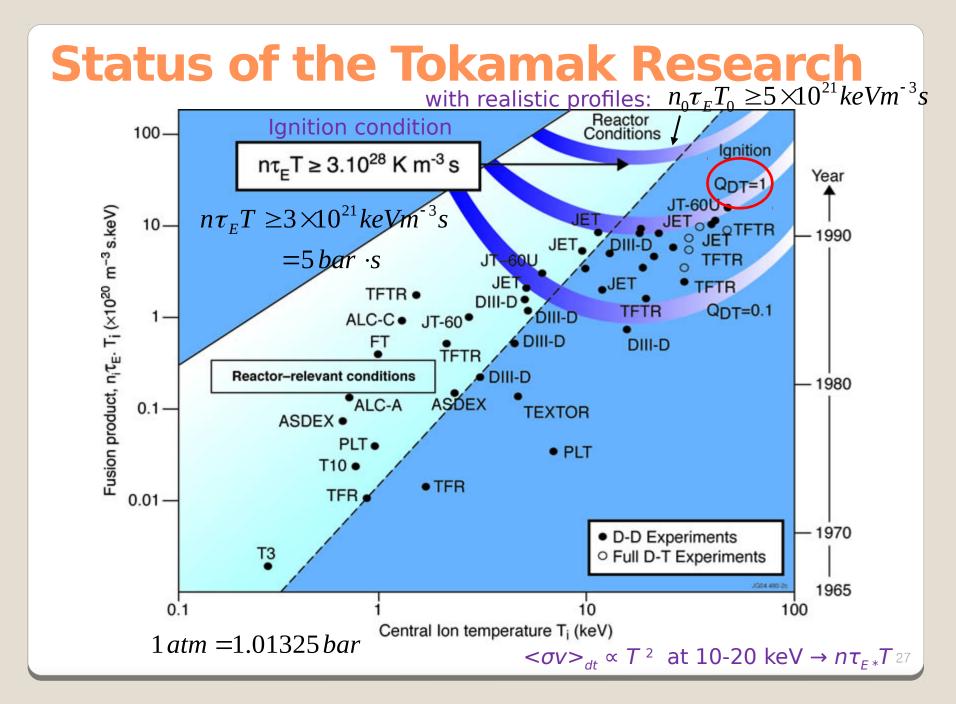
- No particular fusion design was necessary in the derivation of this criterion.
- Although it does not contain all relevant processes such as cyclotron radiation, it is a useful and widely employed criterion.
- For commercial power applications, it would be necessary to exceed the minimum Lawson limit by perhaps a factor of ten or better.

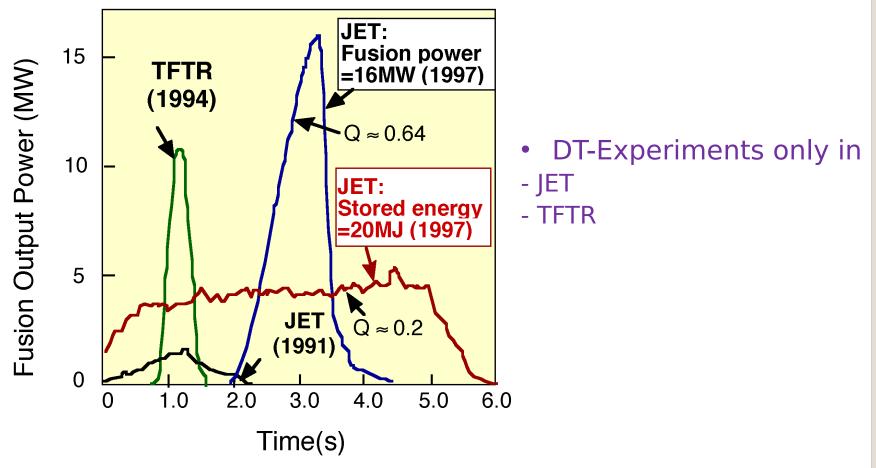


HW. What would be typical  $Q_p$  in KSTAR?

25

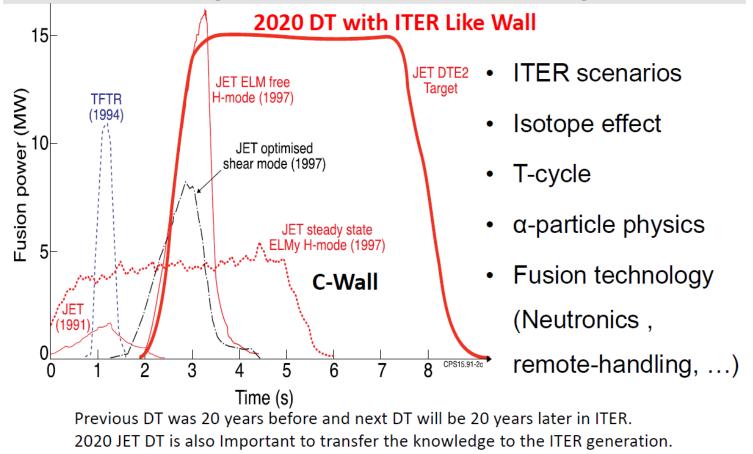






- Present machines produced significant fusion power:
   TFTR (USA) ~10 MW in 1994
  - JET (EU) 16 MW (Q = 0.64) in 1997

Objectives of 2020 JET D-T operation: 15MW fusion power for 5 sec stationary state



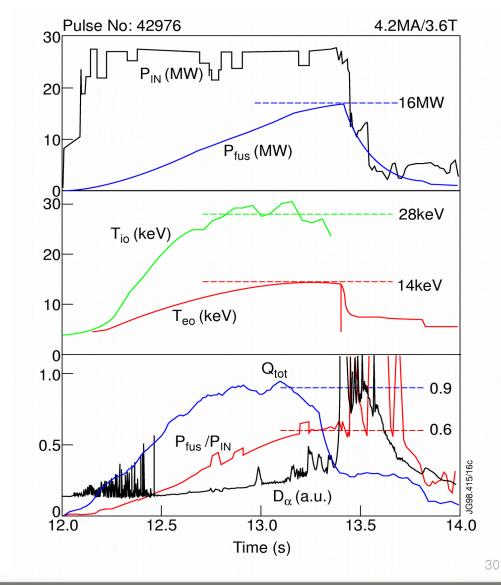


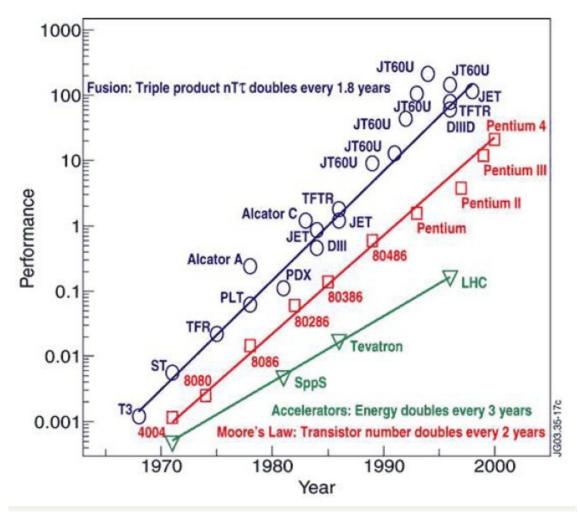
Hyun-Tae Kim | Seminar in Seoul National University | 26 December 2018 | Page 17

- DT-Experiments only inJET
- TFTR
- with world records in JET:

HW.  $Q_{tot}$  ?

- $P_{fusion} = 16 \text{ MW}$
- -Q = 0.65





 Progress in fusion can be compared with the computing power and particle physics accelerator energy.

### Homework

Problems 8.4(submission until the next Thursday)