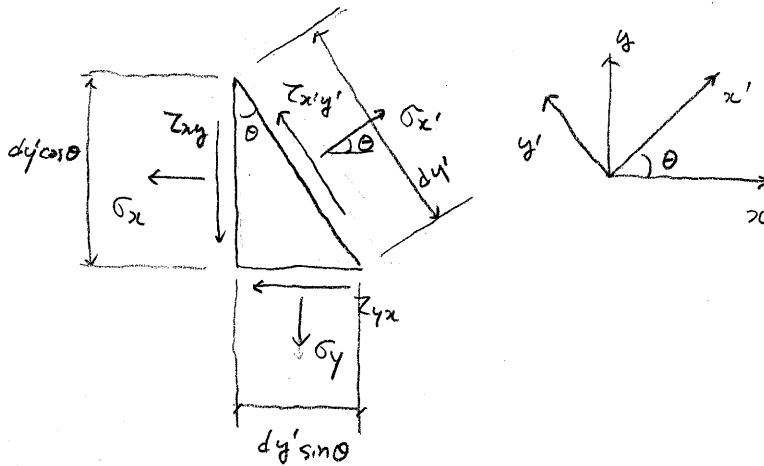
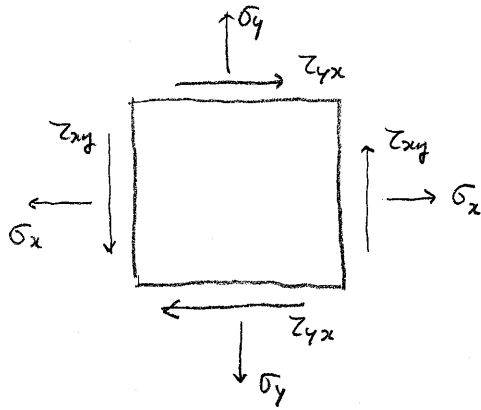


Chapter 2 Plane Stress and Plane Strain

stress - transformation



$$\sum F_{x'} = 0 \Rightarrow \sigma_{x'} dy' - \sigma_x \cos \theta dy' \cos \theta - \sigma_y \sin \theta dy' \sin \theta - \tau_{xy} \sin \theta dy' \cos \theta - \tau_{xy} \cos \theta dy' \sin \theta = 0$$

$$\sigma_{x'} = \sigma_x \cos^2 \theta + \sigma_y \sin^2 \theta + 2 \tau_{xy} \sin \theta \cos \theta$$

$$\sum F_{y'} = 0 \quad \tau_{x'y'} = -(\sigma_x - \sigma_y) \sin \theta \cos \theta + \tau_{xy} (\cos^2 \theta - \sin^2 \theta)$$

$$\sigma_{y'} = \sigma_{x'}(\theta + 90^\circ)$$

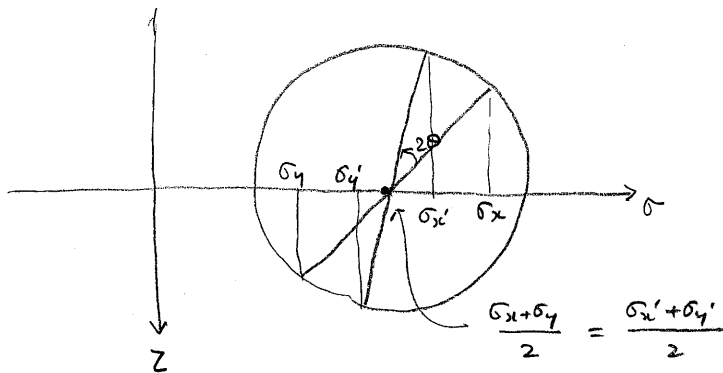
$$= \sigma_x \sin^2 \theta + \sigma_y \cos^2 \theta - 2 \tau_{xy} \sin \theta \cos \theta$$

$$\underline{\sigma}' = \underline{T}_\theta \underline{\sigma}$$

$$\begin{bmatrix} \sigma_x' \\ \sigma_y' \\ \tau_{xy}' \end{bmatrix} = \begin{bmatrix} \cos 2\theta & \sin 2\theta & 2 \sin \theta \cos \theta \\ \sin 2\theta & \cos 2\theta & -2 \sin \theta \cos \theta \\ -\sin \theta \cos \theta & \sin \theta \cos \theta & \cos^2 \theta - \sin^2 \theta \end{bmatrix} \begin{bmatrix} \sigma_x \\ \sigma_y \\ \tau_{xy} \end{bmatrix}$$

$$\underline{T}_\sigma = \underline{T}_\sigma$$

$$\sigma_x' + \sigma_y' = \sigma_x + \sigma_y \Rightarrow \text{invariants}$$



$$\tan 2\theta_p = \frac{2\tau_{xy}}{\sigma_x - \sigma_y} \quad : \text{ principal axes}$$

Strain - displacement relationship in plane

$$\underline{\epsilon} = \begin{bmatrix} \epsilon_x \\ \epsilon_y \\ \gamma_{xy} \end{bmatrix} = \begin{bmatrix} \frac{\partial u}{\partial x} \\ \frac{\partial v}{\partial y} \\ \frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} \end{bmatrix} = \begin{bmatrix} \frac{\partial}{\partial x} & 0 \\ 0 & \frac{\partial}{\partial y} \\ \frac{\partial}{\partial y} & \frac{\partial}{\partial x} \end{bmatrix} \begin{bmatrix} u \\ v \end{bmatrix}$$

$\underline{\epsilon} = \underline{D} \underline{u}$

Axis transformation of strains

Complimentary virtual strain energy densities

$(\delta \underline{\sigma}')^T \underline{\epsilon}' = (\delta \underline{\sigma})^T \underline{\epsilon}$ energy is the same regardless of the definition of axis

$\underline{\sigma}' = \underline{T}_\sigma \underline{\sigma} \Rightarrow \delta \underline{\sigma}' = \underline{T}_\sigma \delta \underline{\sigma}$

$\delta \underline{\sigma}'^T \underline{T}_\sigma^T \underline{\epsilon}' = \delta \underline{\sigma}^T \underline{\epsilon}$

$\underline{T}_\sigma^T \underline{\epsilon}' = \underline{\epsilon} \Rightarrow \underline{\epsilon}' = \underline{T}_\sigma^{-T} \underline{\epsilon}$

$\underline{\epsilon}' = \underline{T}_\epsilon \underline{\epsilon} \quad \underline{T}_\epsilon = \underline{T}_\sigma^{-T}$

$$\begin{bmatrix} \epsilon'_x \\ \epsilon'_y \\ \gamma'_{xy} \end{bmatrix} = \underbrace{\begin{bmatrix} \cos^2\theta & \sin^2\theta & 2\sin\theta\cos\theta \\ \sin^2\theta & \cos^2\theta & -2\sin\theta\cos\theta \\ -2\sin\theta\cos\theta & 2\sin\theta\cos\theta & \cos^2\theta - \sin^2\theta \end{bmatrix}}_{\underline{T}_\epsilon} \begin{bmatrix} \epsilon_x \\ \epsilon_y \\ \gamma_{xy} \end{bmatrix}$$

or

$$\begin{bmatrix} \epsilon'_x \\ \epsilon'_y \\ \epsilon'_{xy} \end{bmatrix} = \underbrace{\begin{bmatrix} \cos^2\theta & \sin^2\theta & 2\sin\theta\cos\theta \\ \sin^2\theta & \cos^2\theta & -2\sin\theta\cos\theta \\ -\sin\theta\cos\theta & \sin\theta\cos\theta & \cos^2\theta - \sin^2\theta \end{bmatrix}}_{\underline{T}'_\epsilon = \underline{T}_\sigma = \underline{T}_\epsilon} \begin{bmatrix} \epsilon_x \\ \epsilon_y \\ \epsilon_{xy} \end{bmatrix}$$

$(\frac{1}{2} \gamma'_{xy})$ $(\frac{1}{2} \gamma_{xy})$

Isotropic material

plane stress (x-y plane)

$$\sigma_z = \tau_{zx} = \tau_{yz} = 0 \quad \gamma_{zx} = \gamma_{yz} = 0, \quad \epsilon_z \neq 0$$

$$\epsilon_x = \frac{1}{E} (\sigma_x - \nu \sigma_y) \quad \epsilon_y = \frac{1}{E} (\sigma_y - \nu \sigma_x) \quad \epsilon_z = \frac{-\nu}{E} (\sigma_x + \sigma_y)$$

$$\gamma_{xy} = \frac{1}{G} \tau_{xy} = \frac{2(1+\nu)}{E} \tau_{xy}$$

$$\underline{\epsilon} = \underline{C} \underline{\sigma} \quad \underline{C} = \frac{1}{E} \begin{bmatrix} 1 & -\nu & 0 \\ -\nu & 1 & 0 \\ 0 & 0 & 2(1+\nu) \end{bmatrix}$$

$$\underline{\sigma} = \underline{E} \underline{\epsilon} \quad \underline{E} = \frac{E}{1-\nu^2} \begin{bmatrix} 1 & \nu & 0 \\ \nu & 1 & 0 \\ 0 & 0 & \lambda \end{bmatrix} \quad \lambda = \frac{1-\nu}{2}$$

plane strain

$$\epsilon_z = \gamma_{zx} = \gamma_{yz} = 0 \quad \tau_{zx} = \tau_{yz} = 0 \quad \sigma_z \neq 0$$

$$\epsilon_x = \frac{1}{E} (\sigma_x - \nu \sigma_y - \nu \sigma_z)$$

$$\epsilon_y = \frac{1}{E} (-\nu \sigma_x + \sigma_y - \nu \sigma_z)$$

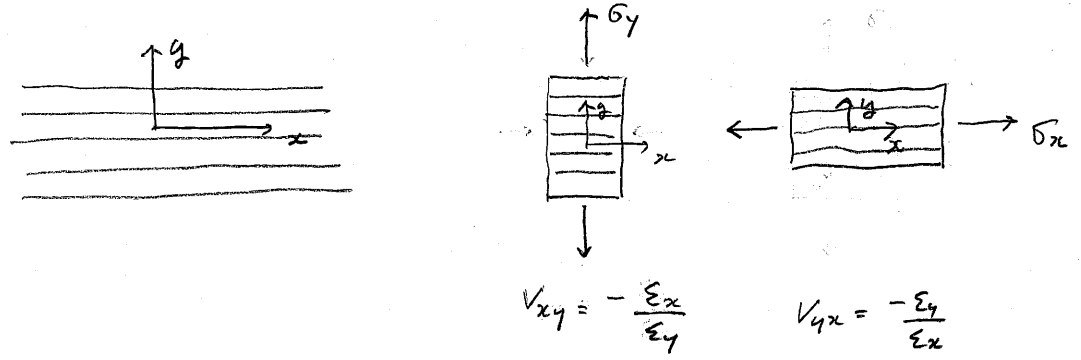
$$\epsilon_z = \frac{1}{E} (-\nu \sigma_x - \nu \sigma_y + \sigma_z) = 0 \Rightarrow \sigma_z = \nu (\sigma_x + \sigma_y)$$

$$\gamma_{xy} = \frac{2(1+\nu)}{E} \tau_{xy}$$

$$\underline{\epsilon} = \underline{C} \underline{\sigma} \quad \underline{C} = \frac{1+\nu}{E} \begin{bmatrix} (1-\nu) & -\nu & 0 \\ -\nu & (1-\nu) & 0 \\ 0 & 0 & 2 \end{bmatrix}$$

$$\underline{\sigma} = \underline{E} \underline{\epsilon} \quad \underline{E} = \frac{E}{(1+\nu)(1-2\nu)} \begin{bmatrix} (1-\nu) & \nu & 0 \\ \nu & (1-\nu) & 0 \\ 0 & 0 & \frac{(1-2\nu)}{2} \end{bmatrix}$$

Orthotropic material



plane stress

$$\epsilon_x = \frac{1}{E_x} \sigma_x - \frac{\nu_{xy}}{E_y} \sigma_y$$

$$\epsilon_y = -\frac{\nu_{yx}}{E_x} \sigma_x + \frac{1}{E_y} \sigma_y$$

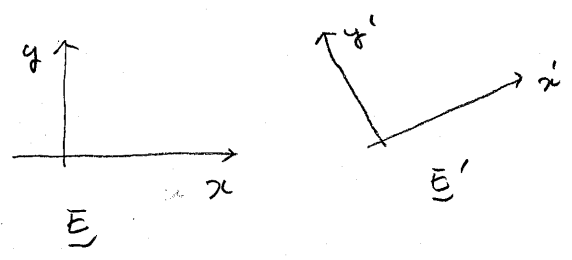
$$\gamma_{xy} = \frac{1}{G_{xy}} \tau_{xy}$$

$$\underline{\epsilon} = \underline{C} \underline{\sigma} \quad \underline{C} = \begin{bmatrix} \frac{1}{E_x} & -\frac{\nu_{xy}}{E_y} & 0 \\ -\frac{\nu_{yx}}{E_x} & \frac{1}{E_y} & 0 \\ 0 & 0 & \frac{1}{G_{xy}} \end{bmatrix}$$

$$\underline{\sigma} = \underline{E} \underline{\epsilon} \quad \underline{E} = \frac{1}{1 - \nu_{xy}\nu_{yx}} \begin{bmatrix} E_x & \nu_{xy} E_x & 0 \\ \nu_{yx} E_y & E_y & 0 \\ 0 & 0 & (-\nu_{xy}\nu_{yx}) G_{xy} \end{bmatrix}$$

Axis - Transformation of Constitutive Matrix

Isotropic material

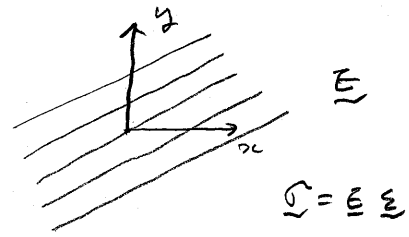
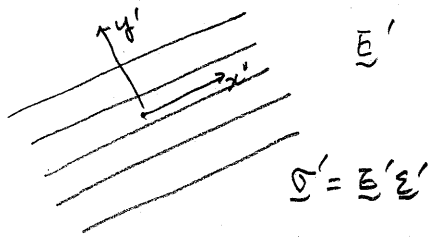


Material properties are the same, regardless of axis - transformation

$\underline{E} = \underline{E}'$ for isotropic material

$\underline{\sigma}' = \underline{E}' \underline{\epsilon}'$

$\underline{\sigma} = \underline{E} \underline{\epsilon}$

Orthotropic material

$$\underline{\sigma}' = \underline{T}_0 \underline{\sigma}$$

$$\underline{\epsilon}' = \underline{T}_0^T \underline{\epsilon}$$

$$\underline{\sigma}' = \underline{E}' \underline{\epsilon}'$$

$$\underline{T}_0 \underline{\sigma} = \underline{E}' \underline{T}_0^T \underline{\epsilon}$$

$$\underline{\sigma} = \underline{T}_0^{-1} \underline{E}' \underline{T}_0^T \underline{\epsilon}$$

$$\underline{T}_0^{-1} = \underline{T}_E$$

$$= \underline{\underbrace{T_E^T \underline{E}' T_E}}_{\underline{E}} \underline{\epsilon}$$

$$\underline{E} = \underline{T_E^T \underline{E}' T_E}$$

for isotropic materials

$$\underline{E} = \underline{T_E^T \underline{E}' T_E} = \underline{E}'$$

2.2 Triangular Elements

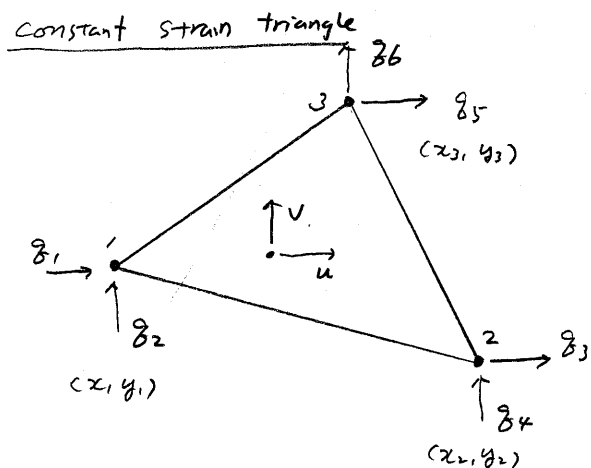
- { Constant strain triangle
- { Linear strain triangle



3-nodes



6-nodes



Generic Displacement $\underline{u} = \begin{bmatrix} u \\ v \end{bmatrix}$

Nodal Displacement $\underline{\delta} = \begin{bmatrix} \delta_1 \\ \delta_2 \\ \delta_3 \\ \delta_4 \\ \delta_5 \\ \delta_6 \end{bmatrix} = \begin{bmatrix} u_1 \\ v_1 \\ u_2 \\ v_2 \\ u_3 \\ v_3 \end{bmatrix}$

Generic - nodal displacement

$$\begin{aligned} u &= c_1 + c_2 x + c_3 y \\ v &= c_4 + c_5 x + c_6 y \end{aligned} \quad \left. \vphantom{\begin{aligned} u \\ v \end{aligned}} \right\} \text{same form}$$

$$u = \underbrace{[1 \quad x \quad y]}_{\underline{\delta}_u} \underbrace{\begin{bmatrix} c_1 \\ c_2 \\ c_3 \end{bmatrix}}_{\underline{c}_u}$$

3. ...

... a number of ...

...

Generally, the assumed displacement function is not the exact displacement that satisfies the force equilibrium.

- 1) However, the principle of virtual displacement results in an acceptable solution, approximately satisfying the force-equilibrium through the energy method.
- 2) Since the displacement function is of a fixed form, the resulting stiffness matrix can be used for any structures with complex conditions as long as a large number of elements are used.
- 3) The fixed form of stiffness automatizes the procedure of FE analysis. Therefore, the FEA procedure can be fitted to analysis using computer.

at node 1 (x_1, y_1) $u = u_1$

$$u_1 = C_1 + C_2 x_1 + C_3 y_1$$

$$u_2 = C_1 + C_2 x_2 + C_3 y_2$$

$$u_3 = C_1 + C_2 x_3 + C_3 y_3$$

$$\underbrace{\begin{bmatrix} u_1 \\ u_2 \\ u_3 \end{bmatrix}}_{\underline{\delta}_u} = \underbrace{\begin{bmatrix} 1 & x_1 & y_1 \\ 1 & x_2 & y_2 \\ 1 & x_3 & y_3 \end{bmatrix}}_{\underline{h}_u} \underbrace{\begin{bmatrix} C_1 \\ C_2 \\ C_3 \end{bmatrix}}_{\underline{c}_u}$$

$$u = \underline{\delta}_u \underline{c}_u = \underline{\delta}_u \underline{h}_u^{-1} \underline{\delta}_u$$

$$= \underline{f}_u \underline{\delta}_u$$

$$\underline{h}_u^{-1} = \frac{1}{\det(\underline{h}_u)} \begin{bmatrix} (x_2 y_3 - x_3 y_2) & (x_3 y_1 - x_1 y_3) & (x_1 y_2 - x_2 y_1) \\ (y_2 - y_3) & (y_3 - y_1) & (y_1 - y_2) \\ (x_3 - x_2) & (x_1 - x_3) & (x_2 - x_1) \end{bmatrix}$$

constants

$$\det(\underline{h}_u) = \begin{vmatrix} 1 & x_1 & y_1 \\ 1 & x_2 & y_2 \\ 1 & x_3 & y_3 \end{vmatrix}$$

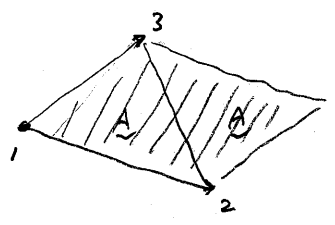
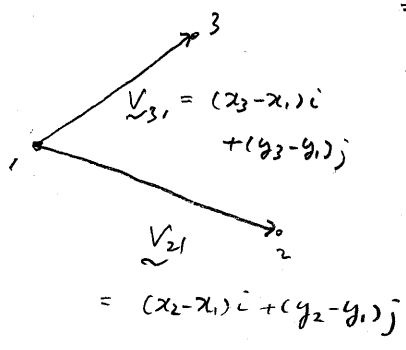
$$= x_2 y_3 - y_2 x_3 - x_1 y_3 + y_1 x_3 + x_1 y_2 - y_1 x_2$$

$$= (x_2 - x_1) y_3 + (x_1 - x_3) y_2 + (x_3 - x_2) y_1$$

$$= (x_2 - x_1) y_3 + (x_1 - x_3) y_2 + (x_3 - x_1) y_1 + (x_1 - x_2) y_1$$

$$= -(x_3 - x_1) (y_2 - y_1) + (x_2 - x_1) (y_3 - y_1)$$

$$= V_{21} \times V_{31} = 2A$$



$$u = f_1 u_1 + f_2 u_2 + f_3 u_3$$

$$\underline{f}_u = \underline{g}_u \underline{h}_u^T = [f_1 \ f_2 \ f_3]$$

$$f_1 = \frac{1}{2A} [(x_2 y_3 - x_3 y_2) + (y_2 - y_3)x + (x_3 - x_2)y]$$

$$f_2 = \frac{1}{2A} [(x_3 y_1 - x_1 y_3) + (y_3 - y_1)x + (x_1 - x_3)y]$$

$$f_3 = \frac{1}{2A} [(x_1 y_2 - x_2 y_1) + (y_1 - y_2)x + (x_2 - x_1)y]$$

$$u = f_1 u_1 + f_2 u_2 + f_3 u_3 = \underline{f}_u \underline{g}_u$$

$$v = f_1 v_1 + f_2 v_2 + f_3 v_3 = \underline{f}_v \underline{g}_v \quad \underline{f}_v = \underline{f}_u$$

$$\underline{\begin{bmatrix} u \\ v \end{bmatrix}} = \underline{\begin{bmatrix} f_1 & 0 & f_2 & 0 & f_3 & 0 \\ 0 & f_1 & 0 & f_2 & 0 & f_3 \end{bmatrix}} \underline{\begin{bmatrix} u_1 \\ v_1 \\ u_2 \\ v_2 \\ u_3 \\ v_3 \end{bmatrix}}$$

\underline{u} \underline{f} \underline{g}

Strain - General displacement

$$\underline{\Sigma} = \underline{\begin{bmatrix} \epsilon_x \\ \epsilon_y \\ \gamma_{xy} \end{bmatrix}} = \underline{\begin{bmatrix} \frac{\partial}{\partial x} & 0 \\ 0 & \frac{\partial}{\partial y} \\ \frac{\partial}{\partial y} & \frac{\partial}{\partial x} \end{bmatrix}} \underline{\begin{bmatrix} u \\ v \end{bmatrix}}$$

\underline{d} \underline{u}

Strain - nodal displacement

$$\underline{\Sigma} = \underline{d} \underline{u} = \underline{d} \underline{f} \underline{g} = \underline{B} \underline{g} \quad \underline{B} = \underline{d} \underline{f}$$

$$\underline{B} = \underline{\begin{bmatrix} \frac{\partial}{\partial x} & 0 \\ 0 & \frac{\partial}{\partial y} \\ \frac{\partial}{\partial y} & \frac{\partial}{\partial x} \end{bmatrix}} \underline{\begin{bmatrix} f_1 & 0 & f_2 & 0 & f_3 & 0 \\ 0 & f_1 & 0 & f_2 & 0 & f_3 \end{bmatrix}}$$

$$\underline{B} = \begin{bmatrix} \frac{\partial f_1}{\partial x} & 0 & \frac{\partial f_2}{\partial x} & 0 & \frac{\partial f_3}{\partial x} & 0 \\ 0 & \frac{\partial f_1}{\partial y} & 0 & \frac{\partial f_2}{\partial y} & 0 & \frac{\partial f_3}{\partial y} \\ \frac{\partial f_1}{\partial y} & \frac{\partial f_1}{\partial x} & \frac{\partial f_2}{\partial y} & \frac{\partial f_2}{\partial x} & \frac{\partial f_3}{\partial y} & \frac{\partial f_3}{\partial x} \end{bmatrix}$$

$$= \frac{1}{2A} \begin{bmatrix} (y_2 - y_3) & 0 & (y_3 - y_1) & 0 & (y_1 - y_2) & 0 \\ 0 & (x_3 - x_2) & 0 & (x_1 - x_3) & 0 & (x_2 - x_1) \\ (x_3 - x_2) & (y_2 - y_3) & (x_1 - x_3) & (y_3 - y_1) & (x_2 - x_1) & (y_1 - y_2) \end{bmatrix}$$

constant

$\underline{\Sigma} = \underline{B}^T \underline{\epsilon} \Rightarrow$ constant strain \Rightarrow constant strain triangle

$$K = \int \underline{B}^T \underline{E} \underline{B} dV$$

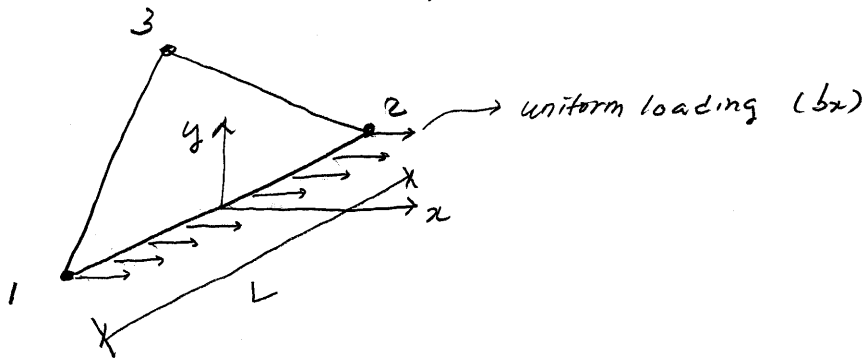
$$= \underline{B}^T \underline{E} \underline{B} \int dV$$

$$= \underline{B}^T \underline{E} \underline{B} A t \Rightarrow \text{constant}$$

\underline{E} $\left\{ \begin{array}{l} \text{plane stress} \\ \text{plane strain} \end{array} \right.$

Equivalent nodal force

uniformly distributed loading



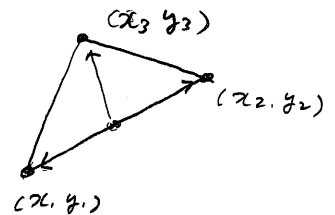
$$\underline{P}_{bx} = \int \underline{f}_y^T b_x dL$$

$$= \int \underline{h}_u^{-T} \underline{g}_u^T b_x dL$$

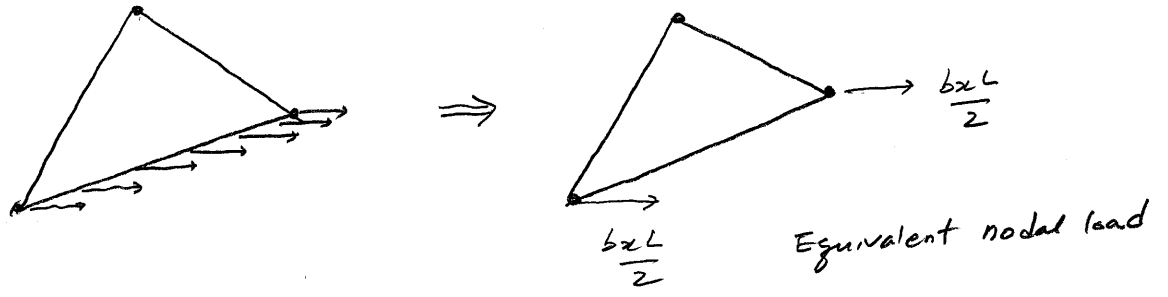
$$= \underline{h}_u^{-T} \int \underline{g}_u^T b_x dL \quad \underline{g}_u = [1 \quad x \quad y]$$

$$= \underline{h}_u^{-T} \begin{bmatrix} \int b_x dL \\ \int x b_x dL \\ \int y b_x dL \end{bmatrix} = \underline{h}_u^{-T} b_x \begin{bmatrix} L \\ 0 \\ 0 \end{bmatrix}$$

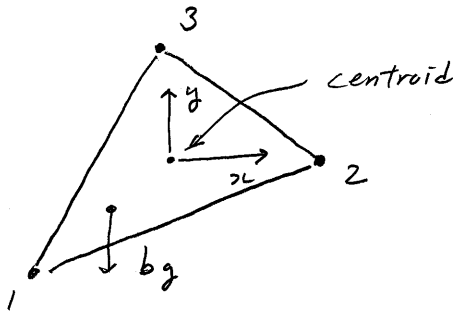
$$= \frac{b_x L}{2A} \begin{bmatrix} (x_2 y_3 - x_3 y_2) \\ (x_3 y_1 - y_3 x_1) \\ (x_1 y_2 - x_2 y_1) \end{bmatrix} \quad \left\{ \frac{A}{2} \times 2 \right.$$



$$= \frac{b_x L}{2A} \begin{bmatrix} A \\ A \\ 0 \end{bmatrix} = \frac{b_x L}{2} \begin{bmatrix} 1 \\ 1 \\ 0 \end{bmatrix}$$



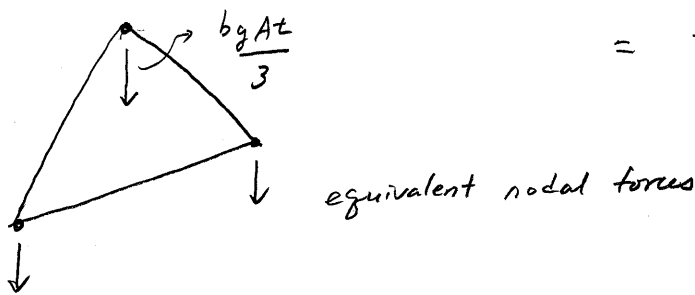
Gravity loading



$$\underline{P}_{bg} = \int f_v^T (-bg) dV = \int \underline{h}_v^{-T} g_v^T (-bg) dV$$

$$= \underline{h}_v^{-T} \begin{bmatrix} \int -bg \, t \, dA \\ \int -x \, bg \, t \, dA \\ \int -y \, bg \, t \, dA \end{bmatrix} = -\underline{h}_v^{-T} bg \, t \begin{bmatrix} A \\ 0 \\ 0 \end{bmatrix}$$

$$= \frac{-bg \, t}{2} \begin{bmatrix} (x_2 y_3 - x_3 y_2) \\ (x_3 y_1 - x_1 y_3) \\ (x_1 y_2 - x_2 y_1) \end{bmatrix} = \frac{-bg \, t}{2} \begin{bmatrix} 2A/3 \\ 2A/3 \\ 2A/3 \end{bmatrix} = \frac{-bg \, t}{2} \begin{bmatrix} 1 \\ 1 \\ 1 \end{bmatrix}$$



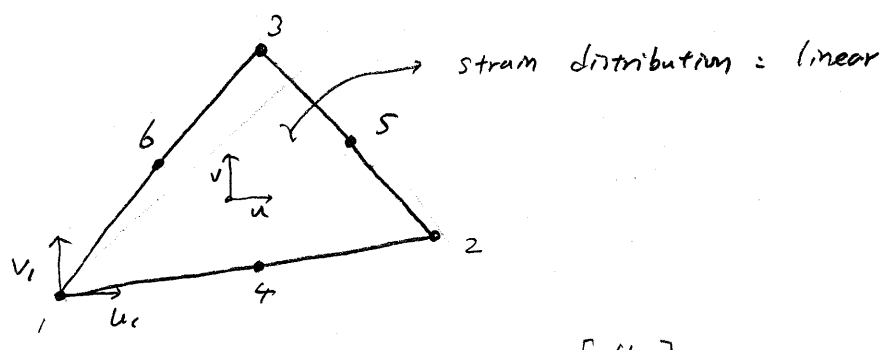
$$= \frac{-bg \, t}{3} \begin{bmatrix} 1 \\ 1 \\ 1 \end{bmatrix}$$

Uniform Temperature change ΔT

$$\underline{\epsilon}_0 = \alpha \Delta T \begin{bmatrix} 1 \\ 1 \\ 0 \end{bmatrix}$$

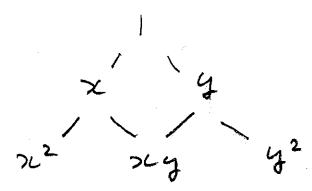
$$\underline{P}_T = \underline{P}_0 = \int_v \underline{B}^T \underline{E} \underline{\epsilon}_0 \, dV = \underline{B}^T \underline{E} \underline{\epsilon}_0 \, A \, t$$

Linear Strain Triangle



nodal displacement $\underline{d} = \begin{bmatrix} u_1 \\ v_1 \\ \vdots \\ u_6 \\ v_6 \end{bmatrix}$

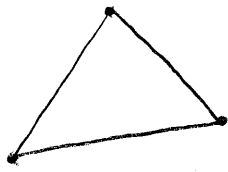
$$\left. \begin{cases} u = C_1 + C_2x + C_3y + C_4x^2 + C_5xy + C_6y^2 \\ v = C_7 + C_8x + C_9y + C_{10}x^2 + C_{11}xy + C_{12}y^2 \end{cases} \right\} 12 \text{ constants}$$



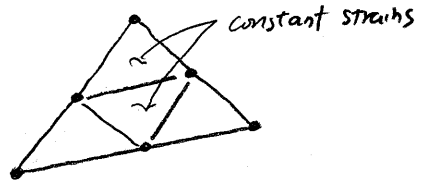
pascal triangle

It is convenient to use the area coordinate system for triangular elements

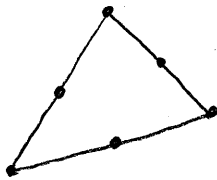
Triangular Element



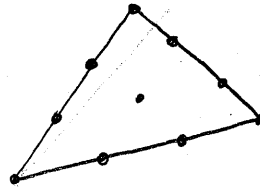
constant strain triangle
(CST) 3 nodes



CST in CST



Linear strain triangle
(LST) 6 nodes



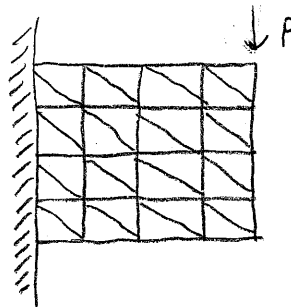
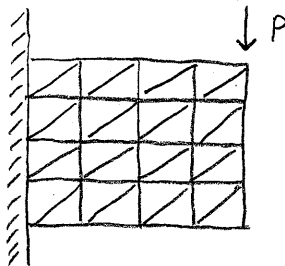
Quadratic strain triangle
10 nodes

Advantages of triangular element

- 1) Simplicity in element stiffness
- 2) usefulness in the description of the geometry of complex structures

Disadvantages

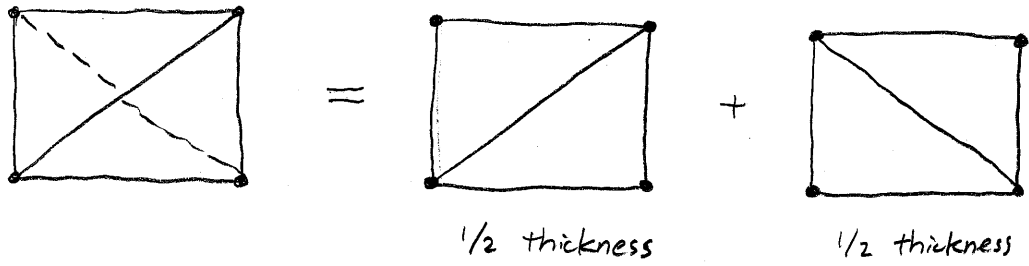
- 1) Lack of geometric isotropy



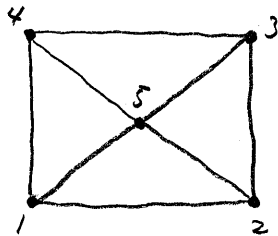
same number of elements
but different solutions

Modification of modeling

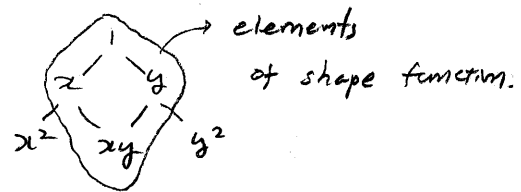
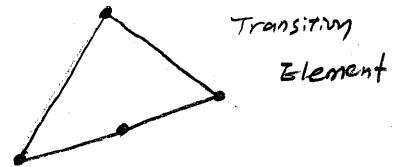
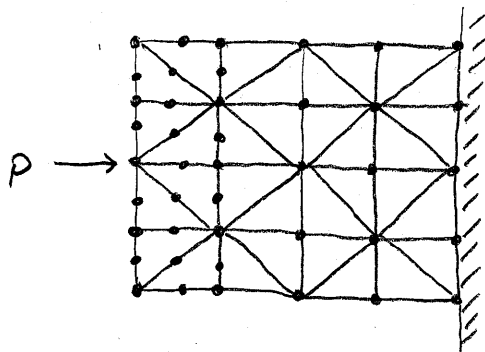
1)



2)

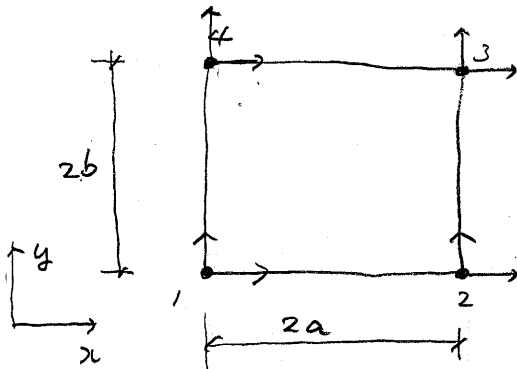


Transition Element

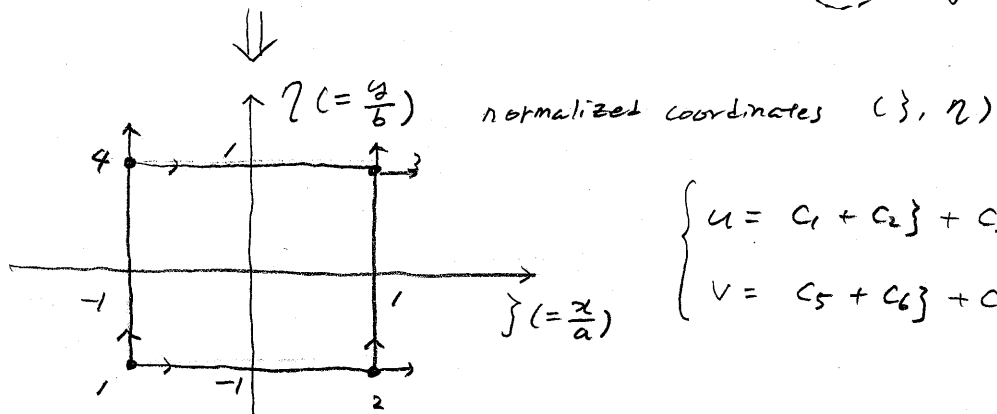
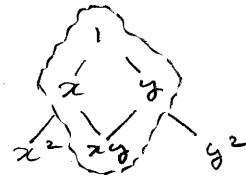


Rectangular Elements

Bilinear Displacement Rectangle by Melosh



$$\begin{cases} u = c_1 + c_2 x + c_3 y + c_4 xy \\ v = c_5 + c_6 x + c_7 y + c_8 xy \end{cases}$$



$$\begin{cases} u = c_1 + c_2 \xi + c_3 \eta + c_4 \xi \eta \\ v = c_5 + c_6 \xi + c_7 \eta + c_8 \xi \eta \end{cases}$$

$$u = c_1 + c_2 \xi + c_3 \eta + c_4 \xi \eta$$

$$= \underbrace{[1 \quad \xi \quad \eta \quad \xi \eta]}_{\underline{\delta}_u} \underbrace{\begin{bmatrix} c_1 \\ c_2 \\ c_3 \\ c_4 \end{bmatrix}}_{\underline{c}_u}$$

using boundary condition at the nodes

$$\underbrace{\begin{bmatrix} u_1 \\ u_2 \\ u_3 \\ u_4 \end{bmatrix}}_{\underline{\delta}_u} = \underbrace{\begin{bmatrix} 1 & -1 & -1 & 1 \\ 1 & 1 & -1 & -1 \\ 1 & 1 & 1 & 1 \\ 1 & -1 & 1 & -1 \end{bmatrix}}_{\underline{b}_u} \underbrace{\begin{bmatrix} c_1 \\ c_2 \\ c_3 \\ c_4 \end{bmatrix}}_{\underline{c}_u}$$

$$h_2^{-1} = \frac{1}{4} \begin{bmatrix} 1 & 1 & 1 & 1 \\ -1 & 1 & 1 & -1 \\ -1 & -1 & 1 & 1 \\ 1 & -1 & 1 & -1 \end{bmatrix}$$

$$\underline{f}_u = \underline{g}_u h_u^{-1}$$

$$\underline{f}_u = [f_1 \quad f_2 \quad f_3 \quad f_4]$$

$$f_1 = \frac{1}{4} (1-\xi)(1-\eta)$$

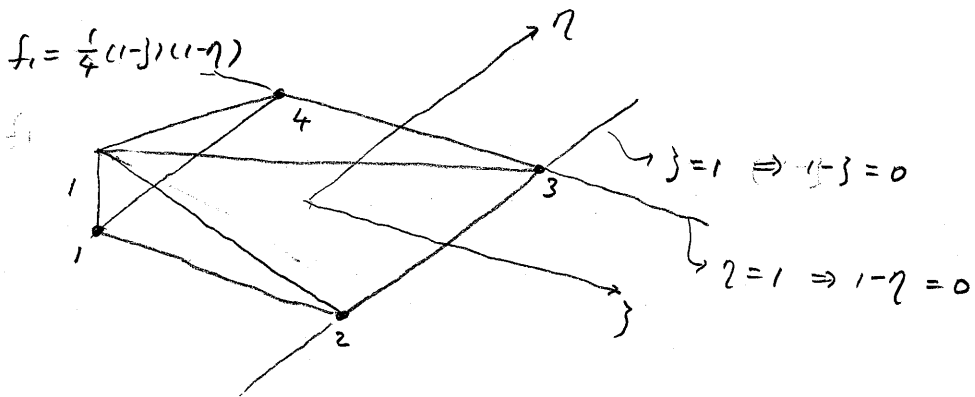
$$f_2 = \frac{1}{4} (1+\xi)(1-\eta)$$

$$f_3 = \frac{1}{4} (1+\xi)(1+\eta)$$

$$f_4 = \frac{1}{4} (1-\xi)(1+\eta)$$

$$f_i = \frac{1}{4} (1+\xi_i)(1+\eta_i)$$

$(\xi_i, \eta_i) \Rightarrow$ coordinate of node i



$$\begin{cases} f_1 = 1 \text{ at node 1 } (\xi_i, \eta_i) = (-1, 1) & a = \frac{1}{4} \\ f_i = 0 \text{ at other nodes} & \Rightarrow f_i = a (1-\xi)(1-\eta) \end{cases}$$

$$u = f_u \underline{z}_u$$

$$v = f_v \underline{z}_v \quad f_y = f_u$$

$$\underline{u} = \begin{bmatrix} u \\ v \end{bmatrix} = \underline{f} \underline{q}$$

$$= \begin{bmatrix} f_1 & 0 & f_2 & 0 & f_3 & 0 & f_4 & 0 \\ 0 & f_1 & 0 & f_2 & 0 & f_3 & 0 & f_4 \end{bmatrix}$$

\underline{f}

$$\begin{bmatrix} u_1 \\ v_1 \\ u_2 \\ v_2 \\ u_3 \\ v_3 \\ u_4 \\ v_4 \end{bmatrix} \underline{q}$$

Strain - Generic Displacement

$$\epsilon_x = \frac{\partial u}{\partial x} = \frac{\partial u}{\partial \xi} \frac{\partial \xi}{\partial x} = \frac{1}{a} \frac{\partial u}{\partial \xi} \quad \xi = \frac{x}{a}$$

$$\epsilon_y = \frac{\partial v}{\partial y} = \frac{\partial v}{\partial \eta} \frac{\partial \eta}{\partial y} = \frac{1}{b} \frac{\partial v}{\partial \eta} \quad \eta = \frac{y}{b}$$

$$\begin{aligned} \gamma_{xy} &= \frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} = \frac{\partial u}{\partial \eta} \frac{\partial \eta}{\partial y} + \frac{\partial v}{\partial \xi} \frac{\partial \xi}{\partial x} \\ &= \frac{1}{b} \frac{\partial u}{\partial \eta} + \frac{1}{a} \frac{\partial v}{\partial \xi} \end{aligned}$$

Strain - matrix displacement

$$\underline{\epsilon} = \begin{bmatrix} \epsilon_x \\ \epsilon_y \\ \gamma_{xy} \end{bmatrix} = \begin{bmatrix} \frac{1}{a} \frac{\partial}{\partial \xi} & 0 \\ 0 & \frac{1}{b} \frac{\partial}{\partial \eta} \\ \frac{1}{b} \frac{\partial}{\partial \eta} & \frac{1}{a} \frac{\partial}{\partial \xi} \end{bmatrix} \begin{bmatrix} u \\ v \end{bmatrix} \underline{u}$$

\underline{d}

Strain - nodal displacement

$$\underline{\underline{\epsilon}} = \underline{\underline{d}} \underline{\underline{u}} = \underline{\underline{d}} \underline{\underline{f}} \underline{\underline{q}} \quad \underline{\underline{B}} = \underline{\underline{d}} \underline{\underline{f}} \quad f_{1,\xi} = \frac{\partial f}{\partial \xi}, \quad f_{1,\eta} = \frac{\partial f}{\partial \eta}$$

3x2 2x8 3x8

$$\underline{\underline{B}} = \begin{bmatrix} \frac{1}{a} f_{1,\xi} & 0 & \frac{1}{a} f_{2,\xi} & 0 & \frac{1}{a} f_{3,\xi} & 0 & \frac{1}{a} f_{4,\xi} & 0 \\ 0 & \frac{1}{b} f_{1,\eta} & 0 & \frac{1}{b} f_{2,\eta} & 0 & \frac{1}{b} f_{3,\eta} & 0 & \frac{1}{b} f_{4,\eta} \\ \frac{1}{b} f_{1,\eta} & \frac{1}{a} f_{1,\xi} & \frac{1}{b} f_{2,\eta} & \frac{1}{a} f_{2,\xi} & \frac{1}{b} f_{3,\eta} & \frac{1}{a} f_{3,\xi} & \frac{1}{b} f_{4,\eta} & \frac{1}{a} f_{4,\xi} \end{bmatrix}$$

$$= \begin{bmatrix} -\frac{1}{a}(1-\eta) & 0 & \frac{1}{a}(1-\eta) & 0 & \frac{1}{a}(1+\eta) & 0 & -\frac{1}{a}(1+\eta) & 0 \\ 0 & -\frac{1}{b}(1-\xi) & 0 & -\frac{1}{b}(1+\xi) & 0 & \frac{1}{b}(1+\xi) & 0 & \frac{1}{b}(1-\xi) \\ -\frac{1}{b}(1-\xi) & -\frac{1}{a}(1-\eta) & -\frac{1}{b}(1+\xi) & \frac{1}{a}(1-\eta) & \frac{1}{b}(1+\xi) & \frac{1}{a}(1+\eta) & \frac{1}{b}(1-\xi) & -\frac{1}{a}(1+\eta) \end{bmatrix}$$

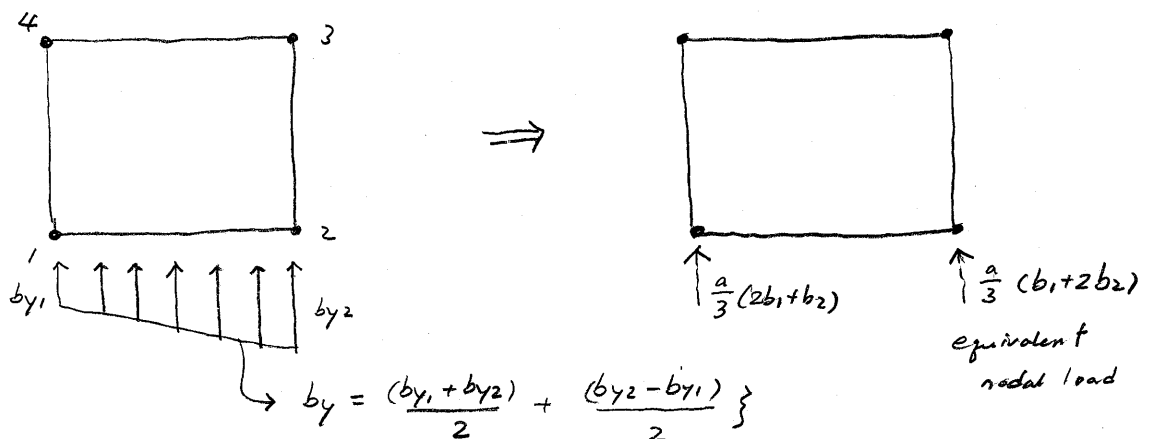
$$\underline{\underline{K}} = \int \underline{\underline{B}}^T \underline{\underline{E}} \underline{\underline{B}} \, dV$$

$$= t \int \underline{\underline{B}}^T \underline{\underline{E}} \underline{\underline{B}} \, dx \, dy \quad dx = a \, d\xi, \quad dy = b \, d\eta$$

$$= abt \int_{-1}^1 \int_{-1}^1 \underline{\underline{B}}^T \underline{\underline{E}} \underline{\underline{B}} \, d\xi \, d\eta$$

$\underline{\underline{E}}$ $\left\{ \begin{array}{l} \text{plane stress} \\ \text{plane strain} \end{array} \right.$

Equivalent nodal force



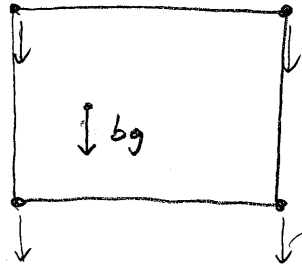
$$\underline{P}_b = \int \underline{f}^T \underline{b} dx$$

$$\underline{f}^T (\eta=+)$$

$$= \begin{bmatrix} \frac{1}{2}(1-\zeta) & 0 \\ 0 & \frac{1}{2}(1-\zeta) \\ \frac{1}{2}(1+\zeta) & 0 \\ 0 & \frac{1}{2}(1+\zeta) \\ 0 & 0 \\ 0 & 0 \\ 0 & 0 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} 0 \\ b_y \end{bmatrix} = \begin{bmatrix} 0 \\ \frac{1}{2}(1-\zeta)b_y \\ 0 \\ \frac{1}{2}(1+\zeta)b_y \\ 0 \\ 0 \\ 0 \\ 0 \end{bmatrix}$$

$$\underline{P}_b = \begin{bmatrix} 0 \\ \int_{-1}^1 \frac{1}{2}(1-\zeta) b_y a d \zeta \\ 0 \\ \int_{-1}^1 \frac{1}{2}(1+\zeta) b_y a d \zeta \\ 0 \\ 0 \\ 0 \\ 0 \end{bmatrix} = \frac{a}{3} \begin{bmatrix} 0 \\ (2b_{y1} + b_{y2}) \\ 0 \\ (b_{y1} + 2b_{y2}) \\ 0 \\ 0 \\ 0 \\ 0 \end{bmatrix}$$

gravity load



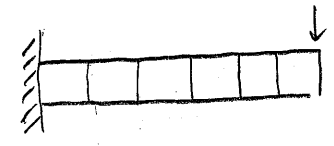
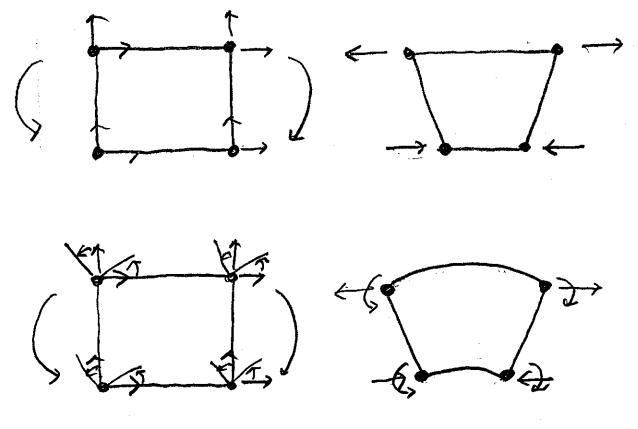
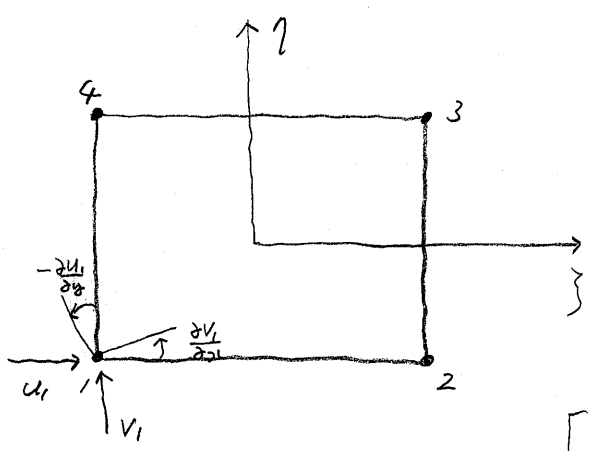
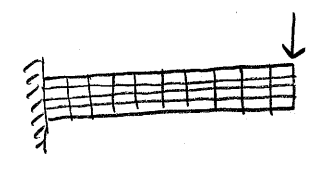
$$\underline{b} = \begin{bmatrix} 0 \\ -b_g \end{bmatrix}$$

equivalent nodal load = $-b_g abt$

$$\underline{P}_b = \int \underline{f}^T \underline{b} dv = abt \int_{-1}^1 \int_{-1}^1 \underline{f}^T \underline{b} d\xi d\eta$$

$$= -b_g abt \begin{bmatrix} 0 \\ \iint \frac{1}{4} (1-\xi)(1-\eta) d\xi d\eta \\ 0 \\ \iint \frac{1}{4} (1+\xi)(1-\eta) d\xi d\eta \\ 0 \\ \iint \frac{1}{4} (1+\xi)(1+\eta) d\xi d\eta \\ 0 \\ \iint \frac{1}{4} (1-\xi)(1+\eta) d\xi d\eta \end{bmatrix} = -b_g abt \begin{bmatrix} 0 \\ 1 \\ 0 \\ 1 \\ 0 \\ 1 \\ 0 \\ 1 \end{bmatrix}$$

Linear cubic Rectangle (Barber)

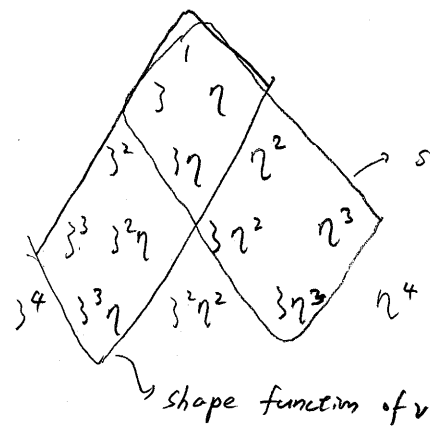


nodal displacement $\delta =$

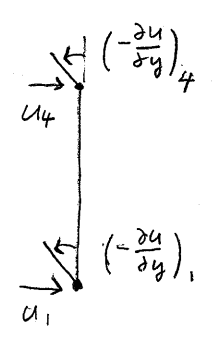
$$\begin{bmatrix} u_1 \\ v_1 \\ (\frac{\partial v}{\partial x})_1 \\ (-\frac{\partial u}{\partial y})_1 \\ \vdots \\ u_4 \\ v_4 \\ (\frac{\partial v}{\partial x})_4 \\ (-\frac{\partial u}{\partial y})_4 \end{bmatrix}$$

$$u = C_1 + C_2 \eta + C_3 \eta^2 + C_4 \eta^3 + C_5 \eta^2 + C_6 \eta^2 + C_7 \eta^3 + C_8 \eta^3 \quad \text{--- (1)}$$

$$v = C_9 + C_{10} \eta + C_{11} \eta^2 + C_{12} \eta^3 + C_{13} \eta^2 + C_{14} \eta^2 + C_{15} \eta^3 + C_{16} \eta^3 \quad \text{--- (2)}$$



in case of u ,



4 conditions

\Rightarrow 3rd order eq. with respect to η

* Condition for selecting polynomial eqns. 2 terms, 3rd order eqns

$$u = [f_1 \ f_2 \ f_3 \ f_4 \ f_5 \ f_6 \ f_7 \ f_8] \begin{bmatrix} u_1 \\ \left(\frac{\partial u}{\partial y}\right)_1 \\ \vdots \\ u_4 \\ \left(\frac{\partial u}{\partial y}\right)_4 \end{bmatrix} \quad u = f_u \underline{q}_u$$

$$v = [f_9 \ f_{10} \ f_{11} \ f_{12} \ f_{13} \ f_{14} \ f_{15} \ f_{16}] \begin{bmatrix} v_1 \\ \left(\frac{\partial v}{\partial x}\right)_1 \\ \vdots \\ v_4 \\ \left(\frac{\partial v}{\partial x}\right)_4 \end{bmatrix} \quad v = f_v \underline{q}_v$$

f_i 's can be defined with C_i 's by using B/C :

at node 1, $u = u_1, \left(\frac{\partial u}{\partial y}\right)_1 = \left(\frac{\partial u}{\partial y}\right)_1$, etc

$$\underline{u} = \begin{bmatrix} u \\ v \end{bmatrix} = \begin{bmatrix} f_1 & 0 & 0 & f_2 & f_3 & 0 & 0 & f_4 & f_5 & 0 & 0 & f_6 & f_7 & 0 & 0 & f_8 \\ 0 & f_9 & f_{10} & 0 & 0 & f_{11} & f_{12} & 0 & 0 & f_{13} & f_{14} & 0 & 0 & f_{15} & f_{16} & 0 \end{bmatrix} \begin{bmatrix} u_1 \\ v_1 \\ \left(\frac{\partial v}{\partial x}\right)_1 \\ \left(\frac{\partial u}{\partial y}\right)_1 \\ \vdots \\ \vdots \end{bmatrix}$$

\underline{f}

$$\underline{\epsilon} = \underline{d} \underline{u} = \underline{d} \underline{f} \underline{q} = \underline{B} \underline{q}$$

$$\underline{B} = \underline{d} \underline{f} \quad \underline{d} = \begin{bmatrix} \frac{\partial}{\partial x} & 0 \\ 0 & \frac{\partial}{\partial y} \\ \frac{\partial}{\partial y} & \frac{\partial}{\partial x} \end{bmatrix}$$

$$\underline{B} = \begin{bmatrix} f_{1,x} & 0 & 0 & f_{2,x} & f_{7,x} & 0 & 0 & f_{8,x} \\ 0 & f_{9,y} & f_{10,y} & 0 & 0 & f_{15,y} & f_{16,y} & 0 \\ f_{1,y} & f_{9,x} & f_{10,x} & f_{2,y} & f_{7,y} & f_{15,x} & f_{16,x} & f_{8,y} \end{bmatrix}$$

$$K = \int_V \underline{B}^T \underline{\epsilon} \underline{B} \, dV$$

$$f_{i,x} = \frac{\partial f_i}{\partial x} \frac{\partial x}{\partial \eta}$$

$$f_{i,y} = \frac{\partial f_i}{\partial y} \frac{\partial y}{\partial \eta}$$

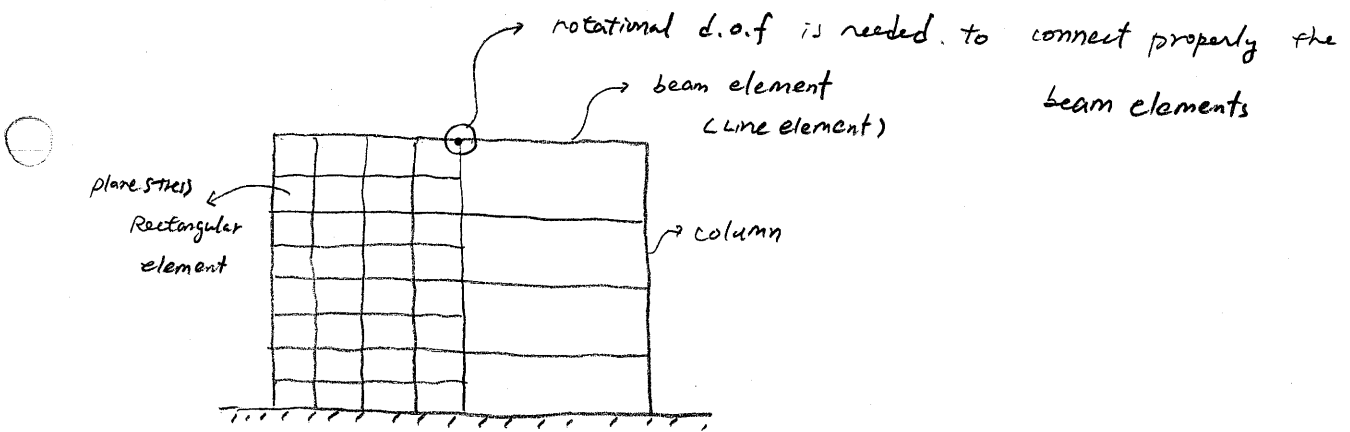
The shape function of u is linear in ξ and cubic in η

The shape function of v is linear in η and cubic in ξ .

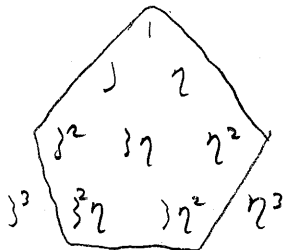
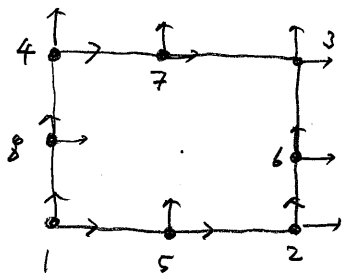
⇒ Linear - cubic Rectangle

The cubic functions allow compatibility between any edge of the rectangle and a bordering flexural element

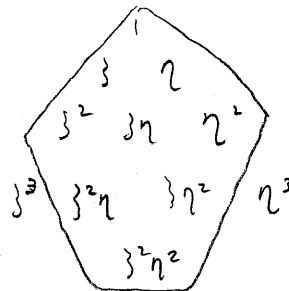
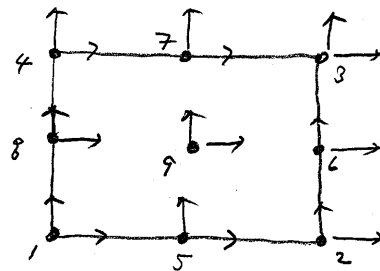
For example

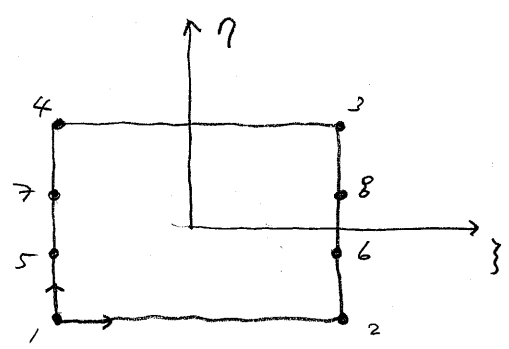


Eight - node Element



nine - node Element





four conditions in η direction
two conditions in ξ direction

$$\left\{ \begin{aligned} u &= C_1 + C_2 \xi + C_3 \eta + C_4 \xi \eta + C_5 \eta^2 + C_6 \xi \eta^2 + C_7 \eta^3 + C_8 \xi \eta^3 \\ v &= C_9 + C_{10} \xi + C_{11} \eta + C_{12} \xi \eta + C_{13} \eta^2 + C_{14} \xi \eta^2 + C_{15} \eta^3 + C_{16} \xi \eta^3 \end{aligned} \right.$$

Nodal displacements should be defined differently in the η direction.

...
...
...
...
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