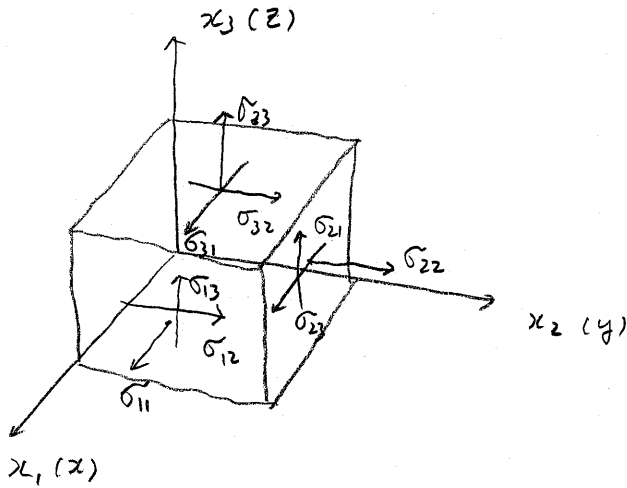
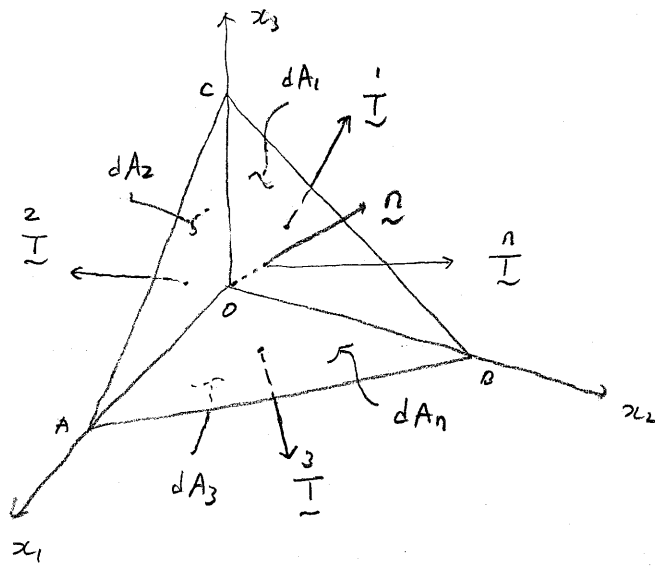


Chapter 4. General Solids



σ_{ij} → direction of stress
 → orientation of face



$$\underline{n} = n_1 \underline{i} + n_2 \underline{j} + n_3 \underline{k}$$

$$\underline{e}_1 = \underline{i}$$

$$\underline{e}_2 = \underline{j}$$

$$\underline{e}_3 = \underline{k}$$

$$\frac{1}{\underline{I}} = \sigma_{11} \underline{i} + \sigma_{12} \underline{j} + \sigma_{13} \underline{k}$$

$$\frac{2}{\underline{I}} = \sigma_{21} \underline{i} + \sigma_{22} \underline{j} + \sigma_{23} \underline{k}$$

$$\frac{3}{\underline{I}} = \sigma_{31} \underline{i} + \sigma_{32} \underline{j} + \sigma_{33} \underline{k}$$

$$\frac{\underline{I}}{\underline{I}} = \sigma_{ij}$$

\underline{I} : surface traction

$$\text{Area of } ABC = dA_n$$

$$OBC = dA_1 = dA_n \cdot \underline{r} \cdot \underline{e}_1 = dA_n n_1$$

$$OCA = dA_2 = dA_n \cdot \underline{r} \cdot \underline{e}_2 = dA_n n_2$$

$$OAB = dA_3 = dA_n \cdot \underline{r} \cdot \underline{e}_3 = dA_n n_3$$

or

$$\frac{dA_1}{dA_n} = n_1, \quad \frac{dA_2}{dA_n} = n_2, \quad \frac{dA_3}{dA_n} = n_3$$

Force Equilibrium

$$\frac{n}{\underline{I}} dA_n = \frac{1}{\underline{I}} dA_1 + \frac{2}{\underline{I}} dA_2 + \frac{3}{\underline{I}} dA_3$$

$$\begin{aligned} \frac{n}{\underline{I}} &= \frac{1}{\underline{I}} \frac{dA_1}{dA_n} + \frac{2}{\underline{I}} \frac{dA_2}{dA_n} + \frac{3}{\underline{I}} \frac{dA_3}{dA_n} \\ &= \frac{1}{\underline{I}} n_1 + \frac{2}{\underline{I}} n_2 + \frac{3}{\underline{I}} n_3 \\ &= \frac{3}{\underline{I}} n_j \end{aligned}$$

or

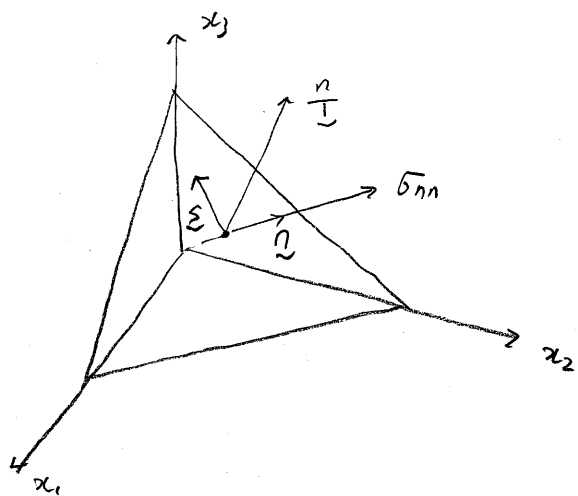
$$\begin{aligned} \frac{n}{T_i} dA_n &= \frac{1}{T_i} dA_1 + \frac{2}{T_i} dA_2 + \frac{3}{T_i} dA_3 \\ &= \sigma_{1i} dA_1 + \sigma_{2i} dA_2 + \sigma_{3i} dA_3 \\ \frac{n}{T_i} &= \sigma_{1i} n_1 + \sigma_{2i} n_2 + \sigma_{3i} n_3 \\ &= \sigma_{ji} n_j \end{aligned}$$

Generally $\frac{n}{T_i} = \sigma_{ji} n_j$ (Cauchy's formula)

or $= \sigma_{ij} n_j$ ($\sigma_{ij} = \sigma_{ji}$)

$$\begin{cases} \frac{n}{l_1} = \sigma_{11} n_1 + \sigma_{21} n_2 + \sigma_{31} n_3 \\ \frac{n}{l_2} = \sigma_{12} n_1 + \sigma_{22} n_2 + \sigma_{32} n_3 \\ \frac{n}{l_3} = \sigma_{13} n_1 + \sigma_{23} n_2 + \sigma_{33} n_3 \end{cases}$$

$$\begin{aligned} \frac{n}{l} &= \begin{bmatrix} \frac{n}{l_1} \\ \frac{n}{l_2} \\ \frac{n}{l_3} \end{bmatrix} = \begin{bmatrix} \sigma_{11} & \sigma_{21} & \sigma_{31} \\ \sigma_{12} & \sigma_{22} & \sigma_{32} \\ \sigma_{13} & \sigma_{23} & \sigma_{33} \end{bmatrix} \begin{bmatrix} n_1 \\ n_2 \\ n_3 \end{bmatrix} \\ &= \begin{bmatrix} \sigma_{11} & \sigma_{12} & \sigma_{13} \\ \sigma_{21} & \sigma_{22} & \sigma_{23} \\ \sigma_{31} & \sigma_{32} & \sigma_{33} \end{bmatrix} \begin{bmatrix} n_1 \\ n_2 \\ n_3 \end{bmatrix} \end{aligned}$$



normal stress

$$\sigma_{nn} = \frac{n_i \sigma_{ij} n_j}{n^2}$$

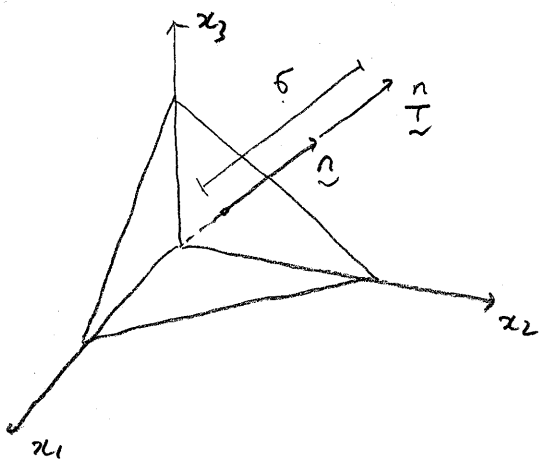
normal stress

$$\begin{aligned}\sigma_{nn} &= \underline{n} \cdot \underline{\sigma} = \underline{n} \cdot \underline{\underline{\sigma}} = \underline{n} \cdot (\sigma_{ij} \underline{e}_j) \\ &= [n_1 \ n_2 \ n_3] \begin{bmatrix} \frac{n_1}{r} \\ \frac{n_2}{r} \\ \frac{n_3}{r} \end{bmatrix} \\ &= [n_1 \ n_2 \ n_3] \begin{bmatrix} \sigma_{11} & \sigma_{12} & \sigma_{13} \\ \sigma_{21} & \sigma_{22} & \sigma_{23} \\ \sigma_{31} & \sigma_{32} & \sigma_{33} \end{bmatrix} \begin{bmatrix} n_1 \\ n_2 \\ n_3 \end{bmatrix}\end{aligned}$$

shear stress

$$\begin{aligned}\sigma_{ns} &= \underline{n} \cdot \underline{s} \\ &= [s_1 \ s_2 \ s_3] \begin{bmatrix} \sigma_{11} & \sigma_{12} & \sigma_{13} \\ \sigma_{21} & \sigma_{22} & \sigma_{23} \\ \sigma_{31} & \sigma_{32} & \sigma_{33} \end{bmatrix} \begin{bmatrix} n_1 \\ n_2 \\ n_3 \end{bmatrix}\end{aligned}$$

principal stress



$$\left. \begin{aligned} \frac{n}{T_i} &= \sigma_{ij} n_j \\ \frac{n}{T_i} &= \sigma n_i \end{aligned} \right\} \rightarrow \sigma n_i = \sigma_{ij} n_j$$

$$\begin{cases} \sigma n_1 = \sigma_{11} n_1 + \sigma_{12} n_2 + \sigma_{13} n_3 \\ \sigma n_2 = \sigma_{21} n_1 + \sigma_{22} n_2 + \sigma_{23} n_3 \\ \sigma n_3 = \sigma_{31} n_1 + \sigma_{32} n_2 + \sigma_{33} n_3 \end{cases}$$

$$\begin{bmatrix} \sigma & 0 & 0 \\ 0 & \sigma & 0 \\ 0 & 0 & \sigma \end{bmatrix} \begin{bmatrix} n_1 \\ n_2 \\ n_3 \end{bmatrix} = \begin{bmatrix} \sigma_{11} & \sigma_{12} & \sigma_{13} \\ \sigma_{21} & \sigma_{22} & \sigma_{23} \\ \sigma_{31} & \sigma_{32} & \sigma_{33} \end{bmatrix} \begin{bmatrix} n_1 \\ n_2 \\ n_3 \end{bmatrix}$$

$$\sigma \underline{I} \underline{n} = \underline{\sigma} \underline{n}$$

$$[\underline{\sigma} - \sigma \underline{I}] \underline{n} = \underline{0}$$

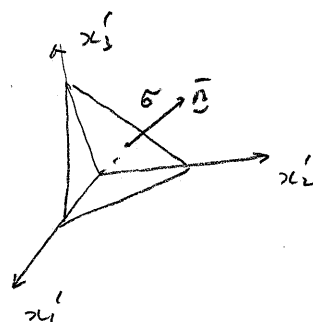
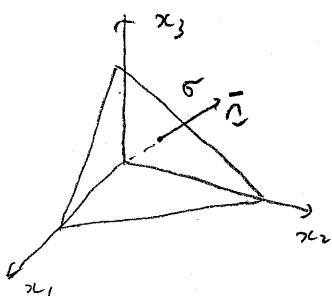
$$\text{or } (\sigma_{ij} - \sigma \delta_{ij}) n_j = 0$$

for non-trivial solution

$$|\underline{\sigma} - \sigma \underline{I}| = 0$$

$$\sigma^3 + I_1 \sigma^2 + I_2 \sigma + I_3 = 0$$

$I_1, I_2, I_3 =$ stress invariants



some principal stresses

$\Rightarrow I_1, I_2, I_3$ are
= invariants

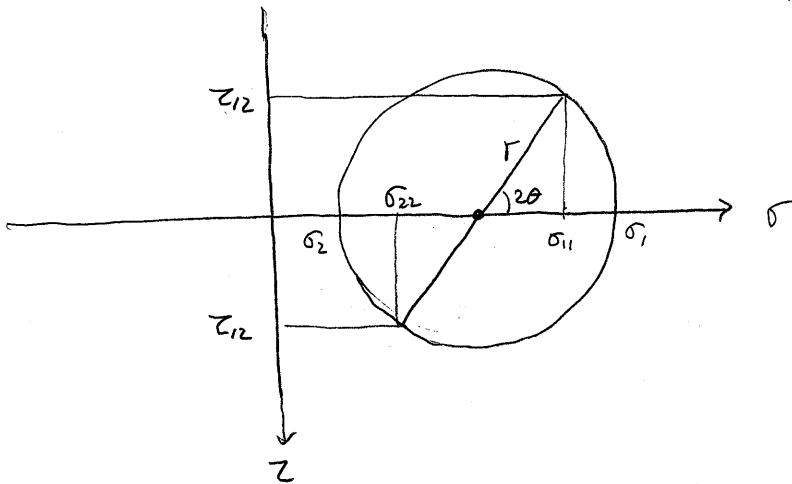
$$I_1 = \sigma_{11} + \sigma_{22} + \sigma_{33} = \sigma_{kk}$$

$$\begin{aligned} I_2 &= \begin{vmatrix} \sigma_{22} & \sigma_{23} \\ \sigma_{32} & \sigma_{33} \end{vmatrix} + \begin{vmatrix} \sigma_{11} & \sigma_{13} \\ \sigma_{31} & \sigma_{33} \end{vmatrix} + \begin{vmatrix} \sigma_{11} & \sigma_{12} \\ \sigma_{21} & \sigma_{22} \end{vmatrix} \\ &= \sigma_{11}\sigma_{22} + \sigma_{22}\sigma_{33} + \sigma_{11}\sigma_{33} - \sigma_{12}^2 - \sigma_{23}^2 - \sigma_{31}^2 \\ &= \frac{1}{2} (\sigma_{ii}\sigma_{jj} - \sigma_{ij}\sigma_{ji}) \end{aligned}$$

$$I_3 = \det(\sigma_{ij})$$

$$= \sigma_{11}\sigma_{22}\sigma_{33} + 2\sigma_{12}\sigma_{23}\sigma_{31} - \sigma_{12}^2\sigma_{33} - \sigma_{23}^2\sigma_{11} - \sigma_{31}^2\sigma_{22}$$

Mohr's circle



stress invariants : $\frac{1}{2} (\sigma_{11} + \sigma_{22})$ or $\frac{1}{2} (\sigma_1 + \sigma_2)$

$$\begin{aligned} r^2 &= \left[\frac{1}{2} (\sigma_{11} - \sigma_{22}) \right]^2 + \tau_{12}^2 \\ &= \frac{1}{4} [\sigma_{11}^2 - 2\sigma_{11}\sigma_{22} + \sigma_{22}^2] + \tau_{12}^2 \\ &= \frac{1}{4} (\sigma_{11} + \sigma_{22})^2 - \sigma_{11}\sigma_{22} + \tau_{12}^2 \\ &= \frac{1}{4} I_1^2 - I_2 \end{aligned}$$

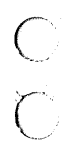
Material Strength Criteria

* J_2 theory - von Mises criterion (For metals)

$$s^3 + J_1 s^2 + J_2 s + J_3 = 0 \qquad s = \sigma - \frac{1}{3} I_1$$

$$J_1 = 0$$

$$\sigma_y = \sqrt{3 J_2}$$

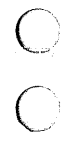


$$\sqrt{3 J_2} = \sqrt{\frac{1}{2} [(\sigma_{11} - \sigma_{22})^2 + (\sigma_{22} - \sigma_{33})^2 + (\sigma_{33} - \sigma_{11})^2] + 6 (\tau_{12}^2 + \tau_{23}^2 + \tau_{31}^2)}$$

$$= \sqrt{\frac{1}{2} [(\sigma_1 - \sigma_2)^2 + (\sigma_2 - \sigma_3)^2 + (\sigma_3 - \sigma_1)^2]}$$

* Elastic - hardening plasticity model (For concrete)

$$A_1 \sigma + A_2 \sqrt{3 J_2} + A_3 = \bar{\sigma} \qquad \sigma = \frac{1}{3} \frac{(\sigma_1 + \sigma_2 + \sigma_3)}{I_1}$$



Why are the existing material failure criteria defined with the stress - invariants ?

⇒ For isotropic materials and similarities,
 the material failure criteria should not be dependent on
 the definition of coordinate system.

$$|\underline{\sigma} - \sigma \underline{I}| = 0$$

solutions = $\sigma_1, \sigma_2, \sigma_3$ Eigenvalues \Rightarrow principal stresses

corresponding vectors $\bar{n}_1, \bar{n}_2, \bar{n}_3$ Eigenvalues

\Rightarrow orientations of principal axes

$$\underline{N} = [\bar{n}_1 \bar{n}_2 \bar{n}_3]$$

$$= \begin{bmatrix} \bar{n}_{11} & \bar{n}_{21} & \bar{n}_{31} \\ \bar{n}_{12} & \bar{n}_{22} & \bar{n}_{32} \\ \bar{n}_{13} & \bar{n}_{23} & \bar{n}_{33} \end{bmatrix}$$

$$\left. \begin{aligned} \sigma_1 \underline{I} \bar{n}_1 &= \underline{\sigma} \bar{n}_1 \\ \sigma_2 \underline{I} \bar{n}_2 &= \underline{\sigma} \bar{n}_2 \\ \sigma_3 \underline{I} \bar{n}_3 &= \underline{\sigma} \bar{n}_3 \end{aligned} \right\} \Rightarrow$$

$$\sigma_1 \bar{n}_1^T \bar{n}_1 = \bar{n}_1^T \underline{\sigma} \bar{n}_1$$

$$\bar{n}_1^T \bar{n}_1 = 1$$

$$\begin{cases} \sigma_1 = \bar{n}_1^T \underline{\sigma} \underline{N} & \bar{n}_i^T \underline{\sigma} \bar{n}_j = 0 \\ \sigma_2 = \bar{n}_2^T \underline{\sigma} \underline{N} \\ \sigma_3 = \bar{n}_3^T \underline{\sigma} \underline{N} \end{cases}$$

$$\underline{\sigma} = \underline{N}^T \underline{\sigma} \underline{N}$$

\underline{N} : orthogonal matrix

$$\underline{N}^T \underline{N} = \underline{I}$$

Orthogonality

$$\sigma \bar{n}_i = \sigma_i \bar{n}_i \quad - \textcircled{1} \quad (i \neq j)$$

$$\sigma \bar{n}_j = \sigma_j \bar{n}_j \quad - \textcircled{2}$$

$$\textcircled{1} - \bar{n}_j^T \sigma \bar{n}_i = \sigma_i \bar{n}_j^T \bar{n}_i$$

$$\textcircled{2} - \bar{n}_i^T \sigma \bar{n}_j = \sigma_j \bar{n}_i^T \bar{n}_j \quad - \textcircled{2}'$$

$$\textcircled{2}' - \bar{n}_j^T \sigma \bar{n}_i = \sigma_j \bar{n}_j^T \bar{n}_i$$

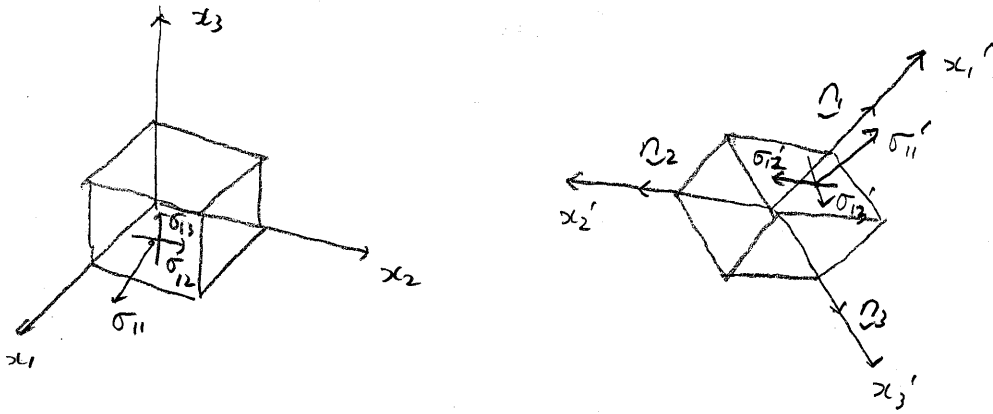
from $\textcircled{1}$ and $\textcircled{2}'$, $\sigma_i \bar{n}_j^T \bar{n}_i = \sigma_j \bar{n}_j^T \bar{n}_i$

$$(\sigma_i - \sigma_j) \bar{n}_j^T \bar{n}_i = 0$$

since $\sigma_i \neq \sigma_j$, $\bar{n}_j^T \bar{n}_i = 0 \quad (i \neq j)$

$$\Rightarrow \bar{n}_j \perp \bar{n}_i \quad (\text{orthogonality})$$

Axis - transformation of stresses



$$\underline{n}_1 = \begin{bmatrix} n_{11} \\ n_{12} \\ n_{13} \end{bmatrix} \quad \underline{n}_2 = \begin{bmatrix} n_{21} \\ n_{22} \\ n_{23} \end{bmatrix} \quad \underline{n}_3 = \begin{bmatrix} n_{31} \\ n_{32} \\ n_{33} \end{bmatrix}$$

$$\sigma_{11}' = \sigma_{n_1 n_1} = [n_{11} \ n_{12} \ n_{13}] \begin{bmatrix} \sigma_{11} & \sigma_{12} & \sigma_{13} \\ \sigma_{21} & \sigma_{22} & \sigma_{23} \\ \sigma_{31} & \sigma_{32} & \sigma_{33} \end{bmatrix} \begin{bmatrix} n_{11} \\ n_{12} \\ n_{13} \end{bmatrix}$$

$$= n_{1l} \sigma_{kl} n_{1k}$$

$$\sigma_{21}' = \sigma_{n_2 n_1} = [n_{21} \ n_{22} \ n_{23}] \begin{bmatrix} \sigma_{11} & \sigma_{12} & \sigma_{13} \\ \sigma_{21} & \sigma_{22} & \sigma_{23} \\ \sigma_{31} & \sigma_{32} & \sigma_{33} \end{bmatrix} \begin{bmatrix} n_{11} \\ n_{12} \\ n_{13} \end{bmatrix}$$

(= σ_{3n})

$$= n_{2l} \sigma_{kl} n_{1k}$$

Generally $\sigma_{ij}' = n_{il} \sigma_{kl} n_{jk}$

$$\begin{matrix} \sigma_{11} & \sigma_{12} & \sigma_{13} \\ \sigma_{21} & \sigma_{22} & \sigma_{23} \\ \sigma_{31} & \sigma_{32} & \sigma_{33} \end{matrix} = \begin{matrix} n_{11} & n_{12} & n_{13} \\ n_{21} & n_{22} & n_{23} \\ n_{31} & n_{32} & n_{33} \end{matrix} \begin{matrix} \sigma_{11} & \sigma_{12} & \sigma_{13} \\ \sigma_{21} & \sigma_{22} & \sigma_{23} \\ \sigma_{31} & \sigma_{32} & \sigma_{33} \end{matrix} \begin{matrix} n_{11} & n_{12} & n_{31} \\ n_{12} & n_{22} & n_{32} \\ n_{13} & n_{23} & n_{33} \end{matrix}$$

$\underbrace{\quad}_{\Sigma'}$
 \quad
 $\underbrace{\quad}_{R^T}$
 \quad
 $\underbrace{\quad}_{\Sigma}$
 \quad
 $\underbrace{\quad}_{R}$

$$\Sigma' = R^T \Sigma R$$

In text book,

$$\underline{R} = \begin{bmatrix} l_1 & l_2 & l_3 \\ m_1 & m_2 & m_3 \\ n_1 & n_2 & n_3 \end{bmatrix}$$

$$\Sigma' = T_{\Sigma} \Sigma$$

$$\begin{matrix} \sigma_{11}' \\ \sigma_{22}' \\ \sigma_{33}' \\ \sigma_{12}' \\ \sigma_{23}' \\ \sigma_{31}' \end{matrix} = \begin{matrix} l_1^2 & m_1^2 & n_1^2 & 2l_1m_1 & 2m_1n_1 & 2n_1l_1 \\ l_2^2 & m_2^2 & n_2^2 & 2l_2m_2 & 2m_2n_2 & 2n_2l_2 \\ l_3^2 & m_3^2 & n_3^2 & 2l_3m_3 & 2m_3n_3 & 2n_3l_3 \\ \hline l_1l_2 & m_1m_2 & n_1n_2 & (l_1m_2+l_2m_1) & (m_1n_2+m_2n_1) & (n_1l_2+n_2l_1) \\ l_2l_3 & m_2m_3 & n_2n_3 & (l_2m_3+l_3m_2) & (m_2n_3+m_3n_2) & (n_2l_3+n_3l_2) \\ l_3l_1 & m_3m_1 & n_3n_1 & (l_3m_1+l_1m_3) & (m_3n_1+m_1n_3) & (n_3l_1+n_1l_3) \end{matrix} \begin{matrix} \sigma_{11} \\ \sigma_{22} \\ \sigma_{33} \\ \sigma_{12} \\ \sigma_{23} \\ \sigma_{31} \end{matrix}$$

$\underbrace{\quad}_{\Sigma'}$
 \quad
 $\underbrace{\quad}_{T_{\Sigma}}$
 \quad
 $\underbrace{\quad}_{\Sigma}$

$$\underline{\underline{\Sigma}} = \begin{bmatrix} \Sigma_{11} & \Sigma_{12} & \Sigma_{13} \\ \Sigma_{21} & \Sigma_{22} & \Sigma_{23} \\ \Sigma_{31} & \Sigma_{32} & \Sigma_{33} \end{bmatrix} \quad \text{where} \quad \Sigma_{12} = \frac{1}{2} \gamma_{12}, \quad \Sigma_{13} = \frac{1}{2} \gamma_{13} \dots$$

$$\underline{\underline{\Sigma}}' = \underline{\underline{R}}^T \underline{\underline{\Sigma}} \underline{\underline{R}}$$

3x3

also,

$$\underline{\underline{\Sigma}}' = \underline{\underline{T}}_{\sigma} \underline{\underline{\Sigma}}$$

6x1

if we use,

$$\underline{\underline{\Sigma}} = [\Sigma_{11} \quad \Sigma_{22} \quad \Sigma_{33} \quad \gamma_{12} \quad \gamma_{23} \quad \gamma_{31}]^T$$

$$\underline{\underline{\Sigma}}' = \underline{\underline{T}}_{\epsilon} \underline{\underline{\Sigma}} \quad \underline{\underline{T}}_{\epsilon} \neq \underline{\underline{T}}_{\sigma}$$

using virtual strain energy density

$$(\delta \underline{\underline{\Sigma}}')^T \underline{\underline{\sigma}}' = \delta \underline{\underline{\Sigma}}^T \underline{\underline{\sigma}}$$

$$\delta \underline{\underline{\Sigma}}^T \underline{\underline{T}}_{\epsilon}^T \underline{\underline{\sigma}}' = \delta \underline{\underline{\Sigma}}^T \underline{\underline{\sigma}}$$

$$\underline{\underline{T}}_{\epsilon}^T \cdot \underline{\underline{T}}_{\sigma} \underline{\underline{\sigma}} = \underline{\underline{\sigma}}$$

$$\underline{\underline{T}}_{\sigma} = \underline{\underline{T}}_{\epsilon}^{-T}$$

Assume

Axis-transformation of constitutive matrix \underline{E}

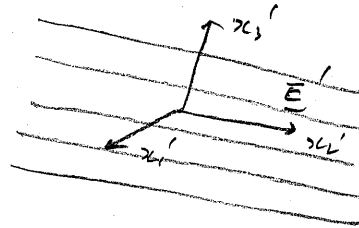
$$\underline{\sigma}' = \underline{E}' \underline{\epsilon}'$$

$$\underline{T}_\sigma \underline{\sigma} = \underline{E}' \underline{T}_\epsilon \underline{\epsilon}$$

$$\underline{\sigma} = \underline{T}_\sigma^{-1} \underline{E}' \underline{T}_\epsilon \underline{\epsilon}$$

$$= \underline{T}_\epsilon^T \underline{E}' \underline{T}_\epsilon \underline{\epsilon}$$

$$= \underline{E} \underline{\epsilon}$$



Therefore, $\underline{E} = \underline{T}_\epsilon^T \underline{E}' \underline{T}_\epsilon$

However, for isotropic material, $\underline{E} = \underline{E}'$

Stress-Strain Relationship

For Iso

$$\underline{\sigma} = \underline{E} \underline{\epsilon}$$

for isotropic material

$$\underline{E} = \frac{E}{(1+\nu)e_2} \begin{bmatrix} e_1 & \nu & \nu & & & \\ & e_1 & \nu & & & \\ & & e_1 & & & \\ & & & \text{sym.} & & \\ & & & & e_3 & \\ & & & & & e_3 \\ & & & & & & e_3 \end{bmatrix}$$

$$e_1 = 1 - \nu, \quad e_2 = 1 - 2\nu, \quad e_3 = \frac{e_2}{2}$$

\underline{E} can be defined with 2 independent constants E and ν

why?

Anisotropic material

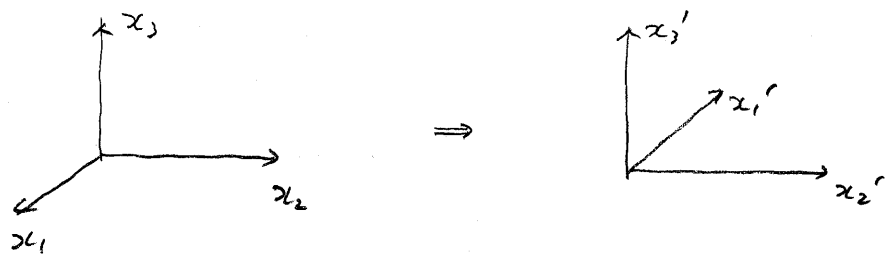
$$\sigma_i = C_{ij} \epsilon_j$$

$$C_{ij} = \begin{bmatrix} C_{11} & C_{12} & C_{13} & C_{14} & C_{15} & C_{16} \\ & C_{22} & C_{23} & C_{24} & C_{25} & C_{26} \\ & & C_{33} & C_{34} & C_{35} & C_{36} \\ & & & C_{44} & C_{45} & C_{46} \\ & & & & C_{55} & C_{56} \\ & & & & & C_{66} \end{bmatrix}$$

Sym.

21 constants

Elastic symmetry with respect to one plane (Monoclinic material)



Transformation matrix

$$\underline{R} = \begin{bmatrix} n_{11} & n_{12} & n_{13} \\ n_{21} & n_{22} & n_{23} \\ n_{31} & n_{32} & n_{33} \end{bmatrix} = \begin{bmatrix} -1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{matrix} x_1' \\ x_2' \\ x_3' \end{matrix}$$

Seek condition of \underline{C} to satisfy $\underline{\sigma} = \underline{C} \underline{\epsilon}$ and $\underline{\sigma}' = \underline{C} \underline{\epsilon}'$,
 where $\underline{\sigma}' = \underline{R}^T \underline{\sigma} \underline{R}$ and $\underline{\epsilon}' = \underline{R} \underline{\epsilon} \underline{R}^T$.

$$\underline{\sigma}' = R^T \underline{\sigma} R$$

$$\Rightarrow \sigma_{11}' = \sigma_{11}, \quad \sigma_{22}' = \sigma_{22}, \quad \sigma_{33}' = \sigma_{33}$$

$$\sigma_{12}' = -\sigma_{12}, \quad \sigma_{23}' = \sigma_{23}, \quad \sigma_{31}' = -\sigma_{31}$$

Also $\underline{\epsilon}' = R^T \underline{\epsilon} R$

$$\Rightarrow \epsilon_{11}' = \epsilon_{11}, \quad \epsilon_{22}' = \epsilon_{22}, \quad \epsilon_{33}' = \epsilon_{33}$$

$$\epsilon_{12}' = -\epsilon_{12}, \quad \epsilon_{23}' = \epsilon_{23}, \quad \epsilon_{31}' = -\epsilon_{31}$$

from $\underline{\sigma} = \underline{C} \underline{\epsilon}$

$$\sigma_{11} = C_{11} \epsilon_{11} + C_{12} \epsilon_{22} + C_{13} \epsilon_{33} + C_{14} \epsilon_{12} + C_{15} \epsilon_{23} + C_{16} \epsilon_{31}$$

from $\underline{\sigma}' = \underline{C}' \underline{\epsilon}'$

$$\sigma_{11}' = C_{11}' \epsilon_{11}' + C_{12}' \epsilon_{22}' + C_{13}' \epsilon_{33}' + C_{14}' \epsilon_{12}' + C_{15}' \epsilon_{23}' + C_{16}' \epsilon_{31}'$$

$$= C_{11} \epsilon_{11} + C_{12} \epsilon_{22} + C_{13} \epsilon_{33} - C_{14} \epsilon_{12} + C_{15} \epsilon_{23} - C_{16} \epsilon_{31}$$

from $\sigma_{11} = \sigma_{11}'$, $C_{14} = C_{16} = 0$

Similarly, $C_{24} = C_{26} = C_{34} = C_{36} = C_{45} = C_{36} = 0$

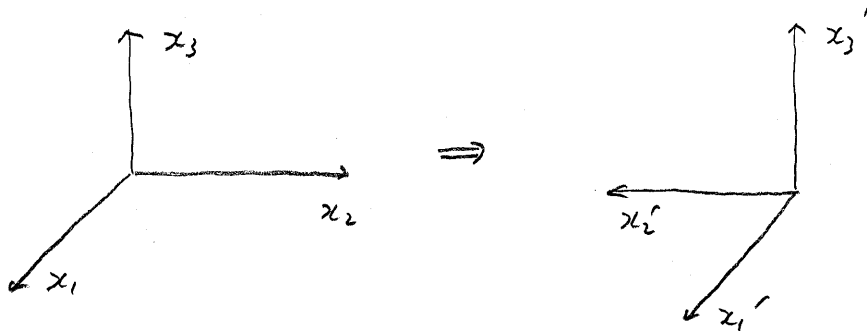
Hence,

$$C_{ij} = \begin{bmatrix} C_{11} & C_{12} & C_{13} & 0 & C_{15} & 0 \\ & C_{22} & C_{23} & 0 & C_{25} & 0 \\ & & C_{33} & 0 & C_{35} & 0 \\ & & & C_{44} & 0 & C_{45} \\ & & & & C_{55} & 0 \\ & & & & & C_{66} \end{bmatrix}$$

13 independent
constants
for monoclinic
material

Elastic symmetry w.r.t two orthogonal planes

(Orthotropic material)



$$(R) = \begin{matrix} & x_1 & x_2 & x_3 \\ \begin{bmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 1 \end{bmatrix} & x_1' \\ & x_2' \\ & x_3' \end{matrix}$$

from $\underline{\epsilon}' = \underline{C} \underline{\epsilon}$, $\sigma_{ii}' = \sigma_{ii}$, $\sigma_{12}' = -\sigma_{12}$, $\sigma_{13}' = -\sigma_{13}$

$\sigma_{23}' = -\sigma_{23}$, $\sigma_{31}' = \sigma_{31}$

$$C_{ij} = \begin{bmatrix} C_{11} & C_{12} & C_{13} & 0 & 0 & 0 \\ & C_{22} & C_{23} & 0 & 0 & 0 \\ & & C_{33} & 0 & 0 & 0 \\ & & & C_{44} & 0 & 0 \\ & & & & C_{55} & 0 \\ & & & & & C_{66} \end{bmatrix}$$

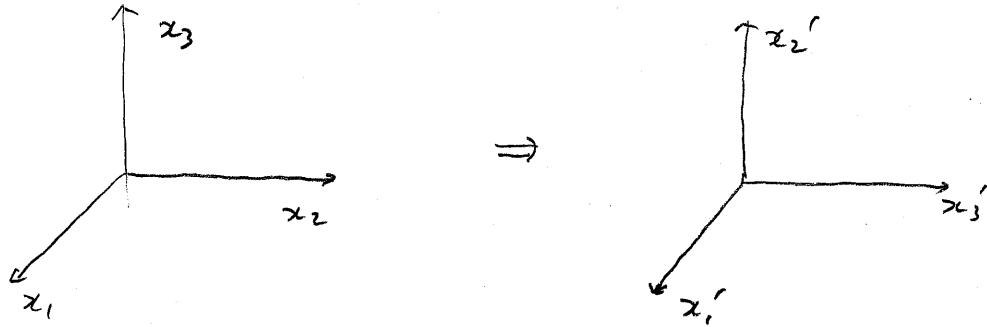
9 constants

for orthotropic material

No coupling between (normal stress - shear strains
normal strain - shear stresses

Directional Independence (Interchange Axes)

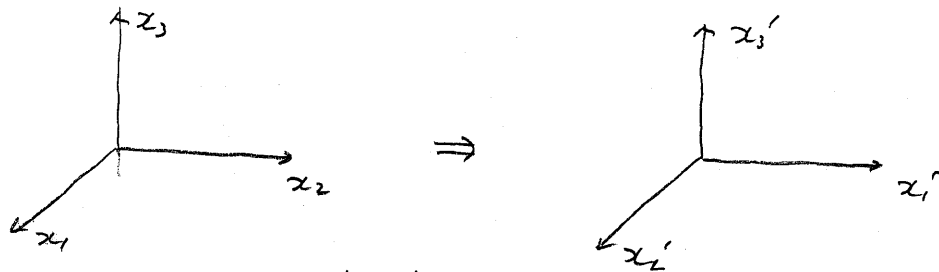
(Cubic material)



$$\underline{R} = \begin{array}{ccc|c} & x_1' & x_2' & x_3' \\ \begin{array}{c} x_1 \\ x_2 \\ x_3 \end{array} & \begin{bmatrix} 1 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{bmatrix} & \begin{array}{c} x_1' \\ x_2' \\ x_3' \end{array} \end{array}$$

from $\underline{\underline{C}}' = \underline{\underline{C}} \underline{\underline{R}}^T$ $\sigma_{11}' = \sigma_{11}$ $\sigma_{22}' = \sigma_{33}$, $\sigma_{33}' = \sigma_{22}$
 $\sigma_{12}' = \sigma_{13}$ $\sigma_{23}' = \sigma_{32}$, $\sigma_{13}' = \sigma_{12}$

$$\Rightarrow C_{22} = C_{33}, \quad C_{44} = C_{66}, \quad C_{12} = C_{13},$$



$$\underline{R} = \begin{array}{ccc|c} & x_1' & x_2' & x_3' \\ \begin{array}{c} x_1 \\ x_2 \\ x_3 \end{array} & \begin{bmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 1 \end{bmatrix} & \begin{array}{c} x_1' \\ x_2' \\ x_3' \end{array} \end{array}$$

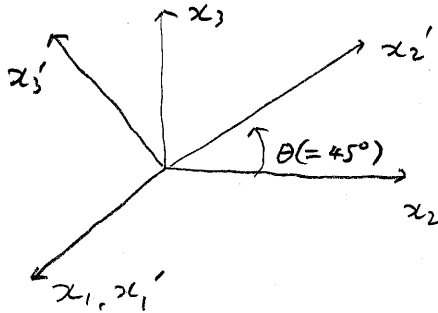
from $\underline{\underline{C}}' = \underline{\underline{C}} \underline{\underline{R}}^T$, $\sigma_{11}' = \sigma_{33}$, $\sigma_{22}' = \sigma_{22}$, $\sigma_{33}' = \sigma_{11}$
 $\sigma_{12}' = \sigma_{12}$, $\sigma_{23}' = \sigma_{31}$, $\sigma_{31}' = \sigma_{23}$

$$\Rightarrow C_{11} = C_{33}, \quad C_{55} = C_{66}, \quad C_{12} = C_{23}$$

$$C_{ij} = \begin{bmatrix} C_{11} & C_{12} & C_{12} & 0 & 0 & 0 \\ & C_{11} & C_{12} & 0 & 0 & 0 \\ & & C_{11} & 0 & 0 & 0 \\ & & & C_{44} & 0 & 0 \\ & \text{sym.} & & & C_{44} & 0 \\ & & & & & C_{44} \end{bmatrix}$$

3 independent constants
for a cubic material

Rotational Independence (Isotropic material)



$$\underline{R} = \begin{bmatrix} x_1 & x_2 & x_3 \\ 1 & 0 & 0 \\ 0 & \cos\theta + \sin\theta & \\ 0 & \sin\theta & \cos\theta \end{bmatrix} \begin{matrix} x_1' \\ x_2' \\ x_3' \end{matrix}$$

from $\underline{\epsilon}' = \underline{\epsilon}$, $\underline{\sigma}' = \underline{R}^T \underline{\sigma} \underline{R}$, $\underline{\epsilon}' = \underline{R}^T \underline{\epsilon} \underline{R}$
 $\underline{\sigma} = \underline{\epsilon} \underline{\epsilon}$, $\underline{\sigma}' = \underline{\epsilon}' \underline{\epsilon}'$

$\Rightarrow C_{11} - C_{12} = C_{44}$

$$C_{ij} = \begin{bmatrix} C_{12} + C_{44} & C_{12} & C_{12} & 0 & 0 & 0 \\ & C_{12} + C_{44} & C_{12} & 0 & 0 & 0 \\ & & C_{12} + C_{44} & 0 & 0 & 0 \\ & & & C_{44} & 0 & 0 \\ & \text{sym.} & & & C_{44} & 0 \\ & & & & & C_{44} \end{bmatrix}$$

set $C_{12} \rightarrow \lambda$, $C_{44} \rightarrow 2\mu$

$$C_{ij} = \begin{bmatrix} \lambda + 2\mu & \lambda & \lambda & & & \\ \lambda & \lambda + 2\mu & \lambda & & & \\ \lambda & \lambda & \lambda + 2\mu & & & \\ & & & 2\mu & & \\ & & & & 2\mu & \\ & & & & & 2\mu \end{bmatrix}$$

$$\lambda = \frac{\nu E}{(1+\nu)(1-2\nu)} \quad \mu = G = \frac{E}{2(1+\nu)}$$

$\lambda, \mu = \text{Lame's constants}$

$$\sigma_{ij} = \lambda \epsilon_{kk} \delta_{ij} + 2\mu \epsilon_{ij}$$

If we define the strains as

$$\underline{\sigma} = \underline{E} \underline{\epsilon}$$

$$\underline{\epsilon} = \begin{bmatrix} \epsilon_{11} \\ \epsilon_{22} \\ \epsilon_{33} \\ \gamma_{12} \\ \gamma_{23} \\ \gamma_{31} \end{bmatrix}$$

$$\underline{E} = \begin{bmatrix} \lambda + 2\mu & \lambda & \lambda & & & \\ \lambda & \lambda + 2\mu & \lambda & & & \\ \lambda & \lambda & \lambda + 2\mu & & & \\ & & & 2\mu & & \\ & & & & 2\mu & \\ & & & & & 2\mu \end{bmatrix}$$

Strain - Displacement Relationship

$$\underline{u} = \begin{bmatrix} u \\ v \\ w \end{bmatrix} \quad \underline{\epsilon} = \begin{bmatrix} \epsilon_{11} \\ \epsilon_{22} \\ \epsilon_{33} \\ \epsilon_{12} \\ \epsilon_{23} \\ \epsilon_{31} \end{bmatrix} \quad \text{or} \quad \begin{bmatrix} \epsilon_{11} \\ \epsilon_{22} \\ \epsilon_{33} \\ \gamma_{12} \\ \gamma_{23} \\ \gamma_{31} \end{bmatrix}$$

$$\epsilon_{11} = \frac{\partial u}{\partial x}, \quad \epsilon_{22} = \frac{\partial v}{\partial y}, \quad \epsilon_{33} = \frac{\partial w}{\partial z}$$

$$\epsilon_{12} = \frac{1}{2} \left(\frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} \right) \quad \gamma_{12} = \frac{\partial u}{\partial y} + \frac{\partial v}{\partial x}$$

$$\epsilon_{23} = \frac{1}{2} \left(\frac{\partial v}{\partial z} + \frac{\partial w}{\partial y} \right) \quad \gamma_{23} = \frac{\partial v}{\partial z} + \frac{\partial w}{\partial y}$$

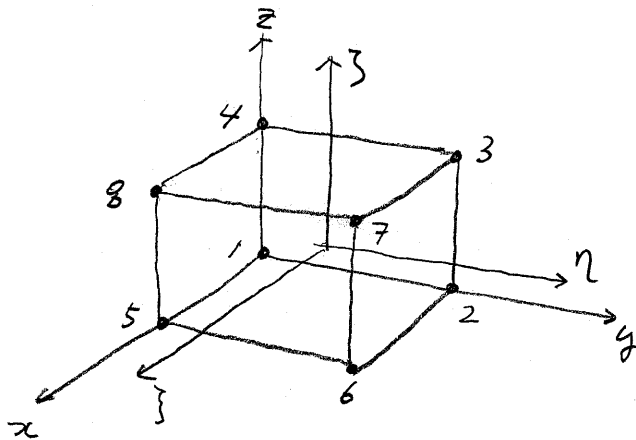
$$\epsilon_{31} = \frac{1}{2} \left(\frac{\partial w}{\partial x} + \frac{\partial u}{\partial z} \right) \quad \gamma_{31} = \frac{\partial w}{\partial x} + \frac{\partial u}{\partial z}$$

$$\begin{bmatrix} \epsilon_{11} \\ \epsilon_{22} \\ \epsilon_{33} \\ \gamma_{12} \\ \gamma_{23} \\ \gamma_{31} \end{bmatrix} = \begin{bmatrix} \frac{\partial}{\partial x} & 0 & 0 \\ 0 & \frac{\partial}{\partial y} & 0 \\ 0 & 0 & \frac{\partial}{\partial z} \\ \frac{\partial}{\partial y} & \frac{\partial}{\partial x} & 0 \\ 0 & \frac{\partial}{\partial z} & \frac{\partial}{\partial y} \\ \frac{\partial}{\partial z} & 0 & \frac{\partial}{\partial x} \end{bmatrix} \begin{bmatrix} u \\ v \\ w \end{bmatrix}$$

$$\underline{\epsilon} = \underline{D} \underline{u}$$

General Solids

8 - node Solid Element (Isoparametric Element)



$$\begin{cases}
 u = f_1 u_1 + f_2 u_2 + f_3 u_3 + f_4 u_4 + \dots + f_8 u_8 = \sum_i^8 f_i u_i \\
 v = \sum f_i v_i \\
 w = \sum f_i w_i \\
 x = \sum f_i x_i \\
 y = \sum f_i y_i \\
 z = \sum f_i z_i
 \end{cases}$$

node 1 (-1 -1 -1) $f_1 = \frac{1}{8} (1-\zeta)(1-\eta)(1-\xi)$

node 2 (-1 1 -1) $f_2 = \frac{1}{8} (1-\zeta)(1+\eta)(1-\xi)$

node 3 (-1 1 1) $f_3 = \frac{1}{8} (1-\zeta)(1+\eta)(1+\xi)$

⋮

node 8 (1 -1 1) $f_8 = \frac{1}{8} (1+\zeta)(1-\eta)(1+\xi)$

$$\Rightarrow f_i = \frac{1}{8} (1+\zeta_i)(1+\eta_i)(1+\xi_i)$$

$$\underline{u} = \begin{bmatrix} u \\ v \\ w \end{bmatrix} = \begin{bmatrix} f_1 & 0 & 0 & f_2 & 0 & 0 & \dots & f_8 & 0 & 0 \\ 0 & f_1 & 0 & 0 & f_2 & 0 & \dots & 0 & f_8 & 0 \\ 0 & 0 & f_1 & 0 & 0 & f_2 & \dots & 0 & 0 & f_8 \end{bmatrix} \begin{bmatrix} u_1 \\ v_1 \\ w_1 \\ \vdots \\ u_8 \\ v_8 \\ w_8 \end{bmatrix}$$

$$\underline{u} = \underline{f} \underline{q}$$

$$\underline{\Sigma} = \underline{B} \underline{\delta}$$

$$\underline{B} = \underline{d} \underline{f}$$

(6x24) (6x3)(3x24)

$$\underline{B} = \begin{bmatrix} \frac{\partial f_1}{\partial x} & 0 & 0 \\ 0 & \frac{\partial f_1}{\partial y} & 0 \\ 0 & 0 & \frac{\partial f_1}{\partial z} \\ \frac{\partial f_2}{\partial x} & \frac{\partial f_2}{\partial y} & 0 \\ 0 & \frac{\partial f_2}{\partial z} & \frac{\partial f_2}{\partial y} \\ \frac{\partial f_3}{\partial x} & 0 & \frac{\partial f_3}{\partial y} \end{bmatrix} \quad \begin{bmatrix} \frac{\partial g}{\partial x} & 0 & 0 \\ 0 & \frac{\partial g}{\partial y} & 0 \\ 0 & 0 & \frac{\partial g}{\partial z} \\ \frac{\partial g}{\partial x} & \frac{\partial g}{\partial z} & 0 \\ 0 & \frac{\partial g}{\partial z} & \frac{\partial g}{\partial y} \\ \frac{\partial g}{\partial z} & 0 & \frac{\partial g}{\partial x} \end{bmatrix}$$

$$\underline{J} = \begin{bmatrix} \frac{\partial x}{\partial \xi} & \frac{\partial y}{\partial \xi} & \frac{\partial z}{\partial \xi} \\ \frac{\partial x}{\partial \eta} & \frac{\partial y}{\partial \eta} & \frac{\partial z}{\partial \eta} \\ \frac{\partial x}{\partial \zeta} & \frac{\partial y}{\partial \zeta} & \frac{\partial z}{\partial \zeta} \end{bmatrix}$$

$$\frac{\partial x}{\partial \xi} = \sum_i \frac{\partial f_i}{\partial \xi} x_i$$

$$\frac{\partial x}{\partial \eta} = \sum_i \frac{\partial f_i}{\partial \eta} x_i$$

$$\frac{\partial x}{\partial \zeta} = \sum_i \frac{\partial f_i}{\partial \zeta} x_i$$

$$\frac{\partial y}{\partial \xi} = \sum_i \frac{\partial f_i}{\partial \xi} y_i$$

$$\frac{\partial y}{\partial \eta} = \sum_i \frac{\partial f_i}{\partial \eta} y_i$$

$$\frac{\partial y}{\partial \zeta} = \sum_i \frac{\partial f_i}{\partial \zeta} y_i$$

$$\frac{\partial z}{\partial \xi} = \sum_i \frac{\partial f_i}{\partial \xi} z_i$$

$$\frac{\partial z}{\partial \eta} = \sum_i \frac{\partial f_i}{\partial \eta} z_i$$

$$\frac{\partial z}{\partial \zeta} = \sum_i \frac{\partial f_i}{\partial \zeta} z_i$$

$$\underline{J}^x = \underline{J}^{-1} = \begin{bmatrix} \frac{\partial \xi}{\partial x} & \frac{\partial \eta}{\partial x} & \frac{\partial \zeta}{\partial x} \\ \frac{\partial \xi}{\partial y} & \frac{\partial \eta}{\partial y} & \frac{\partial \zeta}{\partial y} \\ \frac{\partial \xi}{\partial z} & \frac{\partial \eta}{\partial z} & \frac{\partial \zeta}{\partial z} \end{bmatrix}$$

$$\frac{\partial f_i}{\partial x} = \frac{\partial f_i}{\partial \xi} \frac{\partial \xi}{\partial x} + \frac{\partial f_i}{\partial \eta} \frac{\partial \eta}{\partial x} + \frac{\partial f_i}{\partial \zeta} \frac{\partial \zeta}{\partial x}$$

$$\frac{\partial f_i}{\partial y} = \frac{\partial f_i}{\partial \xi} \frac{\partial \xi}{\partial y} + \frac{\partial f_i}{\partial \eta} \frac{\partial \eta}{\partial y} + \frac{\partial f_i}{\partial \zeta} \frac{\partial \zeta}{\partial y}$$

$$\frac{\partial f_i}{\partial z} = \frac{\partial f_i}{\partial \xi} \frac{\partial \xi}{\partial z} + \frac{\partial f_i}{\partial \eta} \frac{\partial \eta}{\partial z} + \frac{\partial f_i}{\partial \zeta} \frac{\partial \zeta}{\partial z}$$

$$\underline{K} = \int \underline{B}^T \underline{E} \underline{B} dV$$

$$= \iiint_V \underline{B}^T \underline{E} \underline{B} |\underline{J}| d\xi d\eta d\zeta$$

\underline{f}

Using Gaussian Quadrature

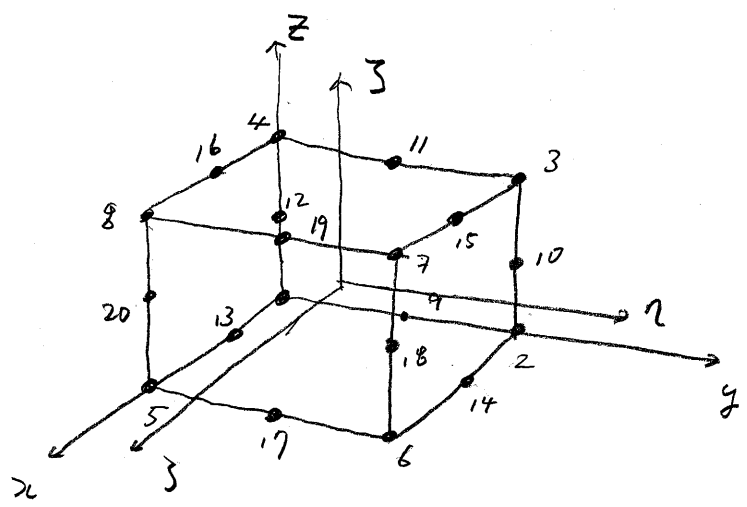
$$\underline{K} = \sum_i^l \sum_j^m \sum_k^n R_i R_j R_k \underline{f}(\xi_i, \eta_j, \zeta_k)$$

Equivalent nodal forces

$$\underline{P}_b = \iiint \underline{f}^T \underline{b} |\underline{J}| d\xi d\eta d\zeta$$

$$\underline{P}_e = \iiint \underline{B}^T \underline{E} \underline{\epsilon}_e |\underline{J}| d\xi d\eta d\zeta$$

20 nodes Solid Element (Isoparametric Element)



see p160 and p161

$$\begin{cases} u = \sum f_i u_i \\ v = \sum f_i v_i \\ w = \sum f_i w_i \end{cases} \quad \begin{cases} x = \sum f_i x_i \\ y = \sum f_i y_i \\ z = \sum f_i z_i \end{cases}$$

$$f_i = \frac{1}{8} (1 + \zeta_i \zeta) (1 + \eta_i \eta) (1 + \xi_i \xi) (\zeta_i \zeta + \eta_i \eta + \xi_i \xi - 2) \quad (i = 1, 2, \dots, 8)$$

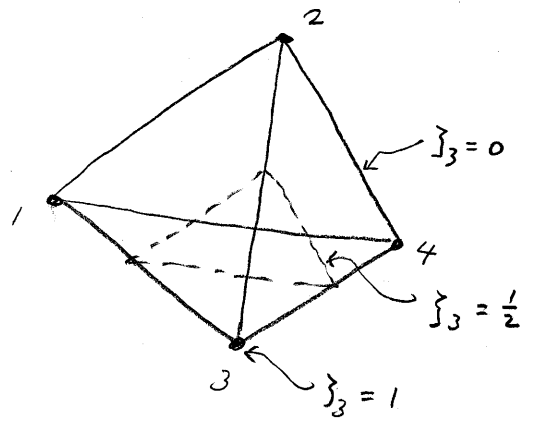
$$f_i = \frac{1}{4} (1 - \zeta^2) (1 + \eta_i \eta) (1 + \xi_i \xi) \quad (i = 9, 11, 17, 19)$$

$$f_i = \frac{1}{4} (1 - \eta^2) (1 + \zeta_i \zeta) (1 + \xi_i \xi) \quad (i = 10, 12, 18, 20)$$

$$f_i = \frac{1}{4} (1 - \xi^2) (1 + \zeta_i \zeta) (1 + \eta_i \eta) \quad (i = 13, 14, 15, 16)$$

4-node Tetrahedron Element

Volume coordinates



$$\xi_1 = \frac{V_1}{V}$$

$$\xi_2 = \frac{V_2}{V}$$

$$\xi_3 = \frac{V_3}{V}$$

$$\xi_4 = \frac{V_4}{V}$$

$$V_1 + V_2 + V_3 + V_4 = V$$

$$\xi_1 + \xi_2 + \xi_3 + \xi_4 = 1$$

$$\begin{cases} x = \sum f_i x_i = \sum \xi_i x_i & f_i = \xi_i \\ y = \sum \xi_i y_i \\ z = \sum \xi_i z_i \end{cases}$$

$$\begin{bmatrix} 1 \\ x \\ y \\ z \end{bmatrix} = \begin{bmatrix} 1 & 1 & 1 & 1 \\ x_1 & x_2 & x_3 & x_4 \\ y_1 & y_2 & y_3 & y_4 \\ z_1 & z_2 & z_3 & z_4 \end{bmatrix} \begin{bmatrix} \xi_1 \\ \xi_2 \\ \xi_3 \\ \xi_4 \end{bmatrix}$$

$$\underline{X} = \underline{H} \underline{\xi}$$

$$\underline{\zeta} = H^T \underline{x}$$

$$\begin{bmatrix} \zeta_1 \\ \zeta_2 \\ \zeta_3 \\ \zeta_4 \end{bmatrix} = \frac{1}{|H|} \begin{bmatrix} H_{11}^* & H_{12}^* & H_{13}^* & H_{14}^* \\ H_{21}^* & H_{22}^* & H_{23}^* & H_{24}^* \\ H_{31}^* & H_{32}^* & H_{33}^* & H_{34}^* \\ H_{41}^* & H_{42}^* & H_{43}^* & H_{44}^* \end{bmatrix} \begin{bmatrix} 1 \\ x \\ y \\ z \end{bmatrix}$$

$$\begin{cases} u = \sum f_i u_i = \sum \zeta_i u_i \\ v = \sum \zeta_i v_i \\ w = \sum \zeta_i w_i \end{cases}$$

$$\underline{\Sigma} = \underline{B} \underline{\delta}$$

$$\underline{B} = \begin{bmatrix} f_{1,x} & 0 & 0 & \dots & f_{4,x} & 0 & 0 \\ 0 & f_{1,y} & 0 & \dots & 0 & f_{4,y} & 0 \\ 0 & 0 & f_{1,z} & \dots & 0 & 0 & f_{4,z} \\ f_{1,y} & f_{1,x} & 0 & \dots & f_{4,y} & f_{4,x} & 0 \\ 0 & f_{1,z} & f_{1,y} & \dots & 0 & f_{4,z} & f_{4,y} \\ f_{1,z} & 0 & f_{1,x} & \dots & f_{4,z} & 0 & f_{4,x} \end{bmatrix}$$

$$\begin{aligned} \frac{\partial f_i}{\partial x} &= \frac{\partial f_i}{\partial \zeta_1} \frac{\partial \zeta_1}{\partial x} + \frac{\partial f_i}{\partial \zeta_2} \frac{\partial \zeta_2}{\partial x} + \frac{\partial f_i}{\partial \zeta_3} \frac{\partial \zeta_3}{\partial x} + \frac{\partial f_i}{\partial \zeta_4} \frac{\partial \zeta_4}{\partial x} \\ &= \frac{\partial f_i}{\partial \zeta_1} \frac{1}{|H|} H_{12}^* + \frac{\partial f_i}{\partial \zeta_2} \frac{1}{|H|} H_{22}^* + \frac{\partial f_i}{\partial \zeta_3} \frac{1}{|H|} H_{32}^* + \frac{\partial f_i}{\partial \zeta_4} \frac{1}{|H|} H_{42}^* \end{aligned}$$

for 4node tetrahedron element. $f_i \approx \zeta_i$

$$\frac{\partial f_i}{\partial x} = \frac{\partial \zeta_i}{\partial \zeta_i} \frac{1}{|H|} H_{i2}^* = \frac{1}{|H|} H_{i2}^*$$

$$\frac{\partial f_i}{\partial y} = \frac{1}{H_1} H_{i3}^* \quad , \quad \frac{\partial f_i}{\partial z} = \frac{1}{H_1} H_{i4}^*$$

$$K = \int B^T E B dV = \underline{B}^T \underline{E} \underline{B} V \quad (\text{constant strain})$$

Generally

$$K = \int_V \left\{ \begin{matrix} a \\ b \\ c \\ d \end{matrix} \right\}_1 \left\{ \begin{matrix} a \\ b \\ c \\ d \end{matrix} \right\}_2 \left\{ \begin{matrix} a \\ b \\ c \\ d \end{matrix} \right\}_3 \left\{ \begin{matrix} a \\ b \\ c \\ d \end{matrix} \right\}_4 dV = \frac{a!b!c!d!}{(a+b+c+d+3)} (6V)$$

or

$$K = \frac{1}{6} \sum_{j=1}^n w_j f(\xi_1, \xi_2, \xi_3, \xi_4) \left| \underline{J}(\xi_1, \xi_2, \xi_3) \right|$$

p156 Table 4.1

$$\underline{J} = \begin{bmatrix} x_{,1} & y_{,1} & z_{,1} \\ x_{,2} & y_{,2} & z_{,2} \\ x_{,3} & y_{,3} & z_{,3} \end{bmatrix}$$

$$x_{,1} = \frac{\partial x}{\partial \xi_1} = \sum \frac{\partial f_i}{\partial \xi_1} x_i$$

$$x = \xi_1 x_1 + \xi_2 x_2 + \xi_3 x_3 + \xi_4 x_4$$

$$= \xi_1 x_1 + \xi_2 x_2 + \xi_3 x_3 + (1 - \xi_1 - \xi_2 - \xi_3) x_4$$

$$x_{,1} = x_1 - x_4 \Rightarrow -\underline{J} \Rightarrow \text{constants}$$

$$|\underline{J}| = 6V$$

Tetrahedron 10 node Element

→ see p. 164

