

Topics in Ship Structural Design (Hull Buckling and Ultimate Strength)

Lecture 3 Plate Bending

Reference : Ship Structural Design Ch.09

NAOE

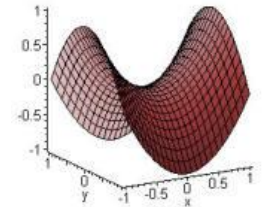
Jang, Beom Seon



"Long" Plates(Cylindrical Bending)

- An equation relating the deflections to the loading can be developed as for the beam.
- Cylindrical Bending:
 - ✓ A plate which is bent about one axis only($a \gg b$)
 - ✓ Transverse deformation does not occur, since such a deformation would required a saddle shape deformation.

Saddle shape



$$\varepsilon_x = \frac{\sigma_x}{E} - \nu \frac{\sigma_y}{E} \quad \leftarrow \quad \Delta \varepsilon_x = \nu \varepsilon_y = \nu \frac{\sigma_y}{E}$$

$$\varepsilon_y = \frac{\sigma_y}{E} - \nu \frac{\sigma_x}{E} = 0 \quad \Rightarrow \quad \sigma_y = \nu \sigma_x$$

$$\varepsilon_x = \frac{\sigma_x (1 - \nu^2)}{E}$$

$$\sigma_x = \frac{E}{1 - \nu^2} \varepsilon_x = E' \varepsilon_x$$

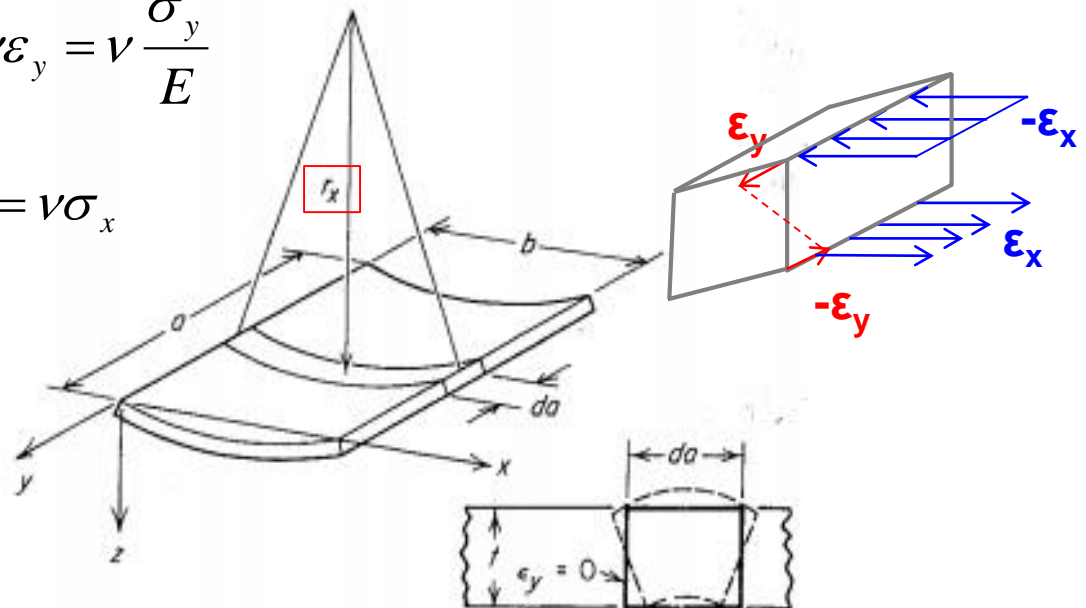
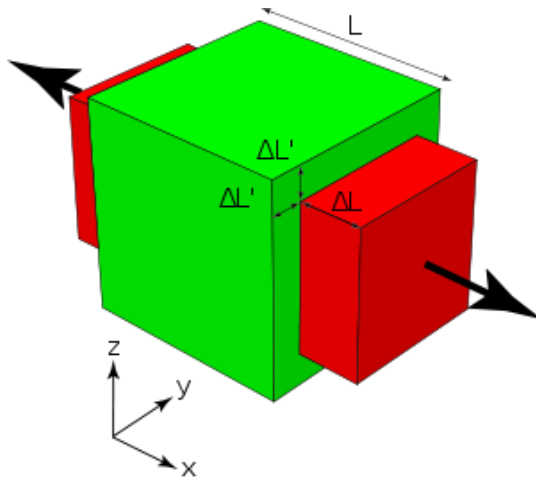


Plate in cylindrical bending

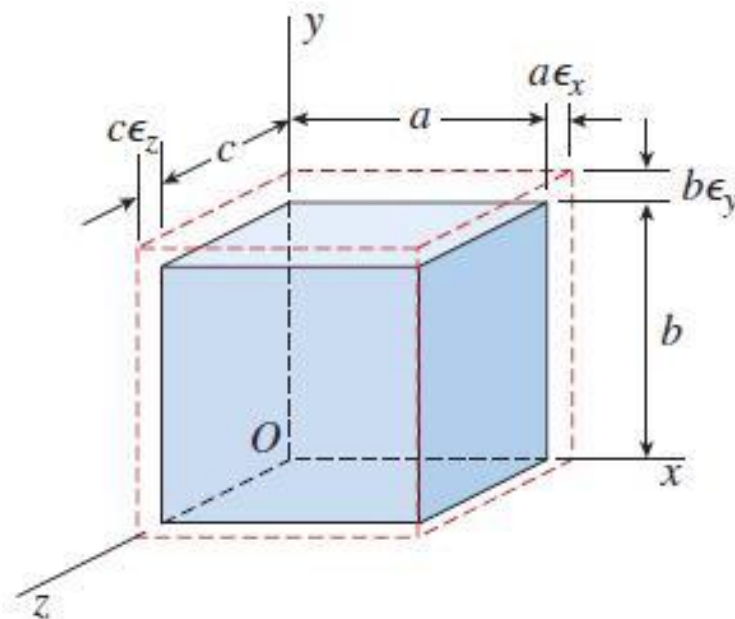
Poisson Ratio



$$\nu = - \frac{\epsilon_{trans}}{\epsilon_{axial}}$$

Strain $\Delta\epsilon_y$ induced by ϵ_y

$$\Delta\epsilon_x = \nu\epsilon_y = \nu \frac{\sigma_y}{E}$$



$$\epsilon_x = \frac{1}{E} (\sigma_x - \nu\sigma_y - \nu\sigma_z)$$

$$\epsilon_y = \frac{1}{E} (\sigma_y - \nu\sigma_z - \nu\sigma_x)$$

$$\epsilon_z = \frac{1}{E} (\sigma_z - \nu\sigma_x - \nu\sigma_y)$$

Formulas between a beam and a long plate

- E' : effective modulus of elasticity, always $> E$.

$$E' = \frac{E}{(1-\nu^2)}$$

→ A plate is always “stiffer” than a row of beams

- For a long prismatically loaded plate, **the extra stiffness** may be fully taken into account for by using E' in place of E .
- For simply supported plates:

$$\omega_{\max} = \frac{5pb^4}{384E'I} = \frac{5pb^4(1-\nu^2)}{32Et^3}$$

- For clamped plates:

$$\omega_{\max} = \frac{pb^4}{384E'I} = \frac{pb^4(1-\nu^2)}{32Et^3}$$

Formulas between a beam and a long plate

- Moment curvature relation may be obtained from
External bending moment = moment of stress force

$$M = \int_{-t/2}^{t/2} \sigma_x z dz = \int_{-t/2}^{t/2} \frac{E}{1-\nu^2} \left(\frac{z^2}{r_x} \right) dz = \frac{Et^3}{12(1-\nu^2)} \left(\frac{1}{r_x} \right)$$

where $\sigma_x = \frac{E}{(1-\nu^2)} \varepsilon_x$ $\varepsilon_x = z/r_x$

$$M = \frac{Et^3}{12(1-\nu^2)} \left(\frac{1}{r_x} \right) = \frac{D}{r_x} \quad D = \frac{Et^3}{12(1-\nu^2)}$$

- D : the flexural rigidity of the plate
the constant of proportionality between moment and curvature
analogous with EI in beam theory

$$M = \frac{EI}{\rho}$$

Formulas between a beam and a long plate

- The radius of curvature of the plate can be expressed in terms of the deflection w of the plate

$$-\frac{\partial^2 w}{\partial x^2} = \frac{1}{r_x} \quad -D \frac{\partial^2 w}{\partial x^2} = M \quad \boxed{-v'' = \frac{M}{EI}}$$

- Also for a unit strip in a long plate the max. stress is $\sigma_{\max} = M_{\max} c / I$ where the section modulus

$$I / c = (t^3 / 12) / (t / 2) = t^2 / 6$$

- For a uniform pressure p , the maximum bending moment is proportional to pb^2 . The stress is expressed in

$$\sigma_{\max} = kp \left(\frac{b}{t} \right)^2$$

- The coefficient k depends on the boundary conditions:

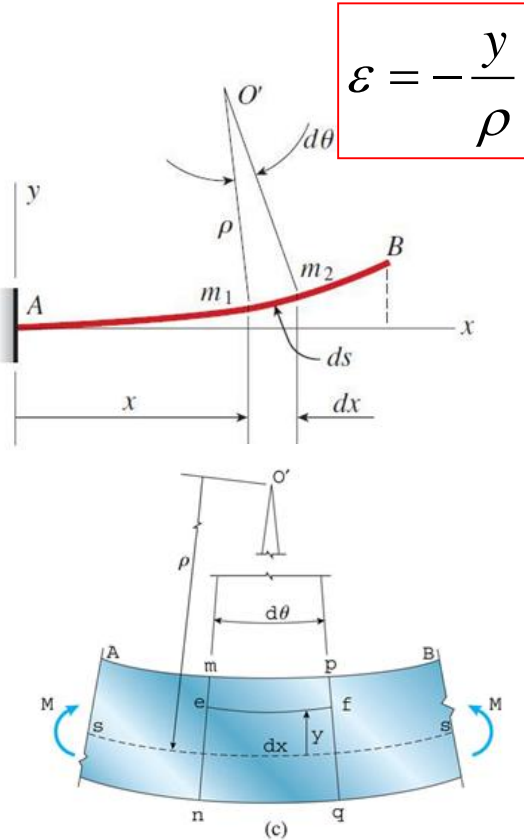
$k=3/4$ for simply supported edges

$k=1/2$ for clamped edges

Beam Theory

stress – strain - curvature - Moment

1) strain-curvature
Geometric relation



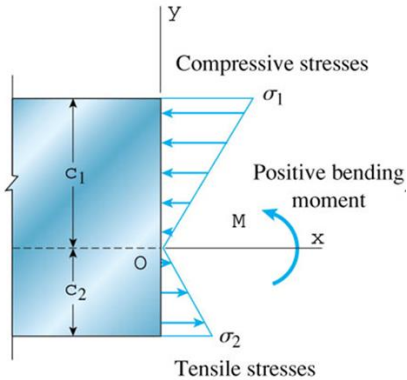
2) strain-stress relation

$\sigma = E\epsilon = -\frac{Ey}{\rho}$

$\frac{1}{\rho} = -\frac{\sigma}{Ey}$

3) moment-curvature relation

$M = -\int_A \sigma_x y dA$



$\therefore K = \frac{1}{\rho} = \frac{M}{EI}$

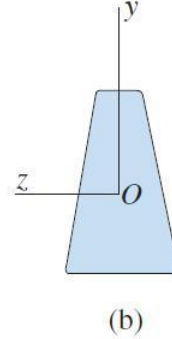
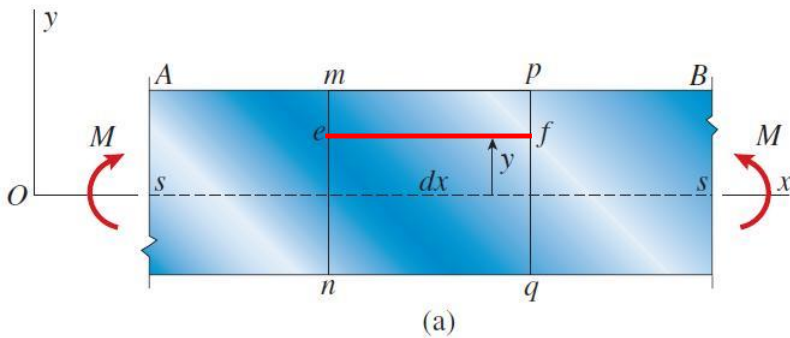
4) Stress-moment relation

$\therefore \sigma_x = -\frac{My}{I}$



Beam theory

1) Strain-Curvature Relation

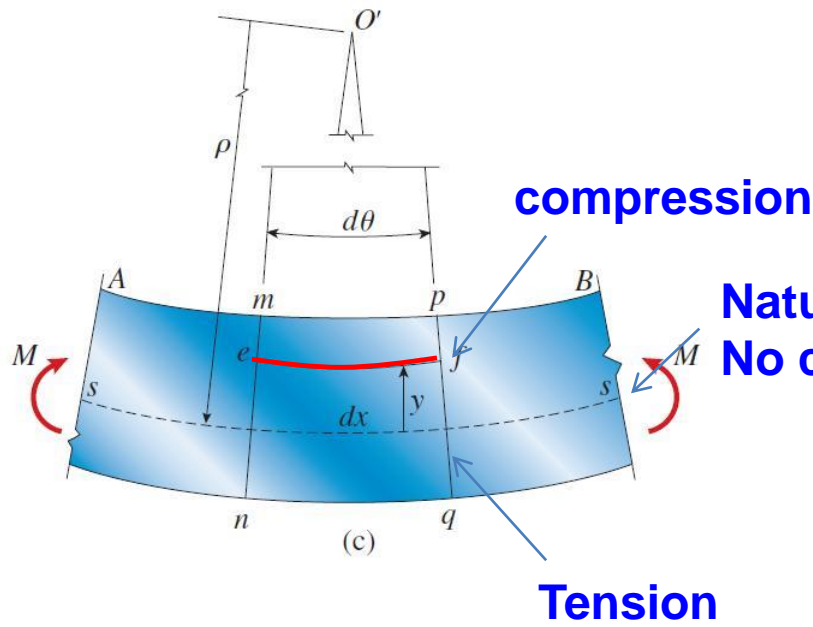


Initial length of Line **e-f** = dx

Length after bending = L_1

$$L_1 = (\rho - y)d\theta = dx - \frac{y}{\rho}dx$$

$$d\theta = \frac{dx}{\rho}$$



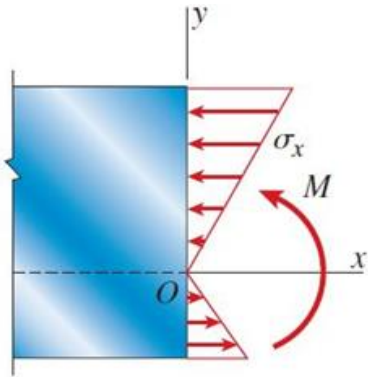
Natural axis

No change in length

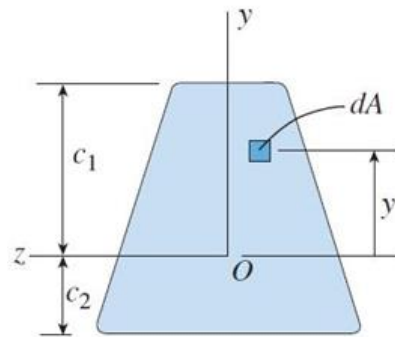
$$\therefore \epsilon_x = \frac{L_1 - dx}{dx} = -\frac{y}{\rho} = -\kappa y$$

Beam theory

3) moment-curvature relation



(a)



(b)

$$dM = -\sigma_x y dA$$

$$M = -\int_A \sigma_x y dA$$

$$M = \int_A \kappa E y^2 dA = \kappa E \int_A y^2 dA = \kappa E I = \frac{1}{\rho} E I \quad \leftarrow \quad \boxed{\sigma = -\kappa E y}$$

$$\boxed{\therefore \kappa = \frac{1}{\rho} = \frac{M}{EI}}$$

: EI = bending rigidity

$$I \equiv \int_A y^2 dA$$

: Moment of Inertia

Beam Theory : Relation between load, shear force, Moment

Force equilibrium

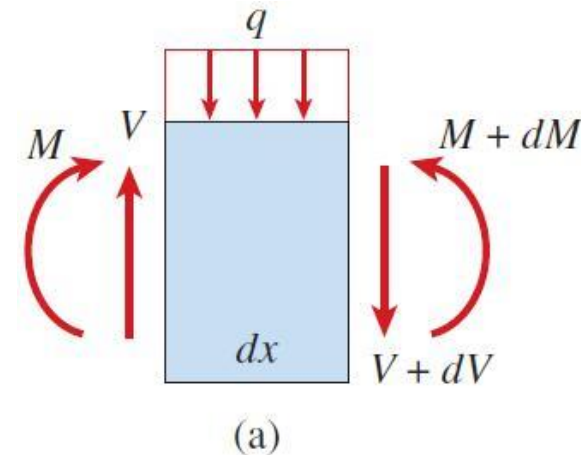
→ Relation between shear force (V) and load (q)

$$\sum F_{vert} = 0 : V - qdx - (V + dV) = 0 \Rightarrow \frac{dV}{dx} = -q$$

Moment equilibrium

→ Relation between Moment (M) and shear force (V)

$$\sum M = 0 : -M - qdx\left(\frac{dx}{2}\right) - (V + dV)dx + (M + dM) = 0 \Rightarrow \frac{dM}{dx} = V$$



Combining two relations

$$\frac{d^2 M}{dx^2} = -q$$

Relation between deflection and distributed load

$$\therefore \frac{d^2 v}{dx^2} = \frac{M}{EI}$$

$$\frac{d^4 v}{dx^4} = -\frac{q}{EI}$$

9.1 Small Deflection Theory

Forces and moments in a plate element

Force equilibrium

→ Relation between shear force (V) and load (q)

$$\frac{\partial q_x}{\partial x} + \frac{\partial q_y}{\partial y} + p = 0 \quad \boxed{\frac{dV}{dx} = -q}$$

Moment equilibrium

$$\frac{\partial m_{xy}}{\partial x} - \frac{\partial m_y}{\partial y} + q_y = 0 \quad \frac{\partial m_{yx}}{\partial y} + \frac{\partial m_x}{\partial x} - q_x = 0 \quad \boxed{\frac{dM}{dx} = V}$$

Combining two relations

$$\frac{\partial^2 m_x}{\partial x^2} - 2 \frac{\partial^2 m_{xy}}{\partial x \partial y} + \frac{\partial^2 m_y}{\partial y^2} + p = 0 \quad \boxed{\frac{d^2 M}{dx^2} = -q}$$

Relation between deflection and distributed load

$$m_x = -D \left(\frac{\partial^2 w}{\partial x^2} + \nu \frac{\partial^2 w}{\partial y^2} \right) \quad m_y = -D \left(\frac{\partial^2 w}{\partial y^2} + \nu \frac{\partial^2 w}{\partial x^2} \right)$$

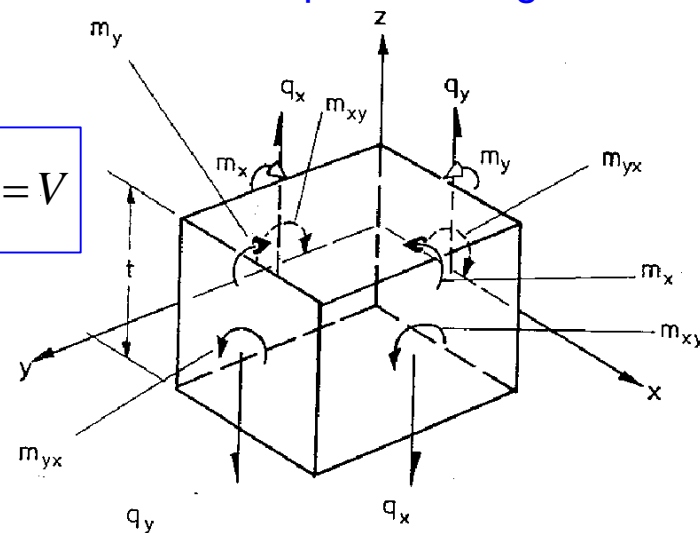
$$m_{xy} = D(1-\nu) \frac{\partial^2 w}{\partial x \partial y}$$



$$\boxed{\frac{\partial^4 w}{\partial x^4} + 2 \frac{\partial^4 w}{\partial x^2 \partial y^2} + \frac{\partial^4 w}{\partial y^4} = \frac{p}{D}}$$

$$\boxed{\frac{d^4 v}{dx^4} = -\frac{q}{EI}}$$

q : shear force per unit length
m : moment per unit length



Derivation of the Plate Bending Equation

- ❖ The previous theory is only applicable if :
 - Plane cross sections remain plane.
 - The deflections of the plate are small (w_{max} not exceeding $3/4 t$)
 - The maximum stress nowhere exceeds the plate yield stress (i.e., the material remains elastic)
- ❖ A panel of plating will have curvature in two directions

$$\varepsilon_x = \frac{\sigma_x}{E} - \frac{\nu \sigma_y}{E} \quad \varepsilon_y = \frac{\sigma_y}{E} - \frac{\nu \sigma_x}{E}$$

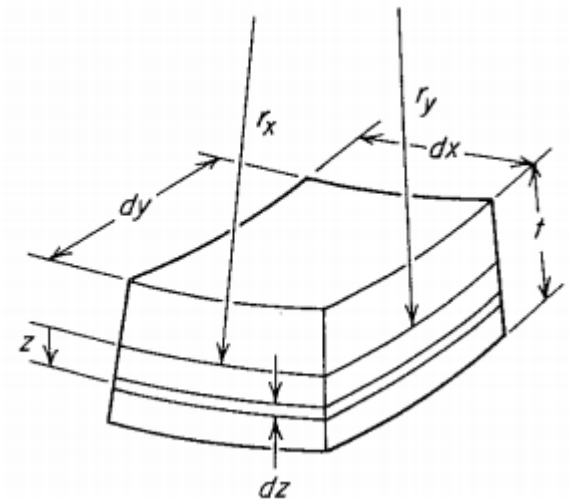
+



$$\nu \varepsilon_y = \frac{\nu \sigma_y}{E} - \frac{\nu^2 \sigma_x}{E}$$

||

$$\frac{\sigma_x}{E} (1 - \nu^2) = \varepsilon_x + \nu \varepsilon_y$$



Differential element of plating

Derivation of the Plate Bending Equation

- Assumptions stated previously give rise to the strain-curvature relations

$$\varepsilon_x = (-z) \left(\frac{\partial^2 w}{\partial x^2} \right) \quad \varepsilon_y = (-z) \left(\frac{\partial^2 w}{\partial y^2} \right) \quad \Rightarrow \quad \frac{\sigma_x}{E} (1 - \nu^2) = \varepsilon_x + \nu \varepsilon_y$$

$$\sigma_x = \frac{E}{1 - \nu^2} (-z) \left(\frac{\partial^2 w}{\partial x^2} + \nu \frac{\partial^2 w}{\partial y^2} \right)$$

$$m_x = \int_{-t/2}^{t/2} \sigma_x z dz = - \int_{-t/2}^{t/2} \frac{E}{1 - \nu^2} \left(\frac{\partial^2 w}{\partial x^2} + \nu \frac{\partial^2 w}{\partial y^2} \right) z^2 dz$$

$$m_x = -D \left(\frac{\partial^2 w}{\partial x^2} + \nu \frac{\partial^2 w}{\partial y^2} \right)$$

$$D = \frac{Et^3}{12(1 - \nu^2)}$$

$$m_y = -D \left(\frac{\partial^2 w}{\partial y^2} + \nu \frac{\partial^2 w}{\partial x^2} \right)$$

A diagram showing a curved beam element in a Cartesian coordinate system. The horizontal axis is labeled x . A red curved beam is shown, with a small segment of length ds highlighted. The beam is at an angle θ to the horizontal at its left end, labeled m_1 . At its right end, labeled m_2 , the angle is $\theta + d\theta$. The vertical distance from the x -axis to the beam's centerline is v at m_1 and $v + dv$ at m_2 . The horizontal distance between the vertical lines through m_1 and m_2 is dx . A line perpendicular to the beam's axis at m_2 is shown, and the angle between this perpendicular and the horizontal is $d\theta$.

$$\rho d\theta = ds$$

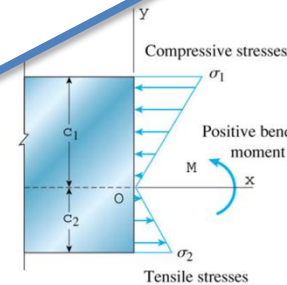
$$\kappa = \frac{1}{\rho} = \frac{d\theta}{ds}$$

$$\frac{dv}{dx} = \tan \theta$$

$$\kappa = \frac{1}{\rho} = \frac{d\theta}{dx}$$

$$\tan \theta = \frac{dv}{dx} \approx \theta$$

$$\therefore \kappa = \frac{1}{\rho} = \frac{d^2 v}{dx^2}$$



$$\underset{\nearrow}{K} = \frac{1}{\rho} = \frac{M}{EI}$$

$$\therefore \frac{d^2 v}{dx^2} = \frac{M}{EI}$$

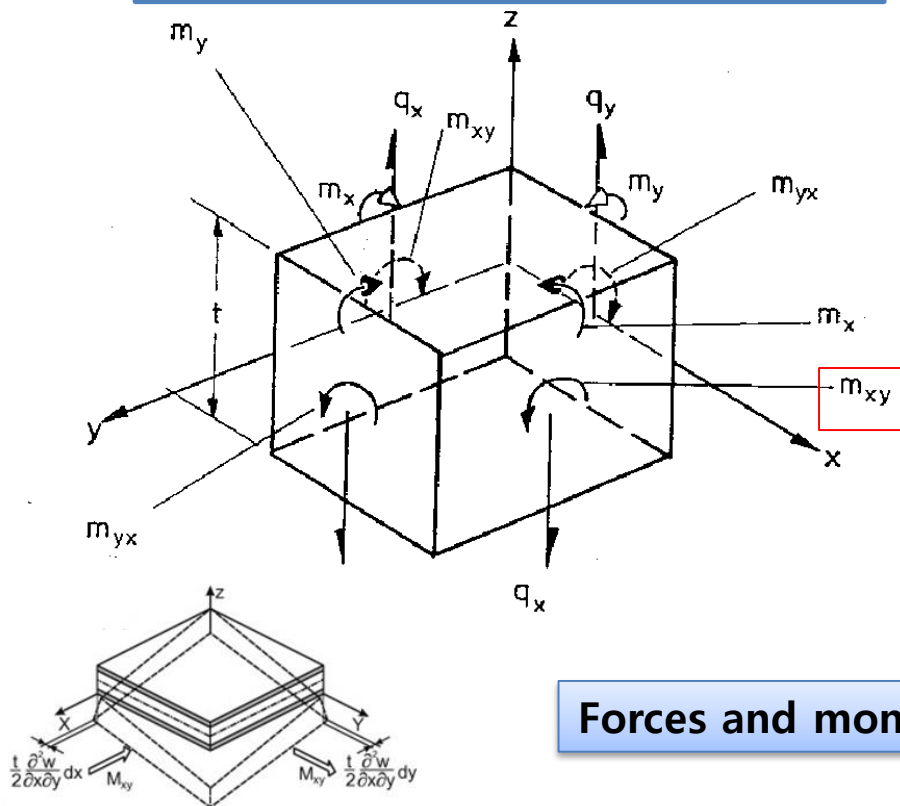
$$\frac{d^4 v}{dx^4} = -\frac{q}{EI}$$

$$\frac{d^2 M}{dx^2} = -q$$

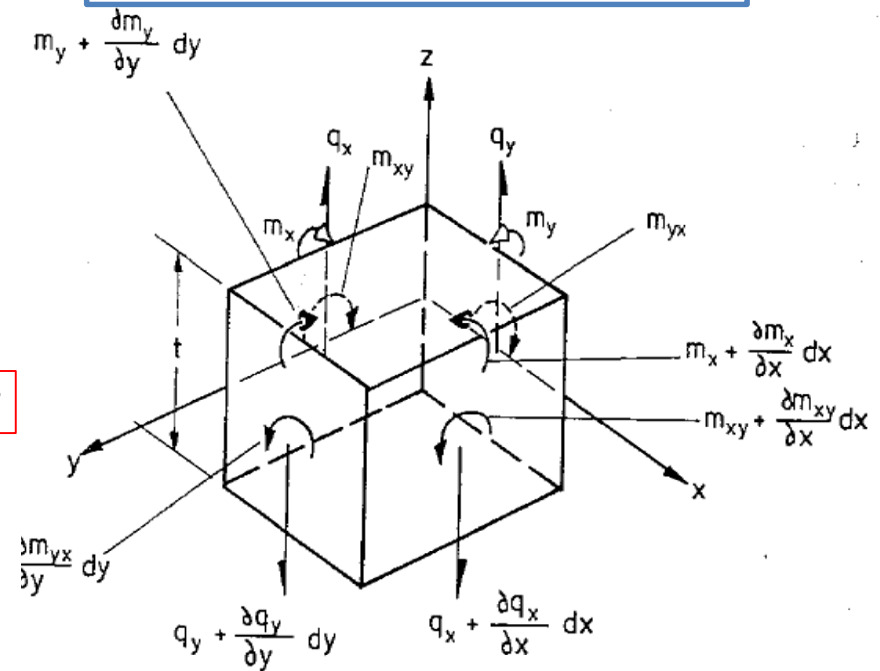
Forces and moments in a plate element

- The lateral load $pdx dy$ is carried by the distributed shear forces q acting on the four edges of the element.
- In the general case of bending, **twisting moments** will also be generated on all the four faces

Without the lateral load $pdx dy$



With the lateral load $pdx dy$

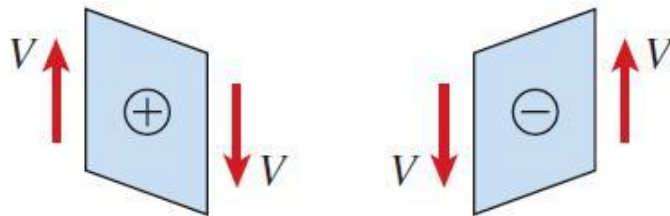


Forces and moments in a plate element

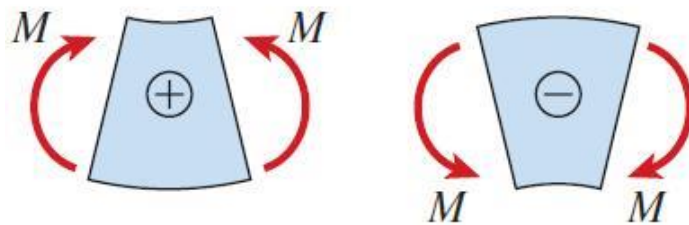
9.1 Small Deflection Theory

Forces and moments in a plate element

- Sine convention in beam theory
 - Determined how deform the material
 - Shear force : rotate an element clockwise (+)
 - Bending moment : shorten the upper skin (+)

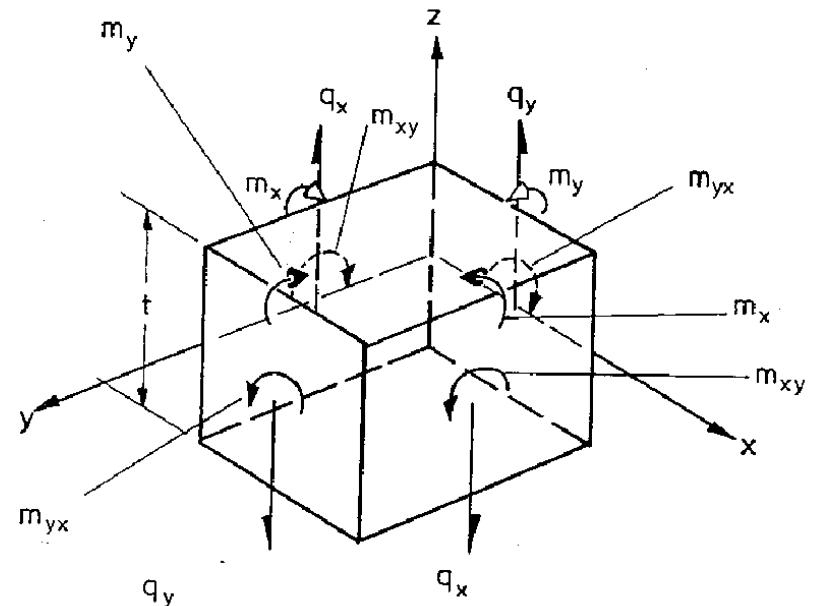


(a)



(b)

Sine convention in beam theory



Sine convention in plate bending

9.1 Small Deflection Theory

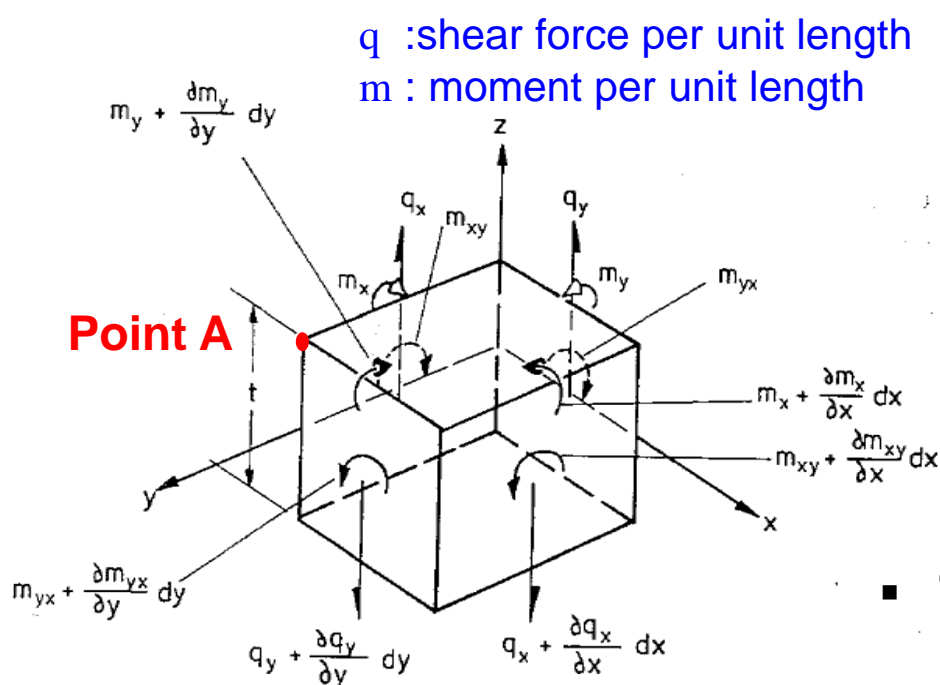
Forces and moments in a plate element

- From equilibrium of vertical forces :

$$\begin{aligned} & (\cancel{q_x} + \frac{\partial q_x}{\partial x} dx) dy - \cancel{q_x} dy \\ & + (\cancel{q_y} + \frac{\partial q_y}{\partial y} dy) dx - \cancel{q_y} dx + p dx dy = 0 \end{aligned} \quad \Rightarrow \quad \frac{\partial q_x}{\partial x} + \frac{\partial q_y}{\partial y} + p = 0$$

- Taking moments parallel to the x-axis, at Point A

q : shear force per unit length
 m : moment per unit length



$$\begin{aligned} & (m_{xy} + \frac{\partial m_{xy}}{\partial x} dx) dy - m_{xy} dy + m_y dx \\ & - (m_y + \frac{\partial m_y}{\partial y} dy) dx + q_y dx dy = 0 \end{aligned}$$

$$\frac{\partial m_{xy}}{\partial x} - \frac{\partial m_y}{\partial y} + q_y = 0$$

- Taking moments parallel to the y-axis

$$\frac{\partial m_{yx}}{\partial y} + \frac{\partial m_x}{\partial x} - q_x = 0$$

Forces and moments in a plate element

- Because of the principle of complementary shear stress, it follows that $m_{yx} = -m_{xy}$, so that

$$\frac{\partial m_{yx}}{\partial y} + \frac{\partial m_x}{\partial x} - q_x = 0 \quad \Rightarrow \quad \boxed{-\frac{\partial m_{xy}}{\partial y} + \frac{\partial m_x}{\partial x} - q_x = 0} \quad \boxed{\frac{\partial m_{xy}}{\partial x} - \frac{\partial m_y}{\partial y} + q_y = 0}$$

- Substituting for q_x and q_y

$$\boxed{\frac{\partial q_x}{\partial x} + \frac{\partial q_y}{\partial y} + p = 0}$$

$$\frac{\partial}{\partial x} \left(-\frac{\partial m_{xy}}{\partial y} + \frac{\partial m_x}{\partial x} \right) + \frac{\partial}{\partial y} \left(\frac{\partial m_y}{\partial y} - \frac{\partial m_{xy}}{\partial x} \right) + p = 0$$

$$\frac{\partial^2 m_x}{\partial x^2} - 2 \frac{\partial^2 m_{xy}}{\partial x \partial y} + \frac{\partial^2 m_y}{\partial y^2} + p = 0$$

Forces and moments in a plate element

- If a point A in the plate a distance z from the neutral surface is displaced a distance v in the y direction,
- The change in slope of the line AB will be

$$\frac{v + \frac{\partial v}{\partial x} dx - v}{dx} = \frac{\partial v}{\partial x}$$

- The change in slope of the line AD

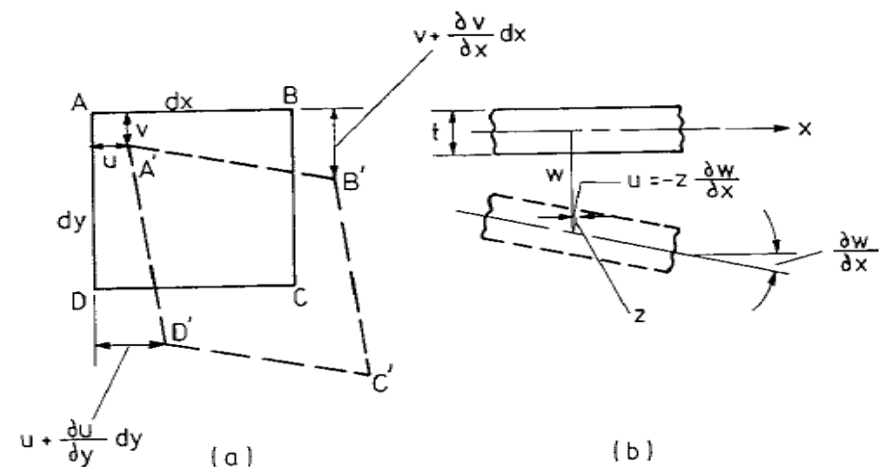
$$\frac{\partial u}{\partial x}$$

- The shear strain is

$$\gamma = \frac{\partial v}{\partial x} + \frac{\partial u}{\partial y}$$

- The shear stress

$$\tau = G \left(\frac{\partial v}{\partial x} + \frac{\partial u}{\partial y} \right)$$



Shear strain and plate deflection

Shear strain is induced by deflection. That is, **the vertical deflection** generates shear strain.

9.1 Small Deflection Theory

Forces and moments in a plate element

- from the geometry of the slope of w in each direction (x and y):

$$u = -z \frac{\partial w}{\partial x} \quad v = -z \frac{\partial w}{\partial y}$$

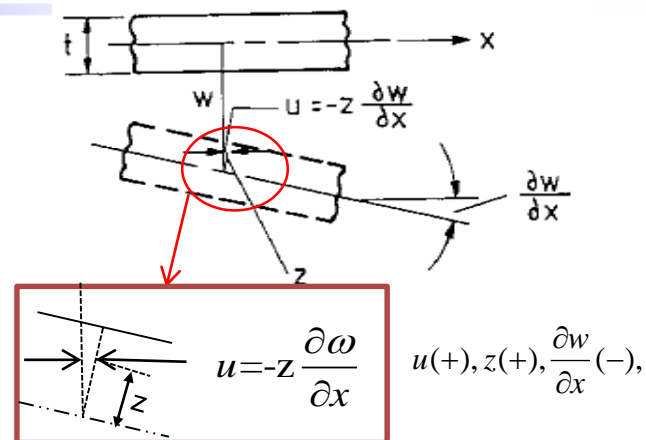
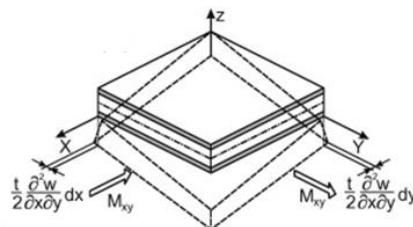
- Making these substitutions gives

$$\tau = -2Gz \frac{\partial^2 w}{\partial x \partial y}$$

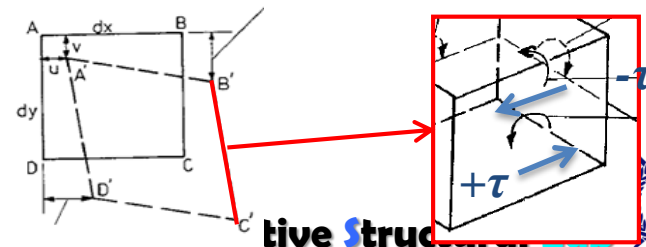
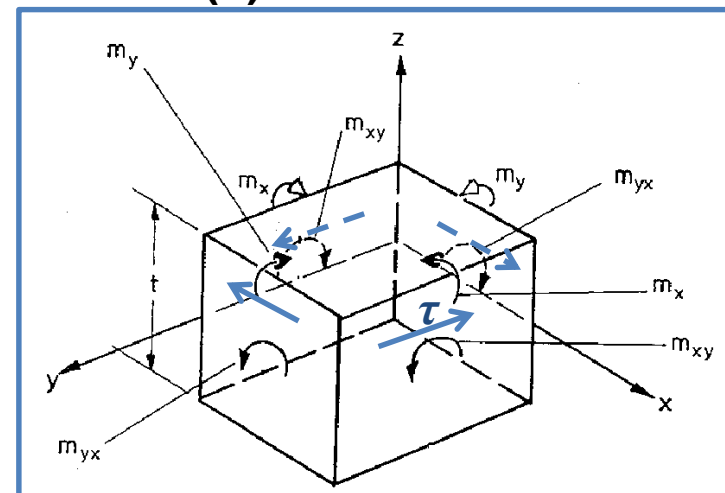
- The twisting moment

$$m_{xy} = - \int_{-t/2}^{t/2} -2G \frac{\partial^2 w}{\partial x \partial y} z^2 dz = \frac{Gt^3}{6} \frac{\partial^2 w}{\partial x \partial y} \quad G = \frac{E}{2(1+\nu)}$$

$$m_{xy} = \frac{Et^3}{12(1+\nu)} \frac{\partial^2 w}{\partial x \partial y} = \frac{Et^3(1-\nu)}{12(1+\nu)(1-\nu)} \frac{\partial^2 w}{\partial x \partial y} = D(1-\nu) \frac{\partial^2 w}{\partial x \partial y}$$



(+) convention



Forces and moments in a plate element

- Substitute m_x , m_y , m_{xy}

$$m_x = -D \left(\frac{\partial^2 w}{\partial x^2} + \nu \frac{\partial^2 w}{\partial y^2} \right)$$

$$m_y = -D \left(\frac{\partial^2 w}{\partial y^2} + \nu \frac{\partial^2 w}{\partial x^2} \right)$$

$$m_{xy} = D(1-\nu) \frac{\partial^2 w}{\partial x \partial y}$$



$$\frac{\partial^2 m_x}{\partial x^2} - 2 \frac{\partial^2 m_{xy}}{\partial x \partial y} + \frac{\partial^2 m_y}{\partial y^2} + p = 0$$

- Finally

$$\frac{\partial^2}{\partial x^2} \left[-D \left(\frac{\partial^2 w}{\partial x^2} + \nu \frac{\partial^2 w}{\partial y^2} \right) \right] - 2 \frac{\partial^2}{\partial x \partial y} \left[D(1-\nu) \frac{\partial^2 w}{\partial x \partial y} \right] + \frac{\partial^2}{\partial y^2} \left[-D \left(\frac{\partial^2 w}{\partial y^2} + \nu \frac{\partial^2 w}{\partial x^2} \right) \right] = -p$$

- Equation of equilibrium for the plate

$$\Rightarrow \frac{\partial^4 w}{\partial x^4} + 2 \frac{\partial^4 w}{\partial x^2 \partial y^2} + \frac{\partial^4 w}{\partial y^4} = \frac{p}{D} \quad \boxed{\nabla^4 w = p / D}$$

- Assumption

- small deflection, max deflection $< \frac{3}{4} t$
- no stretching of the middle plane (no membrane effect)

Boundary Conditions

- The types of restraint around the boundary of a plate can be idealized as follows:

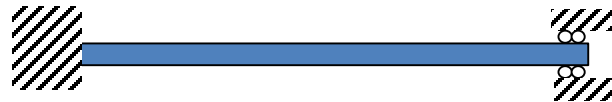
1. Simply supported: edges **free to rotate** and to **move in the plane**



2. Pinned: edges **free to rotate** but **not free to move in the plane**



3. Clamped but free to slide: edges **not free to rotate** but **free to move in the plane of the plate**;



4. Rigidly clamped: edges **not free to rotate** or **move in the plane**



- In most plated frame structures (particularly in vehicles and free-standing structure) : **little restraint against edge pull-in.**
→ BC. 1 and 3 are generally applicable

Simply Supported Plates

- The general expression for the load

$$p(x,y) = \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} A_{mn} \sin \frac{m\pi y}{a} \sin \frac{n\pi x}{b}$$

- The coefficient A_{mn} can be obtained by Fourier analysis for any particular load condition. For a uniform pressure p_0 .

$$A_{mn} = \frac{16p_0}{\pi^2 mn} \quad \text{Homework \#2 Prove it.}$$

- For the biharmonic equation, a sinusoidal load distribution produces a sinusoidal deflection

$$w = \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} B_{mn} \sin \frac{m\pi y}{a} \sin \frac{n\pi x}{b} \quad \Rightarrow \quad \frac{\partial^4 w}{\partial x^4} + 2 \frac{\partial^4 w}{\partial x^2 \partial y^2} + \frac{\partial^4 w}{\partial y^4} = \frac{p}{D}$$

$$B_{mn} \left(\frac{\pi^4 m^4}{a^4} + \frac{2\pi^4 m^2 n^2}{a^2 b^2} + \frac{\pi^4 n^4}{b^4} \right) = \frac{16p_0}{\pi^6 D mn}$$

$$B_{mn} = \frac{16p_0}{\pi^6 D mn \left(\frac{m^2}{a^2} + \frac{n^2}{b^2} \right)^2}$$

9.1 Small Deflection Theory

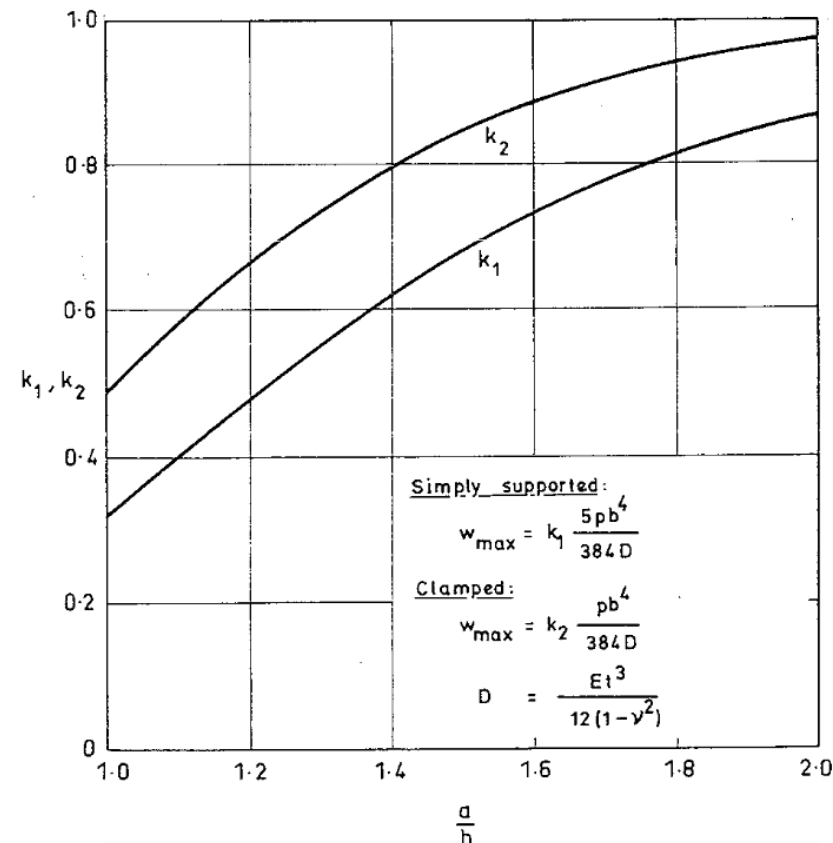
Simply Supported Plates

- Due to the symmetry of the problem, m and n only take odd values.

$$w = \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \frac{16 p_0}{\pi^6 D m n \left(\frac{m^2}{a^2} + \frac{n^2}{b^2} \right)} \sin \frac{m \pi y}{a} \sin \frac{n \pi x}{b}$$

- Effect of boundary condition
Clamped < Simply supported
about 4~5 times.
- Effect of aspect ratio
when $a=b$, the deflection is minimum

Homework #3 Plot the curve using any kinds of software.



Maximum deflection of rectangular plates under uniform pressure

9.1 Small Deflection Theory

Simply Supported Plates

- To determine the bending stress in the plate, calculate the bending stress.

$$\frac{\partial^2 w}{\partial x^2} = - \sum_{m=1}^{m=\infty} \sum_{n=1}^{n=\infty} \frac{16 p_0}{\pi^6 D m n \left(\frac{m^2}{a^2} + \frac{n^2}{b^2} \right)^2} \frac{\pi^2 n^2}{b^2} \sin \frac{m \pi y}{a} \sin \frac{n \pi x}{b} = - \sum_{m=1}^{m=\infty} \sum_{n=1}^{n=\infty} \frac{16 p_0 n}{\pi^4 D m b^2 \left(\frac{m^2}{a^2} + \frac{n^2}{b^2} \right)^2} \sin \frac{m \pi y}{a} \sin \frac{n \pi x}{b}$$

$$\frac{\partial^2 w}{\partial y^2} = - \sum_{m=1}^{m=\infty} \sum_{n=1}^{n=\infty} \frac{16 p_0 m}{\pi^4 D n a^2 \left(\frac{m^2}{a^2} + \frac{n^2}{b^2} \right)^2} \sin \frac{m \pi y}{a} \sin \frac{n \pi x}{b}$$

$$m_x = -D \left(\frac{\partial^2 w}{\partial x^2} + \nu \frac{\partial^2 w}{\partial y^2} \right) = \sum_{m=1}^{m=\infty} \sum_{n=1}^{n=\infty} \frac{16 p_0}{\pi^4 m n \left(\frac{m^2}{a^2} + \frac{n^2}{b^2} \right)^2} \left(\frac{n^2}{b^2} + \nu \frac{m^2}{a^2} \right) \sin \frac{m \pi y}{a} \sin \frac{n \pi x}{b}$$

- The bending moment will always be greater across the shorter span.

Homework #4 Prove this.

- The symbol b is used for shorter dimension, then $a/b > 1$
- m_x is the larger of the two moments and has its maximum value in the center of the plate.

9.1 Small Deflection Theory

Clamped Plates

- The usual methods are the energy(or Ritz) method and the method of Levy
- The energy method, while giving only approximate results
- The Levy type solution is achieved by the superposition of three loading systems applied to a simply supported plate :
 - (1) uniformly distributed along the short edges : deflection w_1
 - (2) moments distributed along the short edges : deflection w_2
 - (3) moments distributed along the long edges : deflection w_3

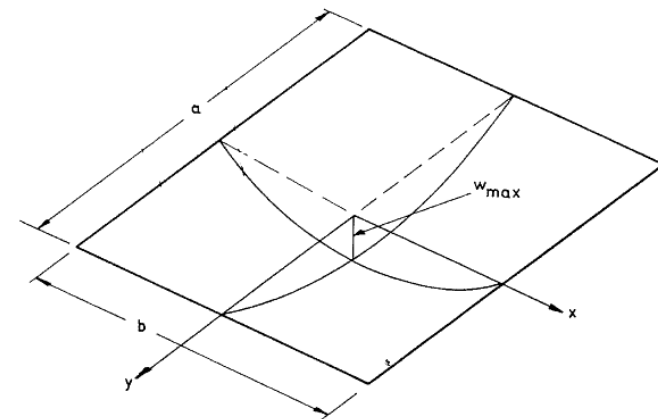
short edges :

$$\left(\frac{\partial \omega_1}{\partial y} \right)_{y=\pm a/2} + \left(\frac{\partial \omega_2}{\partial y} + \frac{\partial \omega_3}{\partial y} \right)_{y=\pm a/2} = 0$$

long edges :

$$\left(\frac{\partial \omega_1}{\partial x} \right)_{x=\pm b/2} + \left(\frac{\partial \omega_2}{\partial x} + \frac{\partial \omega_3}{\partial x} \right)_{x=\pm b/2} = 0$$

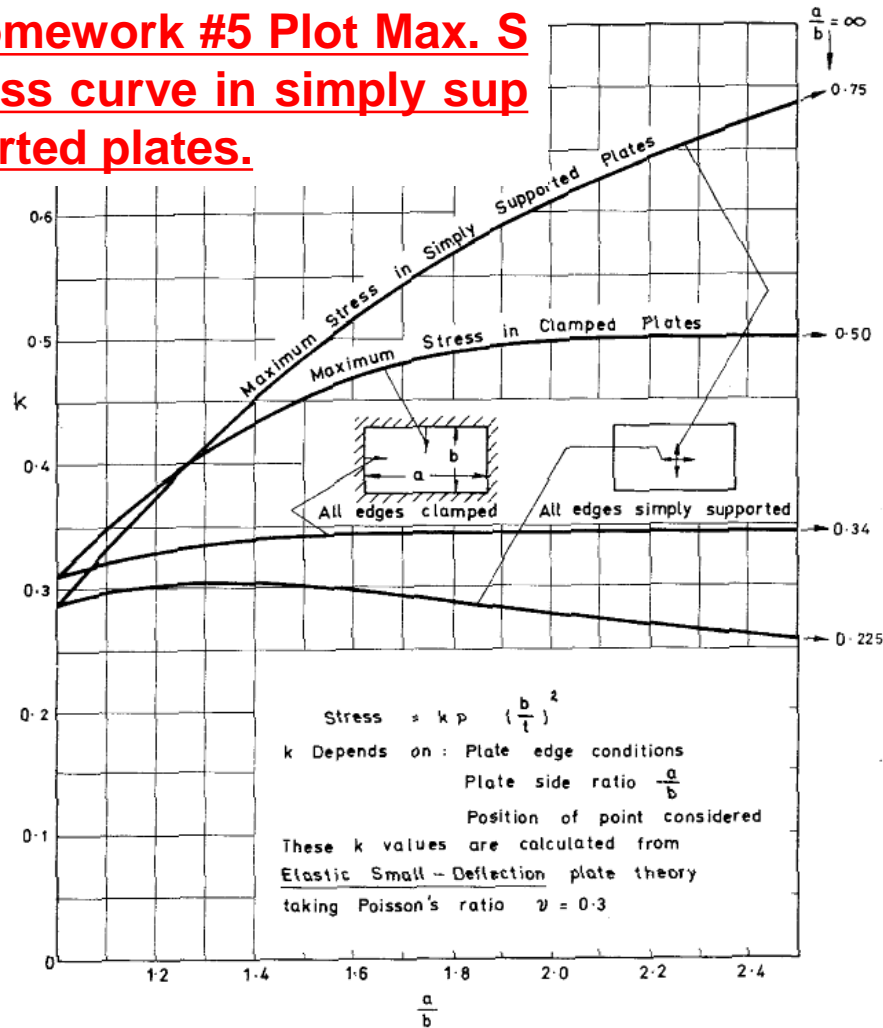
(where $x=0, y=0$ is at the center of the plate)



9.1 Small Deflection Theory

Clamped Plates

Homework #5 Plot Max. Stress curve in simply supported plates.



$$\text{Stress} = kp \left(\frac{b}{t} \right)^2$$

- ❖ k depends on :
 - plate edge conditions
 - Plate side ratio a/b
 - Position of point considered

The effect of aspect ratio is smaller for clamped plates, and beyond about $a/b=2$

such a plate behaves essentially as a clamped strip and the influence of aspect ratio is negligible.

Stresses in rectangular plates under uniform lateral pressure

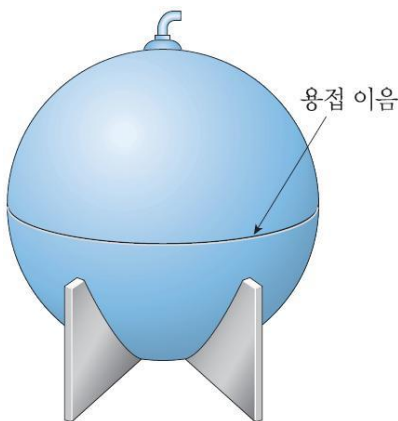
Large-Deflection Plate Theory

- Membrane stress arises when the deflection becomes large and/or when the edges are prevented from pulling in.
- Small-deflection theory fails to allow for membrane stresses.
- As the deflection increases, an increasing prortion of the load is carried by this membrane action.
- The lateral load is supported by both bending and membrane action.
- A more comprehensive plate theory, usually referred to as "large-deflection" plate theory.

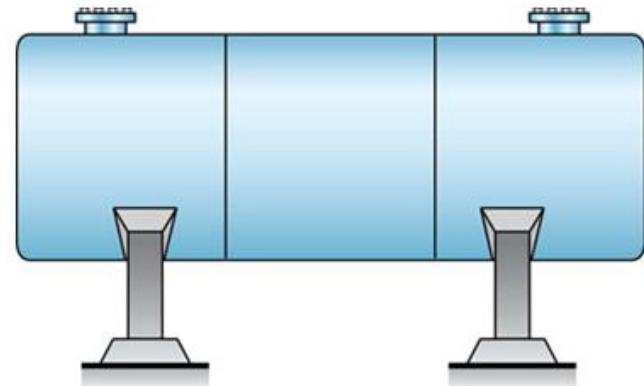
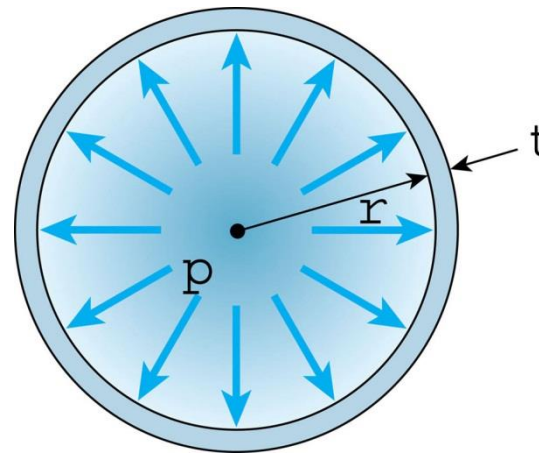
Membrane Tension(Edges Restrained Against Pull-in)

❖ What is membrane tension?

- **Membrane stress** : Stresses that act **tangentially to the curved surface of a shell**
- Pressure vessel : closed structures containing liquids or gases under pressure
- Gage pressure (Internal – external pressure) is resisted by membrane tension



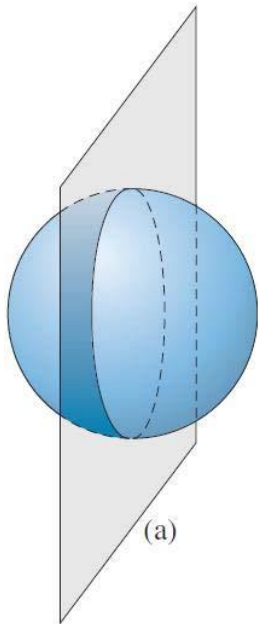
Spherical Pressure vessel



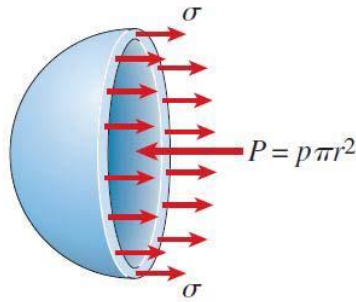
Cylindrical Pressure vessel

Membrane stress

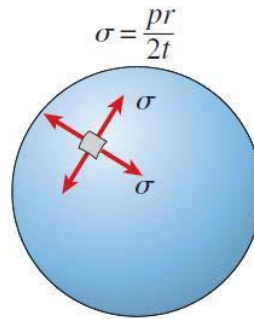
❖ Spherical Pressure vessel



(a)



(b)



(c)

$$\sum F_{horiz} = 0$$

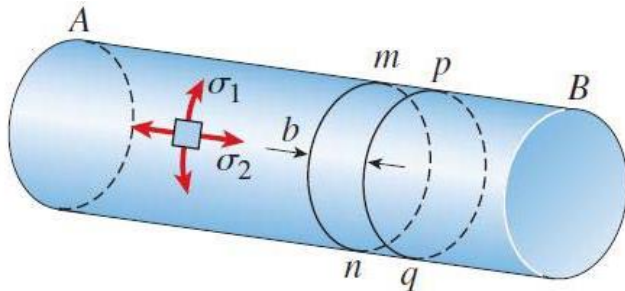
$$\sigma(2\pi r_m t) = p(\pi r^2)$$

$$r_m = r + \frac{t}{2}$$

$$\therefore \sigma = \frac{pr^2}{2r_m t} \approx \frac{pr}{2t}$$

Membrane stress

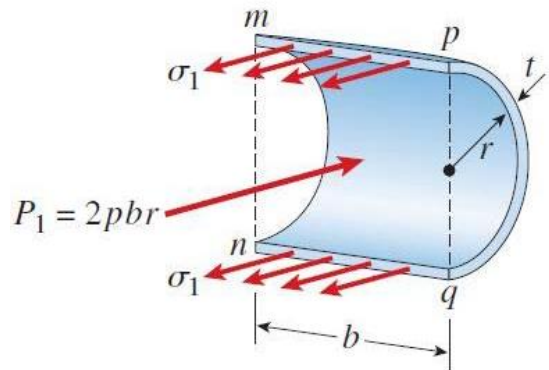
❖ Cylindrical Pressure vessel



σ_1 : **circumferential stress**
hoop stress

$$\sigma_1(2bt) - 2pbr = 0$$

$$\therefore \sigma_1 = \frac{pr}{t}$$



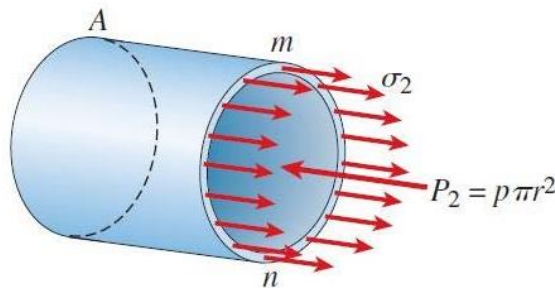
(b)

σ_2 : **longitudinal stress**

$$\sigma_2(2\pi rt) - p\pi r^2 = 0$$

$$\therefore \sigma_2 = \frac{pr}{2t}$$

$$\sigma_1 = 2\sigma_2$$



(c)

Membrane stress

❖ Ship & Offshore Structure

Global Stress

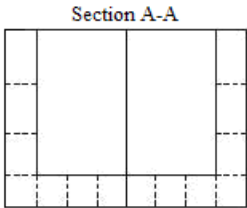
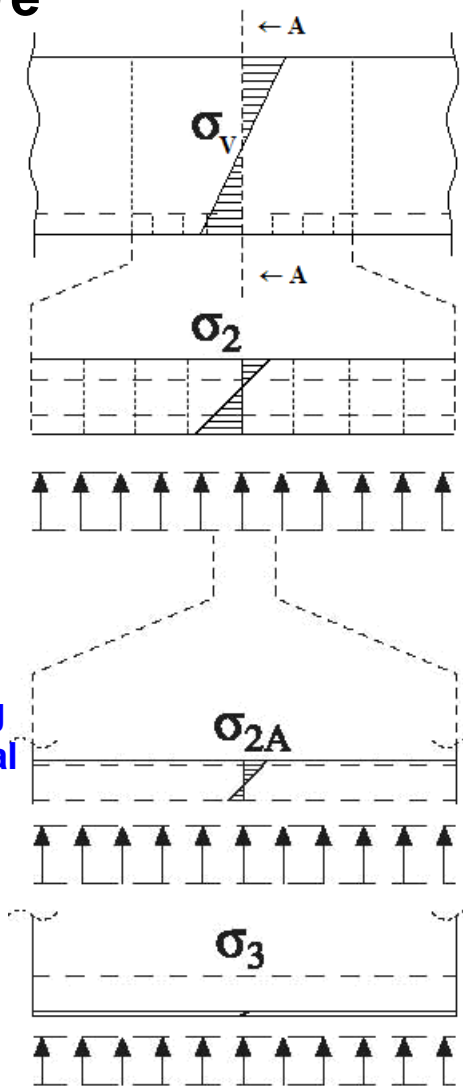
Hull Girder Bending

Double Bottom Bending

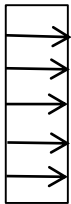
Local Stress

Local Bending of Longitudinal

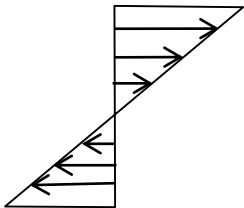
Plate Bending



Membrane stress on plating



Membrane stress



Bending stress

Bending stress on plating

of elements should be more than thee (3) between boundaries.

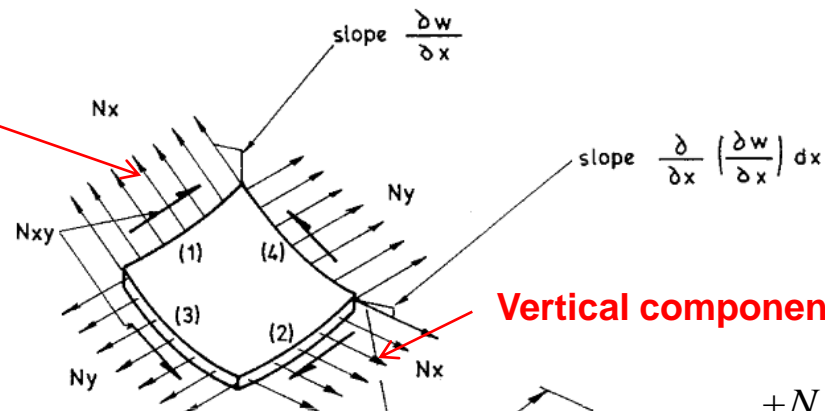


Large-Deflection Plate Theory by von Karman

- N_x, N_y : membrane tension per unit length
- N_{xy} : membrane shear force per unit length
- Vertical component of the tension force at side (1)
 $-N_x dy \sin(\text{slope}) \approx -N_x dy \tan(\text{slope}) = -N_x dy (\partial w / \partial x)$

Vertical component

$$-N_x dy \left(\frac{\partial w}{\partial x} \right)$$



Vertical component

$$+N_x dy \left[\frac{\partial w}{\partial x} + \frac{\partial}{\partial x} \left(\frac{\partial w}{\partial x} \right) dx \right]$$

- Net force in the x-direction

$$-N_x dy \left(\frac{\partial^2 w}{\partial x^2} \right)$$

- Net force in the y-direction

$$-N_y dx \left(\frac{\partial^2 w}{\partial y^2} \right)$$

Large-Deflection Plate Theory by von Karman

- The vertical component of the shear force along sides (3) and (4) is

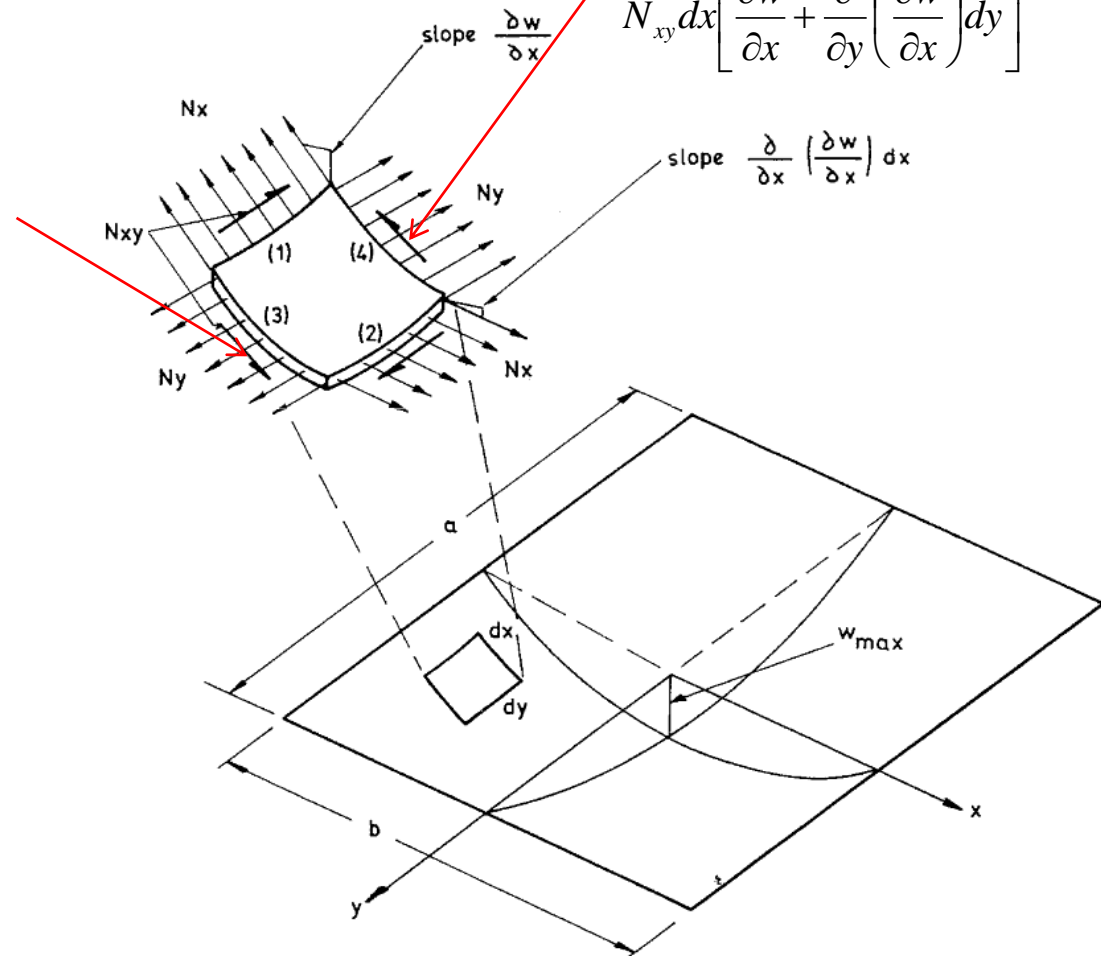
$$N_{xy}dx \left[\frac{\partial w}{\partial x} + \frac{\partial}{\partial y} \left(\frac{\partial w}{\partial x} \right) dy \right] - N_{xy}dx \left(\frac{\partial w}{\partial x} \right)$$

Vertical component

$$N_{xy}dx \left(\frac{\partial w}{\partial x} \right)$$

Vertical component

$$N_{xy}dx \left[\frac{\partial w}{\partial x} + \frac{\partial}{\partial y} \left(\frac{\partial w}{\partial x} \right) dy \right]$$



Large-Deflection Plate Theory by von Karman

- Total vertical forces

$$\left(N_x \frac{\partial^2 w}{\partial x^2} + 2N_{xy} \frac{\partial^2 w}{\partial x \partial y} + N_y \frac{\partial^2 w}{\partial y^2} \right) dx dy$$

- New terms are added

$$\frac{\partial^4 w}{\partial x^4} + 2 \frac{\partial^4 w}{\partial x^2 \partial y^2} + \frac{\partial^4 w}{\partial y^4} = \frac{1}{D} \left(p + N_x \frac{\partial^2 w}{\partial x^2} + 2N_{xy} \frac{\partial^2 w}{\partial x \partial y} + N_y \frac{\partial^2 w}{\partial y^2} \right) \quad \leftarrow \quad \boxed{\frac{\partial^4 w}{\partial x^4} + 2 \frac{\partial^4 w}{\partial x^2 \partial y^2} + \frac{\partial^4 w}{\partial y^4} = \frac{p}{D}}$$

$$\boxed{\nabla^4 w = \frac{1}{D} \left(p + N_x \frac{\partial^2 w}{\partial x^2} + 2N_{xy} \frac{\partial^2 w}{\partial x \partial y} + N_y \frac{\partial^2 w}{\partial y^2} \right)}$$

- N_x , N_y , N_{xy} are functions of x and y .
- Except for a few simple cases, precise mathematical solutions are very difficult.

Membrane Tension(Edges Restrained Against Pull-in)

- The **relative magnitude of membrane effects** depends on
 - : the degree of lateral deflection or "curvature" of the plate surface
 - the degree to which the edges are restrained from pulling in.
- If the edges are restrained and large deflection ($w > 1.5 t$)
the contribution of **membrane tension** > **bending**, vice versa.
- In ship plating : **little restraint against edge pull-in**
for large deflections, supporting stiffeners or beams fails earlier.
- Nevertheless, some situations in which large deflections can be permitted,
→ **the use of membrane tension can give substantial weight savings.**

Membrane Tension(Edges Restrained Against Pull-in)

❖ The case of a unit-width strip

- Laterally loaded
- Edges are prevented from approaching
- the difference between the arc length of the deflected strip and the original straight length

$$\delta = \int_0^b \left(\sqrt{1 + \left(\frac{d\omega}{dx} \right)^2} - 1 \right) dx \quad \delta \cong \int_0^b \frac{1}{2} \left(\frac{d\omega}{dx} \right)^2 dx$$

- Membrane force is unidirectional ($N_y = N_{xy} = 0$) and constant over the length.
- Extension due to the tension

$$T = \frac{Et}{2b} \int_0^b \left(\frac{d\omega}{dx} \right)^2 dx$$

Membrane Tension(Edges Restrained Against Pull-in)

- w_0 initial deflection, w_1 due to the load

$$\omega(x) = (\omega_0 + \omega_1) \sin \frac{\pi x}{b}$$

- The total change in length is

$$\begin{aligned}\delta &= \int_0^b \frac{1}{2} (\omega_0 + \omega_1)^2 \frac{\pi^2}{b^2} \cos^2 \frac{\pi x}{b} dx \\ &= \frac{\pi^2}{4b} (\omega_0 + \omega_1)^2\end{aligned}$$

- The change in length due to the initial deflection

$$\frac{\pi^2 w_0^2}{4b}$$

Membrane Tension(Edges Restrained Against Pull-in)

- The change in length due to the loading is the difference between these two:

$$\delta_1 = \frac{\pi^2}{4b} (2\omega_0\omega_1 + \omega_1^2)$$

- Strain energy due to bending

$$= \int_0^b \frac{E'I}{2} \left(\frac{d^2 w}{dx^2} \right) dx = \int_0^b \frac{E'I}{2} (w_1)^2 \frac{\pi^4}{b^4} \sin^2 \frac{\pi x}{b} dx = \frac{\pi^4 E'I (w_1)^2}{4b^3}$$

- Strain energy due to tension:

$$\frac{1}{2} T \delta_p = \frac{1}{2} \frac{AE}{b} \delta_1^2$$

- The work done

$$W = \int_0^b \frac{1}{2} p w dx = \int_0^b \frac{1}{2} p w_1 \sin\left(\frac{\pi x}{b}\right) dx = \frac{p w_1 b}{\pi}$$

- The work done by the load is equal to the total strain energy, hence:

$$\frac{p w_1 b}{\pi} = \frac{AE \delta_1^2}{2b} + \frac{\pi^4 E'I (w_1^2)}{4b^3} = \frac{\pi^4}{32} \frac{AE}{b^3} (2w_0 w_1 + w_1^2)^2 + \frac{\pi^4 E'I (w_1^2)}{4b^3}$$

Strain Energy of Bending

- When bending moment M varies along its length

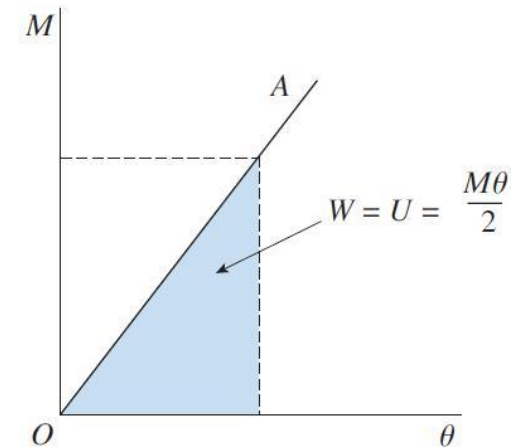
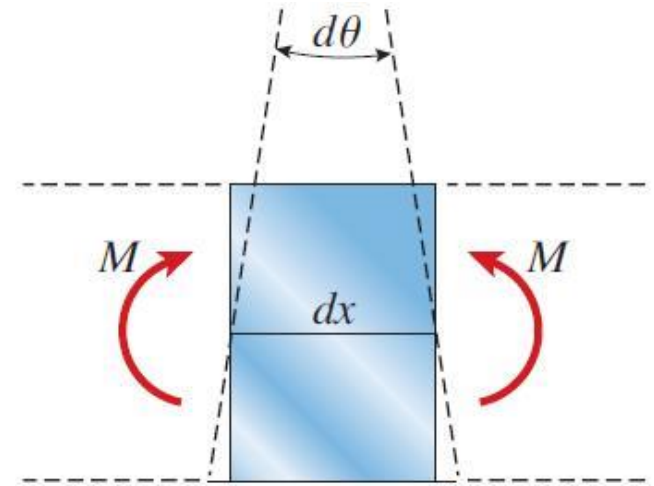
$$d\theta = \kappa dx = \frac{1}{\rho} dx = \frac{d^2v}{dx^2} dx$$

$$dU = \frac{Md\theta}{2} = \frac{M}{2} \frac{d^2v}{dx^2} dx = \frac{M}{2} \frac{M}{EI} dx = \frac{M^2 dx}{2EI} \text{ or}$$

$$dU = \frac{EI}{2dx} \left(\frac{d^2v}{dx^2} dx \right)^2 = \frac{EI}{2} \left(\frac{d^2v}{dx^2} \right)^2 dx$$

- if integrated

$$U = \int \frac{M^2 dx}{2EI} \text{ or } U = \int \frac{EI}{2} \left(\frac{d^2v}{dx^2} \right)^2 dx$$



Membrane Tension(Edges Restrained Against Pull-in)

- The work done by the load is equal to the total strain energy, hence:

$$\frac{\pi^4 AE \omega_1^3}{32b^3} + \frac{\pi^4 AE \omega_0 \omega_1^2}{8b^3} + \left(\frac{\pi^4 AE \omega_0^2}{8b^3} + \frac{\pi^4 E' I}{4b^3} \right) \omega_1 - \frac{pb}{\pi} = 0$$

- $A=t$, $E' = E/(1-\nu)$ and $I=t^3/12$ per unit width

$$\omega_1^3 + 4\omega_0 \omega_1^2 + \left(4\omega_0^2 + \frac{2}{3} \frac{t^2}{1-\nu^2} \right) \omega_1 - \frac{32pb^4}{\pi^5 Et} = 0$$

Effect of Initial Deformation

- If there is no initial deflection:

$$\cancel{\omega_1^3} + \frac{2}{3} \frac{t^2}{1-\nu^2} \omega_1 = \frac{32pb^4}{\pi^5 Et}$$

- For the initial stages of loading the deflection w_1 will be small relative to the thickness, and hence the first term may be neglected:

$$\frac{\omega_1}{t} = \frac{48(1-\nu^2)p}{\pi^5 E} \left(\frac{b}{t}\right)^4 = 0.143 \frac{p}{E} \left(\frac{b}{t}\right)^4$$

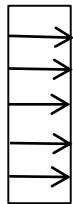
- The result obtained from small-deflection theory, ignoring membrane action:

$$\frac{\omega_1}{t} = \frac{5(1-\nu^2)p}{32E} \left(\frac{b}{t}\right)^4 = 0.142 \frac{p}{E} \left(\frac{b}{t}\right)^4$$

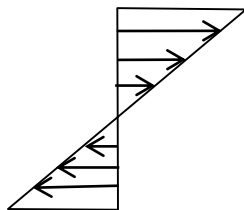
9.2 Combined Bending and Membrane Stresses-Elastic Range

Effect of Initial Deformation

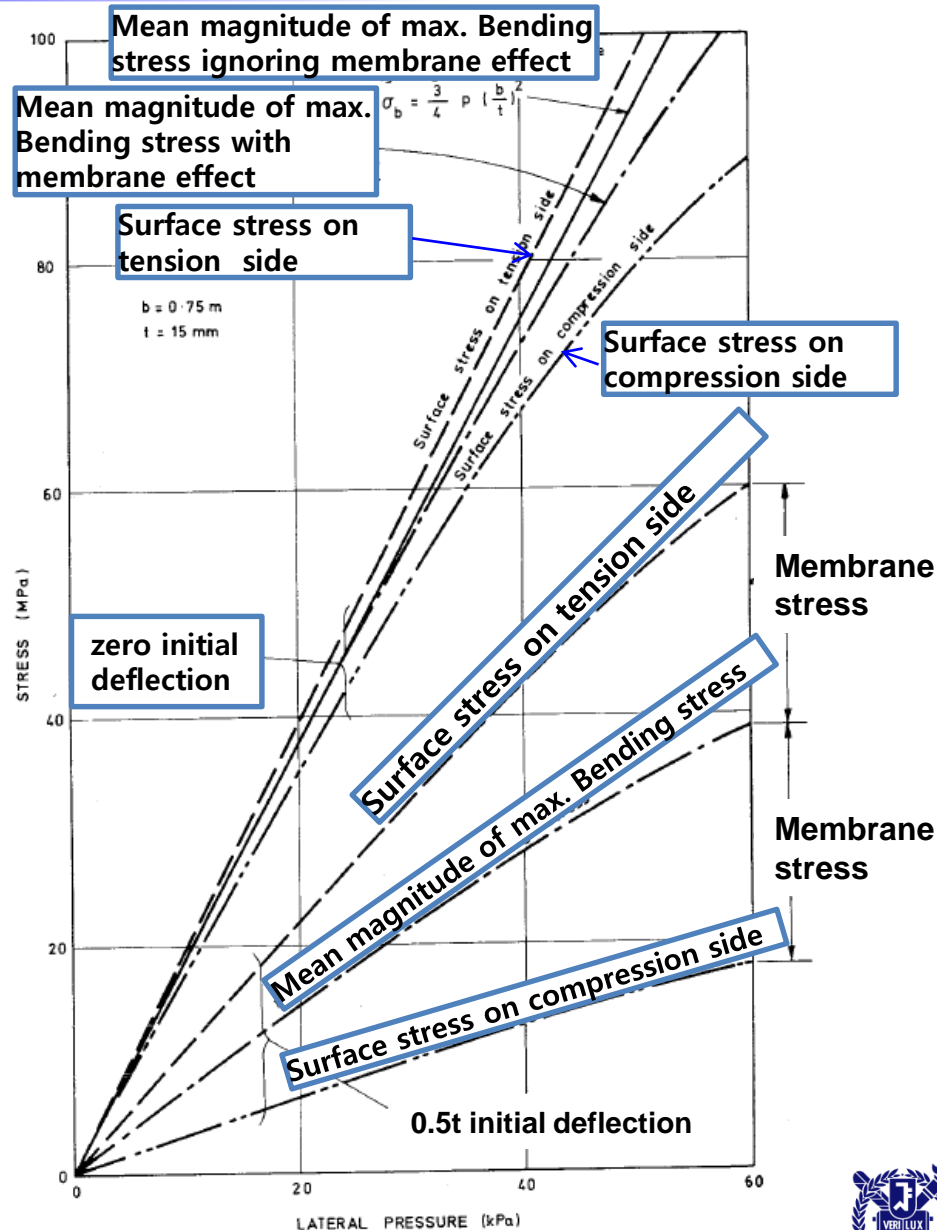
- Membrane action requires **some deflection**, either **initial** or **due to load**
- if there is no initial deflection → membrane action does not become significant until the deflection due to load approaches the plate thickness.
- Welding deformation & permanent set give a beneficial influence on elastic strength of plates.



Membrane stress



Bending stress

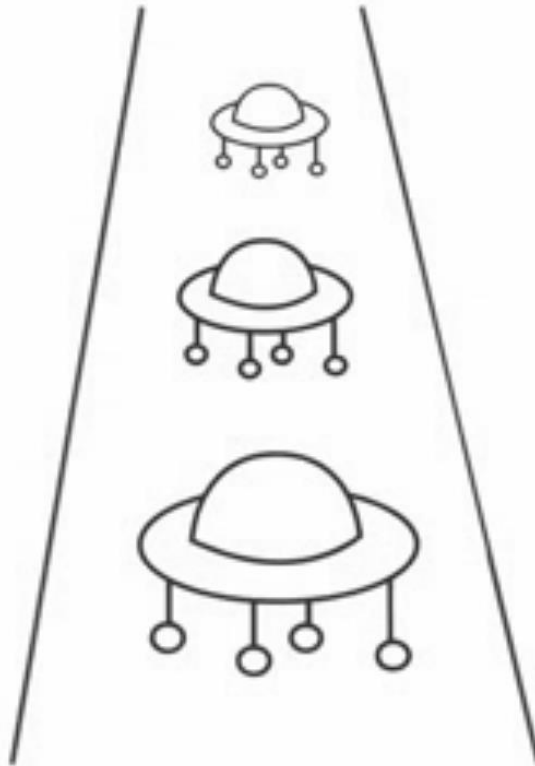


Effect of Initial Deformation

- The beneficial effect of initial deformation.

Type of Initial Deformation	Elastic Strength	Source of Increase in Elastic Strength
Flat plate	1.59	-
Initial deflection (stress-free) equal to plating thickness	2.58	Membrane action
Initially flat plate dished to a permanent set equal to plating thickness	4.50	Membrane action plus residual stresses

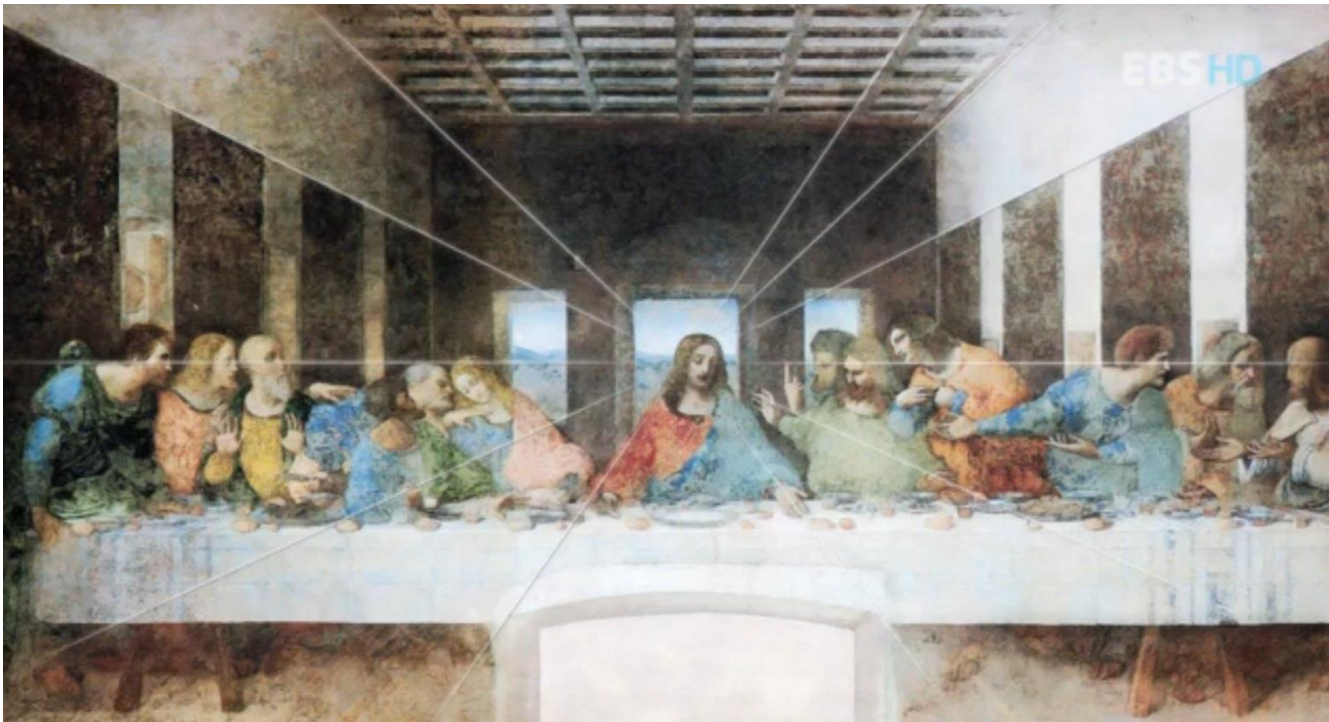
- Initial lateral deformation is beneficial only when the plate edges are **at least partly restrained from pulling in**, thus **allowing the development of in-plane tension**.



Which one is in front?

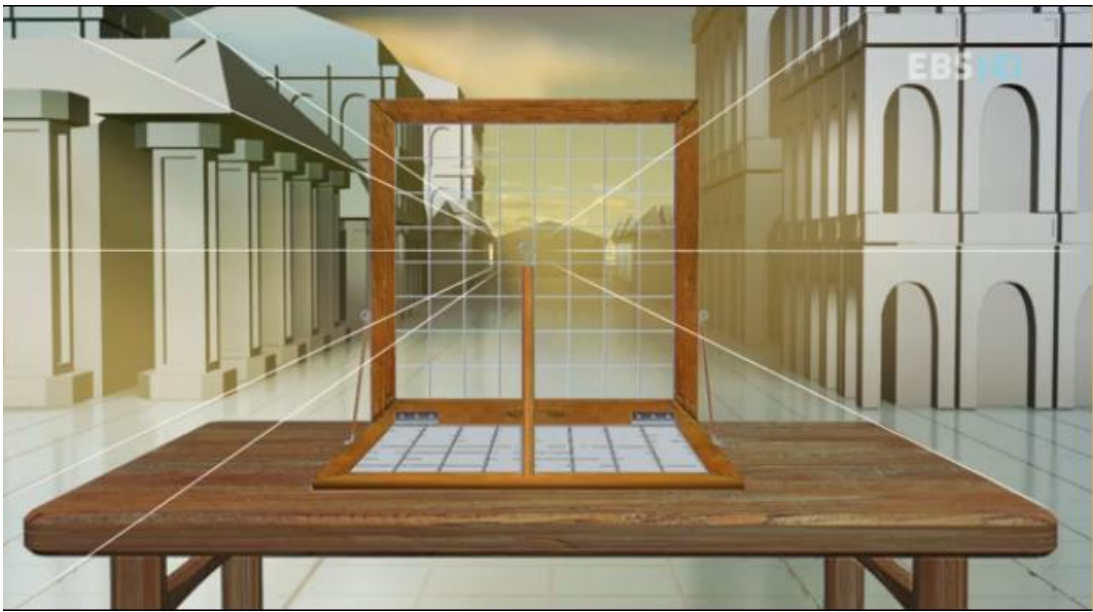
Di Vinci's Last Supper

- Famous for not only artistic value but also fine command of perspective

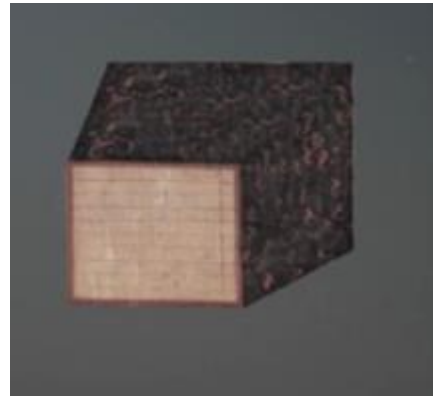


Westerner try to see and Asian try to be

- The law of perspective is one of representative characteristics
- Object has the meaning of observation and all things
- Objective is not subjective.
- I see = I understand
- Seeing is believing, Westerner makes seeing very important.
- From the observer's perspective, farthest thing indicates that in front.



- Closer thing is drawn larger and the farther thing smaller. -> reverse perspective.
- Western artists draw picture seeing an object in front of him.
- Asian artists draw picture after seeing and feeling the object and returning.

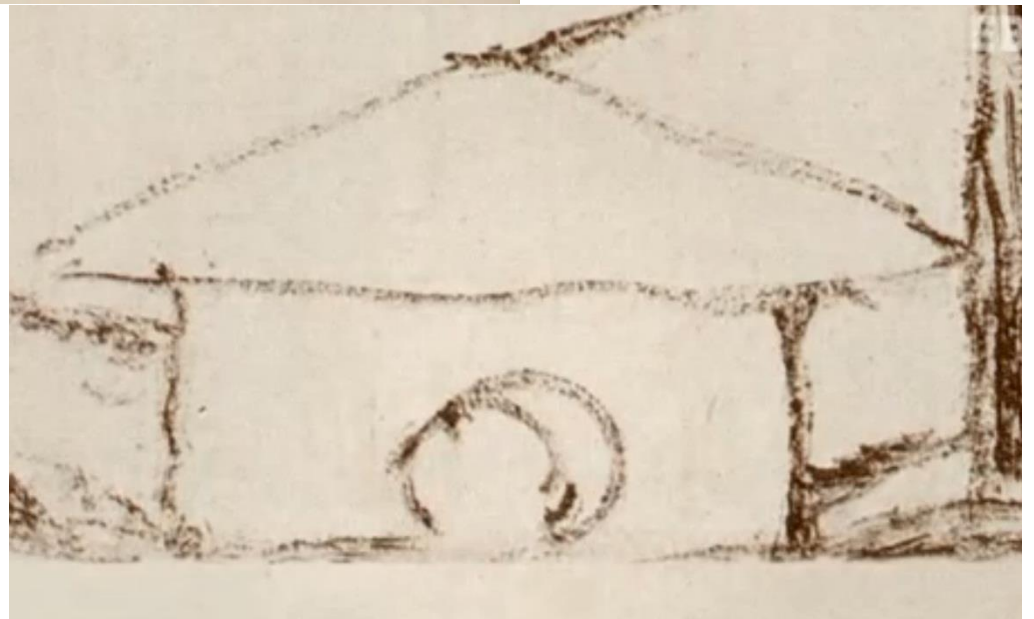


寒山圖

蕭子雲
畫



세한도(국보 180호)



- Universe is covered by Indra net. Each marble is mirroring others



