First Law of Thermodynamics and Energy Equation (1) (Lecture 3)

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(*) Some texts and figures are borrowed from Sonntag & Borgnakke unless noted otherwise.



- 3.1 The Energy Equation
- → Energy of a substance consists of two main categories:

E = *Microscopic Energy* + *Macroscopic Energy* (by molecular dynamics in CM) (by mean position or motion of CM)

$$E = Internal \ Energy + Kinetic \ Energy + Potential \ Energy$$
$$E = U + KE + PE \quad \leftarrow \quad KE = \frac{1}{2}mV^2, \ PE = mgh$$
$$= m(u + \frac{1}{2}V^2 + gh)$$

Examples of internal (microscopic) energy (left: rotational modes, right: vibrational modes)

Examples of external (macroscopic) energy (potential energy vs kinetic energy)



- Then, energy equation for a control mass can be easily expressed as:
 - $\Delta Energy = +(In) (Out) \rightarrow Energy Conservation$



$$dE_{cv} = dU + d(KE) + d(PE) = \delta Q - \delta W$$

→ infinitesimal change

$$\int_{1}^{2} dE_{cv} = \int_{1}^{2} \left[dU + d(KE) + d(PE) \right] = \int_{1}^{2} \delta Q - \int_{1}^{2} \delta W$$

$$E_{cv,2} - E_{cv,1} = U_{2} - U_{1} + \frac{1}{2} m \left(V_{2}^{2} - V_{1}^{2} \right) + mg \left(Z_{2} - Z_{1} \right) = _{1} Q_{2} - _{1} W_{2}$$

finite change from state 1 to state 2



3.2 The First Law of Thermodynamics

→ The first law of thermodynamics states:

"During any cycle a system undergoes, the cyclic integral of the heat is proportional (or equal) to the cyclic integral of the work."

$$\oint \delta Q = \oint \delta W$$



For a thermodynamic cycle, initial gas state should be equal to final gas state.



→ Now consider a control mass that undergoes a change of state. From the first law of thermodynamics,

$$\oint \delta Q = \oint \delta W$$



Case 1:
$$1 \xrightarrow{A} 2 \xrightarrow{B} 1$$

$$\int_{1}^{2} \delta Q_{A} + \int_{2}^{1} \delta Q_{B} = \int_{1}^{2} \delta W_{A} + \int_{2}^{1} \delta W_{B}$$

Case 2: $1 \xrightarrow{A} 2 \xrightarrow{C} 1$

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$$\int_1^2 \delta Q_A + \int_2^1 \delta Q_C = \int_1^2 \delta W_A + \int_2^1 \delta W_C$$



$$\begin{array}{l} \bullet \quad \text{Here,} \qquad \qquad \int_{1}^{2} \delta Q_{A} + \int_{2}^{1} \delta Q_{B} = \int_{1}^{2} \delta W_{A} + \int_{2}^{1} \delta W_{B} \\ \hline & - \underbrace{\Big) \quad \int_{1}^{2} \delta Q_{A} + \int_{2}^{1} \delta Q_{C} = \int_{1}^{2} \delta W_{A} + \int_{2}^{1} \delta W_{C} \\ \hline & \int_{2}^{1} \delta Q_{B} - \int_{2}^{1} \delta Q_{C} = \int_{2}^{1} \delta W_{B} - \int_{2}^{1} \delta W_{C} \\ \hline & \int_{2}^{1} \left(\delta Q_{B} - \delta W_{B} \right) = \int_{2}^{1} \left(\delta Q_{C} - \delta W_{C} \right) \end{array}$$

Then, the value $\delta Q - \delta W$ doesn't depend on the path, but only on the initial and final state of the process. \rightarrow Point function!

→ Define a thermodynamic property of a substance, Energy(E), as:

$$\delta Q - \delta W = dE$$
$$_1 Q_2 - _1 W_2 = E_2 - E_1$$



$$\Delta Energy = +(In) - (Out)$$

→ Same with energy equation



3.3 Definition of Work

- Work is an important way of energy transfer in thermodynamics!
- In general physics, work is defined as a force F acting through a displacement x, where the displacement is in the direction of the force:

$$\delta W = F \cdot dx \quad \rightarrow \quad W = \int_1^2 F \cdot dx$$

 In thermodynamics, the direction of work transfer across the system boundary is important.

outward : work done by a system (output)
inward : work done on a system (input)





Work is the product of force and distance and the unit for work in SI units is joule (J):

$$J = 1 N \cdot m$$

Power is the time rate of doing work and the unit for power in SI units is watt (W):

$$\dot{W} \equiv \frac{\delta W}{dt} \rightarrow 1W = 1 \text{ J/s}$$

→ Specific work is the work per unit mass of the system:

$$w \equiv \frac{W}{m} \left(J/kg \right)$$



- 3.4 Work done at the moving boundary of a simple compressible system
- Let's consider the moving piston in a cylinder where the volume of a simple compressible substance changes in a quasi-equilibrium process.

 $\delta W = F dL$ $\delta W = P A dL$ $\delta W = P dV$ ${}_{1}W_{2} = \int_{1}^{2} F dL = \int_{1}^{2} P dV$ $dV : + \implies W : \text{outward}$

$$dV: - \Rightarrow W:$$
 inward





(Ref.) General Systems that Involve Work





Pressure-Volume (P-V) diagram is widely used to represent the states and processes of a simple compressible substance.

$$W = \int_1^2 \delta W = \int_1^2 P dV$$

=area under curve 1-2

$$dV: + \Rightarrow W:$$
 outward $(2 \rightarrow 1)$

$$dV: - \Rightarrow W:$$
 inward $(1 \rightarrow 2)$

- → Work is path function or, in a mathematical term, inexact differential(e.g. δW).
- → Thermodynamic properties are point functions, or exact differential(e.g. dV).

$$\Delta V = \int_{1}^{2} dV = V_{2} - V_{1}$$





To evaluate the following integral, the relationship between P and V should be given:

$$W = \int_{1}^{2} \delta W = \int_{1}^{2} P dV$$

One common example is a polytropic process, where

 $PV^n = Const. = C$ or $P = \frac{C}{V^n}$ n: polytropic coefficient $(-\infty \sim +\infty)$

$$n \neq 1 \qquad _{1}W_{2} = \int_{1}^{2} \delta W = \int_{1}^{2} P dV = \int_{1}^{2} \frac{C}{V^{n}} dV = \int_{1}^{2} CV^{-n} dV$$
$$= C \cdot \left[\frac{V^{1-n}}{1-n} \right]_{1}^{2} = \frac{C}{1-n} \left(V_{2}^{1-n} - V_{1}^{1-n} \right) = \frac{P_{2}V_{2} - P_{1}V_{1}}{1-n}$$

$$n = 1 \qquad {}_{1}W_{2} = \int_{1}^{2} \delta W = \int_{1}^{2} P dV = \int_{1}^{2} \frac{C}{V} dV = C \ln \frac{V_{2}}{V_{1}} = P_{1}V_{1} \ln \frac{V_{2}}{V_{1}}$$



 For non-equilibrium process in a system, consider external force and volume changes by surroundings.

$$P_{ext} = F_{ext} / A = P_0 + m_p g / A$$
$${}_1W_2 = \int_1^2 P_{ext} dV = P_{ext} (V_2 - V_1)$$

