Cryogenic Engineering

Chapter 3.

Gas Liquefaction System

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Gas – liquefaction systems

The systems that can produce low temperatures required for liquefaction.

System performance parameters

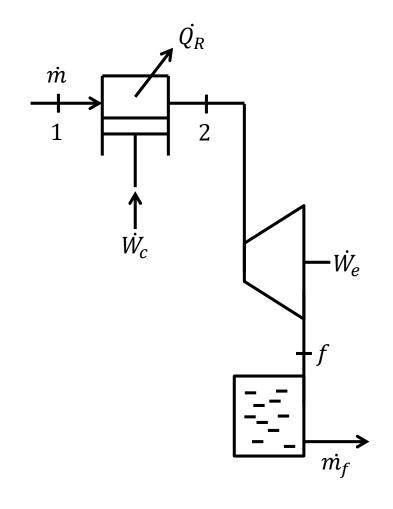
- $\frac{\text{Work required}}{\text{Unit mass of gas compressed}} = -\frac{\dot{W}}{\dot{m}}$
- $\frac{\text{Work required}}{\text{Unit mass of gas liquefied}} = -\frac{\dot{W}}{\dot{m}_{f}}$
- Liquefiecd fraction of the total flow of gas = $y = -\frac{\dot{m_f}}{\dot{m}}$
- Figure of Merit(FoM) = $\frac{\text{Theoretical work}}{\text{Actual work}} = -\frac{\dot{W_i}}{\dot{W}}$ (0~1)

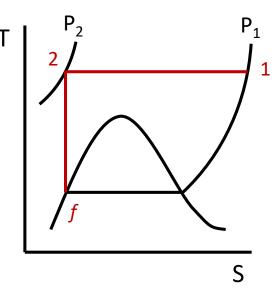
3.1 System Performance Parameters

+ Additional performance parameters

- Compressor adiabatic efficiency
- Expander adiabatic efficiency
- Compressor mechanical efficiency
- Expander mechanical efficiency
- HX effectiveness
- Pressure drop
- Heat transfer to the system

3.2 Thermodynamically Ideal Cycle (Liquefaction system - open cycle)





1-2: Isothermal compression2-f: Isentropic expansion(expander)

In this case, $P_2 = 70 \sim 80$ GPa for N_2 (too high!) Apply 1st and 2nd laws of thermodynamics to the system

1st law of thermodynamics:

$$\begin{split} \dot{Q}_{net} - \dot{W}_{net} &= \sum_{outlets} \dot{m} \left(h + \frac{v^2}{2g_c} + \frac{gz}{g_c} \right) - \sum_{inlets} \dot{m} \left(h + \frac{v^2}{2g_c} + \frac{gz}{g_c} \right) \\ \rightarrow \dot{Q}_{net} - \dot{W}_{net} &= \sum_{outlets} \dot{m}h - \sum_{inlets} \dot{m}h \\ \rightarrow \dot{Q}_R - \dot{W}_i &= \dot{m}(h_f - h_1) = -\dot{m}(h_1 - h_f) \\ (\dot{W}_i &= W_e - W_c) \end{split}$$

3.2 Thermodynamically Ideal Cycle (Liquefaction system - open cycle)

Apply 1st and 2nd laws of thermodynamics to the system

2nd law of thermodynamics:

$$ds = \frac{\delta q}{T} \bigg|_{rev,ideal}$$

$$\rightarrow \delta q = T \cdot ds$$

$$\rightarrow \dot{Q}_{R} = \dot{m}T_{1}(S_{2} - S_{1}) = -\dot{m}T_{1}(S_{1} - S_{f})$$

3.2 Thermodynamically Ideal Cycle (Liquefaction system - open cycle)

Apply 1st and 2nd laws of thermodynamics to the system

From 1st and 2nd laws of thermodynamics, $\dot{Q_R} - \dot{W_i} = \dot{m}(h_f - h_1) = -\dot{m}(h_1 - h_f) \cdots 1$ st law $\dot{Q_R} = \dot{m}T_1(S_2 - S_1) = -\dot{m}T_1(S_1 - S_f) \cdots 2$ nd law

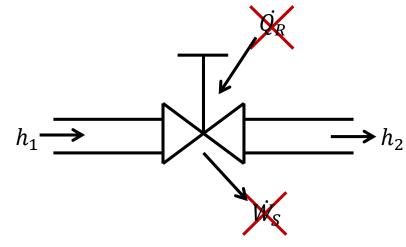
$$\rightarrow -\frac{\dot{W}_i}{\dot{m}} = T_1(S_1 - S_f) - (h_1 - h_f) = -\frac{\dot{W}_i}{\dot{m}_f}$$

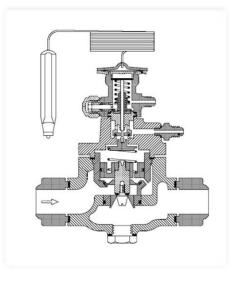
*
$$\dot{m} = \dot{m}_f \rightarrow \text{ liquid yield } \mathbf{y} = \frac{\dot{m}_f}{\dot{m}} = \mathbf{1}$$

Gas	Normal Boiling Point		Ideal Work of Liquefaction, $-\dot{W}_i/\dot{m}_f$	
	K	°R	kJ/kg	Btu/lb _m
Helium-3	3.19	5.74	8 178	3 516
Helium-4	4.21	7.58	6 819	2 931
Hydrogen, H ₂	20,27	36.5	12 019	5 167
Neon, Ne	27.09	48.8	1 335	574
Nitrogen, N ₂	77.36	139.2	768.1	330.2
Air	78.8	142	738.9	317.7
Carbon monoxide, CO	81.6	146.9	768.6	330.4
Argon, A	87.28	157.1	478.6	205.7
Oxygen, O ₂	90.18	162.3	635.6	273.3
Methane, CH ₄	111.7	201.1	1 091	469
Ethane, C_2H_6	184.5	332.1	353.1	151.8
Propane, C ₃ H ₆	231.1	416.0	140.4	60.4
Ammonia, NH ₃	239.8	431.6	359.1	154.4

Table 3.1. Ideal-work requirements for liquefaction of gases beginning at 300 K (80°F) and 101.3 kPa (14.7 psia)







1st law of thermodynamics:

$$\dot{Q}_{R} - \dot{W}_{S} = \sum_{\text{outlets}} \dot{m} \left(h + \frac{v^{2}}{2g_{c}} + \frac{gz}{g_{c}} \right) - \sum_{\text{inlets}} \dot{m} \left(h + \frac{v^{2}}{2g_{c}} + \frac{gz}{g_{c}} \right)$$

$$\rightarrow 0 = \dot{m}(h_{2} - h_{1})$$

$$\rightarrow h_{2} = h_{1}$$

Joule – Thomson coefficient

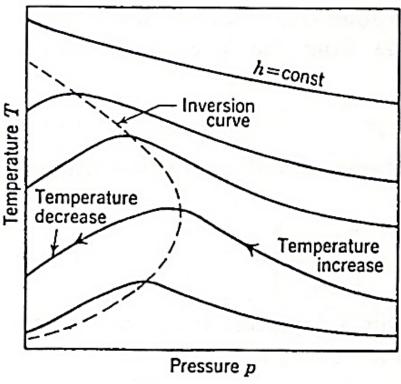
$$\mu_{JT} = \frac{\partial T}{\partial P} \bigg|_{h}$$

 → Change in temperature due to a change in pressure at constant enthalpy (Slope of isenthalpic line)

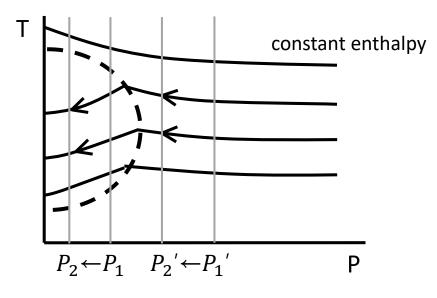
Joule – Thomson coefficient

 $\mu_{JT} = \frac{\partial T}{\partial P} \bigg|_{h}$

 \rightarrow Slope of isenthalpic line



Isenthalpic expansion of a real gas.



 $\begin{array}{ll} P_2 \leftarrow P_1 & : \text{Temperature \downarrow} & \text{if Pressure \downarrow} & \rightarrow \mu_{JT} > 0 \\ \\ P_2 \leftarrow P_1' & : \text{Temperature \uparrow} & \text{if Pressure \downarrow} & \rightarrow \mu_{JT} < 0 \end{array}$

Joule – Thomson coefficient

$$\mu_{JT} = \frac{\partial T}{\partial P} \bigg|_{h} = -\frac{\partial T}{\partial h} \bigg|_{P} \frac{\partial h}{\partial P} \bigg|_{T}$$

from basic thermodynamics (Van Wylen and Sonntag, 1976),

$$dh = \frac{\partial h}{\partial T} \Big|_{P} dT + \frac{\partial h}{\partial P} \Big|_{T} dP = C_{P} dT + \left[v - T \frac{\partial v}{\partial T} \right]_{P} dP$$
$$\rightarrow \frac{\partial h}{\partial T} \Big|_{P} = C_{P}, \qquad \frac{\partial h}{\partial P} \Big|_{T} = v - T \frac{\partial v}{\partial T} \Big|_{P}$$

Joule – Thomson coefficient

$$\mu_{JT} = \frac{\partial T}{\partial P} \bigg|_{h} = -\frac{\partial T}{\partial h} \bigg|_{P} \frac{\partial h}{\partial P} \bigg|_{T}$$
$$= \frac{1}{C_{P}} \left[T \frac{\partial v}{\partial T} \bigg|_{P} - v \right]$$

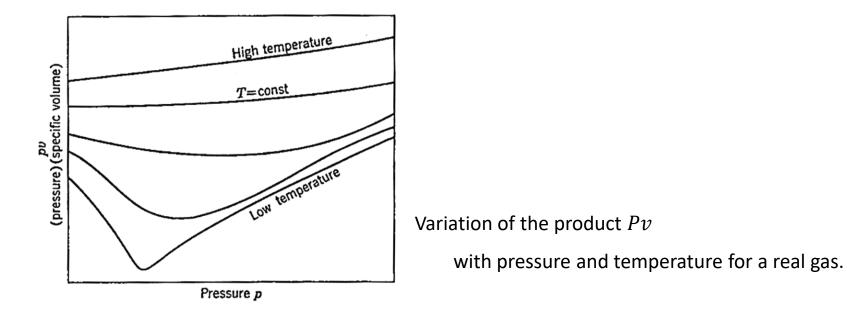
For an ideal gas,

$$\frac{\partial v}{\partial T} \bigg|_{P} = \frac{R}{P} = \frac{v}{T}$$
$$\mu_{JT} = \frac{1}{C_{P}} \left[T \frac{\partial v}{\partial T} \right]_{P} - v = 0$$

Joule – Thomson coefficient

from
$$h = u + Pv$$
,

$$\mu_{JT} = \frac{1}{C_P} \left[T \frac{\partial v}{\partial T} \right]_P - v = -\frac{1}{C_P} \left\{ \frac{\partial u}{\partial P} \right]_T + \frac{\partial (Pv)}{\partial P} \Big]_T$$



Cryogenic Engineering

Van der Waals gas

Equation of state (EOS)

$$\left(P + \frac{a}{v^2}\right)(v - b) = RT$$

Joule – Thomson coefficient for van der Waals gas

$$\mu_{JT} = \frac{(2a/RT)(1 - b/v)^2 - b}{C_P \left[1 - \left(\frac{2a}{vRT}\right)\left(1 - \frac{b}{v}\right)^2\right]}$$

for large value of the specific volume,

$$\mu_{JT} = \frac{1}{C_P} \left(\frac{2a}{RT} - b \right)$$

Inversion curve for van der Waals gas

The inversion curve is represented by all points at which the Joule-Thomson coefficient is zero.

$$\mu_{JT} = \frac{(2a/RT)(1 - b/v)^2 - b}{C_P \left[1 - \left(\frac{2a}{vRT}\right) \left(1 - \frac{b}{v}\right)^2 \right]} = 0 \quad \rightarrow (2a/RT)(1 - b/v)^2 - b = 0$$

Inversion temperature for van der Waals gas

$$T_{i} = \frac{2a}{bR} \left(1 - \frac{b}{v}\right)^{2}$$

$$T_{i} = \frac{2a}{bR} \left(1 - \frac{b}{v}\right)^{2}$$

$$T_{i,max} = \frac{2a}{bR} < T_{room}$$
 : He, H₂, Ne

 \rightarrow Cannot produce low T with expansion valve alone. (Require expander, turbine)

	Maximum Inversion Temperature		
Gas	ĸ	۴R	
Helium-4	45	81	
Hydrogen	205	369	
Neon	250	450	
Nitrogen	621	1118	
Air	603	1085	
Carbon monoxide	652	1174	
Argon	794	1429	
Oxygen	761	1370	
Methane	939	1690	
Carbon dioxide	1500	2700	
Ammonia	1994	3590	

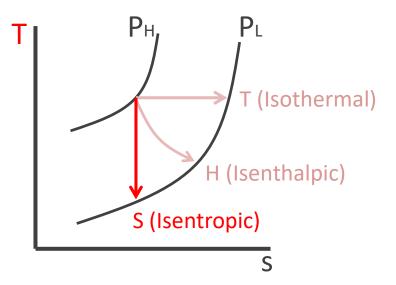
Table 3.2. Maximum inversion

3.4 Adiabatic Expansion

Work producing device : expansion engine (turbine)

Adiabatic Expansion : Most effective means of lowering T of the gas (X Adiabatic + Reversible = Isentropic!)

Problem : 2 Phase Mixture in an expander!



Isentropic expansion coefficient, μ_s

$$\mu_{s} = \left(\frac{\partial T}{\partial P}\right)_{s} = -\left(\frac{\partial T}{\partial s}\right)_{p} \left(\frac{\partial s}{\partial P}\right)_{T} = \frac{T}{C_{p}} \left(\frac{\partial v}{\partial T}\right)_{p}$$

$$C_{p} = \left(\frac{\partial h}{\partial T}\right)_{p} = \left(\frac{\partial Q}{\partial T}\right)_{p}, Q = TdS$$
Maxwell's Relation $-\left(\frac{\partial S}{\partial P}\right)_{T} = \left(\frac{\partial V}{\partial T}\right)_{P}$

$$\left(\frac{\partial T}{\partial V}\right)_{s} = -\left(\frac{\partial p}{\partial S}\right)_{V}$$

$$\left(\frac{\partial T}{\partial p}\right)_{s} = +\left(\frac{\partial V}{\partial S}\right)_{p}$$

$$\left(\frac{\partial S}{\partial V}\right)_{T} = +\left(\frac{\partial p}{\partial T}\right)_{V}$$

• Isentropic expansion coefficient, μ_s

$$\mu_{s} = \left(\frac{\partial T}{\partial P}\right)_{s} = -\left(\frac{\partial T}{\partial s}\right)_{p} \left(\frac{\partial s}{\partial P}\right)_{T} = \frac{T}{C_{p}} \left(\frac{\partial v}{\partial T}\right)_{P}$$

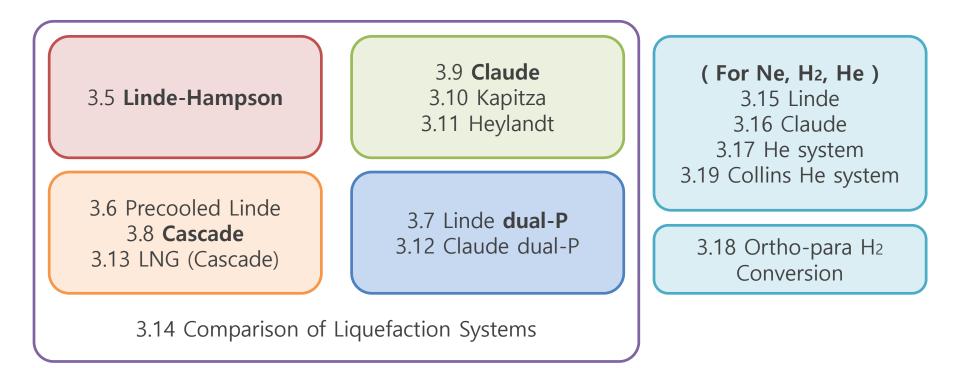
Volume expansion coefficient
$$\beta = \frac{1}{v} \left(\frac{\partial v}{\partial T} \right)_{P} \implies \left[= \frac{T}{C_{p}} \beta v \right]$$

For Ideal gas : Pv = RT $\begin{pmatrix} \frac{\partial v}{\partial T} \end{pmatrix}_{p} = \frac{R}{p} = \frac{v}{T} \quad \Rightarrow \quad = \frac{v}{C_{p}}$ For Van der Waals gas : $(P + \frac{a}{v^{2}})(v - b) = RT$ $= \frac{v(1 - \frac{b}{v})}{C_{p}[1 - (\frac{2a}{vRT})(1 - \frac{b}{v})^{2}]}$

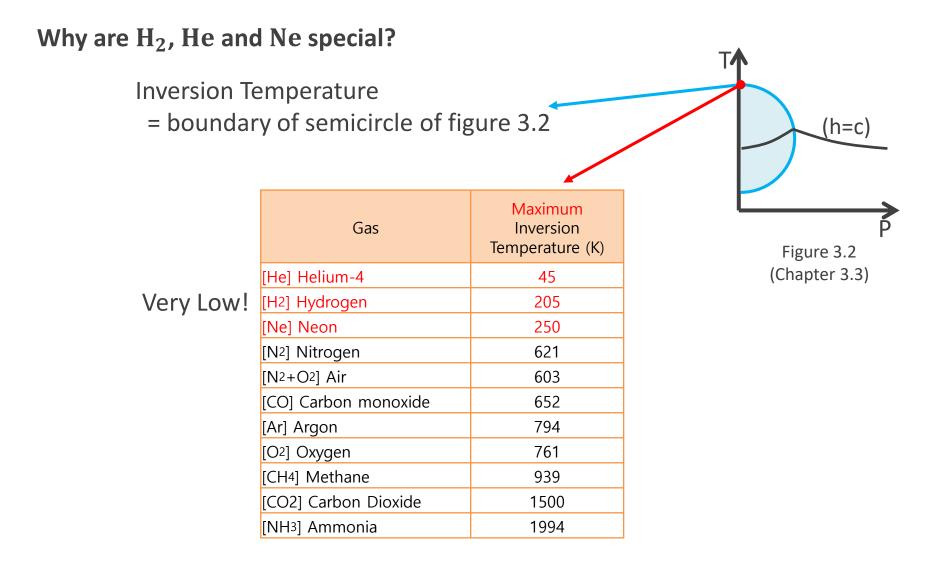
3.4 Adiabatic Expansion

- Methods of cooling
 - External work method : energy is removed as external work
 - Internal work method (Expansion Valve) : do not remove energy from gas but moves molecules farther apart

= A big purpose for learning cryogenic systems! (Korea's energy supply is heavily dependent on import by ships!)



3.4 Adiabatic Expansion



3.4 Adiabatic Expansion

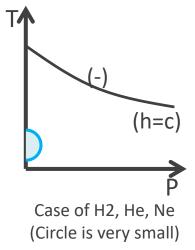
Why are H₂, He and Ne special?

In semicircle, $\mu JT > 0$ = Isenthalpic expansion (J-T Valve) is cooling!

Out of semicircle, $\mu JT < 0$ = Isenthalpic expansion (J-T valve) is heating!

= We cannot make liquid by J-T valve!

= To make liquid, we must applicate expander or dualpressure or pre-cooling.



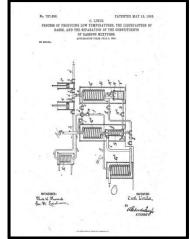
History

The most simple gas liquefaction/separation system.

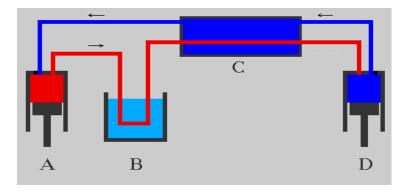
This system is based on the Joule-Thomson effect.

Chapter 3.3. When a (non-ideal) gas expands, it will cold down (below inversion temperature). This is why our whistle is colder than our body temperature!

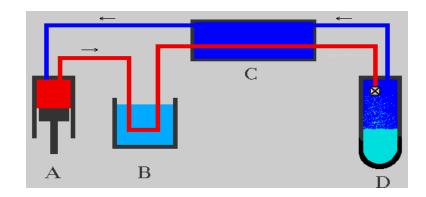
William Hampson and Carl von Linde independently filed for patent of the cycle in 1895.







- A : Compression
- B : Cooling
- C: Cooling (HX)
- D : Expansion
- C: Heating (HX)



- A : Compression
- B: Cooling
- C: Cooling (HX)
- D : Joule-Thompson Orifice

(Reservoir + Expansion Valve)

C : Heating (HX)

History

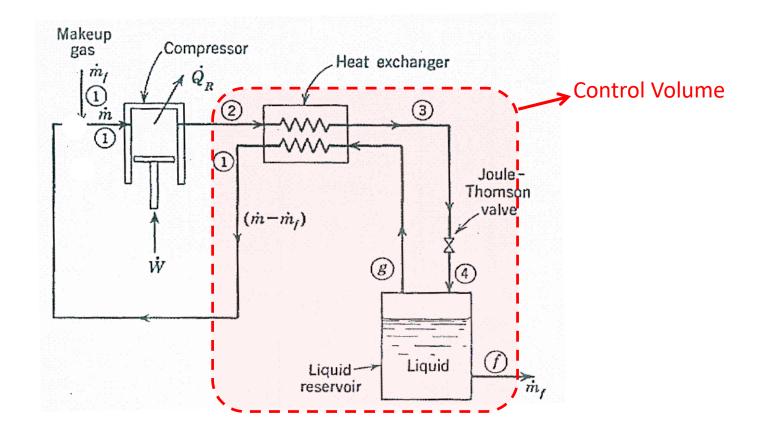
Siemens cycle has poor efficiency and is only efficient for high temperature gases. But Linde-Hampson cycle can easily collect liquid and cool dramatically.

Heike Kamerlingh Onnes made 'liquid helium' (1908) by this Linde-Hampson system and found 'super-conductivity'.

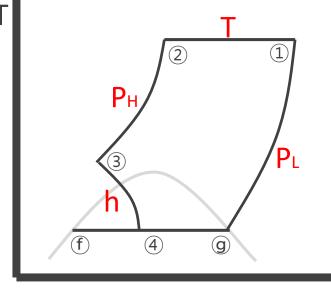
And He got Nobel prize in Physics (1913)!



Diagram



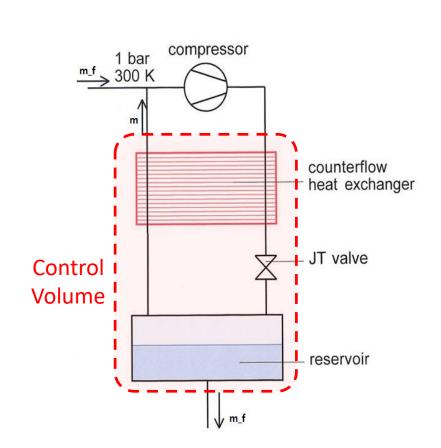
Diagram



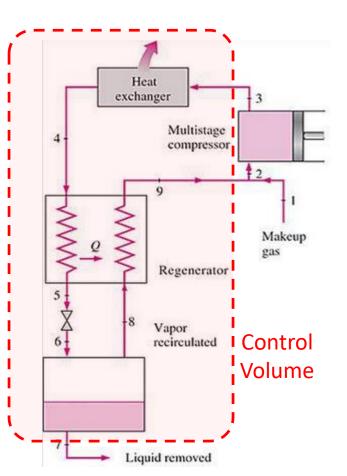
Normal room : T₁=300K, P₁=1bar Liquid N₂ : T₄=77K

Constant T : Isothermal Constant P : Isobaric Constant h : Isenthalpic

S



Another form of diagrams



Diagram

Diagram

Assumption

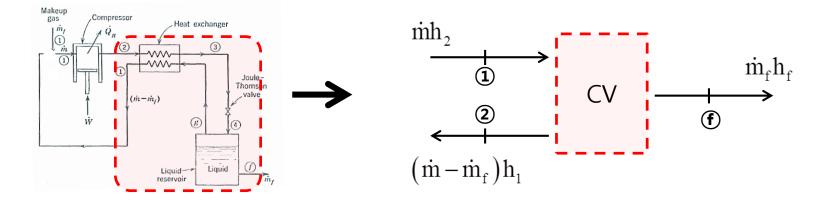
- Reversible pressure drop
- No heat in-leak
 - (Reversible isothermal process in compressor)
- 100% effective heat exchanger

Quality
$$x = \frac{(mass of sat. gas)}{(total mass)}$$

(X Volume of gas is very larger than liquid)

Questions

Q1. We can choose P2 in the system. Which is best P2?



At control volume, $\dot{m}h_2 = (\dot{m} - \dot{m}_f)h_1 + \dot{m}h_f$

Liquid yield
$$y = \frac{m_f}{m} = \frac{h_1 - h_2}{h_1 - h_f}$$

: Fraction of gas flow that is liquefied

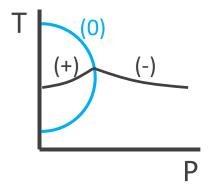
Questions

Best P2 is
$$\frac{\partial y}{\partial P_2} = 0$$

 \rightarrow (h1, hf is fixed value) $\frac{\partial h_2}{\partial P_2} = 0$
 $\rightarrow \frac{\partial h}{\partial P} = -\mu_{JT}C_P$

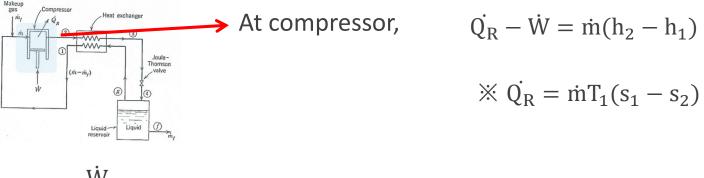
$$\rightarrow$$
 (C_P is positive number) $\mu_{JT} = 0$

 μ_{JT} is slope at T-P diagram.
 So, semicircle in T-P graph means best P₂! (Inversion Curve) (Chapter 3.3)



Questions

Q2. How much works do we need for running this system?



$$-\frac{\dot{W}}{\dot{m}} = T_1(s_1 - s_2) - (h_1 - h_2)$$
$$-\frac{\dot{W}}{\dot{m}_f} = -\frac{\dot{W}}{\dot{m}y} = (\frac{h_1 - h_f}{h_1 - h_2})[T_1(s_1 - s_2) - (h_1 - h_2)]$$

Questions

Q3. What will happen in real system without assumption?

Reversible pressure drop \rightarrow P₂ will be lower.

No heat in-leak Reversible isothermal process in compressor \rightarrow Q_{elec} will be added. And y will be lower.

 $\dot{m}h_2 + \mathbf{Q}_{elec} = (\dot{m} - \dot{m}_f)h_1 + \dot{m}h_f$

3.5 Simple Linde-Hampson System

Questions

100% effective heat exchanger

→ Temperature difference at both side will be lower.
 To keep temperature difference, we need input more pump work.
 (To enlarge mass flow rate)

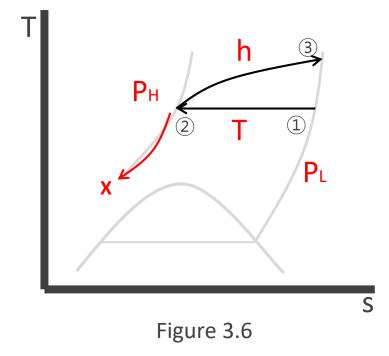


3.5 Simple Linde-Hampson System

Reason 1.

It's not for Ne/H2/He! ----

We will learn later... (Chapter 3.15~3.19)



 Gas
 M.I.T. (K)

 [He] Helium-4
 45

 [H2] Hydrogen
 205

 [Ne] Neon
 250

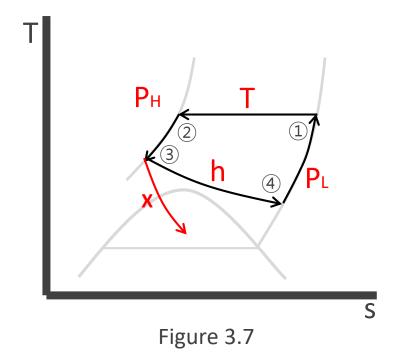
Maximum inversion temp. << room temp.

- \rightarrow Their expansion = heating!
- \rightarrow Gas in HX warmed rather than cooled!

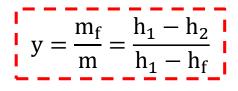
3.5 Simple Linde-Hampson System

It's not for Ne/H2/He! ----

We will learn later... (Chapter 3.15~3.19)

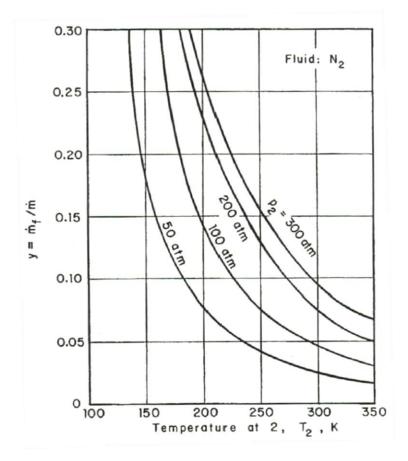


Reason 2.

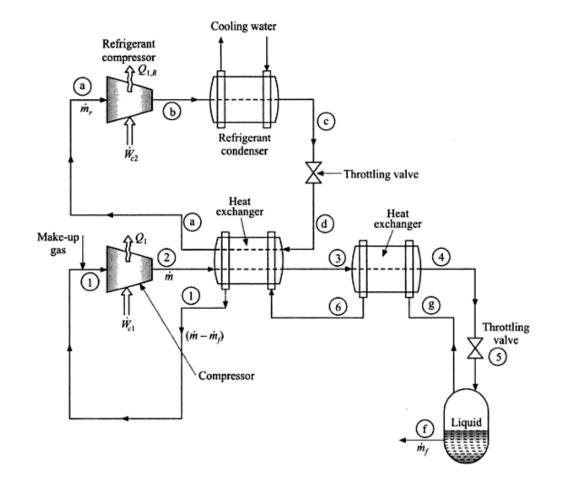


Liquid yield(y) is negative. (h1 < h2)

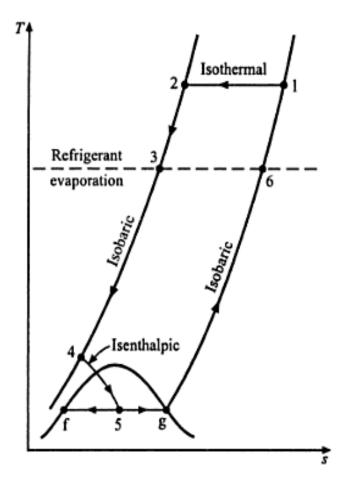
 \rightarrow Even if we could attain low temperature, no gas would be liquefied.



Liquid yield versus compressor temperature for a Linde-Hampson system using nitrogen as the working fluid It is apparent that the performance of a Linde-Hampson system could be improved if the gas entered the heat exchanger at a temperature
 lower than ambient temperature



Precooled Linde-Hampson system



Precooled Linde-Hampson system T-S diagram

Liquid yield

Applying the First Law for steady flow to the heat exchanger, the two liquid receivers, and the two expansion valves.

$$\dot{m}h_{2} + \dot{m}_{r}h_{d} = (\dot{m} - \dot{m}_{f})h_{1} + \dot{m}_{r}h_{a} + \dot{m}_{f}h_{f}$$

$$y = \frac{\dot{m}_{f}}{\dot{m}} = \frac{h_{1} - h_{2}}{h_{1} - h_{f}} + r\frac{h_{a} - h_{d}}{h_{1} - h_{f}}$$

 $r = \dot{m}_r / \dot{m}$

r is the refrigerant mass flow rate ratio $\dot{m_r}$ is the mass flow rate of the auxiliary refrigerant \dot{m} is the total mass flow rate through the high pressure compressor

Liquid yield

The second term of liquid yield represents the improvement in liquid yield that is obtained through the use of precooling.

$$y = \frac{\dot{m}_{f}}{\dot{m}} = \frac{h_{1} - h_{2}}{h_{1} - h_{f}} + r\frac{h_{a} - h_{d}}{h_{1} - h_{f}}$$

Limit of the liquid yield 1

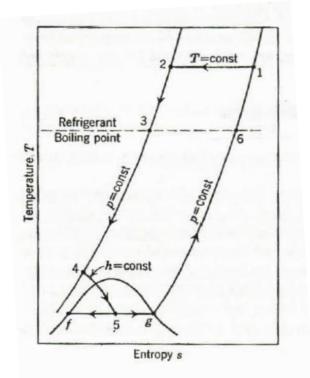
From the Second Law of Thermodynamics, T_3 and T_6 cannot be lower than the boiling point of the auxiliary refrigerant at point d

- Limit of the liquid yield 1
 - : The maximum liquid yield

With a suitable value of the refrigerant flow-rate ratio r, liquid yield could have a value of 1, which means 100 percent for the liquid yield.

$$\mathbf{y}_{\max} = \frac{\mathbf{h}_6 - \mathbf{h}_3}{\mathbf{h}_6 - \mathbf{h}_f}$$

- Limit of the liquid yield 1
 - : The maximum liquid yield

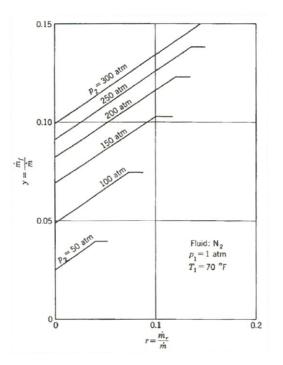


 h_3 and h_6 are taken at the temperature of the boiling refrigerant at point d)

$$\mathbf{y}_{\max} = \frac{\mathbf{h}_6 - \mathbf{h}_3}{\mathbf{h}_6 - \mathbf{h}_f}$$

Precooled Linde-Hampson cycle

Limit of the liquid yield 2



If the refrigerant flow rate ratio were too large, the liquid at point d would not be completely vaporized, and liquid would enter the refrigerant compressor.

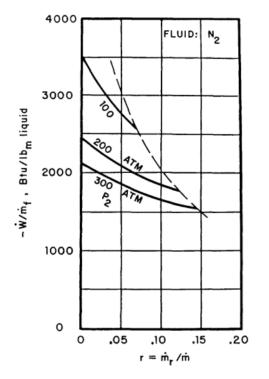
Liquid yield versus refrigerant flow rate ratio for the precooled Linde-Hampson system using nitrogen as the working fluid

The work requirement

- If the main compressor is reversible and isothermal and the auxiliary compressor is reversible and adiabatic.
- The last term represents the additional work requirement for the auxiliary compressor. (usually on the order of 10 percent of the total work)

$$-\frac{W}{\dot{m}} = T_1(s_1 - s_2) - (h_1 - h_2) + r(h_b - h_a)$$

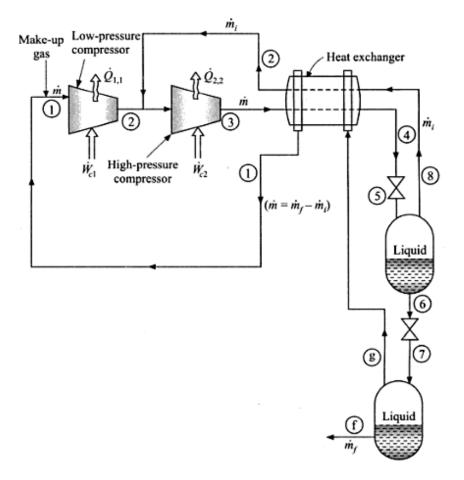
The work requirement



The increase in liquid yield more than offsets the additional work requirement, however, so that the work requirement per unit mass of gas liquefied is actually less for the precooled system than for the simple system.

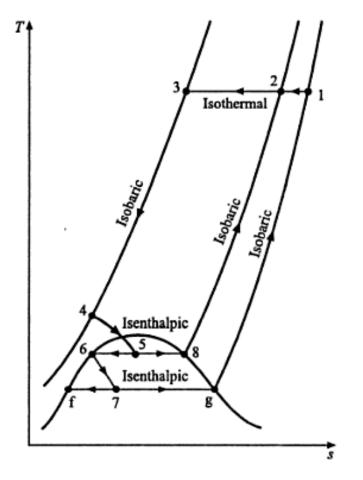
Work required to liquefy a unit mass of nitrogen in a precooled Linde-Hampson system

3.7 Linde Dual-Pressure Cycle



Linde dual-pressure system

3.7 Collins Helium-Liquefaction System



Linde dual-pressure system T-S diagram

Liquid yield

Applying the First Law for steady flow to the heat exchanger, the two liquid receivers, and the two expansion valves.

$$\dot{m}h_3 = \dot{m}_fh_f + \dot{m}_ih_2 + (\dot{m} - \dot{m}_i - \dot{m}_f)h_1$$

$$y = \frac{\dot{m}_{f}}{\dot{m}} = \frac{h_{1} - h_{3}}{h_{1} - h_{f}} - i\frac{h_{1} - h_{2}}{h_{1} - h_{f}}$$

 $i = \dot{m}_i / \dot{m}$

i is the intermediate pressure stream flow rate ratio \dot{m}_i is the mass flow rate of the intermediate pressure stream at point 8 \dot{m} is the total mass flow rate through the high pressure compressor

3.7 Linde Dual-Pressure Cycle

Liquid yield

- This modification reduces the liquid yield somewhat.
- The second term of liquid yield represents the reduction in the liquid yield below that of the simple system because of splitting the flow at the intermediate pressure liquid receiver.

$$y = \frac{\dot{m}_{f}}{\dot{m}} = \frac{h_{1} - h_{3}}{h_{1} - h_{f}} - i\frac{h_{1} - h_{2}}{h_{1} - h_{f}}$$

The work requirement

Applying the First Law for steady flow to the two compressors.

$$(\dot{Q}_{R1} - \dot{W}_{C1}) + (\dot{Q}_{R2} - \dot{W}_{C2}) = \dot{m}h_3 - (\dot{m} - \dot{m}_i)h_1 - \dot{m}_ih_2$$

$$\dot{Q}_{R1} = -(\dot{m} - \dot{m}_i)T_1(s_1 - s_2)$$

 $\dot{Q}_{R2} = -\dot{m}T_1(s_2 - s_3)$

$$-\frac{\dot{W}}{\dot{m}} = \left[T_{1}\left(s_{1}-s_{3}\right)-\left(h_{1}-h_{3}\right)\right]-i\left[T_{1}\left(s_{1}-s_{2}\right)-\left(h_{1}-h_{2}\right)\right]$$

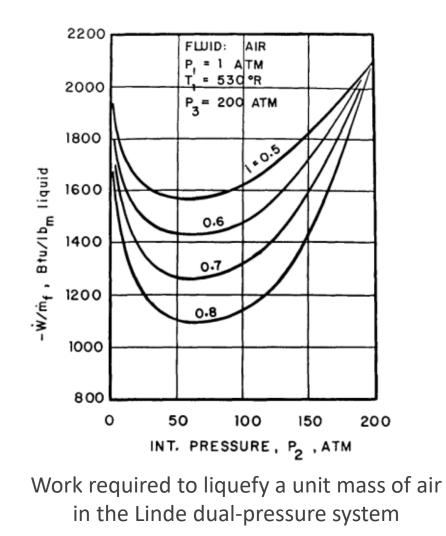
The work requirement

This modification reduces the total work required.

The work requirement is reduced below that of the simple system by the amount given by the second bracketed term.

$$-\frac{W}{\dot{m}} = \left[T_1(s_1 - s_3) - (h_1 - h_3)\right] - i\left[T_1(s_1 - s_2) - (h_1 - h_2)\right]$$

3.7 Linde Dual-Pressure Cycle

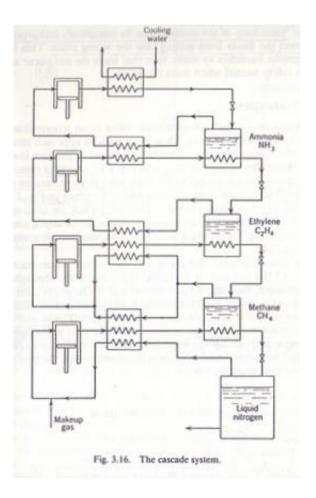


3.7 Linde Dual-Pressure Cycle

Optimal intermediate pressure

- As 'Work required to liquefy a unit mass of air in the Linde dualpressure system' shows, there is an optimum intermediate pressure p₂ for a given intermediate stream mass flow rate ratio, which makes the work requirements a minimum.
- Typical air liquefaction plants operate with i = 0.8, $p_3 = 200$ atm, p_2 between 40 and 50 atm.

3.8 Cascade System



- The cascade system is an extension of the precooled system
- There are refrigeration system chain of ammonia ethylene methane -nitrogen
- From a thermodynamic point of view, the cascade system in very desirable for liquefaction because it approaches the ideal reversible system more closely than any other discussed thus far0

3.8 Cascade System

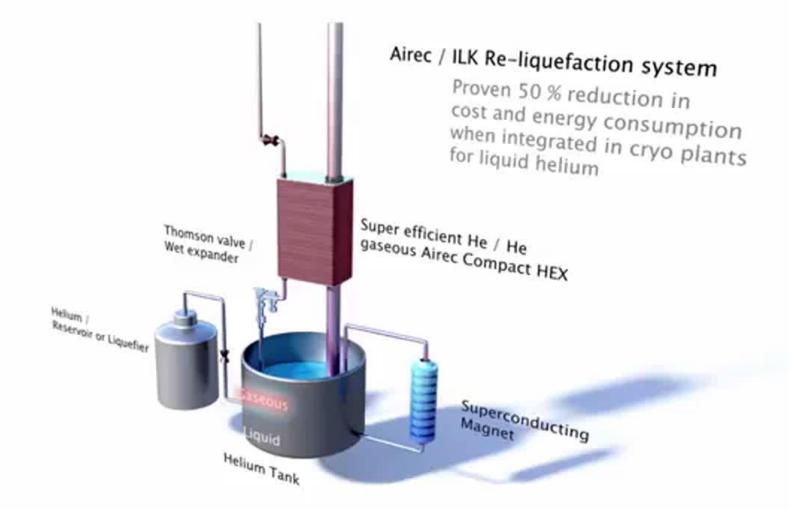
$$\dot{m}_{f}h_{f} + (\dot{m} - \dot{m}_{f})h_{1} - \dot{m}h_{2} + \sum_{i=1}^{n_{comp}} \dot{m}_{x,i}(h_{a,i} - h_{b,i}) + \dot{m}_{x,n}(h_{b,n} - h_{c,n}) = 0$$

Liquid yield :
$$y = \frac{h_1 - h_2}{h_1 - h_f} + \sum_{i=1}^{n_comp} x_i \frac{h_{i,i} - h_{e,i}}{h_1 - h_f} + \dot{m}_{x,n} (h_{b,n} - h_{c,n})$$

$$-\dot{W} = \dot{m}[(h_2 - h_1) - T_1(h_2 - h_1)]h_f + \sum_{i=1}^{n_comp} \dot{m}_{x,i}(h_{e,i} - h_{i,i})$$

Power per mass flow :
$$-\frac{\dot{W}}{\dot{m}} = [(h_2 - h_1) - T_1(h_2 - h_1)]h_f + \sum_{i=1}^{n_comp} x_i(h_{e,i} - h_{i,i})$$

Cryogenic System Video



Cryogenic System Video



Cryogenic System Video

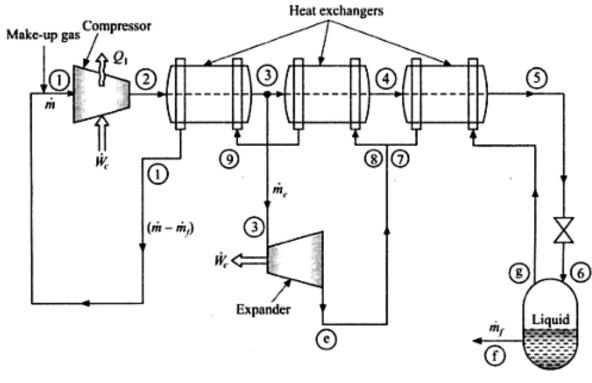
Industrial Refrigeration Basics **Refrigeration** Cycle Where To Find Industrial Refrigeration Systems TheEngineeringMindset.com



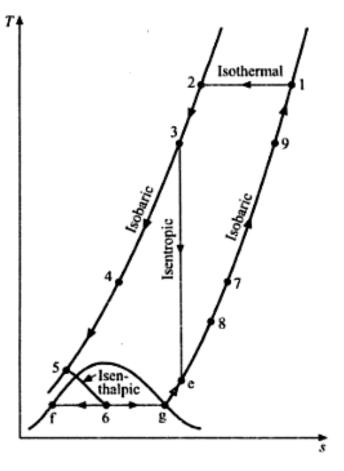


Georges Claude(1870-1960)

In 1902 Claude devised what is now known as the Claude system for liquefying air. The system enabled the production of industrial quantities of liquid nitrogen, oxygen, and argon



Claude system



Claude system T-S diagram

Valve VS Expansion engine or expander

- The expansion through an expansion value is an irreversible process. Thus if we wish to approach closer to the ideal performance, we must seek a better process.
- If the expansion engine is reversible and adiabatic, the expansion process in isentropic, and a much lower temperature is attained than an isenthalpic expansion.

Why we can not eliminate the expansion valve?

The expansion valve could not be eliminated because of the problem of two-phase flow within the engine cylinder or turbine blade flow passages.

Liquid yield

Applying the First Law for steady flow to the heat exchangers, the expansion valve, and the liquid receiver as a unit, for no external heat transfer.

$$\dot{m}h_2 + \dot{m}_eh_e = (\dot{m} - \dot{m}_f)h_1 + \dot{m}_eh_3 + \dot{m}_fh_f$$

$$y = \frac{\dot{m}_{f}}{\dot{m}} = \frac{h_{1} - h_{2}}{h_{1} - h_{f}} + x \frac{h_{3} - h_{e}}{h_{1} - h_{f}}$$

$$x = \dot{m}_e / \dot{m}$$

x is the fraction of the total flow that passes through the expander $\dot{m_e}$ is the mass flow rate of fluid through expander \dot{m} is the total mass flow rate through the high pressure compressor

Liquid yield

The second term of liquid yield represents the improvement in performance over the simple Linde Hampson system.

$$y = \frac{\dot{m}_{f}}{\dot{m}} = \frac{h_{1} - h_{2}}{h_{1} - h_{f}} + x \frac{h_{3} - h_{e}}{h_{1} - h_{f}}$$

The work requirement

The work requirement per unit mass compressed is exactly the same as that for the Linde Hampson system if the expander work is not utilized to help in the compression.

$$-\frac{\dot{W}}{\dot{m}} = T_1(s_1 - s_2) - (h_1 - h_2)$$

The work requirement per unit mass compressed for the Linde Hampson system

The work requirement

If the expander work is used to aid in compression, then the net work requirement is given by

$$-\frac{\dot{W}_{e}}{\dot{m}} = -\frac{\dot{W}_{e}}{\dot{m}} - \frac{\dot{W}_{e}}{\dot{m}}$$
$$\frac{\dot{W}_{e}}{\dot{m}} = x \left(h_{3} - h_{e} \right)$$
$$\frac{\dot{W}_{c}}{\dot{m}} = \left(h_{2} - h_{1} \right) + T_{1} \left(s_{1} - s_{2} \right)$$

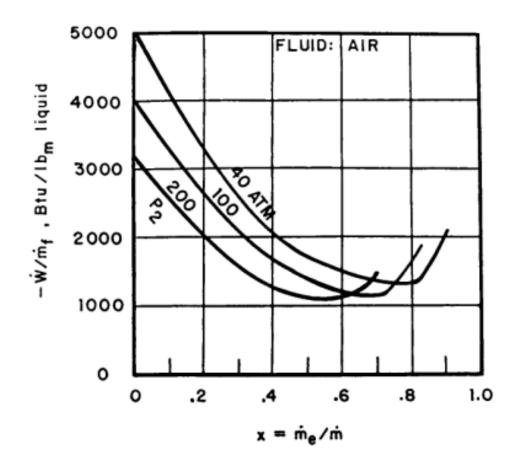
The net work is given by

$$-\frac{\dot{W}}{\dot{m}} = \left[T_{1}(s_{1}-s_{2})-(h_{1}-h_{2})\right]-x(h_{3}-h_{e})$$

The work requirement

The last term is the reduction in energy requirements due to the utilization of the expander work output.

$$-\frac{W}{\dot{m}} = \left[T_{1}(s_{1}-s_{2})-(h_{1}-h_{2})\right]-x(h_{3}-h_{e})$$

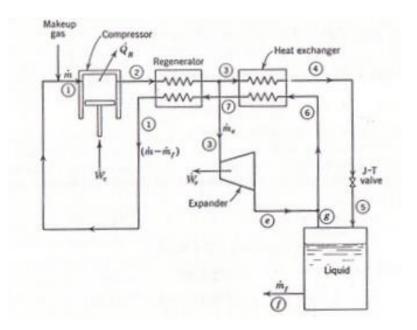


Work required to liquefy a unit mass of air in the Claude system

Smallest work requirement per unit mass liquefied

- There is a finite temperature at point 3 that will yield the smallest work requirement per unit mass liquefied.
- As the high pressure is increased, the minimum work requirement per unit mass liquefied decreases.

3.10 Kapitza System



- Modified Claude system which eliminate the third low temperature heat exchanger
- A rotary expansion engine was used instead of a reciprocating expander
- Kapitza system usually operated at relatively low pressures-on the order of 700kPa

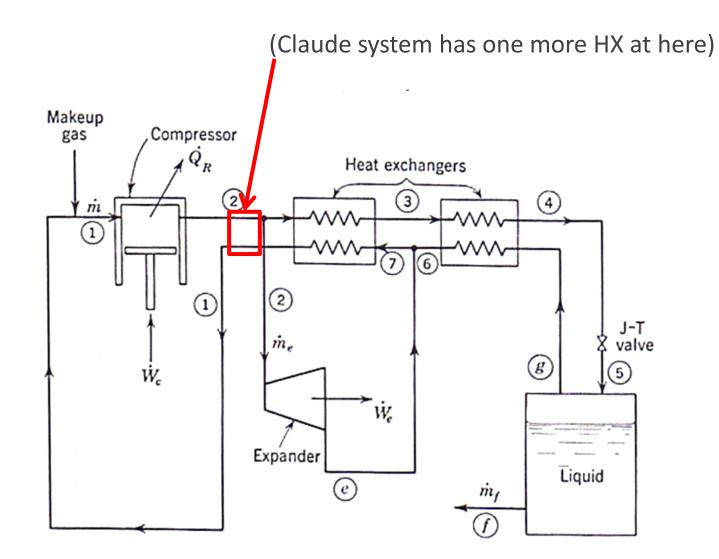
$$\dot{m}_{f}h_{f} + (\dot{m} - \dot{m}_{f})h_{1} - \dot{m}h_{2} + \dot{m}_{x}(h_{3} - h_{e}) = 0$$

$$\label{eq:Liquid yield} \mbox{Liquid yield}: y = \frac{h_1 - h_2}{h_1 - h_f} + x \frac{h_3 - h_e}{h_1 - h_f}$$

$$-\dot{W} = \dot{m}[(h_2 - h_1) - T_1(h_2 - h_1)]h_f + \dot{m}_x(h_e - h_3)$$

Power per mass flow : $-\frac{\dot{W}}{\dot{m}} = [(h_2 - h_1) - T_1(h_2 - h_1)]h_f + x(h_e - h_3)$

3.11 Heylandt System



For a high pressure (Approximately 20 MPa ≒ 200 atm) For an expansion-engine (flow-rate ratio of approximately 0.6)

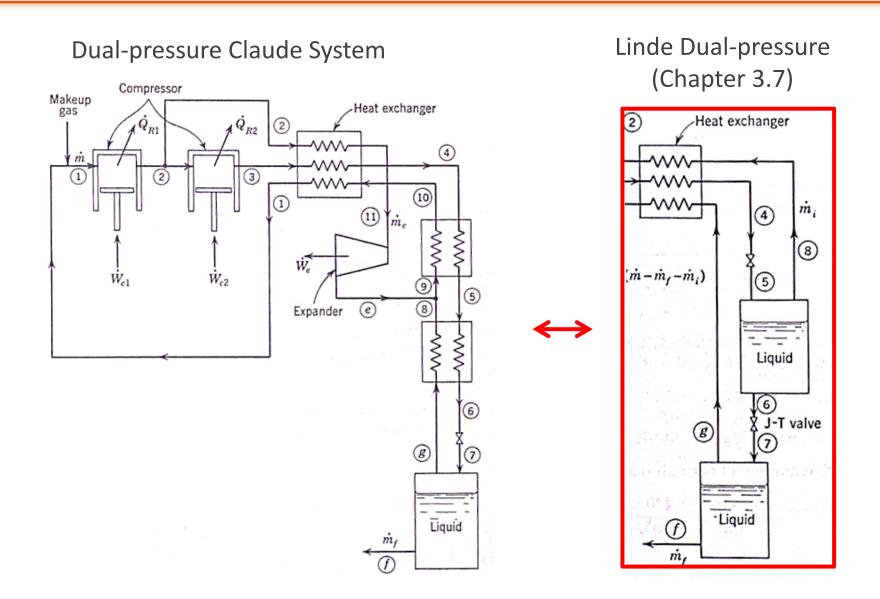
The optimum value of T before expansion = near ambient T
 → So, it can eliminate first HX in the Claude system by compressing!
 (∴ Modified Claude system)

Advantage

→ The lubrication problems in the expander are easy to solve! (Because T of expander is very low)

※ Contribution of expander and expansion valve is nearly equal.(At original Claude system, expander makes more contribution)

3.12 Other Liquefaction Systems Using Expanders



... It is similar to the Linde dual-pressure system. (Chapter 3.7) (A reservoir is replaced by expander and two HX)

Advantage

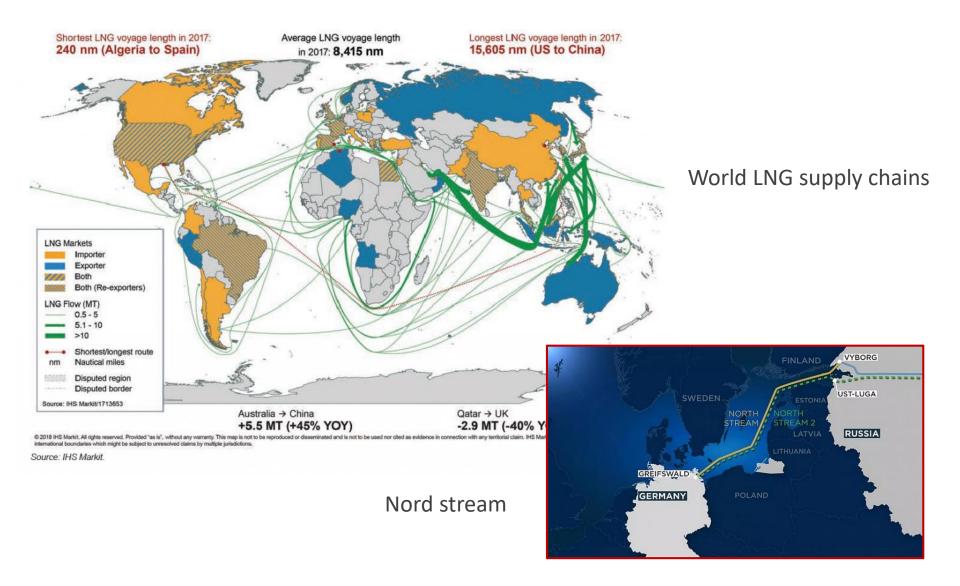
Gas through expander is compressed to some intermediate P. \rightarrow Work requirement per unit mass of gas liquid is reduced.

X If nitrogen compressed from 1 atm to 35 atm, optimum performance is attained when 75 percent of flow diverted through the expander.

3.13 Liquefaction Systems for LNG



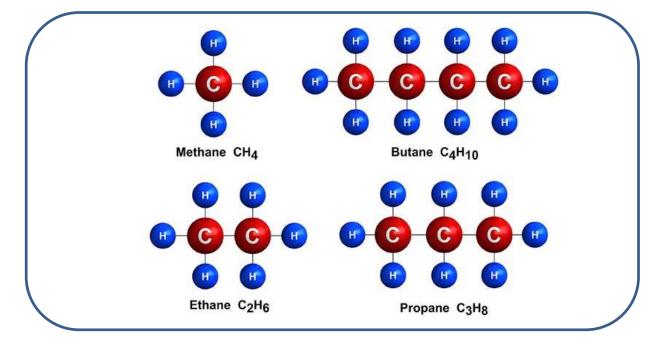
3.13 Liquefaction Systems for LNG



3.13 Liquefaction Systems for LNG

LNG liquefaction

Typical 4 Natural Gases

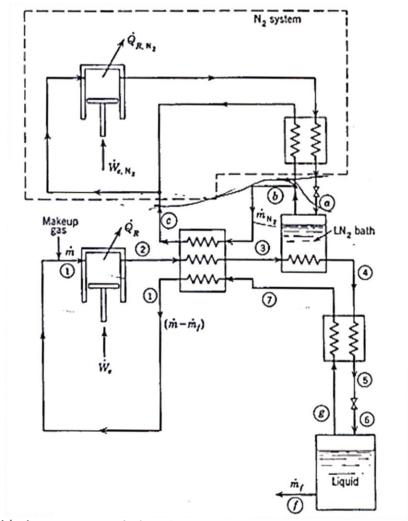


- Condense at different temperature levels.
- LNG deposits are normally found deeper position than oil
- Deepest deposits can be made of pure LNG
- Require mixed refrigerant liquefaction system.

The Linde-Hampson system is desirable for small-scale liquefaction plants.

However, the basic Linde-Hampson system with no precooling would not work for neon, hydrogen, or helium

Because the maximum inversion temperature for these gases is below ambient temperature. So, it is normally using for air separation



Liquid-nitrogen-precooled Linde-Hampson system for neon or Hydrogen

$$y = \frac{h_7 - h_4}{h_7 - h_f}$$

$$0 = \dot{m}_{N_2} h_c + (\dot{m} - \dot{m}_f) h_1 + \dot{m}_f h_f - \dot{m}_{N_2} h_a - \dot{m} h_2$$

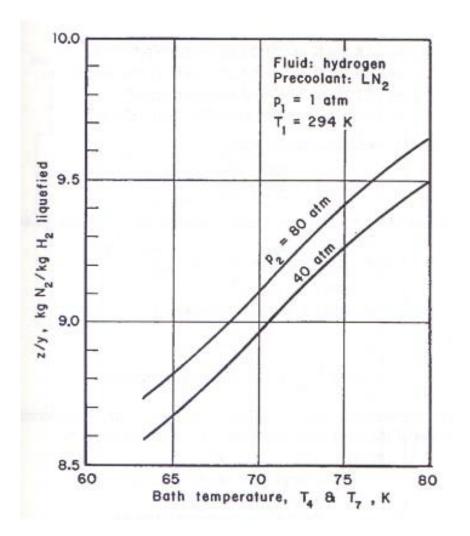
- \dot{m}_{N_2} : the mass flow rate of liquid nitrogen boiled away to precool the incoming hydrogen or neon
- m : the mass flow rate of hydrogen or neon through the compressor
- $\dot{m}_{\rm f}~$: the mass flow rate of hydrogen or neon which is liquefied

The nitrogen boil-off rate per unit mass of hydrogen or neon compressed

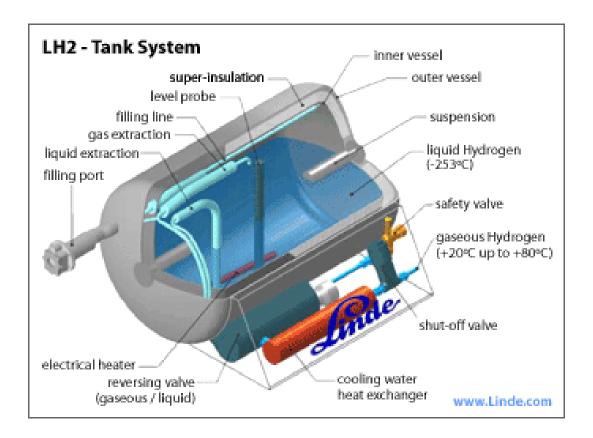
$$z = \frac{\dot{m}_{N_2}}{\dot{m}}$$

$$z = \frac{h_2 - h_1}{h_c - h_a} + y \frac{h_1 - h_f}{h_c - h_a}$$

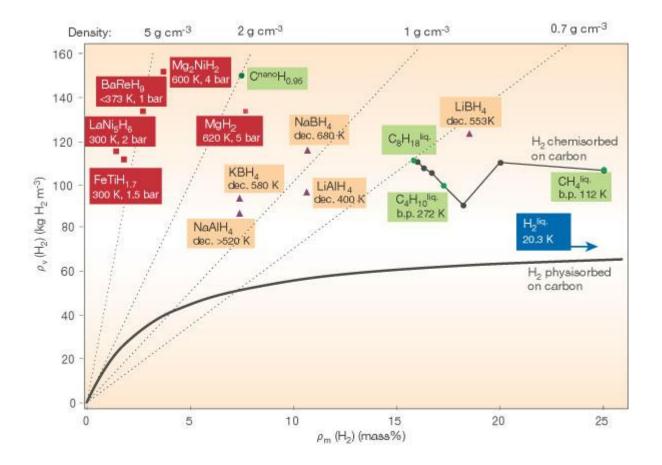
$$\frac{\dot{m}_{N_2}}{\dot{m}_f} = \frac{\dot{m}_{N_2} / \dot{m}}{\dot{m}_f / \dot{m}} = \frac{z}{y}$$



Nitrogen boil-off per unit mass of hydrogen produced for the liquid-nitrogen-precooled Linde-Hampson system as a function of the liquid-nitrogen bath temperature.

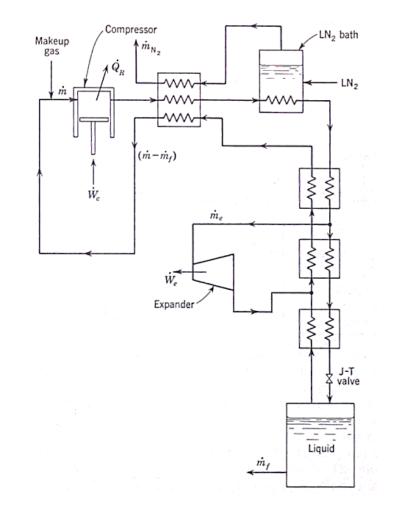


Liquid hydrogen storage tank system, horizontal mounted with double gasket and dual seal



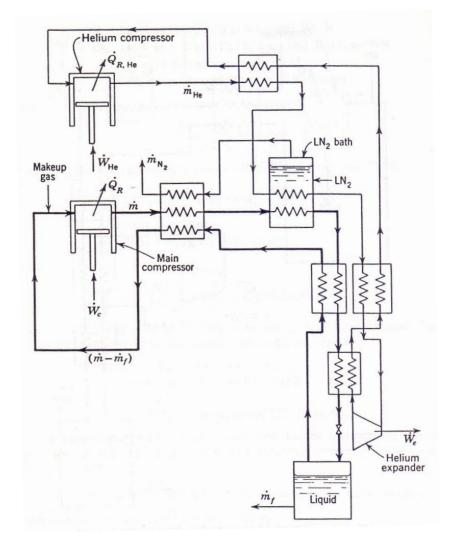
A complete survey plot of hydrogen storage in metal hydrides and carbon-based materials

3.16 Claude System for Hydrogen or Neon



Precooled Claude system for hydrogen or neon

3.17 Helium-Refrigerated Hydrogen-Liquefaction System



Helium-refrigerated hydrogen liquefaction system

3.17 Helium-Refrigerated Hydrogen-Liquefaction System

- Advantage
 - Relatively low pressures can be used
 - The compressor size can be reduced

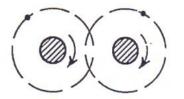
(although two compressors are required)

- The pipe thickness can be reduced
- The hydrogen or neon need be compressed only to a pressure high enough to overcome the irreversible pressure drops through the heat exchangers and piping in an actual system

N.B.P. = 20.3K

Types of hydrogen molecules

ortho $- H_2$ (Spins aligned, high energy) para $- H_2$ (Spins aligned, high energy)



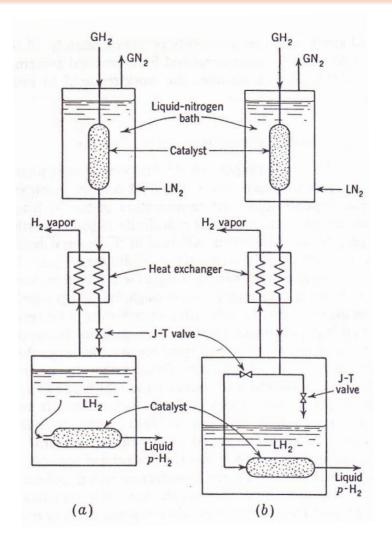
Ortho-hydrogen

Para-hydrogen

ortho – $H_2 \rightarrow para - H_2 + \Delta \alpha$ (heat of conversion) 70.3 kJ/kg Latent heat 44.3 kJ/kg

At high temperature is a mixture of 75% ortho - H₂ and 25% para - H₂ As temperature is cooled to the normal boiling point of hydrogen, the ortho - H₂ concentration decreases from 75 to 0.2% The changeover from ortho – to para – hydrogeninvolves a heat of conversion that is greater than the heat of vaporization of para – hydrogen

Therefore serious boil-off losses will result unless measures are taken to prevent it.



Ortho-para-hydrogen conversion arrangements

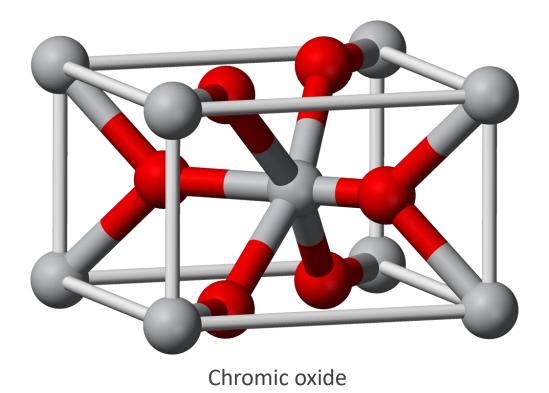
Catalysts

hydrous ferric oxide



Nickel-based catalyst

Catalysts

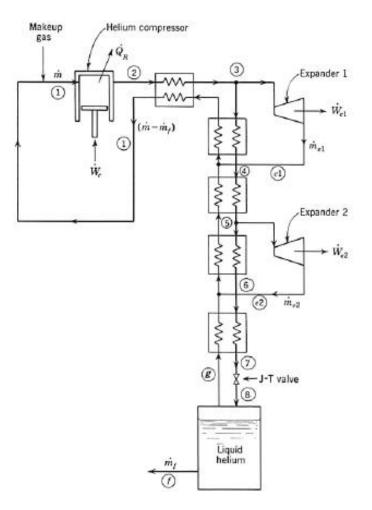




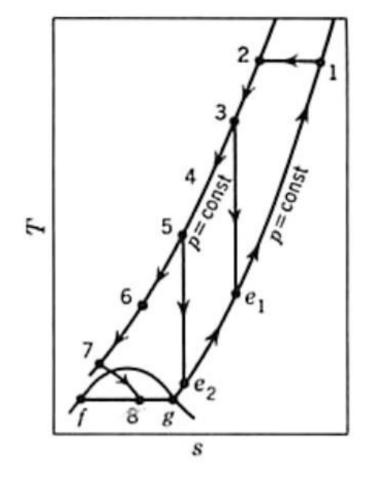
Helium refrigerator-liquefier at MIT

Samuel C. Collins (MIT) designed Helium liquefier (1952)

It is an extension of the Claude system



Collins helium-liquefaction system



Collins helium-liquefaction system T-S diagram

Liquid yield

Consisting of all components except the helium compressor and the expansion engines

$$y = \frac{h_1 - h_2}{h_1 - h_f} + x_1 \frac{\Delta h_{e_i}}{h_1 - h_f} + x_2 \frac{\Delta h_{e_2}}{h_1 - h_f}$$

$$\begin{split} x_1 &= \dot{m}_{e_1}/\dot{m} \\ x_2 &= \dot{m}_{e_2}/\dot{m} \\ \Delta h_{e_i} &= \text{enthalpy change of fluid passing through expander 1} \\ \Delta h_{e_2} &= \text{enthalpy change of fluid passing through expander 2} \\ \dot{m}_{e_1}, \dot{m}_{e_2} &= \text{mass flow rates of fluid through expander 1 and 2, respectively} \end{split}$$

Liquid yield

For more than two engines?

→ Additional term similar to the second term for each expander would be added!

Increase the Collins liquefier performance?

- By using liquid-nitrogen precooling bath
 - \rightarrow Cool-down time would be reduced from 4h to 2h
 - ightarrow Liquid yield can be tripled

But,

Pre coolant bath is not required because it doesn't dependent solely on the Joule-Thomson effect!

3.20 Simon Helium-Liquefaction System

System Characteristics

- Liquefy small quantities of helium
- Not operate as a steady-flow system
- It produces helium in a batch process

3.20 Simon Helium-Liquefaction System

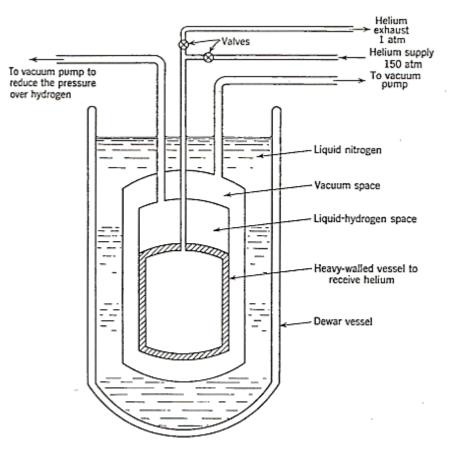
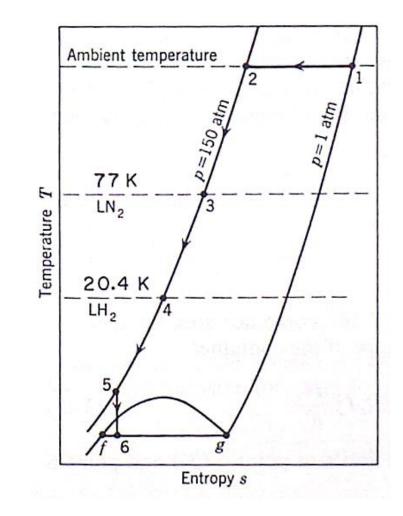


Fig. 3.30. Simon helium-liquefaction system.

Simon helium-liquefaction system

3.20 Simon Helium-Liquefaction System



Simon helium-liquefaction system T-S diagram

Simon helium-liquefaction system process

Process 1-2 : Helium gas is introduced into the heavy-wall container. (Pressure : 10~15 MPa, Ambient temperature)

Process 2-3 : • Liquid Nitrogen is introduced into the enclosing bath.

- Cools the entire container and contents (77K)
- Vacuum space is filled with helium gas. (act as heat transfer medium)
- At the completion of the process, inner vessel is thermally isolated

Simon helium-liquefaction system process

- Process 3-4 : Liquid hydrogen is introduced into the upper part of the inner container. Inner container and contents are further cooled to liquid-hydrogen temperature. (20.4 K)
- Process 4-5 : The pressure above the liquid hydrogen is reduced to approximately 0.23 kPa.
 - Liquid hydrogen boils as the pressure is lowered until the hydrogen solidifies.
 - At the end of the process, the solid hydrogen is at approximately 10 K.

Simon helium-liquefaction system process

- Process 5-6 : The pressure of the gaseous helium is reduced to atmosphere pressure by allowing the helium to be released to a gas holder external to the system.
 - The helium that finally remains in the inner container does work against the helium gas that is discharged.
 - The discharged gas removes energy from the system, thereby lowering the temperature of the remaining helium.
 - At the end of this process, 75~100% is liquid helium.

If heat is transferred reversibly,

$$dQ = m_c T_c ds_c = m_c c_c dT_c$$

 $m_c = mass of container$ $T_c = container temperature$ $c_c = container specific heat$

Substituting the Debye expression for the container specific heat,

$$ds_{c} = 233.78R_{c}T_{c}^{2}\frac{dT_{c}}{\theta_{D}^{3}}$$

Integration,

$$\Delta s_{c} = -\frac{77.93R_{c}}{\theta_{D}^{3}}(T_{5}^{3} - T_{f}^{3})$$

There is no external heat transfer from the surroundings to the inner vessel,

$$(m_6 - m_5)s_5 - 0 = (m_6s_6 - m_5s_5) + m_c\Delta s_c$$

 $m_6 = total mass of helium in the vessel after expansion <math>m_5 = total mass of helium in the vessel before the expansion <math>s_5, s_6 = specific entropy values at points 5 and 6, respectively$

Solving for s_6 ,

$$s_6 = s_5 - (\frac{m_c}{m_6})\Delta s_c$$

Write the entropy s₆ in terms of the liquid yield

 $y = m_f/m_6$,

$$\mathbf{s}_6 = \mathbf{y}\mathbf{s}_f + (1 - \mathbf{y})\mathbf{s}_g$$

 $s_{\rm f}=$ entropy of the saturated liquid at the final condition $s_{\rm g}=$ entropy of the saturated vapor at the final condition

We can solve for the liquid yield,

$$y = \frac{\left(s_g - s_5\right) + \left(\frac{m_c}{m_6}\right)\Delta s_c}{s_g - s_f}$$

The helium specific volume,

$$v_6 = yv_f + (1 - y)v_g = v_g - y(v_g - v_f) = V/m_6$$

V = container internal volume

Fraction volume that is filled with liquid helium,

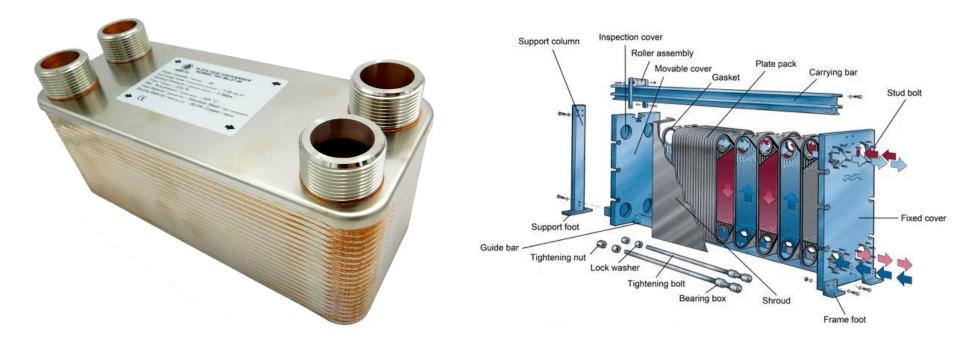
$$\frac{V_f}{V} = \frac{m_f v_f}{m_6 v_6} = \frac{y v_f}{v_6}$$

Substituting for the specific volume v_6 we can obtain,

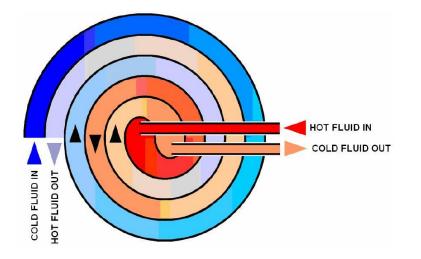
$$\frac{V_{f}}{V} = \frac{y}{\left(\frac{V_{g}}{V_{f}}\right) - y\left[\left(\frac{V_{g}}{V_{f}}\right) - 1\right]}$$

Single tube Linde exchanger	High pressure stream for inner circle because of its high density	High-pressure stream Low-pressure stream (a)
Multiple tube Linde exchanger	For multi heat exchanging ex) Linde dual pressure system	H-P stream (b)
Single tube Linde exchanger with a wire spacer	wire spacer causes turbulent flow which increases heat exchange rate while also increases pressure drop	H-P stream L-P stream (c)
Bundle type heat exchanger	bundle of circular tube with an ensure of good thermal contact	H-P stream Solder (d)

1.Plate and Frame Exchangers

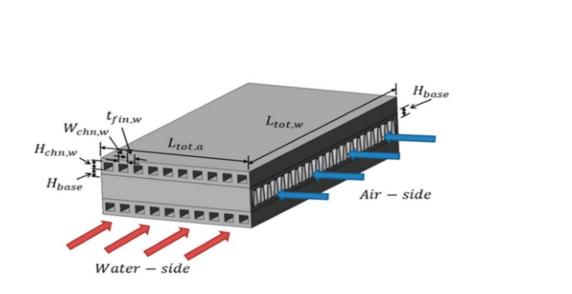


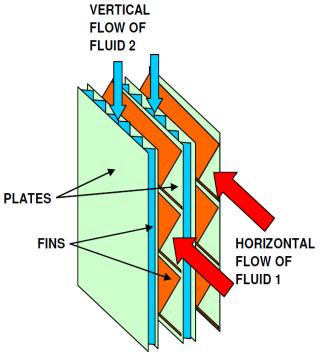
2. Spiral Heat Exchangers



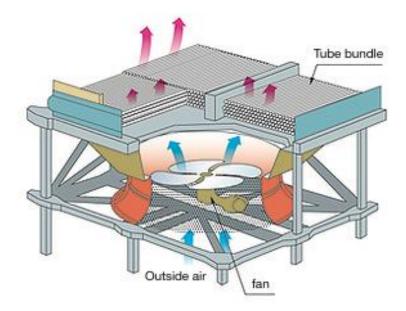


3. Plate-Fin Heat Exchangers





4. Air coolers



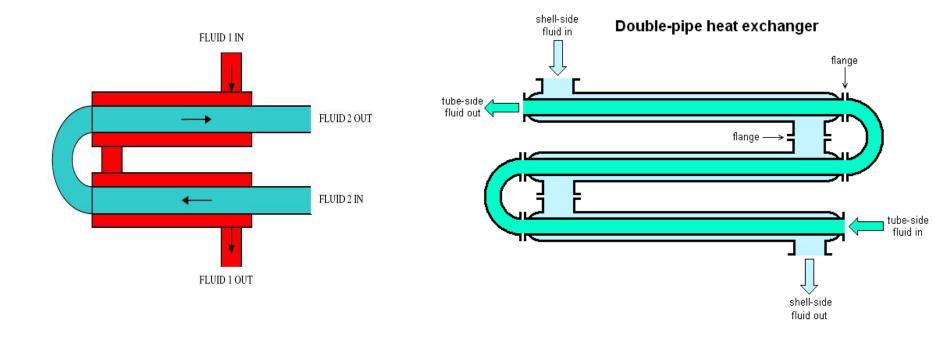


Reference: Dry Cooler

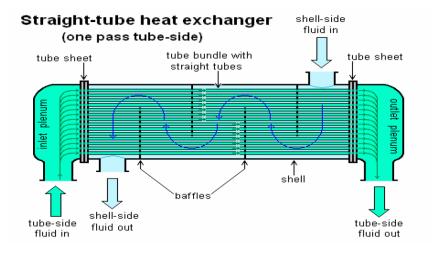
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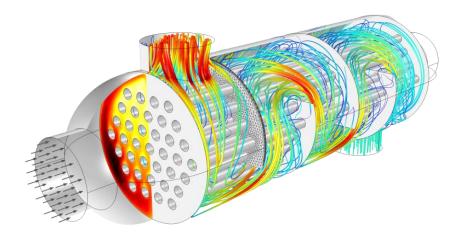
Cryogenic Engineering

5. Double Pipe Exchangers



6. Shell and Tube Exchangers





3.22 Heat-Transfer Coefficients and Friction Factors

• Heat-transfer data are usually presented in terms of dimensionless moduli, including the following parameters.

1. Prandtl number :
$$N_{Pr} = {{}^{\mu C_p}/_{k_t}}$$

2. Nusselt number :
$$N_{Nu} = \frac{h_c D_e}{k_t}$$

3. Colburn J-factor :
$$j_H = (\frac{h_c}{GC_p}) N_{Pr}^{2/3}$$

4. Reynolds number :
$$N_{Re} = \frac{D_e G}{\mu}$$

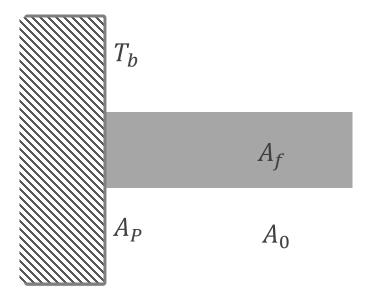
5. Friction factor :
$$f = \left(\frac{\Delta p}{L}\right) \left(\frac{G^2}{2g\rho D_e}\right)$$

3.23 Fin Effectiveness

Fin effectiveness

$\eta_{f} = \frac{actual \ \dot{Q}_{f}}{\dot{Q}_{f} \ when \ fin \ is \ maintained \ at \ T_{b}}$

3.23 Fin Effectiveness



 $A_0 = total \ surface \ area$ $A_f = fin \ surface \ area$ $A_p = primary \ surface \ area$ $\eta_0 = surface \ effectiveness$

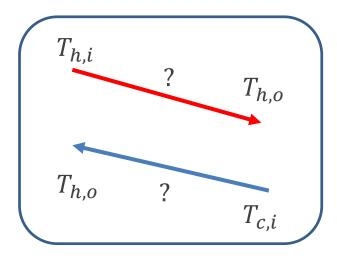
 $\eta_0 = 1 - (A_f/A_0)(1 - \eta_f)$

Surface effectiveness

 $\eta_0 A_0 = (1)A_0 + \eta_f A_f = A_0 - A_f + \eta_f A_f$

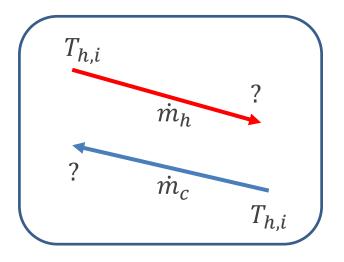
Overall heat transfer coefficient

$$\dot{Q} = UA_0, \qquad \frac{1}{UA} = \frac{1}{A_0h_0} + \frac{1}{A_ih_i} + \frac{t}{kA} + R$$



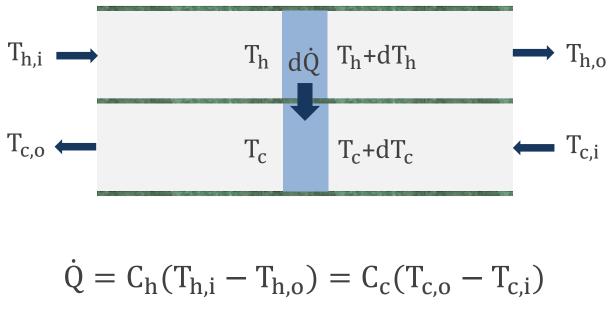
LMTD method

When inlet and outlet temperatures are known



e-NTU method

When inlet conditions and mass flows are known



$$d\dot{Q} = -C_h dT_h = -C_c dT_c = U dA\Delta T$$

3.24 LMTD Method

$$d(\Delta T) = dT_{h} - dT_{c} = -(\frac{UdA}{C_{h}}\Delta T - \frac{UdA}{C_{c}}\Delta T)$$

$$\frac{\mathrm{d}(\Delta \mathrm{T})}{\Delta \mathrm{T}} = -(\frac{1}{\mathrm{C}_{\mathrm{h}}} - \frac{1}{\mathrm{C}_{\mathrm{c}}})\mathrm{U}\mathrm{d}\mathrm{A}$$

$$\rightarrow d(\Delta T) = dT_{h} - dT_{c} = -(\frac{UdA}{C_{h}}\Delta T - \frac{UdA}{C_{c}}\Delta T)$$

$$\rightarrow \frac{d(\Delta T)}{\Delta T} = -(\frac{1}{C_{h}} - \frac{1}{C_{c}})UdA$$

$$\rightarrow \ln\left(\frac{(\Delta T)_2}{(\Delta T)_1}\right) = \frac{(\Delta T)_2 - (\Delta T)_1}{\dot{Q}} UA$$

3.24 LMTD Method

$$\dot{Q} = \frac{\ln((\Delta T)_2 / (\Delta T)_1)}{(\Delta T)_2 - (\Delta T)_1} UA$$

co-flow

$$(\Delta T)_1 = T_{h,i} - T_{c,i}$$

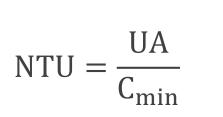
 $(\Delta T)_2 = T_{h,o} - T_{c,o}$

counter flow

$$(\Delta T)_1 = T_{h,i} - T_{c,o}$$

 $(\Delta T)_2 = T_{h,o} - T_{c,i}$

$$\epsilon = \frac{q}{q_{\text{max}}} = \frac{C_{\text{c}}(T_{\text{c},\text{o}} - T_{\text{c},\text{i}})}{C_{\text{min}}(T_{\text{h},\text{i}} - T_{\text{c},\text{i}})} = \frac{C_{\text{h}}(T_{\text{h},\text{i}} - T_{\text{h},\text{o}})}{C_{\text{min}}(T_{\text{h},\text{i}} - T_{\text{c},\text{i}})}$$



 $C_{\min} = C_C$ for most cryonic system

$$\epsilon = \frac{(T_{c,o} - T_{c,i})}{(T_{h,i} - T_{c,i})} \qquad \text{NTU} = \frac{\text{UA}}{\text{C}_{\text{C}}}$$

$$\frac{(\Delta T)_2}{(\Delta T)_1} = \frac{UA}{C_C} \left(1 - \frac{C_C}{C_h} \right) \rightarrow \frac{T_{h,o} - T_{c,i}}{T_{h,i} - T_{c,o}} = \exp\left(\frac{UA}{C_{\min}} \left(1 - \frac{C_{\min}}{C_{\max}} \right) \right)$$

$$\frac{T_{h,o} - T_{c,i}}{T_{h,i} - T_{c,o}} = \frac{T_{h,o} - T_{c,i}}{T_{h,i} - T_{c,o}} = \frac{T_{h,o} - T_{h,i} + T_{h,i} - T_{c,i}}{T_{h,i} - T_{c,i} + T_{c,i} - T_{c,o}} = \frac{1 - \frac{T_{h,i} - T_{h,o}}{T_{h,i} - T_{c,i}}}{1 - \frac{T_{c,o} - T_{c,i}}{T_{h,i} - T_{c,i}}} = \frac{1 - C_R \epsilon}{1 - \epsilon} = \exp(\text{NTU}(1 - C_R))$$

for counter flow when
$$C_{\min} = C_C$$
: $\epsilon = \frac{1 - \exp[-NTU(1 + C_r)]}{1 - C_r \exp[-NTU(1 - C_r)]}$

3.24 e-NTU Method

Co-flow:

$$\epsilon = \frac{1 - \exp[-NTU(1 + C_r)]}{1 + C_r}$$
Counter-flow:

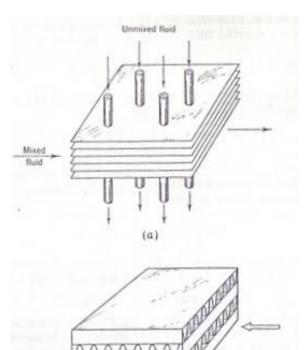
$$\epsilon = \frac{1 - \exp[-NTU(1 + C_r)]}{1 - C_r \exp[-NTU(1 - C_r)]}$$

$$\epsilon = \frac{NTU}{1 + NTU}$$

$$C_r = 1$$

$$\epsilon = 1 - \exp(-NTU)C_r = 0$$

3.24 e-NTU Method



(6)

both unmixed:

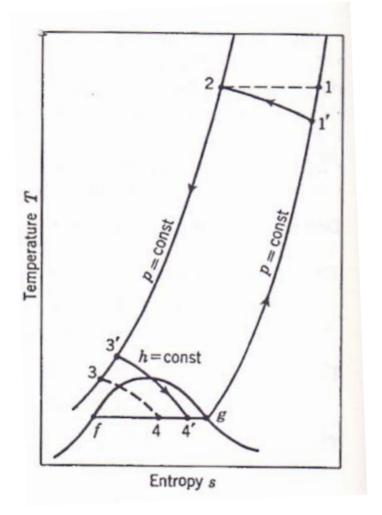
$$\epsilon = 1 - \exp\left[\left(\frac{1}{C_r}\right)(NTU)^{0.22} \{\exp\left[-C_r(NTU)^{0.78}\right] - 1\}\right]$$

C_{max} mixed:

 $\epsilon = \left(\frac{1}{C_r}\right)(1 - \exp(-C_r^{-1}\{1 - \exp[-C_r(NTU)]\})$ C_{in} mixed:

$$\epsilon = (1 - \exp(-C_r^{-1}\{1 - \exp[-C_r(NTU)]\}))$$

3.25 Effect of ε on System Performance



Linde-Hampson cycle with a heat exchanger with an effectiveness less than 100 percent. Points 1, 2, 3, and 4 are the state points for the ideal Linde-Hampson system. Point 1', 2', 3', and 4' are the state points for the system with the less effective heat exchanger.

3.25 Effect of ε on System Performance

Mean specific heat

$$\overline{C}_{p,h} = \frac{h_2 - h_3}{T_2 - T_3} > \overline{C}_{p,c} = \frac{h_1 - h_g}{T_1 - T_g}$$

In Heat Exchanger

$$C_{h} = \dot{m}\overline{C}_{p,h} > C_{c} = \dot{m}(1-y)\overline{C}_{p,c}$$
$$C_{min} = C_{c}$$

HX Effectiveness

$$\varepsilon = \frac{C_{c} \left(T_{1}^{'} - T_{g}^{'} \right)}{C_{min} \left(T_{1}^{'} - T_{g}^{'} \right)} = \frac{h_{1}^{'} - h_{g}^{'}}{h_{1}^{'} - h_{g}^{'}}$$

3.25 Effect of ε on System Performance

Liquid yield

$$y = \frac{h_{1} - h_{2}}{h_{1} - h_{f}} = \frac{h_{1} - h_{2} + (h_{1} - h_{1})}{h_{1} - h_{f} + (h_{1} - h_{1})} \quad (\because h_{1} - h_{1} = (1 - \varepsilon)(h_{1} - h_{g}))$$
$$= \frac{h_{1} - h_{2} + (1 - \varepsilon)(h_{1} - h_{g})}{h_{1} - h_{f} + (1 - \varepsilon)(h_{1} - h_{g})} < \frac{h_{1} - h_{2}}{h_{1} - h_{f}}$$

➔ liquid yield is decreased

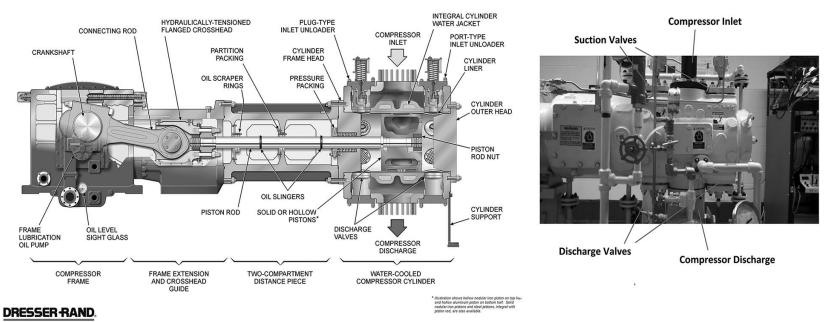
If y = 0

$$\varepsilon' = 1 - \frac{h_1 - h_2}{h_1 - h_g} = 0.869$$

HX is very critical in the liquefaction system

$$\Delta \mathbf{W} = \dot{\mathbf{m}} \left(\mathbf{h}_{1} - \mathbf{h}_{1}^{'} \right)$$
$$= \dot{\mathbf{m}} \left(1 - \varepsilon \right) \left(\mathbf{h}_{1} - \mathbf{h}_{g} \right) \quad : \text{ increased work}$$

- There are two broad classes of compressors and expanders used in liquefaction systems
 - 1. Reciprocating

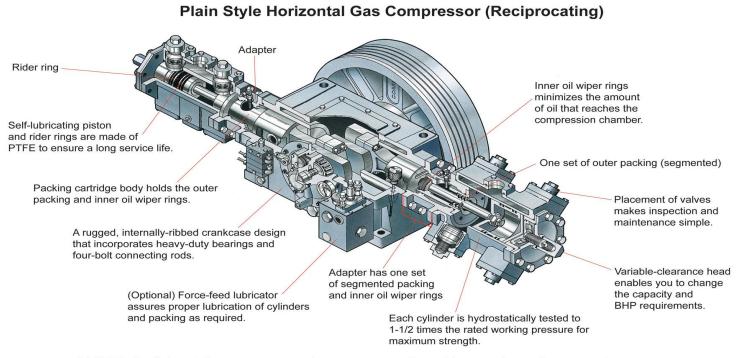


Reciprocating Process Compressor Components

1-1. Reciprocating Compressor-IC engine type

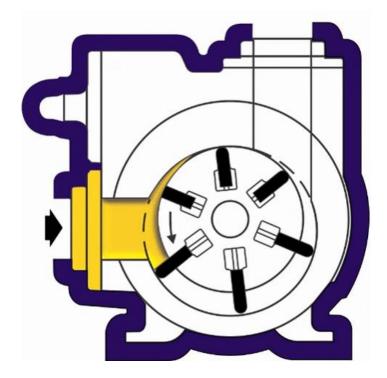


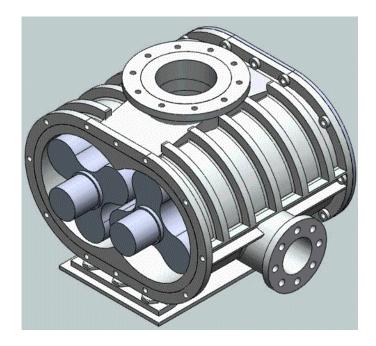
1-2. Reciprocating Compressor-steam engine type



NOTE: A plain style compressor has one set of packing and no distance piece.

2. Rotary Compressor





Reference: 1) gascompressor.co.uk 2)Giphy

Cryogenic Engineering

2-2. Rotary Expander

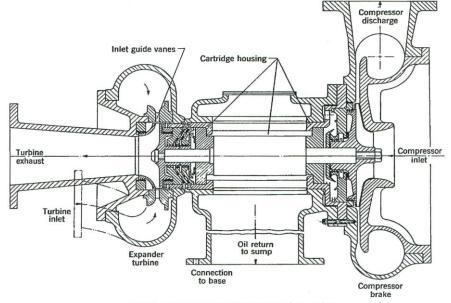
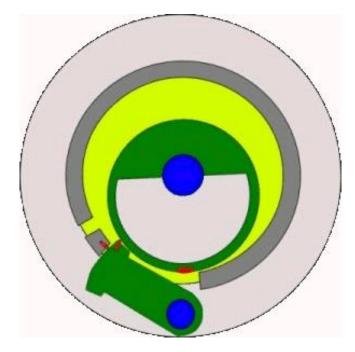
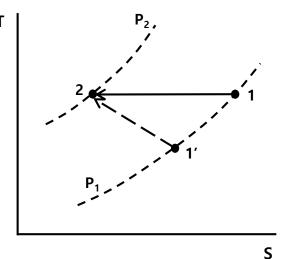


Fig. 3.44. Rotary expander (Worthington Corporation).



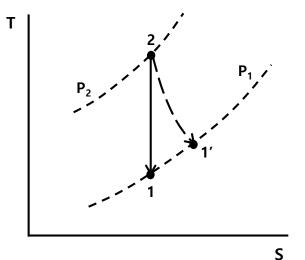
1. Isothermal efficiency of a compressor

$$\eta_{T=c} = \frac{\left(\dot{W} / \dot{m}\right)_{T=c}}{\left(\dot{W} / \dot{m}\right)_{indicated}}$$



2. Adiabatic efficiency of an expander

$$\eta_{ad} = \frac{\left(h_1 - h_2\right)_{indicated}}{\left(h_1 - h_2\right)_{ad}}$$



3. Mechanical efficiency of a compressor

$$\eta_{c,m} = \frac{\left(\dot{W} / \dot{m}\right)_{indicated}}{\left(\dot{W} / \dot{m}\right)_{actual}} = 1 - \frac{\dot{W}_{friction}}{\dot{W}_{actual}}$$

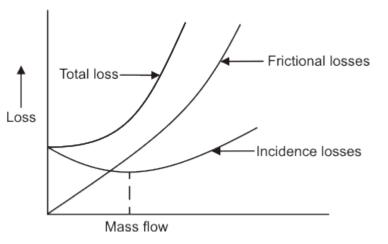
4. Mechanical efficiency of an expander

$$\eta_{e,m} = \frac{\left(\dot{W}/\dot{m}\right)_{actual}}{\left(h_1 - h_2\right)_{actual}} = \frac{\dot{W}_{actual}}{\dot{W}_{actual} + \dot{W}_{friction}}$$

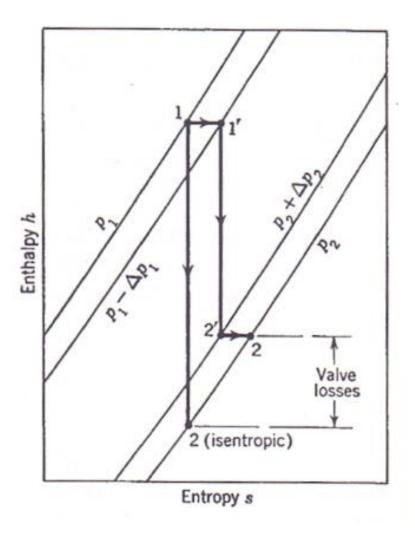
5. Overall efficiency

 $\eta_{o}=\eta_{T=c}\eta_{c,m}\ : \text{compressor}$

 $\eta_{o}=\eta_{ad}\eta_{e,m} \quad : \text{expander}$



Dependence of various losses with mass flow in a centrifugal compressor



The effect of inlet and exhaust valve losses in an expander.

Reciprocating compressors and expanders, the following factors contribute to the efficiency being less than 100 percent:

- **1. Inlet-valve and outlet-valve losses.** Pressure drops occur across the inlet and exhaust valves.
- 2. Incomplete expansion. By shortening the cutoff, the friction losses are reduced at the expense of the small amount of work output at the end of the stroke.
- 3. Heat transfer.

$$\dot{\mathbf{Q}} - \dot{\mathbf{W}} = \dot{\mathbf{m}} (\mathbf{h}_2 - \mathbf{h}_1)_{\text{actual}}$$

4. Piston friction

$$\eta_{ad} = 1 - \frac{\dot{Q}/\dot{m}}{(h_1 - h_2)_{ad}}$$

 Work for adiabatic expansion (the valve losses were the only irreversible losses)

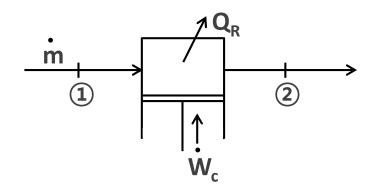
$$\frac{-\dot{W}}{\dot{m}} = \frac{\gamma R T_1}{\gamma - 1} \left\{ \left[\frac{p_2 \left(1 + \Delta p_2 / p_2 \right)}{p_1 \left(1 + \Delta p_1 / p_1 \right)} \right]^{(\gamma - 1)/\gamma} - 1 \right\}$$

• Work for an isothermal compressor (the valve losses were the only irreversible losses)

$$\frac{-\dot{W}}{\dot{m}} = RT_1 \ln\left[\frac{p_2(1+\Delta p_2/p_2)}{p_1(1+\Delta p_1/p_1)}\right]$$

3.28 Effect of Compressor and Expander Efficiency on Performance

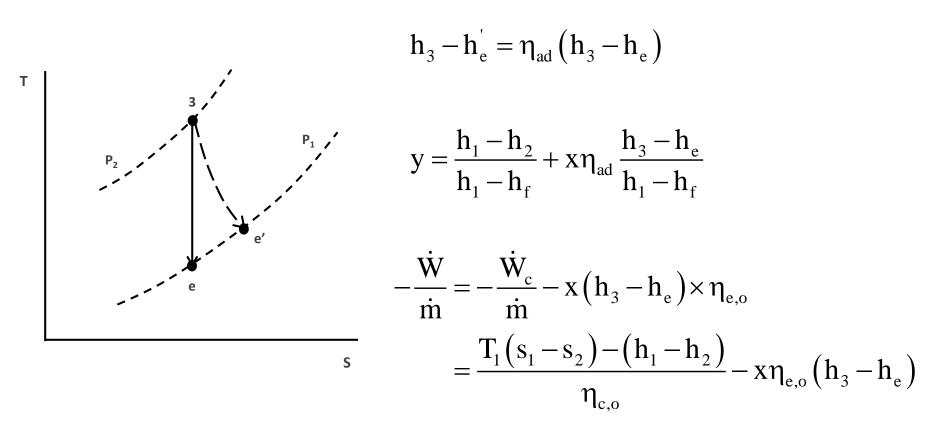
For compressor,



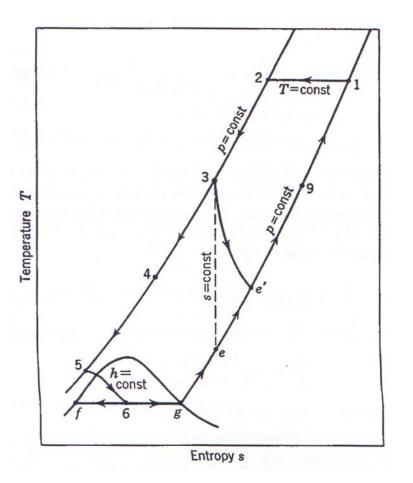
$$-\frac{\dot{W}_{c}}{\dot{m}} = \left[T_{1}\left(s_{1}-s_{2}\right)-\left(h_{1}-h_{2}\right)\right] \times \frac{1}{\eta_{c,o}}$$

3.28 Effect of Compressor and Expander Efficiency on Performance

For expander,



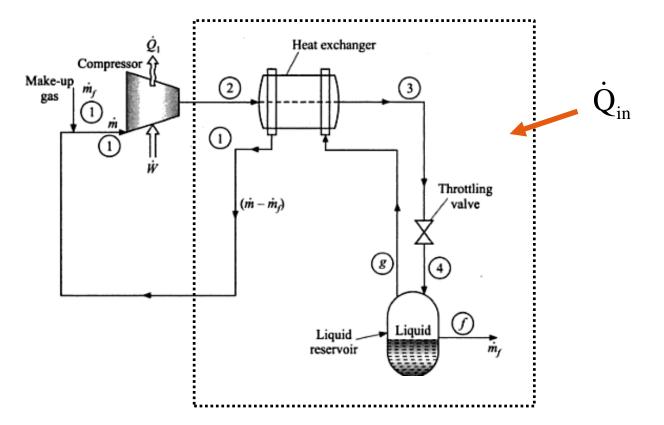
3.28 Effect of Compressor and Expander Efficiency on Performance



Claude system with an expander having an adiabatic efficiency less than unity. The gas leaves the expander at point e' in this case, instead of point e in the case of a reversible expander. (See Fig.3.18.)

3.29 Effect of Heat Transfer to the System

What happens when heat is coming into system?



$$\dot{\mathbf{Q}}_{\mathrm{in}} = \left(\dot{\mathbf{m}} - \dot{\mathbf{m}}_{\mathrm{f}}\right)\mathbf{h}_{1} + \dot{\mathbf{m}}_{\mathrm{f}}\mathbf{h}_{\mathrm{f}} - \dot{\mathbf{m}}\mathbf{h}_{2}$$

$$=\dot{m}(1-y)h_1+\dot{m}yh_f-\dot{m}h_2$$

$$y = \frac{h_1 - h_2}{h_1 - h_f} - \frac{\dot{Q}_{in}/\dot{m}}{h_1 - h_f}$$

The heat transfer rate per unit mass compressed is the important parameter!