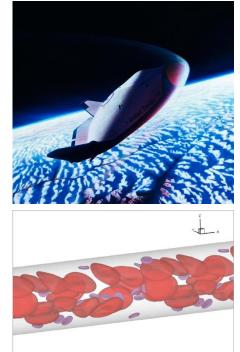
원자로 열유체 실험

Department of Nuclear Engineering, Seoul National Univ. Hyoung Kyu Cho

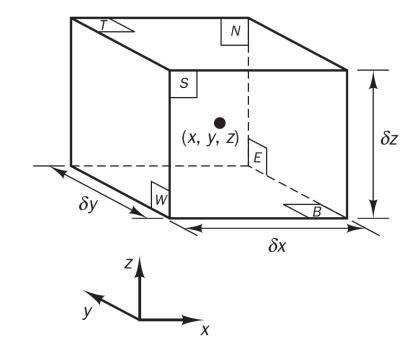
- The governing equations of fluid flow represent mathematical statements of the conservation laws of physics.
 - ✓ The mass of fluid is conserved.
 - ✓ The rate of change of momentum equals the sum of the forces on a fluid particle
 - Newton's second law
 - ✓ The rate of change of energy is equal to the sum of the rate of heat addition to and the rate of work done on a fluid particle
 - First law of thermodynamics

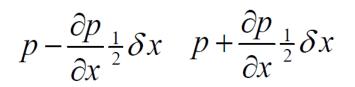
✓ Continuum assumption

- Fluid flows at macroscopic length scales > 1 μm
- The molecular structure and motions may be ignored.
- Macroscopic properties
 - Velocity, pressure, density, temperature
 - Averaged over suitably large numbers of molecules
- Fluid particle
 - The smallest possible element of fluid whose macroscopic properties are not influenced by individual molecules.



- Control volume
 - ✓ Six faces: N, S, E, W, T, B
 - \checkmark The center of the element: (x, y, z)
- Properties at the volume center
 - $\rho = \rho (x, y, z, t)$ p = p (x, y, z, t)T = T (x, y, z, t)u = u (x, y, z, t)
- Fluid properties at faces are approximated by means of the two terms of the Taylor series.
 The pressure at the W and E faces





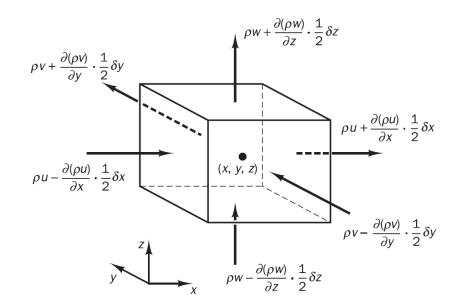
Mass Conservation in Three Dimensions

✓ Rate of increase of mass

$$\frac{\partial}{\partial t}(\rho\,\delta x\,\delta y\,\delta z) = \frac{\partial\rho}{\partial t}\,\delta x\,\delta y\,\delta z$$

✓ Net rate of flow of mass into the element

$$\left(\rho u - \frac{\partial(\rho u)}{\partial x}\frac{1}{2}\delta x\right)\delta y \delta z - \left(\rho u + \frac{\partial(\rho u)}{\partial x}\frac{1}{2}\delta x\right)\delta y \delta z$$
$$+ \left(\rho v - \frac{\partial(\rho v)}{\partial y}\frac{1}{2}\delta y\right)\delta x \delta z - \left(\rho v + \frac{\partial(\rho v)}{\partial y}\frac{1}{2}\delta y\right)\delta x \delta z$$
$$+ \left(\rho w - \frac{\partial(\rho w)}{\partial z}\frac{1}{2}\delta z\right)\delta x \delta y - \left(\rho w + \frac{\partial(\rho w)}{\partial z}\frac{1}{2}\delta z\right)\delta x \delta y$$



$$\frac{\partial \rho}{\partial t} + \frac{\partial (\rho u)}{\partial x} + \frac{\partial (\rho v)}{\partial y} + \frac{\partial (\rho w)}{\partial z} = 0$$

$$\frac{\partial \rho}{\partial t} + \operatorname{div}(\rho \mathbf{u}) = 0$$

Mass conservation/ Continuity Eq.

Mass Conservation in Three Dimensions

$$\frac{\partial \rho}{\partial t} + \operatorname{div}(\rho \mathbf{u}) = 0$$

 \checkmark For an incompressible fluid, the density ρ is constant.

div
$$\mathbf{u} = 0$$
 $\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} = 0$

- Rates of change following a fluid particle and for a fluid element
 - ✓ Lagrangian approach/ changes of properties of a fluid particle
 - Total or substantial derivative of ϕ
 - $\checkmark \phi$: Function of the position (x,y,z), property per unit mass

$$\frac{D\phi}{Dt} = \frac{\partial\phi}{\partial t} + \frac{\partial\phi}{\partial x}\frac{dx}{dt} + \frac{\partial\phi}{\partial y}\frac{dy}{dt} + \frac{\partial\phi}{\partial z}\frac{dz}{dt}$$

- ✓ A fluid particle follows the flow, so
 - dx / dt = udy / dt = vdz / dt = w
- \checkmark Hence, the substantive derivative of ϕ is given by

$$\frac{D\phi}{Dt} = \frac{\partial\phi}{\partial t} + u\frac{\partial\phi}{\partial x} + v\frac{\partial\phi}{\partial y} + w\frac{\partial\phi}{\partial z} = \frac{\partial\phi}{\partial t} + \mathbf{u} \text{ . grad } \phi$$

- Rates of change following a fluid particle and for a fluid element
 - $\checkmark \quad \frac{D\phi}{Dt} \quad \text{defines the rate of change of property } \phi \text{ per unit mass.}$
 - $\checkmark \rho \frac{D\phi}{Dt}$ The rate of change of property ϕ per unit volume

$$\frac{D\phi}{Dt} = \frac{\partial\phi}{\partial t} + u\frac{\partial\phi}{\partial x} + v\frac{\partial\phi}{\partial y} + w\frac{\partial\phi}{\partial z} = \frac{\partial\phi}{\partial t} + \mathbf{u} \text{ grad } \phi$$

$$\rho \frac{D\phi}{Dt} = \rho \left(\frac{\partial\phi}{\partial t} + \mathbf{u} \text{ grad } \phi\right)$$

✓ Eulerian approach/ changes of properties in a fluid element

- Far more common than Lagrangian approach
- Develop equations for collections of fluid elements making up a region fixed in space

- Rates of change following a fluid particle and for a fluid element
 - $\checkmark\,$ LHS of the mass conservation equation

$$\frac{\partial \rho}{\partial t} + \operatorname{div}(\rho \mathbf{u})$$

 \checkmark The generalization of these terms for an arbitrary conserved property

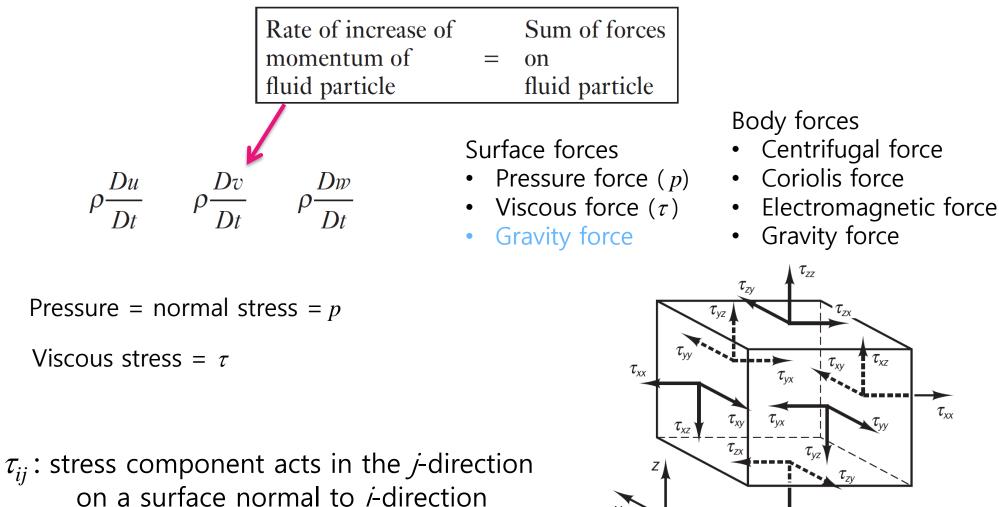
$$\frac{\partial(\rho\phi)}{\partial t} + \operatorname{div}(\rho\phi\mathbf{u}) \qquad \begin{pmatrix} \text{Rate of increase} \\ \text{of }\phi \text{ per unit volume} \end{pmatrix} + \begin{pmatrix} \text{Net rate of flow of }\phi \\ \text{out of fluid element} \\ \text{per unit volume} \end{pmatrix}$$
$$\frac{\partial(\rho\phi)}{\partial t} + \operatorname{div}(\rho\phi\mathbf{u}) = \rho \left[\frac{\partial\phi}{\partial t} + \mathbf{u} \cdot \operatorname{grad} \phi \right] + \phi \left[\frac{\partial\rho}{\partial t} + \operatorname{div}(\rho\mathbf{u}) \right]$$
$$\rho \frac{D\phi}{Dt} = \rho \left(\frac{\partial\phi}{\partial t} + \mathbf{u} \cdot \operatorname{grad} \phi \right) \qquad \frac{\partial\rho}{\partial t} + \operatorname{div}(\rho\mathbf{u}) = 0$$

$\partial(\rho\phi)$	Dφ	Rate of increase		Net rate of flow		Rate of increase	
$\frac{\partial(\rho\phi)}{\partial t} + \operatorname{div}(\rho\phi\mathbf{u})$	$= \rho \frac{D \psi}{d \phi}$	of ϕ of fluid	+	of ϕ out of	=	of ϕ for a	
∂t	^{P}Dt	element		fluid element		fluid particle	

- Rates of change following a fluid particle and for a fluid element
 - \checkmark Relevant entries of ϕ for momentum and energy equations

<i>x</i> -momentum	U	$\rho \frac{Du}{Dt}$	$\frac{\partial(\rho u)}{\partial t} + \operatorname{div}(\rho u \mathbf{u})$
y-momentum	υ	$\rho \frac{Dv}{Dt}$	$\frac{\partial(\rho v)}{\partial t} + \operatorname{div}(\rho v \mathbf{u})$
z-momentum	W	$\rho \frac{Dw}{Dt}$	$\frac{\partial(\rho w)}{\partial t} + \operatorname{div}(\rho w \mathbf{u})$
energy	E	$\rho \frac{DE}{Dt}$	$\frac{\partial(\rho E)}{\partial t} + \operatorname{div}(\rho E \mathbf{u})$

- Momentum equation in three dimensions
 - ✓ Newton's second law



 au_{zz}

Momentum equation in three dimensions

 \checkmark *x*-component of the forces due to pressure and viscous stress

On the pair of faces (E,W)

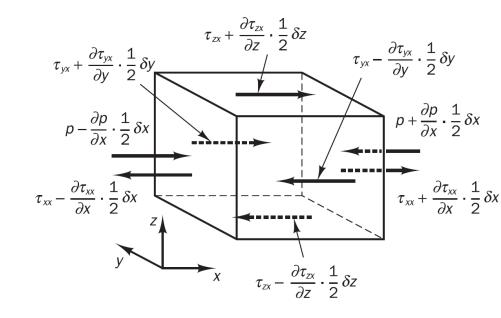
$$\left[\left(p - \frac{\partial p}{\partial x}\frac{1}{2}\delta x\right) - \left(\tau_{xx} - \frac{\partial \tau_{xx}}{\partial x}\frac{1}{2}\delta x\right)\right]\delta y \delta z + \left[-\left(p + \frac{\partial p}{\partial x}\frac{1}{2}\delta x\right) + \left(\tau_{xx} + \frac{\partial \tau_{xx}}{\partial x}\frac{1}{2}\delta x\right)\right]\delta y \delta z = \left(-\frac{\partial p}{\partial x} + \frac{\partial \tau_{xx}}{\partial x}\right)\delta x \delta y \delta z$$

On the pair of faces (N,S)

$$-\left(\tau_{yx} - \frac{\partial \tau_{yx}}{\partial y} \frac{1}{2} \delta y\right) \delta x \, \delta z + \left(\tau_{yx} + \frac{\partial \tau_{yx}}{\partial y} \frac{1}{2} \delta y\right) \delta x \, \delta z = \frac{\partial \tau_{yx}}{\partial y} \, \delta x \, \delta y \, \delta z$$

• On the pair of faces (T,B)

$$-\left(\tau_{zx} - \frac{\partial \tau_{zx}}{\partial z}\frac{1}{2}\delta z\right)\delta x\,\delta y + \left(\tau_{zx} + \frac{\partial \tau_{zx}}{\partial z}\frac{1}{2}\delta z\right)\delta x\,\delta y = \frac{\partial \tau_{zx}}{\partial z}\delta x\,\delta y\,\delta z$$



- Momentum equation in three dimensions
 - \checkmark *x*-component of the forces due to pressure and viscous stress
 - Total surface force per unit volume

$$\frac{\partial(-p+\tau_{xx})}{\partial x} + \frac{\partial\tau_{yx}}{\partial y} + \frac{\partial\tau_{zx}}{\partial z}$$

\checkmark *x*-component of the momentum equation

Rate of increase of		Sum of forces
momentum of	=	on
fluid particle		fluid particle

$$\rho \frac{Du}{Dt} = \frac{\partial (-p + \tau_{xx})}{\partial x} + \frac{\partial \tau_{yx}}{\partial y} + \frac{\partial \tau_{zx}}{\partial z} + S_{Mx}$$

Momentum equation in three dimensions

Rate of increase of		Sum of forces
momentum of	=	on
fluid particle		fluid particle

 \checkmark *x*-component of the momentum equation

$$\rho \frac{Du}{Dt} = \frac{\partial (-p + \tau_{xx})}{\partial x} + \frac{\partial \tau_{yx}}{\partial y} + \frac{\partial \tau_{zx}}{\partial z} + S_{Mx}$$

 \checkmark *y*-component of the momentum equation

$$\rho \frac{Dv}{Dt} = \frac{\partial \tau_{xy}}{\partial x} + \frac{\partial (-p + \tau_{yy})}{\partial y} + \frac{\partial \tau_{zy}}{\partial z} + S_{My}$$

Body force $S_{Mx} = 0$

$$S_{My} = 0$$

$$S_{Mz} = -\rho g$$

✓ z-component of the momentum equation

$$\begin{bmatrix}
\rho \frac{Dw}{Dt} = \frac{\partial \tau_{xz}}{\partial x} + \frac{\partial \tau_{yz}}{\partial y} + \frac{\partial (-p + \tau_{zz})}{\partial z} + S_{Mz}
\end{bmatrix}$$

- Energy equation in three dimensions
 - ✓ The first law of thermodynamics

$\rho \frac{DE}{Dt}$	Rate of increase of energy of fluid particle	=	Net rate of heat added to fluid particle	+	Net rate of work done on fluid particle

- $\checkmark\,$ Rate of work done by surface forces
 - $(F_{\text{surface forces}})(V)$
- \checkmark In *x*-direction,

$$\frac{\partial(-p+\tau_{xx})}{\partial x} + \frac{\partial\tau_{yx}}{\partial y} + \frac{\partial\tau_{zx}}{\partial z} \longrightarrow \left[\frac{\partial(u(-p+\tau_{xx}))}{\partial x} + \frac{\partial(u\tau_{yx})}{\partial y} + \frac{\partial(u\tau_{zx})}{\partial z}\right] \delta x \delta y \delta z$$

 \checkmark In y and z-directions,

$$\left[\frac{\partial(v\tau_{xy})}{\partial x} + \frac{\partial(v(-p + \tau_{yy}))}{\partial y} + \frac{\partial(v\tau_{zy})}{\partial z}\right] \delta x \delta y \delta z \quad \left[\frac{\partial(w\tau_{xz})}{\partial x} + \frac{\partial(w\tau_{yz})}{\partial y} + \frac{\partial(w(-p + \tau_{zz}))}{\partial z}\right] \delta x \delta y \delta z$$

- Energy equation in three dimensions
 - ✓ Total rate of work done on the fluid particle by surface stresses

$$\begin{bmatrix} \frac{\partial(u(-p+\tau_{xx}))}{\partial x} + \frac{\partial(u\tau_{yx})}{\partial y} + \frac{\partial(u\tau_{zx})}{\partial z} \end{bmatrix} \delta x \, \delta y \, \delta z$$

+
$$\begin{bmatrix} \frac{\partial(v\tau_{xy})}{\partial x} + \frac{\partial(v(-p+\tau_{yy}))}{\partial y} + \frac{\partial(v\tau_{zy})}{\partial z} \end{bmatrix} \delta x \, \delta y \, \delta z$$

+
$$\begin{bmatrix} \frac{\partial(m\tau_{xz})}{\partial x} + \frac{\partial(m\tau_{yz})}{\partial y} + \frac{\partial(m(-p+\tau_{zz}))}{\partial z} \end{bmatrix} \delta x \, \delta y \, \delta z$$

$$\begin{bmatrix} -\operatorname{div}(p\mathbf{u}) \end{bmatrix} + \begin{bmatrix} \frac{\partial(u\tau_{xx})}{\partial x} + \frac{\partial(u\tau_{yx})}{\partial y} + \frac{\partial(u\tau_{zx})}{\partial z} + \frac{\partial(v\tau_{xy})}{\partial x} + \frac{\partial(v\tau_{yy})}{\partial y} \end{bmatrix}$$
$$+ \frac{\partial(v\tau_{zy})}{\partial z} + \frac{\partial(w\tau_{xz})}{\partial x} + \frac{\partial(w\tau_{yz})}{\partial y} + \frac{\partial(w\tau_{zz})}{\partial z} \end{bmatrix}$$

Energy equation in three dimensions

Rate of increase	Net rate of	Net rate of work
of energy of =	heat added to	+ done on
fluid particle	fluid particle	fluid particle

✓ Net rate of heat transfer to the fluid particle

In x-direction,

$$\begin{bmatrix} \left(q_x - \frac{\partial q_x}{\partial x} \frac{1}{2} \delta x\right) - \left(q_x + \frac{\partial q_x}{\partial x} \frac{1}{2} \delta x\right) \end{bmatrix} \delta y \delta z = -\frac{\partial q_x}{\partial x} \delta x \delta y \delta z$$

= In y and z-directions,

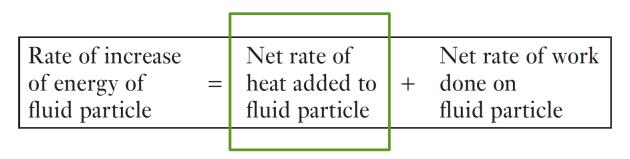
$$-\frac{\partial q_y}{\partial y} \delta x \delta y \delta z = -\frac{\partial q_z}{\partial z} \delta x \delta y \delta z$$

$$q_x + \frac{\partial q_y}{\partial y} \cdot \frac{1}{2} \delta y$$

$$q_x - \frac{\partial q_x}{\partial x} \cdot \frac{1}{2} \delta x$$

$$q_x - \frac{\partial q_x}{\partial x} \cdot \frac{1}{2} \delta x$$

Energy equation in three dimensions



✓ Total rate of heat added to the fluid particle per unit volume

$$-\frac{\partial q_x}{\partial x} - \frac{\partial q_y}{\partial y} - \frac{\partial q_z}{\partial z} = -\text{div } \mathbf{q}$$

✓ Fourier's law of heat conduction

$$q_x = -k \frac{\partial T}{\partial x}$$
 $q_y = -k \frac{\partial T}{\partial y}$ $q_z = -k \frac{\partial T}{\partial z}$ $\mathbf{q} = -k \operatorname{grad} T$ in vector form

 $-\operatorname{div} \mathbf{q} = \operatorname{div}(k \operatorname{grad} T)$

Energy equation in three dimensions *

Rate of increase of energy of fluid particle	=	Net rate of heat added to fluid particle	+	Net rate of work done on fluid particle		Energy Source
--	---	--	---	---	--	---------------

$$\rho \frac{DE}{Dt} \qquad \qquad \left[-\operatorname{div}(p\mathbf{u})\right] + \left[\frac{\partial(u\tau_{xx})}{\partial x} + \frac{\partial(u\tau_{yx})}{\partial y} + \frac{\partial(u\tau_{zx})}{\partial z} + \frac{\partial(v\tau_{xy})}{\partial x} + \frac{\partial(v\tau_{yy})}{\partial y} + \frac{\partial(v\tau_{yy})}{\partial y} + \frac{\partial(v\tau_{yy})}{\partial z} + \frac{\partial(v\tau_{yy$$

✓ Energy equation

DD

$$E = i + \frac{1}{2}(u^2 + v^2 + w^2)$$

• *E*: Sum of internal energy and kinetic energy

$$\rho \frac{DE}{Dt} = -\operatorname{div}(\rho \mathbf{u}) + \left[\frac{\partial (u\tau_{xx})}{\partial x} + \frac{\partial (u\tau_{yx})}{\partial y} + \frac{\partial (u\tau_{zx})}{\partial z} + \frac{\partial (v\tau_{xy})}{\partial z} \right]$$
$$+ \frac{\partial (v\tau_{yy})}{\partial y} + \frac{\partial (v\tau_{zy})}{\partial z} + \frac{\partial (w\tau_{xz})}{\partial x} + \frac{\partial (w\tau_{yz})}{\partial y} + \frac{\partial (w\tau_{zz})}{\partial z} \right]$$
$$+ \operatorname{div}(k \text{ grad } T) + S_E$$

- Energy equation in three dimensions
 - ✓ Kinetic energy equation

$$\rho \frac{Du}{Dt} = \frac{\partial (-p + \tau_{xx})}{\partial x} + \frac{\partial \tau_{yx}}{\partial y} + \frac{\partial \tau_{zx}}{\partial z} + S_{Mx} \qquad \times u$$

$$\rho \frac{Dv}{Dt} = \frac{\partial \tau_{xy}}{\partial x} + \frac{\partial (-p + \tau_{yy})}{\partial y} + \frac{\partial \tau_{zy}}{\partial z} + S_{My} \qquad \times v$$

$$\rho \frac{Dw}{Dt} = \frac{\partial \tau_{xz}}{\partial x} + \frac{\partial \tau_{yz}}{\partial y} + \frac{\partial (-p + \tau_{zz})}{\partial z} + S_{Mz} \qquad \times w$$

$$\rho \frac{D[\frac{1}{2}(u^{2} + v^{2} + w^{2})]}{Dt} = -\mathbf{u} \cdot \operatorname{grad} p + u \left(\frac{\partial \tau_{xx}}{\partial x} + \frac{\partial \tau_{yx}}{\partial y} + \frac{\partial \tau_{zx}}{\partial z} \right) + v \left(\frac{\partial \tau_{xy}}{\partial x} + \frac{\partial \tau_{yy}}{\partial y} + \frac{\partial \tau_{zy}}{\partial z} \right) + w \left(\frac{\partial \tau_{xz}}{\partial x} + \frac{\partial \tau_{yz}}{\partial y} + \frac{\partial \tau_{zz}}{\partial z} \right) + \mathbf{u} \cdot \mathbf{S}_{M}$$

- Energy equation in three dimensions
 - ✓ Total energy equation kinetic energy equation

$$\rho \frac{DE}{Dt} = -\operatorname{div}(\rho \mathbf{u}) + \left[\frac{\partial(u\tau_{xx})}{\partial x} + \frac{\partial(u\tau_{yx})}{\partial y} + \frac{\partial(u\tau_{zx})}{\partial z} + \frac{\partial(v\tau_{xy})}{\partial x} + \frac{\partial(v\tau_{xy})}{\partial x} + \frac{\partial(v\tau_{xy})}{\partial z} + \frac{\partial(w\tau_{xz})}{\partial x} + \frac{\partial(w\tau_{yz})}{\partial y} + \frac{\partial(w\tau_{zz})}{\partial z} \right] + \operatorname{div}(k \operatorname{grad} T) + S_E$$

$$E = i + \frac{1}{2}(u^2 + v^2 + w^2)$$

$$\rho \frac{D[\frac{1}{2}(u^2 + v^2 + m^2)]}{Dt} = -\mathbf{u} \cdot \operatorname{grad} p + u \left(\frac{\partial \tau_{xx}}{\partial x} + \frac{\partial \tau_{yx}}{\partial y} + \frac{\partial \tau_{zx}}{\partial z} \right) + v \left(\frac{\partial \tau_{xy}}{\partial x} + \frac{\partial \tau_{yy}}{\partial y} + \frac{\partial \tau_{zy}}{\partial z} \right) + m \left(\frac{\partial \tau_{xz}}{\partial x} + \frac{\partial \tau_{yz}}{\partial y} + \frac{\partial \tau_{zz}}{\partial z} \right) + \mathbf{u} \cdot \mathbf{S}_{M}$$

✓ Internal energy equation

$$\rho \frac{Di}{Dt} = -p \text{ div } \mathbf{u} + \text{ div}(k \text{ grad } T) + \tau_{xx} \frac{\partial u}{\partial x} + \tau_{yx} \frac{\partial u}{\partial y} + \tau_{zx} \frac{\partial u}{\partial z}$$
$$+ \tau_{xy} \frac{\partial v}{\partial x} + \tau_{yy} \frac{\partial v}{\partial y} + \tau_{zy} \frac{\partial v}{\partial z}$$
$$+ \tau_{xz} \frac{\partial w}{\partial x} + \tau_{yz} \frac{\partial w}{\partial y} + \tau_{zz} \frac{\partial w}{\partial z} + S_i$$

- Energy equation in three dimensions
 - \checkmark For the special case of an incompressible fluid, temperature equation

$$i = cT$$
 $div \mathbf{u} = 0$

$$\rho c \frac{DT}{Dt} = \operatorname{div}(k \operatorname{grad} T) + \tau_{xx} \frac{\partial u}{\partial x} + \tau_{yx} \frac{\partial u}{\partial y} + \tau_{zx} \frac{\partial u}{\partial z} + \tau_{xy} \frac{\partial v}{\partial x}$$
$$+ \tau_{yy} \frac{\partial v}{\partial y} + \tau_{zy} \frac{\partial v}{\partial z} + \tau_{xz} \frac{\partial w}{\partial x} + \tau_{yz} \frac{\partial w}{\partial y} + \tau_{zz} \frac{\partial w}{\partial z} + S_i$$

- Energy equation in three dimensions
 - ✓ Enthalpy equation

 $h = i + p/\rho$ $h_0 = h + \frac{1}{2}(u^2 + v^2 + w^2)$ Total enthalpy

$$h_0 = i + p/\rho + \frac{1}{2}(u^2 + v^2 + w^2) = E + p/\rho$$

✓ Total enthalpy equation

$$\rho \frac{DE}{Dt} = -\operatorname{div}(\rho \mathbf{u}) + \left[\frac{\partial(u\tau_{xx})}{\partial x} + \frac{\partial(u\tau_{yx})}{\partial y} + \frac{\partial(u\tau_{zx})}{\partial z} + \frac{\partial(v\tau_{xy})}{\partial x} + \frac{\partial(v\tau_{xz})}{\partial z} + \frac{\partial(v\tau_{zz})}{\partial y} + \frac{\partial(v\tau_{zz})}{\partial z} \right] + \operatorname{div}(k \operatorname{grad} T) + S_E + \operatorname{div}(\rho E \mathbf{u}) = \frac{\partial(\rho h_0)}{\partial t} - \frac{\partial p}{\partial t} + \operatorname{div}(\rho h_0 \mathbf{u}) - \operatorname{div}(\rho \mathbf{u}) + \operatorname{div}(\rho h_0 \mathbf{u}) = \operatorname{div}(k \operatorname{grad} T) + \frac{\partial p}{\partial t} + \operatorname{div}(\rho h_0 \mathbf{u}) = \operatorname{div}(k \operatorname{grad} T) + \frac{\partial p}{\partial t} + \operatorname{div}(\rho h_0 \mathbf{u}) = \operatorname{div}(k \operatorname{grad} T) + \frac{\partial p}{\partial t} + \frac{\partial(u\tau_{xx})}{\partial x} + \frac{\partial(u\tau_{xx})}{\partial y} + \frac{\partial(u\tau_{xx})}{\partial z} + \frac{\partial(u\tau_{xx})}{\partial z} + \frac{\partial(v\tau_{xx})}{\partial y} + \frac{\partial(v\tau_{xx})}{\partial z} + \frac{\partial(v\tau_{xx})}{\partial z} + \frac{\partial(v\tau_{xx})}{\partial y} + \frac{\partial(v\tau_{xx})}{\partial z} + \frac{\partial(v\tau_{xx})}{\partial z} + \frac{\partial(v\tau_{xx})}{\partial z} + \frac{\partial(v\tau_{xx})}{\partial y} + \frac{\partial(v\tau_{xx})}{\partial z} + \frac{\partial(v\tau_{xx})}{\partial$$

Summary

Governing Equations of Fluid Flow and Heat Transfer

$$\sqrt[]{\frac{\partial \rho}{\partial t}} + \operatorname{div}(\rho \mathbf{u}) = 0$$

$$\rho \frac{Du}{Dt} = \frac{\partial (-p + \tau_{xx})}{\partial x} + \frac{\partial \tau_{yx}}{\partial y} + \frac{\partial \tau_{zx}}{\partial z} + S_{Mx}$$

$$\rho \frac{Dv}{Dt} = \frac{\partial \tau_{xy}}{\partial x} + \frac{\partial (-p + \tau_{yy})}{\partial y} + \frac{\partial \tau_{zy}}{\partial z} + S_{My}$$

$$\rho \frac{Dw}{Dt} = \frac{\partial \tau_{xz}}{\partial x} + \frac{\partial \tau_{yz}}{\partial y} + \frac{\partial (-p + \tau_{zz})}{\partial z} + S_{Mz}$$

 τ_{ij} : stress component acts in the *j*-direction on a surface normal to *i*-direction

$$\rho \frac{DE}{Dt} = -\operatorname{div}(p\mathbf{u}) + \left[\frac{\partial(u\tau_{xx})}{\partial x} + \frac{\partial(u\tau_{yx})}{\partial y} + \frac{\partial(u\tau_{zx})}{\partial z} + \frac{\partial(v\tau_{xy})}{\partial z} + \frac{\partial(v\tau_{xy})}{\partial x} \right] \rho \frac{Di}{Dt} = -p \operatorname{div} \mathbf{u} + \operatorname{div}(k \operatorname{grad} T) + \tau_{xx} \frac{\partial u}{\partial x} + \tau_{yx} \frac{\partial u}{\partial y} + \tau_{zx} \frac{\partial u}{\partial z} + \tau_{yx} \frac{\partial u}{\partial z} + \tau_{yx} \frac{\partial u}{\partial z} + \tau_{xx} \frac{\partial u}{\partial z} + \tau_{$$

Equations of State

Thermodynamic variables

 ρ , p, i and T

- ✓ Assumption of thermodynamic equilibrium
- Equations of the state
 - ✓ Relate two state variables to the other variables

 $p = p(\rho, T)$ $i = i(\rho, T)$

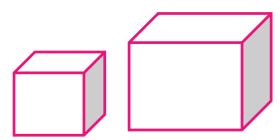
- Compressible fluids
 - ✓ EOS provides the linkage between the energy equation and other governing equations.
- Incompressible fluids
 - ✓ No linkage between the energy equation and the others.
 - ✓ The flow field can be solved by considering mass and momentum equations.

- * Viscous stresses τ_{ij} in momentum and energy equations
 - ✓ Viscous stresses can be expressed as functions of the local deformation rate (or strain rate).
 - $\checkmark\,$ In 3D flows the local rate of deformation is composed of
 - the linear deformation rate
 - the volumetric deformation rate.
 - ✓ All gases and many liquids are **isotropic.**
- The rate of linear deformation of a fluid element
 - ✓ Nine components in 3D
 - ✓ Linear elongating deformation $s_{xx} = \frac{\partial u}{\partial x} + s_{yy} = \frac{\partial v}{\partial y} + s_{zz} = \frac{\partial w}{\partial z}$ ←
 - ✓ Shearing linear deformation components

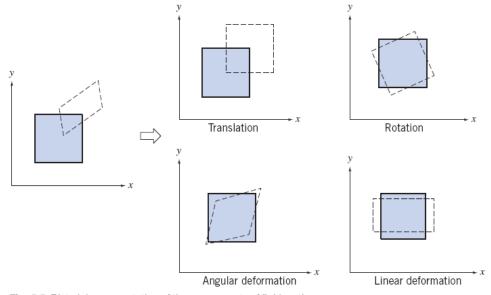
$$s_{xy} = s_{yx} = \frac{1}{2} \left(\frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} \right) \qquad s_{xz} = s_{zx} = \frac{1}{2} \left(\frac{\partial u}{\partial z} + \frac{\partial w}{\partial x} \right) \qquad s_{yz} = s_{zy} = \frac{1}{2} \left(\frac{\partial v}{\partial z} + \frac{\partial w}{\partial y} \right)$$

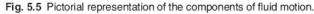
The rate of volume deformation of a fluid element

$$e_{xx} + e_{yy} + e_{zz} = \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} = \text{div } \mathbf{u}$$



• Viscous stresses τ_{ij} in momentum and energy equations





All gases and many liquids are isotropic.

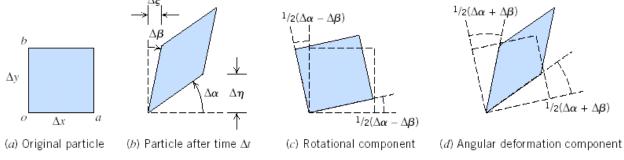
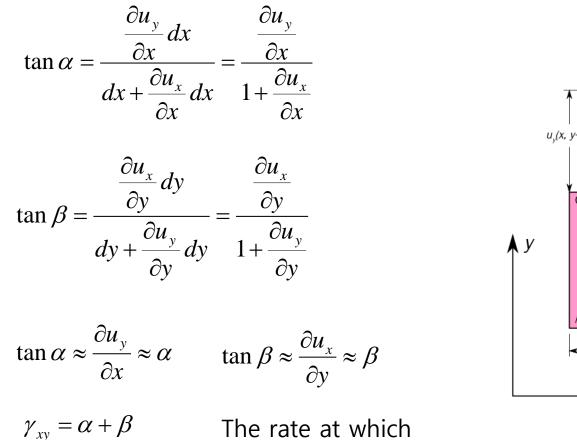
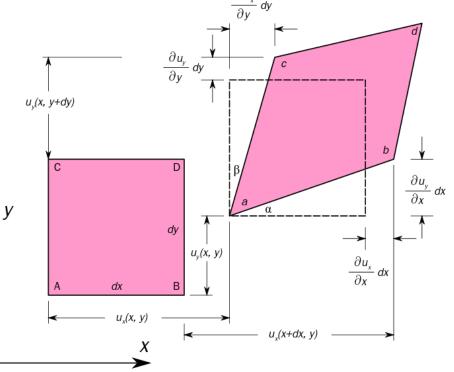


Fig. 5.7 Rotation and angular deformation of perpendicular line segments in a two-dimensional flow.

• Viscous stresses τ_{ii} in momentum and energy equations





The rate at which two sides close toward each other

 $s_{xy} = \frac{1}{2} \gamma_{xy}$

- Newtonian fluid
 - ✓ Viscous stresses are proportional to the rates of deformation.

✓ Two constants of proportionality

- Dynamic viscosity (μ): to relate stresses to linear deformations
- Second viscosity (λ): to relate stresses to volumetric deformation

$$(\tau_{ij})_{linear} = 2 \mu s_{ij}$$
$$(\tau_{ii})_{volume} = \lambda (e_{xx} + e_{yy} + e_{zz})$$

$$e_{xx} + e_{yy} + e_{zz} = \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} = \text{div } \mathbf{u}$$

Viscous stress components

$$\tau_{xx} = 2\mu \frac{\partial u}{\partial x} + \lambda \operatorname{div} \mathbf{u}$$
 $\tau_{yy} = 2\mu \frac{\partial v}{\partial y} + \lambda \operatorname{div} \mathbf{u}$ $\tau_{zz} = 2\mu \frac{\partial w}{\partial z} + \lambda \operatorname{div} \mathbf{u}$

$$\tau_{xy} = \tau_{yx} = \mu \left(\frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} \right) \quad \tau_{xz} = \tau_{zx} = \mu \left(\frac{\partial u}{\partial z} + \frac{\partial w}{\partial x} \right) \quad \tau_{yz} = \tau_{zy} = \mu \left(\frac{\partial v}{\partial z} + \frac{\partial w}{\partial y} \right)$$

✓ Second viscosity

- For gases: $\lambda = -\frac{2}{3}\mu$
- For liquid: $div \mathbf{u} = 0$

Momentum equations

$$\rho \frac{Du}{Dt} = \frac{\partial (-p + \tau_{xx})}{\partial x} + \frac{\partial \tau_{yx}}{\partial y} + \frac{\partial \tau_{zx}}{\partial z} + S_{Mx} \qquad \tau_{xx} = 2\mu \frac{\partial u}{\partial x} + \lambda \operatorname{div} \mathbf{u} \qquad \tau_{xy} = \tau_{yx} = \mu \left(\frac{\partial u}{\partial y} + \frac{\partial v}{\partial x}\right) \qquad \tau_{xz} = \tau_{zx} = \mu \left(\frac{\partial u}{\partial z} + \frac{\partial w}{\partial x}\right)$$

$$\rho \frac{Du}{Dt} = -\frac{\partial p}{\partial x} + \frac{\partial}{\partial x} \left[2\mu \frac{\partial u}{\partial x} + \lambda \operatorname{div} \mathbf{u} \right] + \frac{\partial}{\partial y} \left[\mu \left(\frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} \right) \right] + \frac{\partial}{\partial z} \left[\mu \left(\frac{\partial u}{\partial z} + \frac{\partial w}{\partial x} \right) \right] + S_{Mx}$$

$$\rho \frac{Dv}{Dt} = -\frac{\partial p}{\partial y} + \frac{\partial}{\partial x} \left[\mu \left(\frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} \right) \right] + \frac{\partial}{\partial y} \left[2\mu \frac{\partial v}{\partial y} + \lambda \operatorname{div} \mathbf{u} \right] \qquad \rho \frac{Dw}{Dt} = -\frac{\partial p}{\partial z} + \frac{\partial}{\partial x} \left[\mu \left(\frac{\partial u}{\partial z} + \frac{\partial w}{\partial x} \right) \right] + \frac{\partial}{\partial y} \left[\mu \left(\frac{\partial v}{\partial z} + \frac{\partial w}{\partial y} \right) \right] + \frac{\partial}{\partial z} \left[\mu \left(\frac{\partial v}{\partial z} + \frac{\partial w}{\partial y} \right) \right] + S_{My} \qquad \qquad + \frac{\partial}{\partial z} \left[2\mu \frac{\partial w}{\partial z} + \lambda \operatorname{div} \mathbf{u} \right] + S_{Mz}$$

Rearrangement

$$\frac{\partial}{\partial x} \left[2\mu \frac{\partial u}{\partial x} + \lambda \operatorname{div} \mathbf{u} \right] + \frac{\partial}{\partial y} \left[\mu \left(\frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} \right) \right] + \frac{\partial}{\partial z} \left[\mu \left(\frac{\partial u}{\partial z} + \frac{\partial w}{\partial x} \right) \right]$$
$$= \frac{\partial}{\partial x} \left(\mu \frac{\partial u}{\partial x} \right) + \frac{\partial}{\partial y} \left(\mu \frac{\partial u}{\partial y} \right) + \frac{\partial}{\partial z} \left(\mu \frac{\partial u}{\partial z} \right) + \left[\frac{\partial}{\partial x} \left(\mu \frac{\partial u}{\partial x} \right) + \frac{\partial}{\partial y} \left(\mu \frac{\partial v}{\partial x} \right) + \frac{\partial}{\partial z} \left(\mu \frac{\partial w}{\partial x} \right) + \frac{\partial}{\partial x} (\lambda \operatorname{div} \mathbf{u}) \right]$$

 $= \operatorname{div}(\mu \operatorname{grad} u) + [s_{Mx}]$

• N.-S. equations can be written as follows with modified source terms; $S_M = S_M + [s_M]$

$$\rho \frac{Du}{Dt} = -\frac{\partial p}{\partial x} + \operatorname{div}(\mu \operatorname{grad} u) + S_{Mx}$$
$$\rho \frac{Dv}{Dt} = -\frac{\partial p}{\partial y} + \operatorname{div}(\mu \operatorname{grad} v) + S_{My}$$
$$\rho \frac{Dw}{Dt} = -\frac{\partial p}{\partial z} + \operatorname{div}(\mu \operatorname{grad} w) + S_{Mz}$$

\diamond For incompressible fluids with constant μ

$$[s_{Mx}] = \left[\frac{\partial}{\partial x}\left(\mu\frac{\partial u}{\partial x}\right) + \frac{\partial}{\partial y}\left(\mu\frac{\partial v}{\partial x}\right) + \frac{\partial}{\partial z}\left(\mu\frac{\partial w}{\partial x}\right) + \frac{\partial}{\partial z}\left(\lambda \operatorname{div} \mathbf{u}\right)\right]$$
$$= \mu \left[\frac{\partial}{\partial x}\left(\frac{\partial u}{\partial x}\right) + \frac{\partial}{\partial y}\left(\frac{\partial v}{\partial x}\right) + \frac{\partial}{\partial z}\left(\frac{\partial w}{\partial x}\right)\right] = \mu \left[\frac{\partial}{\partial x}\left(\frac{\partial u}{\partial x}\right) + \frac{\partial}{\partial x}\left(\frac{\partial w}{\partial z}\right)\right]$$
$$= \mu \left[\frac{\partial}{\partial x}\left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z}\right)\right] = 0$$

Internal energy equation $\rho \frac{Di}{Dt} = -p \text{ div } \mathbf{u} + \text{div}(k \text{ grad } T) + \tau_{xx} \frac{\partial u}{\partial x} + \tau_{yx} \frac{\partial u}{\partial y} + \tau_{zx} \frac{\partial u}{\partial z}$ $+ \tau_{xy} \frac{\partial v}{\partial x} + \tau_{yy} \frac{\partial v}{\partial y} + \tau_{zy} \frac{\partial v}{\partial z}$ $\tau_{xx} = 2\mu \frac{\partial u}{\partial x} + \lambda \text{ div } \mathbf{u} \quad \tau_{yy} = 2\mu \frac{\partial v}{\partial y} + \lambda \text{ div } \mathbf{u} \quad \tau_{zz} = 2\mu \frac{\partial w}{\partial z} + \lambda \text{ div } \mathbf{u}$ $+ \tau_{xz} \frac{\partial w}{\partial x} + \tau_{yz} \frac{\partial w}{\partial y} + \tau_{zz} \frac{\partial w}{\partial z} + S_i$ $\tau_{xy} = \tau_{yx} = \mu \left(\frac{\partial u}{\partial y} + \frac{\partial v}{\partial x}\right) \quad \tau_{xz} = \tau_{zx} = \mu \left(\frac{\partial u}{\partial z} + \frac{\partial w}{\partial x}\right) \quad \tau_{yz} = \tau_{zy} = \mu \left(\frac{\partial v}{\partial z} + \frac{\partial w}{\partial y}\right)$

$$\rho \frac{Di}{Dt} = -p \operatorname{div} \mathbf{u} + \operatorname{div}(k \operatorname{grad} T) + \Phi + S_i$$

\checkmark Dissipation function Φ

- Always positive
- Source of internal energy due to deformation work on the fluid particle.
- Mechanical energy is converted into internal energy or heat.

$$\Phi = \mu \left\{ 2 \left[\left(\frac{\partial u}{\partial x} \right)^2 + \left(\frac{\partial v}{\partial y} \right)^2 + \left(\frac{\partial w}{\partial z} \right)^2 \right] + \left(\frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} \right)^2 + \left(\frac{\partial u}{\partial z} + \frac{\partial w}{\partial x} \right)^2 + \left(\frac{\partial v}{\partial z} + \frac{\partial w}{\partial y} \right)^2 \right\} + \lambda (\operatorname{div} \mathbf{u})^2$$

Conservative form of the governing equations of fluid flow

 $\frac{\partial \rho}{\partial t} + div(\rho \mathbf{u}) = 0$ Mass $\frac{\partial(\rho u)}{\partial t} + div(\rho u \mathbf{u}) = -\frac{\partial p}{\partial r} + div(\mu \operatorname{grad} u) + S_{Mx}$ x-momentum $\frac{\partial(\rho v)}{\partial t} + div(\rho v \mathbf{u}) = -\frac{\partial p}{\partial v} + div(\mu \operatorname{grad} v) + S_{My}$ y-momentum $\frac{\partial(\rho w)}{\partial t} + div(\rho w \mathbf{u}) = -\frac{\partial p}{\partial \tau} + div(\mu \operatorname{grad} w) + S_{Mz}$ *z*-momentum $\frac{\partial(\rho i)}{\partial t} + div(\rho i \mathbf{u}) = -p \, div \, \mathbf{u} + div(k \, grad \, T) + \Phi + S_i$ Internal energy u, v, w, p, i, ρ, T + EOS

This system is mathematically closed!