

Nonlinear Optical Engineering

Polarization Effects (1)
(NFO 5th ed: 6.1~6.2)

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Nonlinear Birefringence (1)

Origin of nonlinear birefringence:

Electric field: $\mathbf{E}(\mathbf{r}, t) = \frac{1}{2}(\hat{x}E_x + \hat{y}E_y) \exp(-i\omega_0 t) + c.c.$

Third-order susceptibility: How many elements for $\chi^{(3)}_{ijkl}$? $\rightarrow 81$ elements

For isotropic media: $\rightarrow 21$ nonzero elements

$$\chi^{(3)}_{xxxx} = \chi^{(3)}_{yyyy} = \chi^{(3)}_{zzzz},$$

$$\chi^{(3)}_{xxyy} = \chi^{(3)}_{xxzz} = \chi^{(3)}_{yyxx} = \chi^{(3)}_{yyzz} = \chi^{(3)}_{zzxx} = \chi^{(3)}_{zzyy},$$

$$\chi^{(3)}_{xyxy} = \chi^{(3)}_{xzzz} = \chi^{(3)}_{yzyz} = \chi^{(3)}_{yxyx} = \chi^{(3)}_{zxzx} = \chi^{(3)}_{zyzy}, \quad \leftarrow \text{Centrosymmetry}$$

$$\chi^{(3)}_{xyyx} = \chi^{(3)}_{xzzz} = \chi^{(3)}_{yxyx} = \chi^{(3)}_{yzyz} = \chi^{(3)}_{zxxz} = \chi^{(3)}_{zyyz}.$$

$$\rightarrow \chi^{(3)}_{xxxx} = \chi^{(3)}_{xxyy} + \chi^{(3)}_{xyxy} + \chi^{(3)}_{xyyx}$$

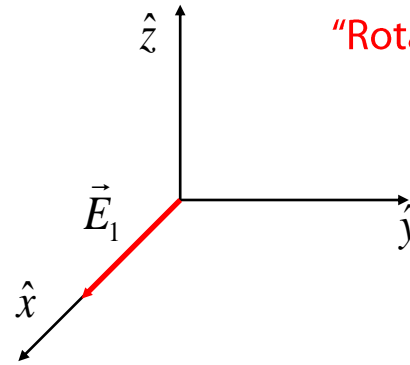
In the compact form:

$$\rightarrow \chi^{(3)}_{ijkl} = \chi^{(3)}_{xxyy} \delta_{ij} \delta_{kl} + \chi^{(3)}_{xyxy} \delta_{ik} \delta_{jl} + \chi^{(3)}_{xyyx} \delta_{il} \delta_{jk}$$

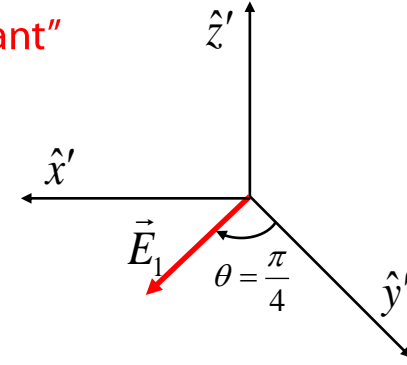
Nonlinear Birefringence (2)

Third-order susceptibility:

$$\vec{P}_{NL} = \hat{x} \frac{3}{4} \varepsilon_0 \chi_{xxxx}^{(3)} |E_1|^2 E_1 = \vec{P}'_{NL}$$



"Rotation invariant"



$$\vec{E}'_1 = \hat{x}' \frac{1}{\sqrt{2}} E_1 + \hat{y}' \frac{1}{\sqrt{2}} E_1 = \hat{x}' E'_{1,x} + \hat{y}' E'_{1,y}$$

$$\begin{aligned} P'_{NL,x'} &= \frac{3}{4} \varepsilon_0 (\chi_{xxxx}^{(3)} |E'_{1,x}|^2 E'_{1,x} + \chi_{xxyy}^{(3)} |E'_{1,y}|^2 E'_{1,x} + \chi_{xyxy}^{(3)} |E'_{1,y}|^2 E'_{1,x} + \chi_{xyyx}^{(3)} |E'_{1,y}|^2 E'_{1,x}) \\ &= \frac{3}{4} \varepsilon_0 |E_1|^2 E_1 \frac{1}{2\sqrt{2}} (\chi_{xxxx}^{(3)} + \chi_{xxyy}^{(3)} + \chi_{xyxy}^{(3)} + \chi_{xyyx}^{(3)}) \end{aligned}$$

$$\begin{aligned} P'_{NL,y'} &= \frac{3}{4} \varepsilon_0 (\chi_{yyyy}^{(3)} |E'_{1,y}|^2 E'_{1,y} + \chi_{yyxx}^{(3)} |E'_{1,x}|^2 E'_{1,y} + \chi_{yxyx}^{(3)} |E'_{1,x}|^2 E'_{1,y} + \chi_{yyxx}^{(3)} |E'_{1,x}|^2 E'_{1,y}) \\ &= \frac{3}{4} \varepsilon_0 |E_1|^2 E_1 \frac{1}{2\sqrt{2}} (\chi_{xxxx}^{(3)} + \chi_{xxyy}^{(3)} + \chi_{xyxy}^{(3)} + \chi_{xyyx}^{(3)}) \end{aligned}$$

$$\vec{P}'_{NL} = \hat{x}' P'_{NL,x'} + \hat{y}' P'_{NL,y'}$$

$$\rightarrow \chi_{xxxx}^{(3)} = \chi_{xxyy}^{(3)} + \chi_{xyxy}^{(3)} + \chi_{xyyx}^{(3)}$$

→

$$P'_{NL} = \sqrt{2} P'_{NL,x'} = \frac{3}{4} \varepsilon_0 |E_1|^2 E_1 \frac{1}{2} (\chi_{xxxx}^{(3)} + \chi_{xxyy}^{(3)} + \chi_{xyxy}^{(3)} + \chi_{xyyx}^{(3)})$$

Nonlinear Birefringence (3)

Nonlinear polarization:

$$\rightarrow \mathbf{P}'_{NL}(\mathbf{r}, t) = \varepsilon_0 \chi^{(3)} : \mathbf{E}(\mathbf{r}, t) \mathbf{E}(\mathbf{r}, t) \mathbf{E}(\mathbf{r}, t)$$

$$\rightarrow \mathbf{P}_{NL}(\mathbf{r}, t) = \frac{1}{2} (\hat{x} P_x + \hat{y} P_y) \exp(-i\omega_0 t) + c.c. \leftarrow \text{Oscillating at } \sim \omega_0$$

$$\rightarrow P_i = \frac{3\varepsilon_0}{4} \sum_j \left(\chi_{xxyy}^{(3)} E_i E_j E_j^* + \chi_{xyxy}^{(3)} E_j E_i E_j^* + \chi_{xyyx}^{(3)} E_j E_j E_i^* \right)$$

$$\leftarrow \chi_{ijkl}^{(3)} = \chi_{xxyy}^{(3)} \delta_{ij} \delta_{kl} + \chi_{xyxy}^{(3)} \delta_{ik} \delta_{jl} + \chi_{xyyx}^{(3)} \delta_{il} \delta_{jk}$$

$$\rightarrow P_x = \frac{3\varepsilon_0}{4} \chi_{xxxx}^{(3)} \left[\left(|E_x|^2 + \frac{2}{3} |E_y|^2 \right) E_x + \frac{1}{3} (E_x^* E_y) E_y \right]$$

$$\leftarrow \chi_{xxyy}^{(3)} \approx \chi_{xyxy}^{(3)} \approx \chi_{xyyx}^{(3)}$$

“Mostly electronic”

$$\rightarrow P_y = \frac{3\varepsilon_0}{4} \chi_{xxxx}^{(3)} \left[\left(|E_y|^2 + \frac{2}{3} |E_x|^2 \right) E_y + \frac{1}{3} (E_y^* E_x) E_x \right]$$

Nonlinear refractive indices:

$$\rightarrow P_j = \varepsilon_0 \varepsilon_j^{NL} E_j \rightarrow \varepsilon_j = \varepsilon_j^L + \varepsilon_j^{NL} = (n_j^L + \Delta n_j)^2$$

$$\rightarrow \Delta n_x = n_2 \left(|E_x|^2 + \frac{2}{3} |E_y|^2 \right)$$

$$\rightarrow \Delta n_y = n_2 \left(|E_y|^2 + \frac{2}{3} |E_x|^2 \right)$$

Nonlinear Birefringence (4)

Coupled-mode equations:

Electric fields: $\rightarrow E_j(\mathbf{r}, t) = F(x, y)A_j(z, t) \exp(i\beta_{0j}z)$

Coupled NLSE in terms of x and y pols.:

$$\begin{aligned} \rightarrow \frac{\partial A_x}{\partial z} + \beta_{1x} \frac{\partial A_x}{\partial t} + \frac{i\beta_2}{2} \frac{\partial^2 A_x}{\partial t^2} + \frac{\alpha}{2} A_x \\ = i\gamma \left(|A_x|^2 + \frac{2}{3} |A_y|^2 \right) A_x + \frac{i\gamma}{3} A_x^* A_y^2 \exp(-2i\Delta\beta z) \end{aligned}$$

$$\begin{aligned} \rightarrow \frac{\partial A_y}{\partial z} + \beta_{1y} \frac{\partial A_y}{\partial t} + \frac{i\beta_2}{2} \frac{\partial^2 A_y}{\partial t^2} + \frac{\alpha}{2} A_y \\ = i\gamma \left(|A_y|^2 + \frac{2}{3} |A_x|^2 \right) A_y + \frac{i\gamma}{3} A_y^* A_x^2 \exp(2i\Delta\beta z) \end{aligned}$$

$$\leftarrow \Delta\beta = \beta_{0x} - \beta_{0y} = \frac{2\pi}{\lambda} B_m = \frac{2\pi}{L_B}$$

Nonlinear Birefringence (5)

Coupled-mode equations:

Coupled NLSE in terms of circularly polarized components:

$$\rightarrow A_+ = (\bar{A}_x + i\bar{A}_y) / \sqrt{2}, \quad A_- = (\bar{A}_x - i\bar{A}_y) / \sqrt{2}$$

$$\begin{aligned} \rightarrow \frac{\partial A_+}{\partial z} + \beta_1 \frac{\partial A_+}{\partial t} + \frac{i\beta_2}{2} \frac{\partial^2 A_+}{\partial t^2} + \frac{\alpha}{2} A_+ \\ = \frac{i}{2} (\Delta\beta) A_- + \frac{2i\gamma}{3} (|A_+|^2 + 2|A_-|^2) A_+ \end{aligned}$$

$$\begin{aligned} \rightarrow \frac{\partial A_-}{\partial z} + \beta_1 \frac{\partial A_-}{\partial t} + \frac{i\beta_2}{2} \frac{\partial^2 A_-}{\partial t^2} + \frac{\alpha}{2} A_- \\ = \frac{i}{2} (\Delta\beta) A_+ + \frac{2i\gamma}{3} (|A_-|^2 + 2|A_+|^2) A_- \end{aligned}$$

$\leftarrow \beta_{1x} \approx \beta_{1y} = \beta_1 \leftarrow$ Low birefringence

Nonlinear Birefringence (6)

Elliptically birefringent fibers:

Electric field: $\rightarrow \mathbf{E}(\mathbf{r}, t) = \frac{1}{2} (\hat{e}_x E_x + \hat{e}_y E_y) \exp(-i\omega_0 t) + c.c.$

$$\leftarrow \hat{e}_x = \frac{\hat{x} + ir\hat{y}}{\sqrt{1+r^2}}, \quad \hat{e}_y = \frac{r\hat{x} - i\hat{y}}{\sqrt{1+r^2}} \leftarrow \text{Orthogonal pol. eigenvectors}$$

Coupled NLSE:

$\leftarrow r = \tan(\theta / 2)$ $\leftarrow \theta$: Ellipticity angle

$$\begin{aligned} \rightarrow \frac{\partial A_x}{\partial z} + \beta_{1x} \frac{\partial A_x}{\partial t} + \frac{i\beta_2}{2} \frac{\partial^2 A_x}{\partial t^2} + \frac{\alpha}{2} A_x & \quad \rightarrow \frac{\partial A_y}{\partial z} + \beta_{1y} \frac{\partial A_x}{\partial t} + \frac{i\beta_2}{2} \frac{\partial^2 A_y}{\partial t^2} + \frac{\alpha}{2} A_y \\ = i\gamma \left[\left(|A_x|^2 + B |A_y|^2 \right) A_x + C A_x^* A_y^2 e^{-2i\Delta\beta z} \right] & \quad = i\gamma \left[\left(|A_y|^2 + B |A_x|^2 \right) A_y + C A_y^* A_x^2 e^{2i\Delta\beta z} \right] \\ + i\gamma D \left[A_y^* A_x^2 e^{i\Delta\beta z} + \left(|A_y|^2 + 2|A_x|^2 \right) A_y e^{-i\Delta\beta z} \right] & \quad + i\gamma D \left[A_x^* A_y^2 e^{-i\Delta\beta z} + \left(|A_x|^2 + 2|A_y|^2 \right) A_x e^{i\Delta\beta z} \right] \\ \leftarrow B = \frac{2 + 2\sin^2 \theta}{2 + \cos^2 \theta}, \quad C = \frac{\cos^2 \theta}{2 + \cos^2 \theta}, \quad D = \frac{\sin \theta \cos \theta}{2 + \cos^2 \theta} \end{aligned}$$

$\rightarrow L \gg L_B$

$$\rightarrow \frac{\partial A_x}{\partial z} + \beta_{1x} \frac{\partial A_x}{\partial t} + \frac{i\beta_2}{2} \frac{\partial^2 A_x}{\partial t^2} + \frac{\alpha}{2} A_x = i\gamma \left(|A_x|^2 + B |A_y|^2 \right) A_x$$

$$\rightarrow \frac{\partial A_y}{\partial z} + \beta_{1y} \frac{\partial A_x}{\partial t} + \frac{i\beta_2}{2} \frac{\partial^2 A_y}{\partial t^2} + \frac{\alpha}{2} A_y = i\gamma \left(|A_y|^2 + B |A_x|^2 \right) A_y$$

Nonlinear Phase Shift (1)

Nondispersive XPM:

Coupled NLSE: $\rightarrow \frac{\partial A_x}{\partial z} + \frac{\alpha}{2} A_x = i\gamma \left(|A_x|^2 + B |A_y|^2 \right) A_x$

$$\rightarrow \frac{\partial A_y}{\partial z} + \frac{\alpha}{2} A_y = i\gamma \left(|A_y|^2 + B |A_x|^2 \right) A_y$$

← CW radiation

$$\leftarrow A_x = \sqrt{P_x} e^{-\alpha z/2} e^{i\phi_x}, \quad A_y = \sqrt{P_y} e^{-\alpha z/2} e^{i\phi_y}$$

$$\rightarrow \frac{\partial \phi_x}{\partial z} = \gamma e^{-\alpha z} (P_x + B P_y), \quad \frac{\partial \phi_y}{\partial z} = \gamma e^{-\alpha z} (P_y + B P_x)$$

$$\rightarrow \phi_x = \gamma (P_x + B P_y) L_{\text{eff}}, \quad \phi_y = \gamma (P_y + B P_x) L_{\text{eff}} \quad \leftarrow L_{\text{eff}} = [1 - \exp(-\alpha L)] / \alpha$$

Relative phase difference:

$$\rightarrow \Delta \phi_{NL} = \phi_x - \phi_y = \gamma L_{\text{eff}} (1 - B)(P_x - P_y)$$

$$\leftarrow B = \frac{2}{3}, \quad P_x = P_0 \cos^2 \theta, \quad P_y = P_0 \sin^2 \theta \quad (\theta: \text{Pol. angle to } \hat{x})$$

$$\rightarrow \Delta \phi_{NL} = (\gamma P_0 L_{\text{eff}} / 3) \cos(2\theta)$$

Nonlinear Phase Shift (2)

Optical Kerr effect:

Relative phase difference:

$$\rightarrow \Delta\phi = (2\pi / \lambda)(\tilde{n}_x - \tilde{n}_y)L$$

$$\leftarrow \tilde{n}_x = n_x + \Delta n_x, \quad \tilde{n}_y = n_y + \Delta n_y$$

By pump:

$$\rightarrow \Delta n_x = 2n_2 |E_p|^2 \leftarrow \mathbf{E}_p = E_p \hat{x}$$

$$\rightarrow \Delta n_y = 2n_2 b |E_p|^2 \leftarrow b = \chi_{xyxy}^{(3)} / \chi_{xxxx}^{(3)} \leftarrow b = \frac{1}{3} \leftarrow \text{Purely electronic}$$

$$\rightarrow \Delta\phi \equiv \Delta\phi_L + \Delta\phi_{NL} = (2\pi L / \lambda)(\Delta n_L + n_{2B} |E_p|^2) \leftarrow n_{2B} = 2n_2(1-b)$$

→ Kerr shutter

→ Pulse shaping