

Introduction to Photonics

Optics of Solids (1)

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Macroscopic Fields and Maxwell's Equations

Maxwell's equations:

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$

$$\nabla \times \mathbf{H} = \mathbf{J} + \frac{\partial \mathbf{D}}{\partial t}$$

$$\nabla \cdot \mathbf{D} = \rho$$

$$\nabla \cdot \mathbf{B} = 0$$

Constitutive relations:

$$\mathbf{D} = \epsilon \mathbf{E} = \epsilon_0 \mathbf{E} + \mathbf{P} \quad \leftarrow \quad \mathbf{P} = (\epsilon - \epsilon_0) \mathbf{E} = \chi \epsilon_0 \mathbf{E}$$

$$\mathbf{B} = \mu \mathbf{H} = \mu_0 \mathbf{H} + \mu_0 \mathbf{M}$$



Electric susceptibility

Current equation (Ohm's law):

$$\mathbf{J} = \sigma \mathbf{E}$$

General Wave Equation

For nonmagnetic & electrically neutral media:

$$\nabla \times \mathbf{E} = -\mu_0 \frac{\partial \mathbf{H}}{\partial t}$$

$$\nabla \times \mathbf{H} = \varepsilon_0 \frac{\partial \mathbf{E}}{\partial t} + \frac{\partial \mathbf{P}}{\partial t} + \mathbf{J}$$

$$\nabla \cdot \mathbf{E} = -\frac{1}{\varepsilon_0} \nabla \cdot \mathbf{P}$$

$$\nabla \cdot \mathbf{H} = 0$$

General wave equation:

$$\nabla \times (\nabla \times \mathbf{E}) + \frac{1}{c^2} \frac{\partial^2 \mathbf{E}}{\partial t^2} = -\mu_0 \frac{\partial^2 \mathbf{P}}{\partial t^2} - \mu_0 \frac{\partial \mathbf{J}}{\partial t}$$

Ohmic loss and metallic nature



Material characteristics:

Dispersion, absorption, birefringence, optical activity, etc.

Propagation of Light in Isotropic Dielectrics (1)

Macroscopic polarization:

$$\mathbf{P} = -Ner$$

For a static electric field:

$$-e\mathbf{E} = K\mathbf{r} \quad \leftarrow \text{Restoring force for "bound" charges}$$

Static polarization:

$$\rightarrow \mathbf{P} = \frac{Ne^2}{K} \mathbf{E}$$

Equation of motion:

$$m \frac{d^2 \mathbf{r}}{dt^2} + m\gamma \frac{\partial \mathbf{r}}{\partial t} + K\mathbf{r} = -e\mathbf{E}$$



Frictional damping force

Effective resonance frequency:

For a monochromatic electric field: $\rightarrow E \sim e^{-i\omega t}$

$$\leftarrow \omega_0 = \sqrt{\frac{K}{m}}$$

$$\rightarrow \mathbf{P} = \frac{Ne^2}{-m\omega^2 - i\omega m\gamma + K} \mathbf{E} = \frac{Ne^2 / m}{\omega_0^2 - \omega^2 - i\gamma\omega} \mathbf{E}$$

Propagation of Light in Isotropic Dielectrics (2)

Wave equation for linear dielectrics:

$$\nabla \times (\nabla \times \mathbf{E}) + \frac{1}{c^2} \frac{\partial^2 \mathbf{E}}{\partial t^2} = -\mu_0 \frac{\partial^2 \mathbf{P}}{\partial t^2} - \mu_0 \frac{\partial \mathbf{J}}{\partial t}$$

$$\nabla(\nabla \cdot \mathbf{E}) - \nabla^2 \mathbf{E} \rightarrow \nabla^2 \mathbf{E} = \frac{1}{c^2} \left(1 + \frac{Ne^2}{m\epsilon_0} \cdot \frac{1}{\omega_0^2 - \omega^2 - i\gamma\omega} \right) \frac{\partial^2 \mathbf{E}}{\partial t^2}$$

Homogeneous plane harmonic wave solution:

$$\rightarrow \mathbf{E} = \mathbf{E}_0 e^{i(\mathcal{K}z - \omega t)}$$

Dispersion relation:

$$\rightarrow \mathcal{K}^2 = \frac{\omega^2}{c^2} \left(1 + \frac{Ne^2}{m\epsilon_0} \cdot \frac{1}{\omega_0^2 - \omega^2 - i\gamma\omega} \right)$$

$$\rightarrow \mathcal{K} = k + i\alpha = \frac{\omega}{c} \mathcal{N}$$

$$\leftarrow \mathcal{N} = n + i\kappa \leftarrow \alpha = \frac{\omega}{c} \kappa$$

$$\rightarrow \mathbf{E} = \mathbf{E}_0 e^{-\alpha z} e^{i(kz - \omega t)}$$

Propagation of Light in Isotropic Dielectrics (3)

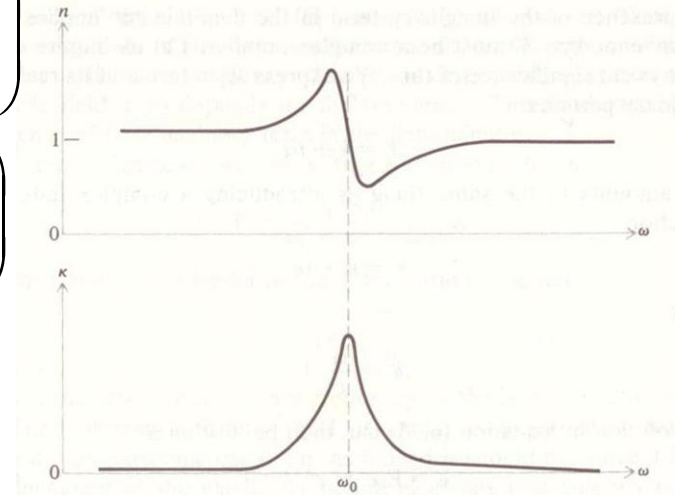
Complex index of refraction:

$$\mathcal{N}^2 = (n + i\kappa)^2 = 1 + \frac{Ne^2}{m\epsilon_0} \left(\frac{1}{\omega_0^2 - \omega^2 - i\gamma\omega} \right)$$

$$\rightarrow n^2 - \kappa^2 = 1 + \frac{Ne^2}{m\epsilon_0} \left(\frac{\omega_0^2 - \omega^2}{(\omega_0^2 - \omega^2)^2 + \gamma^2\omega^2} \right)$$

$$\rightarrow 2n\kappa = \frac{Ne^2}{m\epsilon_0} \left(\frac{\gamma\omega}{(\omega_0^2 - \omega^2)^2 + \gamma^2\omega^2} \right)$$

→ Kramers-Kronig relations (H.W.)

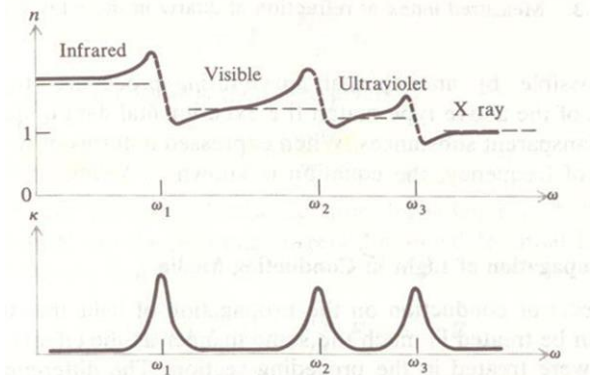


G. R. Fowles, Introduction to Modern Optics, 1975.

Complex index of refraction in general:

$$\rightarrow \mathcal{N}^2 = 1 + \frac{Ne^2}{m\epsilon_0} \sum_j \left(\frac{f_j}{\omega_j^2 - \omega^2 - i\gamma_j\omega} \right)$$

$$\rightarrow n^2 = 1 + \frac{Ne^2}{m\epsilon_0} \sum_j \left(\frac{f_j}{\omega_j^2 - \omega^2} \right) \rightarrow \text{c.f. Sellmeier's formula}$$



G. R. Fowles, Introduction to Modern Optics, 1975.

Propagation of Light in Conducting Media (1)

Equation of motion:

$$m \frac{d\mathbf{v}}{\partial t} + m\tau^{-1}\mathbf{v} = -e\mathbf{E} \quad \leftarrow \text{No restoring force}$$

$$\leftarrow \mathbf{J} = -Ne\mathbf{v}$$

$$\rightarrow \frac{d\mathbf{J}}{\partial t} + \tau^{-1}\mathbf{J} = \frac{Ne^2}{m}\mathbf{E}$$

For a static electric field:

$$\rightarrow \mathbf{J} = \frac{Ne^2}{m}\tau\mathbf{E} = \sigma\mathbf{E}$$

For a monochromatic electric field:

$$\rightarrow \mathbf{J} = \frac{\sigma}{1-i\omega\tau}\mathbf{E}$$

Wave equation for linear conducting media:

$$\nabla \times (\nabla \times \mathbf{E}) + \frac{1}{c^2} \frac{\partial^2 \mathbf{E}}{\partial t^2} = -\mu_0 \frac{\partial^2 \mathbf{P}}{\partial t^2} - \mu_0 \frac{\partial \mathbf{J}}{\partial t}$$

$$\rightarrow \nabla^2 \mathbf{E} = \frac{1}{c^2} \frac{\partial^2 \mathbf{E}}{\partial t^2} + \frac{\mu_0 \sigma}{1-i\omega\tau} \frac{\partial \mathbf{E}}{\partial t}$$

Propagation of Light in Conducting Media (2)

Homogeneous plane harmonic wave solution:

$$\rightarrow \mathbf{E} = \mathbf{E}_0 e^{i(\mathcal{K}z - \omega t)}$$

Dispersion relation:

$$\rightarrow \mathcal{K}^2 = \frac{\omega^2}{c^2} + \frac{i\omega\mu_0\sigma}{1 - i\omega\tau}$$

For very low frequencies:

$$\rightarrow \mathcal{K}^2 \approx i\omega\mu_0\sigma$$

$$\rightarrow \mathcal{K} \approx \sqrt{i\omega\mu_0\sigma} = (1+i)\sqrt{\omega\mu_0\sigma/2} = k + i\alpha$$

$$\rightarrow k \approx \alpha \approx \sqrt{\frac{\omega\mu_0\sigma}{2}} \quad \rightarrow \quad \delta \approx \frac{1}{\alpha} \approx \sqrt{\frac{2}{\omega\mu_0\sigma}} \quad \leftarrow \text{Skin depth}$$

$$\rightarrow n \approx \kappa \approx \sqrt{\frac{\sigma}{2\omega\epsilon_0}}$$

Propagation of Light in Conducting Media (3)

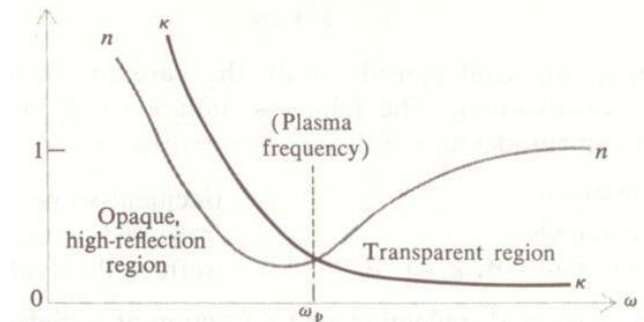
Complex index of refraction in general:

$$\rightarrow \mathcal{N}^2 = 1 - \frac{\omega_p^2}{\omega + i\omega\tau^{-1}} \quad \leftarrow \quad \omega_p = \sqrt{\frac{\mu_0 \sigma c^2}{\tau}} = \sqrt{\frac{Ne^2}{m\epsilon_0}}$$

$$\rightarrow n^2 - \kappa^2 = 1 - \frac{\omega_p^2}{\omega^2 + \tau^{-2}}$$

$$\rightarrow 2n\kappa = \frac{\omega_p^2}{\omega^2 + \tau^{-2}} \left(\frac{1}{\omega\tau} \right)$$

← Plasma frequency



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For poor conductors and semiconductors:

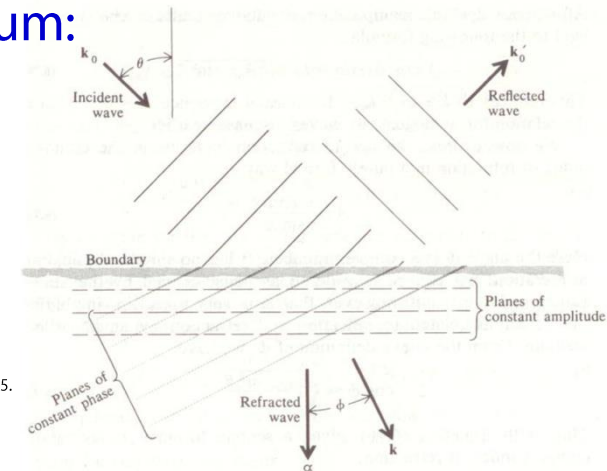
$$\rightarrow \mathcal{N}^2 = 1 - \frac{\omega_p^2}{\omega + i\omega\tau^{-1}} + \frac{Ne^2}{m\epsilon_0} \sum_j \left(\frac{f_j}{\omega_j^2 - \omega^2 - i\gamma_j\omega} \right)$$

Reflection and Refraction: Absorbing Media (1)

Complex index of refraction of an absorbing medium:

$$\mathcal{N} = n + i\kappa$$

$$\rightarrow \mathcal{H} = \mathbf{k} + i\boldsymbol{\alpha}$$



G. R. Fowles, Introduction to Modern Optics, 1975.

Boundary conditions:

$$\rightarrow \mathbf{k}_0 \cdot \mathbf{r} = \mathbf{k}'_0 \cdot \mathbf{r}$$

$$\rightarrow \mathbf{k}_0 \cdot \mathbf{r} = \mathcal{H} \cdot \mathbf{r} = (\mathbf{k} + i\boldsymbol{\alpha}) \cdot \mathbf{r} \rightarrow \mathbf{k}_0 \cdot \mathbf{r} = \mathbf{k} \cdot \mathbf{r} \rightarrow k_0 \sin \theta = k \sin \phi$$

$$\rightarrow \boldsymbol{\alpha} \cdot \mathbf{r} = 0$$

Wave equation:

$$\rightarrow \nabla^2 \mathbf{E} = \frac{\mathcal{N}^2}{c^2} \frac{\partial^2 \mathbf{E}}{\partial t^2} \rightarrow (\mathbf{k} + i\boldsymbol{\alpha}) \cdot (\mathbf{k} + i\boldsymbol{\alpha}) = (n + i\kappa)^2 k_0^2$$

$$\rightarrow k^2 - \alpha^2 = (n^2 - \kappa^2) k_0^2$$

$$\rightarrow \mathbf{k} \cdot \boldsymbol{\alpha} = k\alpha \cos \phi = n\kappa k_0^2$$

Complex index of refraction:

$$\rightarrow k \cos \phi + i\alpha = k_0 \sqrt{\mathcal{N}^2 - \sin^2 \theta}$$

$$\rightarrow \mathcal{N} \sin \phi = \sin \theta \rightarrow \cos \phi = \sqrt{1 - \frac{\sin^2 \theta}{\mathcal{N}^2}} \rightarrow \mathcal{N} = \frac{k \cos \phi + i\alpha}{k_0 \cos \phi}$$

\rightarrow c.f. Snell's law

Reflection and Refraction: Absorbing Media (2)

Electric and magnetic fields:

$$\mathbf{E}, \mathbf{H} = \frac{1}{\mu_0 \omega} \mathbf{k}_0 \times \mathbf{E} \quad \leftarrow \text{Incident}$$

$$\mathbf{E}', \mathbf{H}' = \frac{1}{\mu_0 \omega} \mathbf{k}'_0 \times \mathbf{E}' \quad \leftarrow \text{Reflected}$$

$$\mathbf{E}'', \mathbf{H}'' = \frac{1}{\mu_0 \omega} \mathcal{N} \times \mathbf{E}'' = \frac{1}{\mu_0 \omega} (\mathbf{k} \times \mathbf{E}'' + i\alpha \times \mathbf{E}'') \quad \leftarrow \text{Refracted}$$

Boundary conditions (TE):

$$E + E' = E''$$

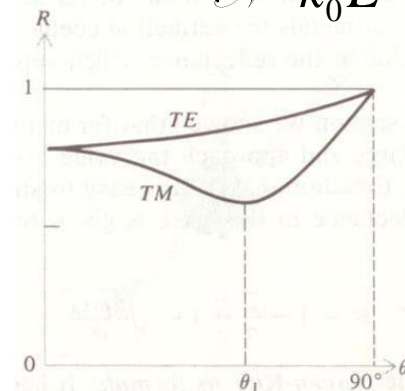
$$H_t + H'_t = H''_t \rightarrow -k_0 E \cos \theta + k_0 E' \cos \theta = -k E'' \cos \phi + i\alpha E''$$

$$= -\mathcal{N} k_0 E'' \cos \phi$$

Reflection coefficients:

$$\rightarrow r_s = \frac{\cos \theta - \mathcal{N} \cos \phi}{\cos \theta + \mathcal{N} \cos \phi} \quad (TE)$$

$$\rightarrow r_p = \frac{-\mathcal{N} \cos \theta + \cos \phi}{\mathcal{N} \cos \theta + \cos \phi} \quad (TM)$$



G. R. Fowles, Introduction to Modern Optics, 1975.