XRD-4 Peak broadening

Size broadening, Strain broadening Size/strain broadening

> XRD-4 Read

Cullity Chapter 5-1, 5-2, 5-4

Jenkins & Snyder chap 3.9.2; 3.9.3 (p89~p94)

Hammond chap 9.3

Cullity Chapter 5-5, 5-6

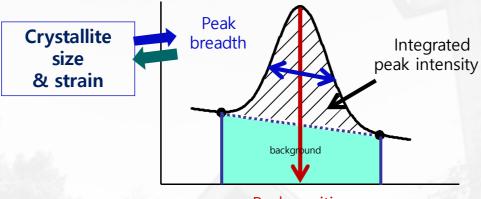
Krawitz chap 11.6 (p343~p346)

Cullity Chapter 14-1. 14-2, 14-3, 14-4, 14-6

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Shape of Peak

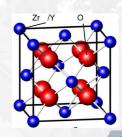


Peak position

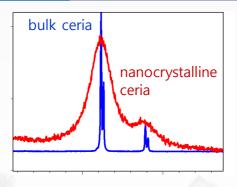
Crystal Structure of "cubic" "ZrO2"

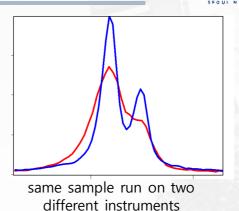
Space Group Fm3m (225) → cubic Lattice Parameter a=5.11

Atom	х	у	Z	B _{iso}	occupancy
Zr	0	0	0	1,14	1
0	0.25	0.25	0.25	2.4	1



Peak Broadening





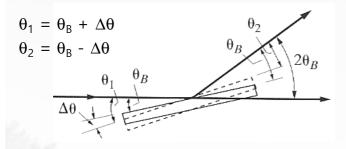
- ➤ Peak broadening ←
 - ✓ Small crystallite size
 - ✓ Stacking faults, Microstrain, and other Defects in the crystal structure
 - ✓ An inhomogeneous composition in a solid solution or alloy
- > Different instrument configurations can change the peak width, too.
 - ← Instrument contribution

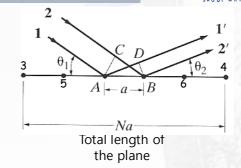
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Scott A. Speakman

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Geometrical factor - 1 of Lorenz Factor



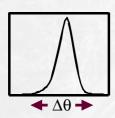


$$\delta_{1'2'} = 2a\Delta\theta \sin\theta_{B_s}$$

Path difference b/w rays scattered by atoms at $2Na~\Delta\theta~\sin~\theta_B$ Path difference by wild associated either end of the plane (3 & 4)

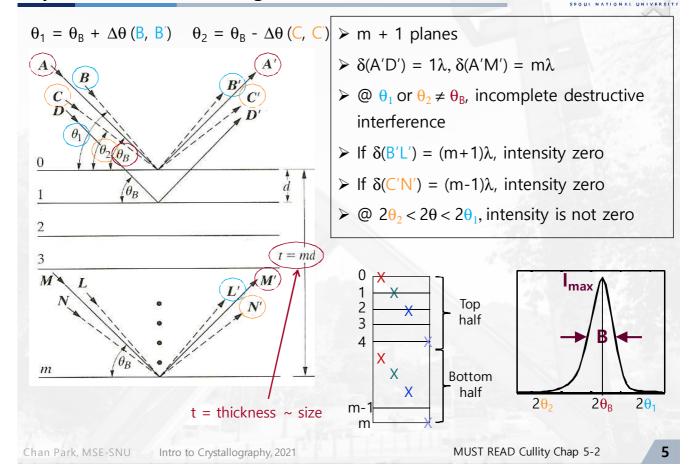
Diffracted intensity = zero when $2Na \Delta\theta \sin \theta_B = \lambda$

$$2Na\ \Delta\theta\ \sin\ \theta_B = \lambda$$



$$\Delta \theta = \frac{\lambda}{2 Na \sin \theta_B} \quad I_{\text{max}} \propto 1/\sin \theta$$

Max angular range of crystal rotation over which appreciable energy can be diffracted in the direction $2\theta_B$



Crystallite size broadening

 $\theta_1 = \theta_B + \Delta\theta \, (B, B)$ $\theta_2 = \theta_B - \Delta\theta \, (C, C)$ > Ray B; $\theta_1 = \theta_B + \Delta\theta$, $\delta(B'L') = (m+1)\lambda$ > Compare when m = 10 and m = 10,000 > $m = 10 \rightarrow (B'L') = \delta(0m) = 11\lambda$, $\delta(01) = 1.1\lambda$ > $m = 10,000 \rightarrow (B'L') = \delta(0m) = 10,001\lambda$, $\delta(01) = 1.0001\lambda$ > $\theta_1 = (m+1)\lambda$ > $\theta_2 \downarrow$ as $m \downarrow$ > $\theta_1 \uparrow$ as $m \downarrow$ > $\theta_2 \downarrow$ as $m \downarrow$ > $\theta_1 \uparrow$ as $m \downarrow$ > $\theta_2 \downarrow$ as $m \downarrow$ > $\theta_2 \downarrow$ as $m \downarrow$ > $\theta_3 \uparrow$ as thickness $\phi_3 \uparrow$ > $\phi_3 \uparrow$ as $\phi_3 \uparrow$ > $\phi_3 \uparrow$ = $\phi_3 \uparrow$ > $\phi_3 \uparrow$ = $\phi_3 \uparrow$ > $\phi_3 \uparrow$ > $\phi_3 \uparrow$ = $\phi_3 \uparrow$ > $\phi_3 \uparrow$

 $t(\sin\theta_1 - \sin\theta_2) = \lambda$

 $2\theta_2$

 $2\theta_B$ $2\theta_1$

 $2t\sin\,\theta_2=(m-1)\lambda$

Assume diffraction line is triangular in shape

$$B = \frac{1}{2}(2\theta_1 - 2\theta_2) = \theta_1 - \theta_2.$$

$$2t \sin \theta_1 = (m+1)\lambda$$

$$2t \sin \theta_2 = (m-1)\lambda$$

$$t\left(\sin\,\theta_1-\sin\,\theta_2\right)=\lambda,$$

$$2t\cos\left(\frac{\theta_1+\theta_2}{2}\right)\sin\left(\frac{\theta_1-\theta_2}{2}\right)=\lambda$$

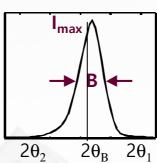
$$\theta_1 + \theta_2 = 2\theta_B \text{(approx.)}$$

$$\sin\left(\frac{\theta_1-\theta_2}{2}\right) = \left(\frac{\theta_1-\theta_2}{2}\right) \text{ (approx.)}$$

$$2t\left(\frac{\theta_1-\theta_2}{2}\right)\cos\,\theta_B=\lambda$$

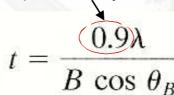
$$t = \frac{\lambda}{B \cos \theta_B}$$

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B; an angular width, in terms of 2θ (not a linear width)

Shape factor; depends on the shape of the crystallites



Scherrer equation

Cullity p169

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plane

2nd plane

3rd

plane

Crystallite size broadening

- \triangleright @ θ_B ; ABC = λ , DEF = $2\lambda \rightarrow$ diffraction peak
- \triangleright ABC = 0.5 λ , DEF = 1 λ \rightarrow no diffraction peak
- \triangleright ABC = 1.1 λ , DEF = 2.2 λ

 \rightarrow PD (path diff.) in 6th plane = 5.5 $\lambda \rightarrow$ 1' & 6' out of phase \rightarrow no net diffraction

- \triangleright ABC = 1.001 $\lambda \rightarrow 1'$ & 501' out of phase; ABC = 1.00001 $\lambda \rightarrow 1'$ & 50001' out of phase $\rightarrow \rightarrow$ Sharp diffraction peak @ θ_B
- ➤ When crystal is only 100nm in size, 5000' or 50000' are not present
- \triangleright Peak begins to show intensity at a lower θ and ends at a higher θ than $\theta_B \rightarrow$ particle size broadening
- ➤ Crystallites smaller than 1um can cause broadening → size can be determined using the peak width (← incomplete destructive interference)





size broadening; degree of being "out-of-phase" that can be tolerate

$$\blacktriangleright$$
 In case λ = 1.5 Å, d= 1.0 Å, θ = 49 °, $t = \frac{0.9\lambda}{B \cos \theta_B}$

- \rightarrow 1mm(millimeter) diameter crystal \rightarrow 10⁷ parallel lattice planes, ~10⁻⁷ radian*, $\sim 10^{-5}$ degree \rightarrow too small to observe
- ➤ 500 Å diameter crystal → 500 parallel lattice planes, ~10⁻³ radian, ~0.2 degree → measurable
- ➤ Non-parallel incident beam, non-monochromatic incident beam → diffraction @ angles not exactly satisfying Bragg's law → line broadening

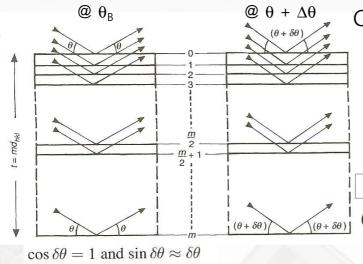
* B = $(0.9 \times 1.5 \times 10^{-10})/(10^{-3} \times \cos 49^{\circ}) \sim 2 \times 10^{-7} \text{ rad}$

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Cullity, p170

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Crystallite size broadening

between planes 0 & (m/2)

Constructive interference at angle θ $(m/2)\lambda = (m/2)2d_{hkl}\sin\theta$

Destructive interference at angle $\theta + \delta\theta$

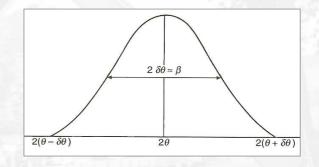
$$(m/2)\lambda + \lambda/2 = (m/2)2d_{hkl}\sin(\theta + \delta\theta)$$
$$\delta(1, (m/2)+1) = 0.5\lambda$$

 $(m/2)\lambda + \lambda/2 = (m/2)2d_{hkl}\sin\theta + (m/2)2d_{hkl}\cos\theta\,\delta\theta$

$$md_{hkl} = t$$
$$2 \, \delta\theta = \frac{\lambda}{t \cos \theta} = \beta$$

$$B = \beta = \frac{\lambda}{t \cos \theta} = \frac{\lambda \sec \theta}{t}$$

Scherrer equation



Scherrer equation

$$t = \frac{0.9\lambda}{B \cos \theta_B}$$

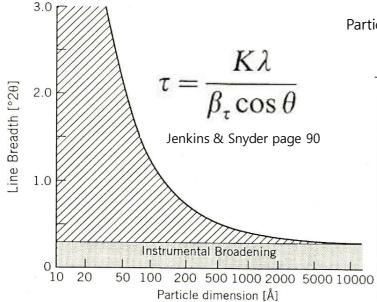
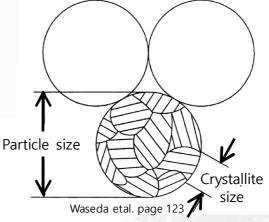
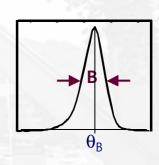


Figure 3.21. Line width as a function of particle dimension.





Waseda & Matsubara, X-ray diffraction Crystallography, Springer, 2011

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Particle Size vs. Crystallite Size

Particles can be individual crystallites.



Particle size = crystallite size

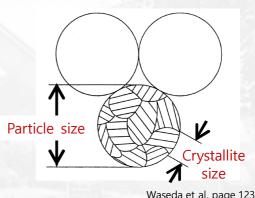
Particles may be imperfect single crystals.



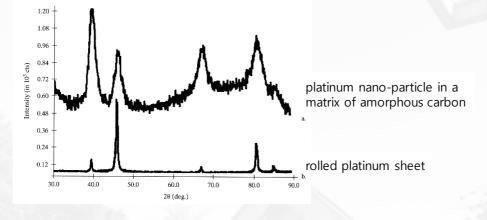


- Individual crystallites are perfect.
- **Boundaries**
 - Dislocations
 - Twin walls
 - Anti-phase walls
 - Stacking faults

From presentation of Dr. Mark Rodriguez @ DXC 2017 "What usually causes trouble?"



Waseda et al. page 123



hkl	FWHM (°2θ)	t (Å)
111	1.9	50
200	1.7	55
220	2.1	50
311	2.5	45-50

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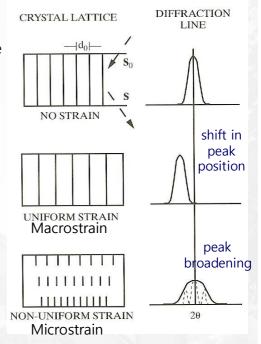
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Strain broadening

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Strain/Stress

- ➤ Macrostrain/Macrostress → shift in peak position
 - ✓ Stress is uniformly compressive or tensile over large distances.
 ← lattice parameter measurement
- ➤ Microstrain/Microstress → peak broadening
 - ✓ Distribution of both tensile & compressive stress → distribution of d-values
 - ✓ Can come from dislocations, vacancies, defects, shear planes, thermal expansion/contraction, etc.



Cullity page 176

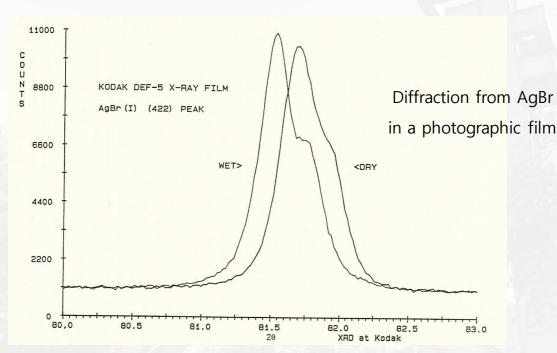
Jenkins & Snyder, page 91~93

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Peak shift ← macrostrain



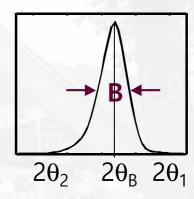
Differential expansion between the film substrate & AgBr causes macrostrain → changes lattice parameter → peak shift

Intro to Crystallography, 2021

Jenkins & Snyder, page 92

Strain broadening

- $> \lambda = 2d \sin\theta$
- $> 0 = 2d \cos\theta \delta\theta + 2 \sin\theta \delta d$
- $\triangleright \Delta(2\theta) = -2(\delta d/d)\tan\theta = B$; extra broadening produced by microstrain
- $> \beta_{\varepsilon} = 4\varepsilon \tan\theta$ (Jenkins & Snyder p93)
- $> \beta = -2\varepsilon \tan\theta$ (Hammond p266)



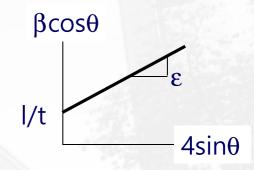
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Size & Strain broadening

 \triangleright β_{size} = λ /(tcosθ), β_{strain} = 4ε tanθ, $\underline{\beta}_{instrument}$

- \triangleright $\beta(total) = \lambda/(tcos\theta) + 4\epsilon tan\theta + \beta_{instrument}$
- $> \beta = \lambda/(t\cos\theta) + 4\epsilon (\sin\theta/\cos\theta)$
- \triangleright $\beta \cos\theta = \lambda/t + 4\epsilon \sin\theta$
- \triangleright $\beta \cos\theta/\lambda = 1/t + (4\epsilon \sin\theta)/\lambda$



LaB₆ (SRM 660c)

 \triangleright plot βcosθ/λ vs sinθ/λ (Williamson-Hall plot) → can separate size & strain contributions to line broadening --- semi-quantitative

Broadening

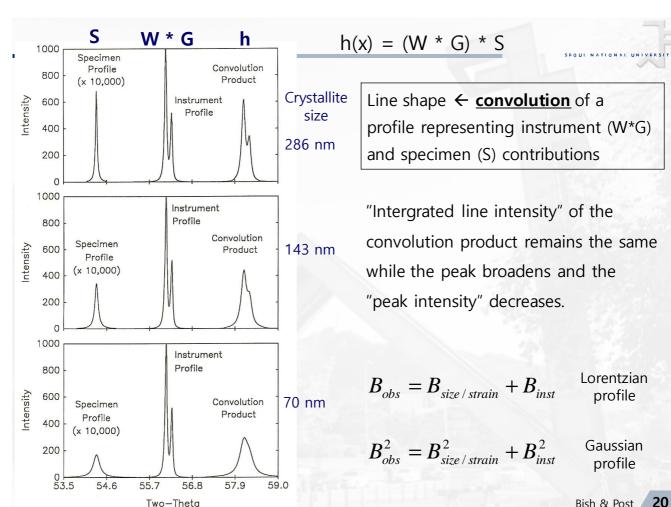
- > Darwin width
 - ✓ Incident photon is confined to certain volume.
 - ✓ Result of uncertainty principle ($\Delta p \Delta x = h$) --- Location of the photon in a xtal is restricted to a certain volume.
 - ✓ Δp must be finite. $\rightarrow \Delta \lambda$ must be finite. \rightarrow finite width of diffraction peak
- > Specimen contribution (S)
- > Spectral distribution (radiation source contribution) (W)
- > Instrumental contribution (G)
- ➤ (W * G) ← X-ray source image, flat specimen, axial divergence of incident beam, specimen transparency, receiving slit, etc.
- > (W * G); fixed for a particular instrument/target system → instrumental profile g(x)
- \triangleright Overall line profile h(x) = (W * G) * S + background = g(x) * S + BKG

LaB₆ SRM

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Standard Reference Materials (SRMs)

- ➤ Powder Line Position + Line Shape Std for Powder Dif
 - ✓ Silicon (SRM 640f); \$745/7.5q
- ➤ Line position Fluorophlogopite mica (SRM 675); \$809/7.5g
- > Line profile LaB₆ (SRM 660c); \$907/6q

No broadening from size & strain

- > Intensity
 - ✓ ZnO, TiO₂ (rutile), Cr₂O₃, CeO₂ (SRM 674b); out of stock
- > Quantitative phase analysis
 - ✓ Al₂O₃ (SRM 676a); out of stock, Silicon Nitride (SRM 656); \$580/20g
- > Instrument Response Std
 - ✓ Alumina plate (SRM 1976c); \$721/1 disc

Prices; 2021-06-17 www.nist.gov/srm/index.cfm

Gold \$58.66 / gram (2021-06-17)goldprice.org

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Conventional Williamson-Hall Plot

➤ Size + Strain

$$\frac{B\cos\theta}{\lambda} = \frac{0.9}{d} - 2\frac{\Delta d}{d} \frac{\sin\theta}{\lambda}$$

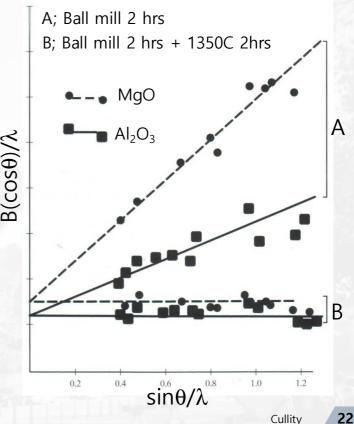
$$y = a + bx$$



✓ Horizontal line

➤ Size << strain

✓ Linear function



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The effect of dislocation contrast on x-ray line broadening: A new approach to line profile analysis

T. Ungára) and A. Borbély Institute for General Physics, Eötvös University Budapest, H-1445 Múzeum krt. 6-8, Budapest VIII, P.O.B. 323, Hungary

(Received 29 May 1996; accepted for publication 6 September 1996)

> explained strain broadening by dislocations.

$$\frac{B\cos\theta}{\lambda} = \frac{0.9}{d} + \Delta K^D \qquad y = a + X$$

Classical
$$X = -2\frac{\Delta d}{d} \frac{\sin \theta}{\lambda}$$

Modified
$$X = A(\rho^*)^{1/2} + A'(Q^*)^{1/2}$$

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Modified Williamson-Hall Plot

$$\frac{B\cos\theta}{\lambda} = \frac{0.9}{d} + A(\rho^*)^{1/2} + A'(Q^*)^{1/2}$$

- ho * : (formal) dislocation density
- $ightharpoonup Q^*$: (formal) two-particle correlations in the dislocation ensemble
- $\triangleright A.A'$: parameter determined by dislocations
- > True values of dislocation density, correlation factor

$$\rho^* = \rho (\pi g^2 b^2 \overline{C})/2$$
 $Q^* = Q(\pi g^2 b^2 \overline{C})^2/4$

- $\frac{C}{C} \text{ :average contrast factor of dislocation}$ $\frac{b}{A} \text{ :Burgers vector of dislocation}$ $\frac{b}{A} \text{ Particular reflection}$ $g = \frac{2\sin\theta}{\lambda}$

$$y = a + X$$

Conventional

$$X = -2\frac{\Delta d}{d} \frac{\sin \theta}{\frac{\lambda}{x}}$$

$$= bx$$

$$K = 2\frac{\sin\theta}{\lambda}$$

Modified

$$X = (\pi Ab^{2}/2)^{1/2} \rho^{1/2} \frac{2\sin\theta}{\lambda} C^{1/2}$$

$$+ (\pi A'b^{2}/2)Q^{1/2} \left(\frac{2\sin\theta}{\lambda} C^{1/2}\right)^{2}$$

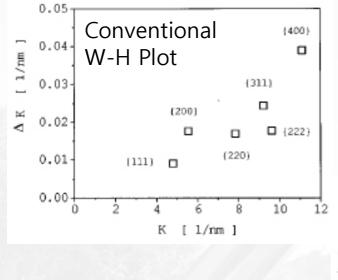
$$= b'x' + b''x'^{2}$$

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Appl. Phys. Lett. 69 (21),3173 (1996)

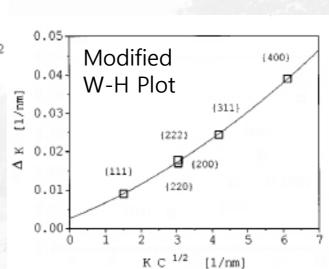
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Conventional vs. Modified W-H Plot



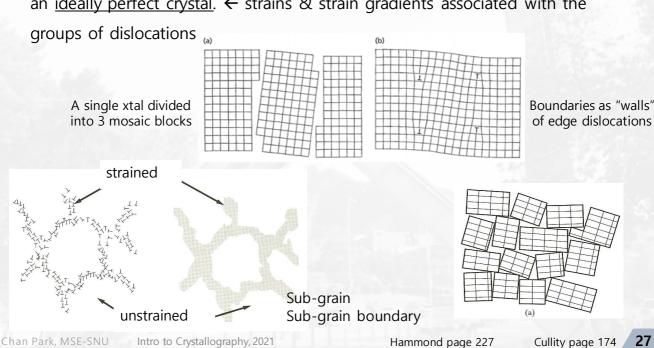
Appl. Phys. Lett. 69 (21),3173 (1996)

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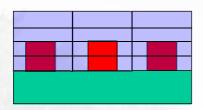
Mosaic structure, mosaic blocks

- ➤ Angle of disorientation between the tiny blocks is ε. → diffraction occurs at all angles between θ_B and θ_B + ε.
- ➤ Increases the integrated intensity relative to that obtained (or calculated) for an ideally perfect crystal. ← strains & strain gradients associated with the groups of dislocations



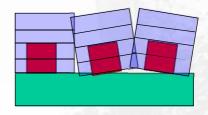
Mosaic spread

➤ Mosaicity is created by slight misorientations of different crystals as they nucleate and grow on the substrate. When the crystals join, they form boundaries.



In an ideal case, each nuclei (red) is perfectly oriented.

When the crystals grow and meet, there is perfect bounding between the crystallites → no boundary.

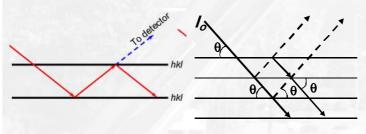


If the nuclei (red) are slightly misaligned, then boundaries will be formed.

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Ideally imperfect crystal

- ➤ Diffracted intensity; perfect xtal << ideally imperfect xtal
- ➤ Decrease in intensity as the crystal becomes more perfect (large mosaic blocks)
- > Ideally imperfect crystal consists of very small mosaic blocks, uniformly disoriented. → no extinction
- Kinematical theory vs. dynamical theory
- > Powder specimens should be ground as fine as possible.
- ➤ Grinding → reduce crystal size, increase # of diffraction cones, decrease mosaic block size, disorient mosaic blocks, strain the crystals non-uniformly



Primary Extinction

- Does not kill the reflection but lower intensity.
- ➤ How to avoid? give some stress (increase mosaicity by e.g. LN2 quenching, heat & quenching, etc.)

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Cullity page 177, 360

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todos

> XRD-4 Read

Cullity Chapter 5-1, 5-2, 5-4

Jenkins & Snyder chap 3.9.2; 3.9.3 (p89~p94)

Hammond chap 9.3

Cullity Chapter 5-5, 5-6

Krawitz chap 11.6 (p343~p346)

Cullity Chapter 14-1. 14-2, 14-3, 14-4, 14-6

> XRD-4 Homework (due in 1 week) Cullity 5-1; 5-3