

# XRD – 6

## Diffraction

### XRD-6 READ

Sherwood & Cooper, Chap 4.13, 5.6, 6.1~6.7, 6.14

Sherwood & Cooper, Chap 6.8~6.14

Ott, Chapter 13.4

Sherwood & Cooper, Chap 4.1~4.3, 4.11, 5.1~5.3

Sherwood & Cooper, Chap 4.4~4.10, 5.4, 5.7~5.10

## Waves

### ➤ Plane wave

- ✓ direction of propagation is left to right
- ✓ wave fronts are a series of parallel lines

### ➤ Circular wave

- ✓ direction of propagation is radial
- ✓ wave fronts are a set of concentric circles

### ➤ Transverse wave

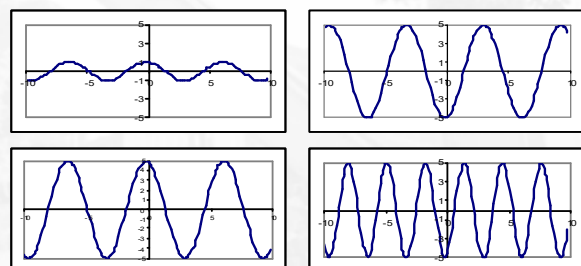
- ✓ local motion is  $\perp$  to the direction of propagation
- ✓ water waves, waves on a string, X-ray

### ➤ Longitudinal wave

- ✓ local motion is  $//$  to the direction of propagation
- ✓ sound wave

$$f(x) = A \cos(kx + \phi)$$

Amplitude    Frequency    Phase



## ➤ 1-D

Wave equations

$$\frac{\partial^2 \psi(x, t)}{\partial x^2} = \frac{1}{v^2} \frac{\partial^2 \psi(x, t)}{\partial t^2}$$

$$\psi(x, t) = \sum_n \psi_n \cos(k_n x - \omega_n t + \phi_n) + \sum_m \psi_m \cos(k_m x - \omega_m t + \phi_m)$$

$$\psi(x, t) = \psi_0 \cos(kx - \omega t)$$

## ➤ 3-D

$$\psi(\mathbf{r}, t) = \sum_n \psi_n \cos(\mathbf{k}_n \cdot \mathbf{r} - \omega_n t + \phi_n) + \sum_m \psi_m \cos(\mathbf{k}_m \cdot \mathbf{r} - \omega_m t + \phi_m)$$

$$\psi(x, y, z, t) = \psi_0 \cos(k_x x + k_y y + k_z z - \omega t)$$

$$\psi(\mathbf{r}, t) = \psi_0 \cos(\mathbf{k} \cdot \mathbf{r} - \omega t)$$

Wave vector  $|\mathbf{k}| = k = \frac{2\pi}{\lambda}$

Phase velocity  $v = \frac{\omega}{k}$

Wave frequency  $\omega = \frac{2\pi}{\tau} = 2\pi\nu$

Group velocity  $g = \frac{d\omega}{dk}$

Intensity  $I = |\psi(\mathbf{r}, t)|^2$

## Fourier transform

## ➤ Joseph Fourier

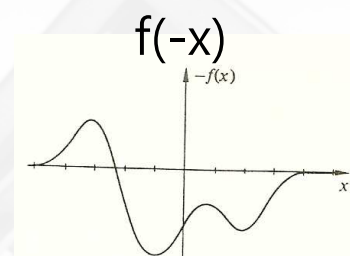
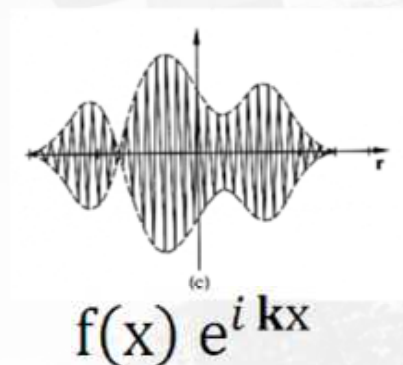
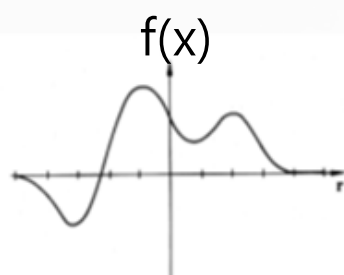
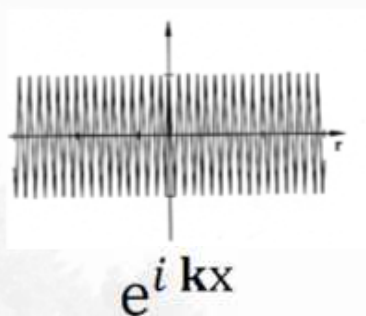
- ✓ Joined Napoleon's army as a scientific adviser
- ✓ Developed the theory on heat propagation using expanded series of sinusoids
- ✓ A periodic function can be described as the sum of simple sine and cosine functions that have wavelengths as integrals of the function



## ➤ Fourier transform; integral of a set of periodic (imaginary) primitive functions

## ➤ Fourier transforms are used in:

- ✓ crystallography; image analysis; (NMR) spectroscopy
- ✓ electron microscopy; optics & light-microscopy
- ✓ communication



$$F(k) = \int_{-\infty}^{\infty} f(x)e^{ikx} dx$$

## Fourier transform

### ➤ Fourier transform of $f(x)$

$$F(k) = \int_{\text{all } r} f(r)e^{i\mathbf{k}\cdot\mathbf{r}} dr = T f(r)$$

$$F(k) = \int_{-\infty}^{\infty} f(x)e^{ikx} dx$$

$$F(k) = T f(x)$$

### ➤ Inverse Fourier transform of $F(k)$

$$f(r) = \int_{\text{all } k} F(k)e^{-i\mathbf{k}\cdot\mathbf{r}} dk = T^{-1} F(k)$$

$$f(x) = \int_{-\infty}^{\infty} F(k)e^{-ikx} dk$$

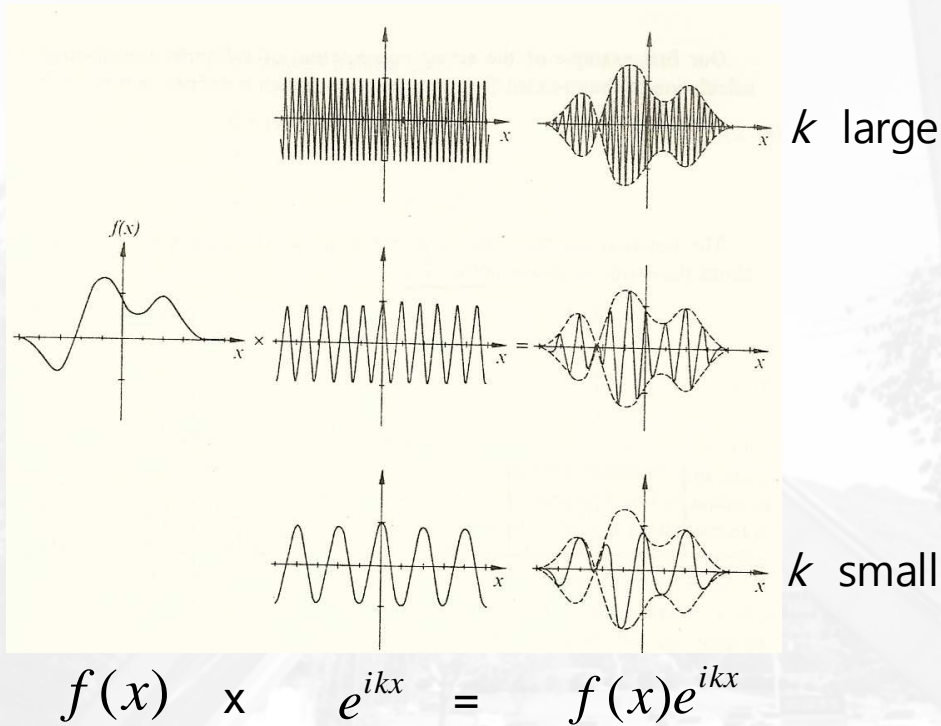
$$f(x) = T^{-1} F(k)$$

Fourier mates  $f(x) \leftrightarrow F(k)$

- Space defined by  $r \rightarrow$  real space
- Space defined by  $k \rightarrow k$  space, reciprocal space, Fourier space

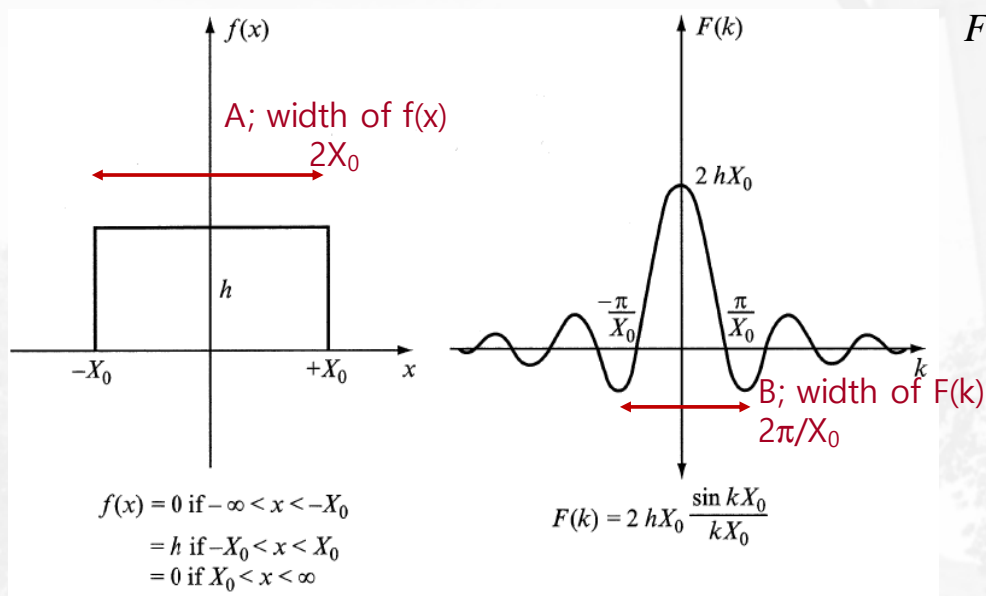
## Fourier transform

$$F(k) = \int_{-\infty}^{\infty} f(x) e^{ikx} dx$$



## Fourier transform

$$F(k) = \int_{-\infty}^{\infty} f(x) e^{ikx} dx$$



As  $A \nearrow$ ,  $B \searrow$

The narrower the  $f(x)$ , the wider the  $F(k)$

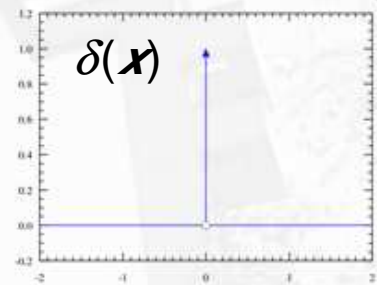
## Dirac delta function

- Unit impulse function

- $\delta(x-x_0) = \infty$ , when  $x = x_0$   
 $= 0$ , elsewhere

$$\int_0^{\infty} \delta(x-x_0) dx = 1$$

an infinitely high, infinitely thin spike at the origin, with total area one under the spike, and physically represents an idealized point mass or point charge.



Crystal Lattice; 3D array of points in space

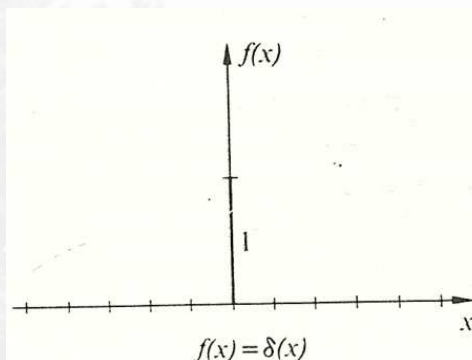
Crystal Lattice can be described using an array of delta function

- Lattice point  $\mathbf{r} = p\mathbf{a} + q\mathbf{b} + r\mathbf{c}$
- Lattice = array of  $\delta(\mathbf{r} - [p\mathbf{a} + q\mathbf{b} + r\mathbf{c}])$
- Lattice function  $l(\mathbf{r}) = \sum_{\text{all } p,q,r} \delta(\mathbf{r} - [p\mathbf{a} + q\mathbf{b} + r\mathbf{c}])$

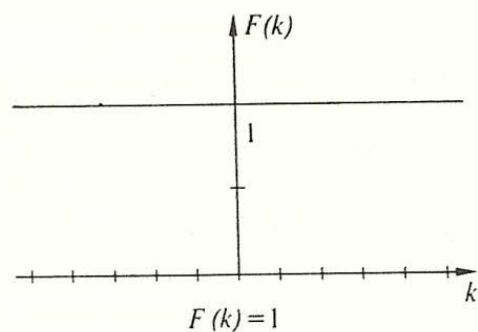
## Fourier transform > 1 delta function

$$\begin{aligned} F(k) &= \int_{-\infty}^{\infty} \delta(x) e^{ikx} dx = [e^{ikx}]_{x=0} \\ &= e^0 \\ &= 1 \end{aligned}$$

$$F(k) = \int_{-\infty}^{\infty} f(x) e^{ikx} dx$$



$$f(x) = \delta(x)$$



$$F(k) = 1$$

## Fourier transform > 2 delta functions

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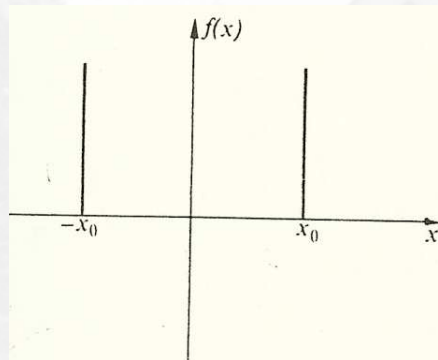
$$F(k) = \int_{-\infty}^{\infty} f(x) e^{ikx} dx$$

$$= \int_{-\infty}^{\infty} \delta(x+x_0) e^{ikx} dx + \int_{-\infty}^{\infty} \delta(x-x_0) e^{ikx} dx$$

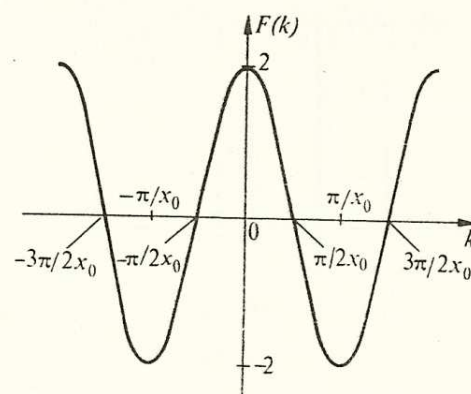
$$= e^{-ikx_0} + e^{ikx_0}$$

$$\delta = kx_0$$

$$F(x) = 2 \cos kx_0$$



$$f(x) = \delta(x+x_0) + \delta(x-x_0)$$



$$F(k) = 2 \cos kx_0$$

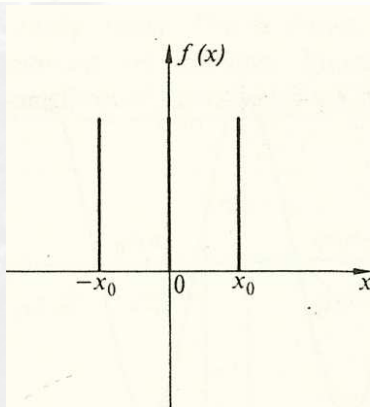
## Fourier transform > 3 delta functions

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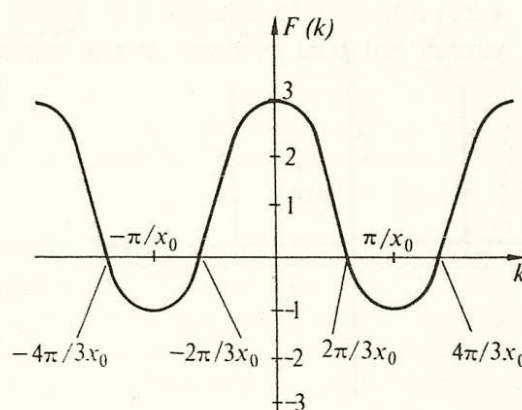
$$F(k) = \int_{-\infty}^{\infty} f(x) e^{ikx} dx$$

$$= \int_{-\infty}^{\infty} \delta(x+x_0) e^{ikx} dx + \int_{-\infty}^{\infty} \delta(x) e^{ikx} dx + \int_{-\infty}^{\infty} \delta(x-x_0) e^{ikx} dx$$

$$= e^{-ikx_0} + 1 + e^{ikx_0} = 1 + 2 \cos kx_0$$



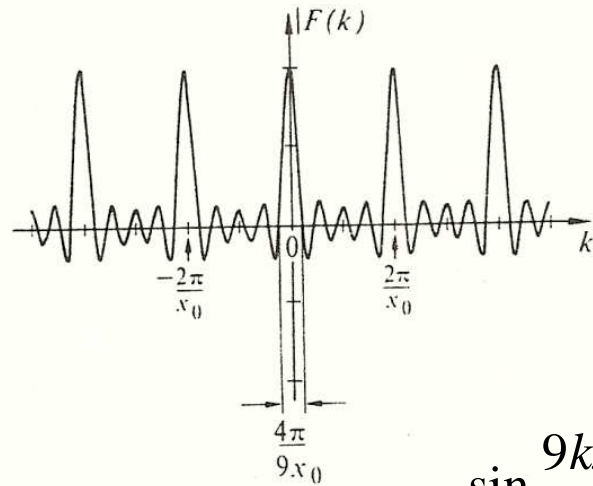
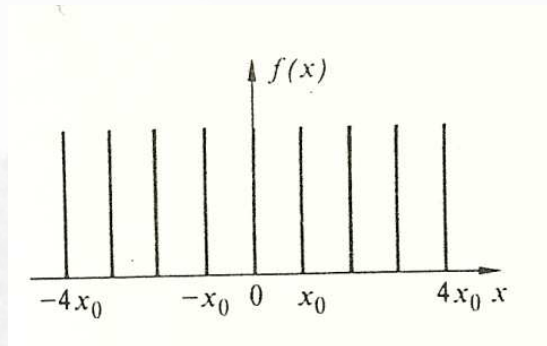
$$f(x) = \delta(x+x_0) + \delta(x) + \delta(x-x_0)$$



$$F(k) = 1 + 2 \cos kx_0$$

## Fourier transform > 9 delta functions

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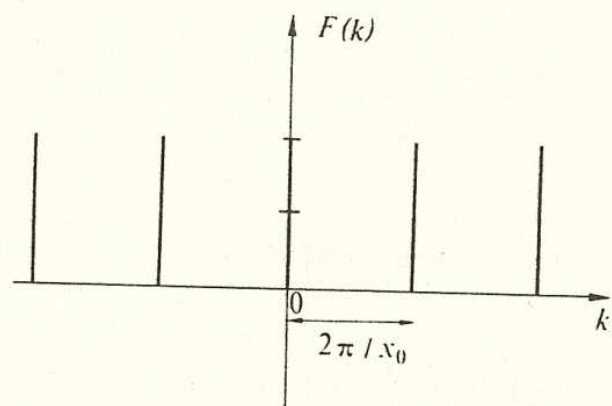
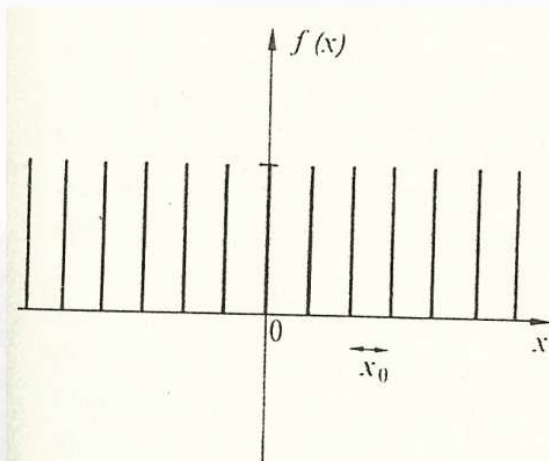


$$f(x) = \sum_{n=-4}^{n=4} \delta(x - nx_0)$$

$$F(k) = \frac{\sin \frac{9kx_0}{2}}{\sin \frac{kx_0}{2}}$$

## Fourier transform > infinite number of delta functions

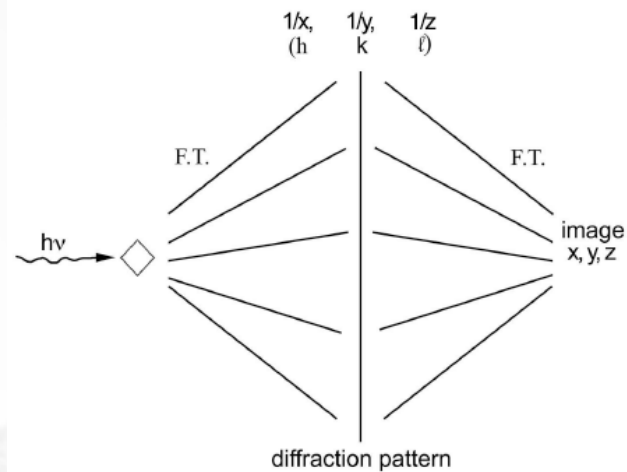
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$$f(x) = \sum_{n=-\infty}^{n=\infty} \delta(x - nx_0)$$

$$F(k) = \sum_{n=-\infty}^{n=\infty} \delta(k - 2n \frac{\pi}{x_0})$$

# How can we describe diffraction pattern of a xtal?



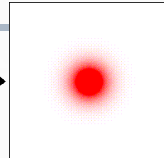
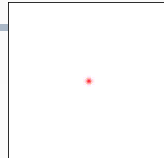
- The Fourier transform decomposes or separates a periodic function into sinusoids of different frequency which sum to the original function.
- Decompose  $f(x)$  into a series of cosine waves. That when summed, reconstruct  $f(x)$ .

$$F(\mathbf{k}) = \int_{\text{all } \mathbf{r}} f(\mathbf{r}) e^{i\mathbf{k}\cdot\mathbf{r}} d\mathbf{r} = \mathcal{T}f(\mathbf{r})$$

## Fourier transform

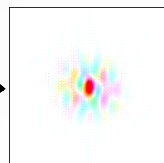
An atom

Fourier Transform



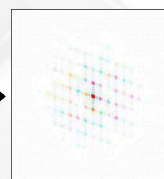
An atom is a sharp feature, whereas its transform is a broad smooth function.

A molecule

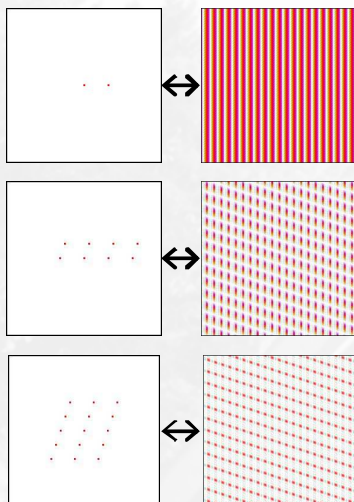


A molecule can be considered as the convolution of the *point atom structure* and the *atomic shape*. → Its transform is the product of the *point atom transform* and the *atomic transform*.

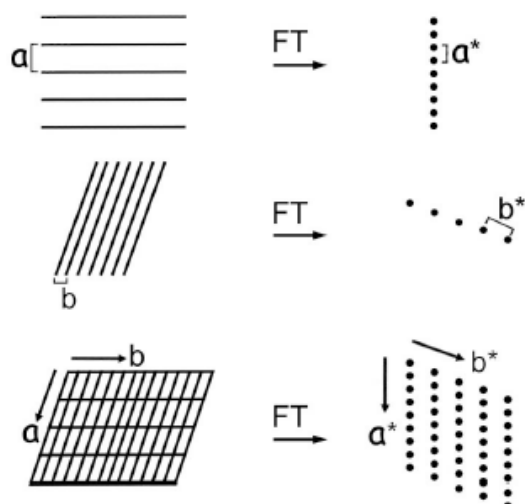
Molecules within a grid



A crystal can be built up by convoluting the *molecule with the grid*. → crystal structure  
The Fourier transform of the crystal is the product of the *molecular transform* and the *reciprocal lattice*. → *diffraction pattern*



The Fourier transform of a grid is a grid with reciprocal *directions* and *spacings*. → the origin of the *reciprocal lattice*



- Diffraction pattern of a set of lines is a row of dots perpendicular to the lines

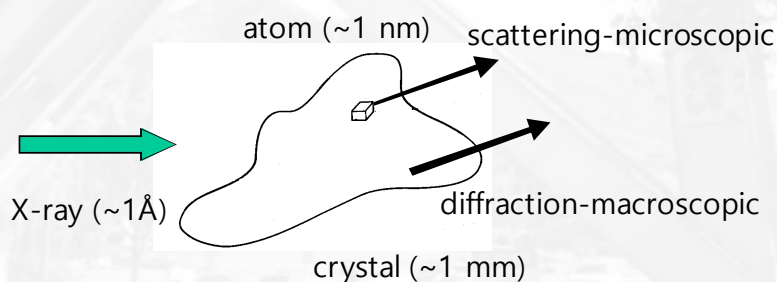
- The separation between the dots is proportional to the inverse of the separation between the lines

Form lattice by multiplying the two functions

The diffraction pattern of the lattice is a convolution of the diffraction patterns of the two sets of lines

## Interaction of waves with obstacles

- Infinite plane wave with wave vector  $\mathbf{k}$  and frequency  $\omega$ ;  $\psi = \psi_0 e^{i(\mathbf{k} \cdot \mathbf{r} - \omega t)}$
- What happens when a wave motion interacts with an obstacle placed in its path?
- How is the wave equation modified to take account of the interaction of the wave with the obstacle?
- **Scattering**; wave-obstacle interaction such that the dimensions of obstacles and wavelength are comparable
- **Diffraction**; wave-obstacle interaction such that the dimensions of obstacles are much larger than the wavelength



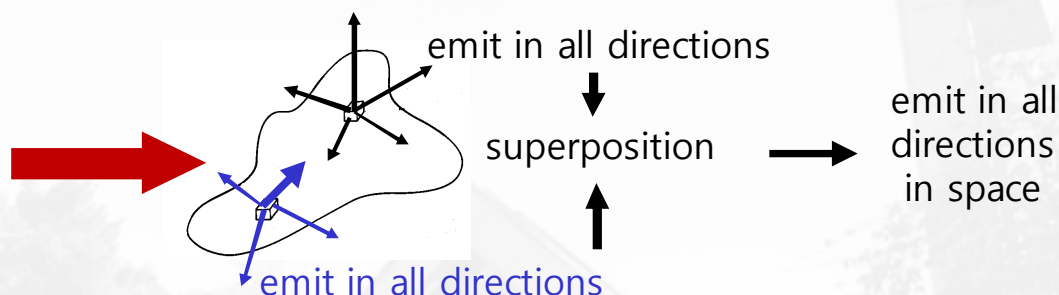
plane electromagnetic waves of a single frequency  $\omega$  & wavelength  $\lambda$



- In any medium, velocity of electromagnetic radiation is constant  $\rightarrow \lambda$  unchanged.  $\rightarrow$  diffracted waves have the same  $\omega$  and  $\lambda$  as the incident waves.
- Electrons behave as if it were free, when  $\omega$  of orbital motion of  $e^- \ll \omega$  of incident wave (See page 194, 195 of Sherwood)
  - ✓  $\omega$  (electron) =  $10^{16}\text{Hz}$ ,  $\omega$  (X-ray) =  $10^{19}\text{Hz}$

## Mathematics of Diffraction (continued)

- $\omega$  &  $\lambda$  unchanged  $\rightarrow$  there should be some other effect of diffracting obstacle.



- plane wave  $\rightarrow$  spatially different wave
- Information of the wave is somehow contained in the way in which they are spread out in space.  $\rightarrow$  direction  $\rightarrow$  only  $\mathbf{k}$  has spatial info among  $\psi_0$ ,  $\omega$ ,  $\mathbf{k}$ .

- significance of the wave vector  $\mathbf{k}$

$$|\mathbf{k}| = k = \frac{2\pi}{\lambda}$$

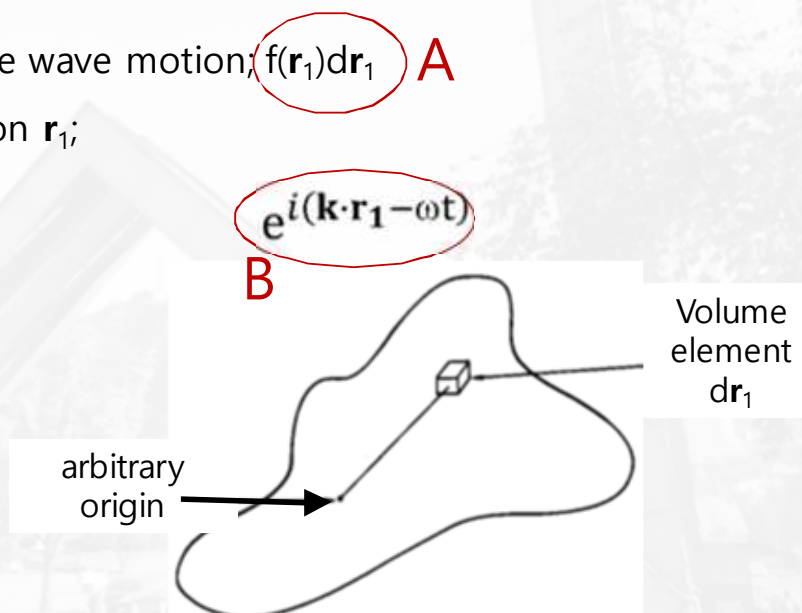
$$\psi(\mathbf{r}, t) = \psi_0 \exp i(\mathbf{k} \cdot \mathbf{r} - \omega t)$$

- $\mathbf{k}$  of diffracted wave has the same magnitude as the incident wave, but has different direction.
- Total set of diffracted waves = set of wave vectors all of which have same magnitude (equal to that of incident wave), but different directions.

# Mathematics of Diffraction (continued)

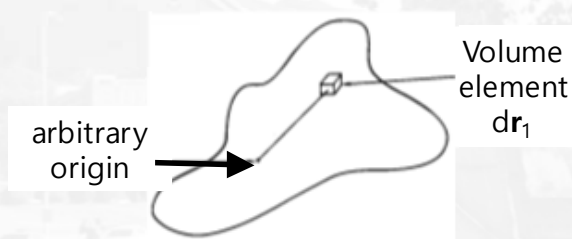
- mathematical description
  - ✓ Waves are propagated thru obstacle.
  - ✓ Obstacle perturbs these waves.
- Perturbing effect;  $f(\mathbf{r})$
- The effect of  $d\mathbf{r}_1$  on the wave motion;  $f(\mathbf{r}_1)d\mathbf{r}_1$  **A**
- Wave in  $d\mathbf{r}_1$  centered on  $\mathbf{r}_1$ ;

Diffracted wave must be represented by some mathematical combination of A & B.



- The effect of  $d\mathbf{r}_1$  on the wave motion;  $f(\mathbf{r}_1)d\mathbf{r}_1$  <sup>A</sup>
- Wave in  $d\mathbf{r}_1$  centered on  $\mathbf{r}_1$ ;  $e^{i(\mathbf{k} \cdot \mathbf{r}_1 - \omega t)}$  <sup>B</sup>
- Wave diffracted from  $d\mathbf{r}_1 = AB = f(\mathbf{r}_1)e^{i(\mathbf{k} \cdot \mathbf{r}_1 - \omega t)}d\mathbf{r}_1$
- Wave diffracted from  $d\mathbf{r}_2 = f(\mathbf{r}_2)e^{i(\mathbf{k} \cdot \mathbf{r}_2 - \omega t)}d\mathbf{r}_2$

➤ Diffraction pattern  $= \int_V f(\mathbf{r})e^{i(\mathbf{k} \cdot \mathbf{r} - \omega t)}d\mathbf{r}$



- Diffraction pattern  $= \int_V f(\mathbf{r})e^{i(\mathbf{k} \cdot \mathbf{r} - \omega t)}d\mathbf{r} = e^{-i\omega t} \int_V f(\mathbf{r})e^{i(\mathbf{k} \cdot \mathbf{r})}d\mathbf{r}$
- Intensity, not amplitude, is recorded.
- XRD experiment measures time average of the I of the diffracted wave.  
→ result obtained is independent of time.
- We need info on space, not info on time. →

Diffraction pattern  $= \int_V f(\mathbf{r})e^{i(\mathbf{k} \cdot \mathbf{r})}d\mathbf{r} = \int_{\text{all } \mathbf{r}} f(\mathbf{r})e^{i(\mathbf{k} \cdot \mathbf{r})}d\mathbf{r}$

Fourier transform of  $f(\mathbf{r})$

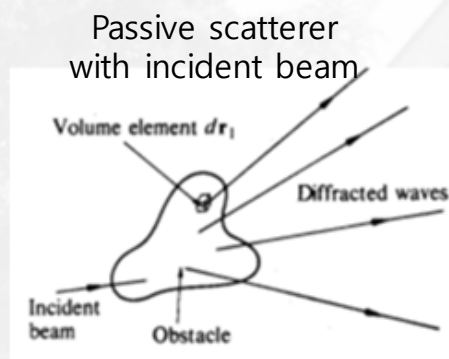
➤ Diffraction pattern =  $\int_V f(\mathbf{r}) e^{i(\mathbf{k} \cdot \mathbf{r})} d\mathbf{r} = \int_{\text{all } \mathbf{r}} f(\mathbf{r}) e^{i\mathbf{k} \cdot \mathbf{r}} d\mathbf{r}$

Fourier transform of  $f(\mathbf{r})$

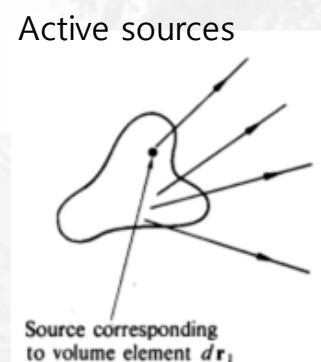
- To compute the diffraction pattern of any obstacle for which  $f(\mathbf{r})$  is known, we can compute the Fourier transform of  $f(\mathbf{r})$ .
- $f(\mathbf{r})$ ; amplitude function
- Diffraction pattern of any obstacle = Fourier transform of amplitude function

## Significance of Fourier Transform

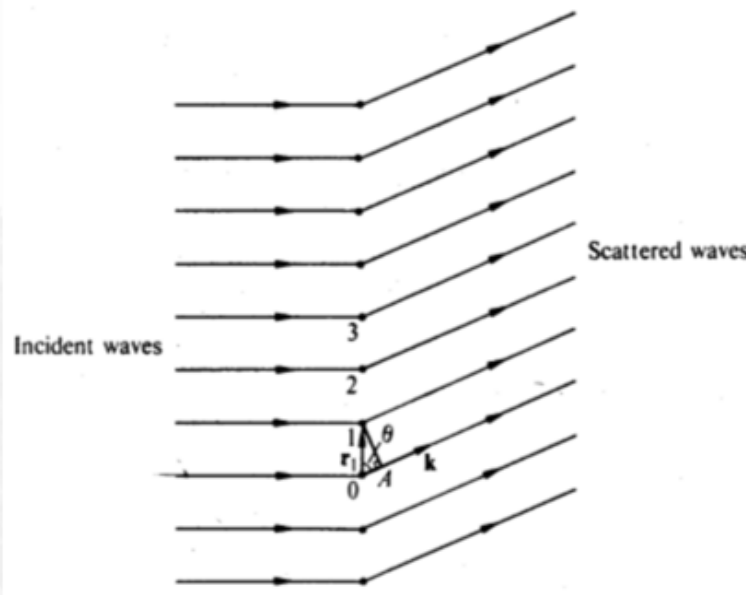
- Diffraction pattern =  $F(\mathbf{k}) = \int_{\text{all } \mathbf{r}} f(\mathbf{r}) e^{i\mathbf{k} \cdot \mathbf{r}} d\mathbf{r}$
- Info in the diffraction pattern is essentially spatial.
- The only way waves can contain spatial info is by means of their wave vector  $\mathbf{k}$ .



Identical effect



- Phase relationships b/w waves scattered by various scattering centers in an obstacle

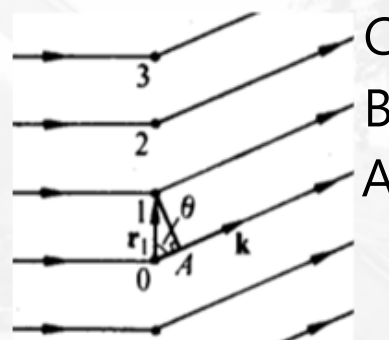


## Fourier Transform and phase

- Path difference b/w A & B =  $r_1 \cos \theta$
- Phase difference b/w A & B =  $(2\pi/\lambda) r_1 \cos \theta = \mathbf{k} \cdot \mathbf{r}_1$
- Phase difference b/w A & C =  $(2\pi/\lambda) r_2 \cos \theta = \mathbf{k} \cdot \mathbf{r}_2$
- Wave scattered from point 0 =  $f(\mathbf{r}_0) d\mathbf{r}_0$
- Wave scattered from point 1 =  $f(\mathbf{r}_1) e^{i\mathbf{k} \cdot \mathbf{r}_1} d\mathbf{r}_1$
- Wave scattered from point 2 =  $f(\mathbf{r}_2) e^{i\mathbf{k} \cdot \mathbf{r}_2} d\mathbf{r}_2$
- Total scattered wave

$$= \int_{\text{obstacle}} f(\mathbf{r}) e^{i\mathbf{k} \cdot \mathbf{r}} d\mathbf{r}$$

$$= \int_{\text{all } \mathbf{r}} f(\mathbf{r}) e^{i\mathbf{k} \cdot \mathbf{r}} d\mathbf{r}$$



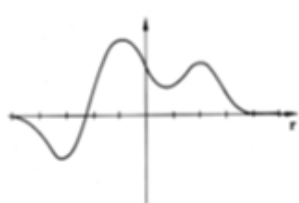
## Fourier Transform and Information

- How is information transmitted in general? see chap 6.11

$$\text{Diffraction pattern} = F(\mathbf{k}) = \int_{\text{all } \mathbf{r}} f(\mathbf{r}) e^{i\mathbf{k} \cdot \mathbf{r}} d\mathbf{r}$$

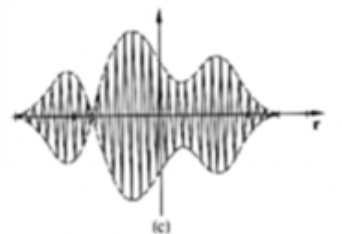


$e^{i\mathbf{k} \cdot \mathbf{r}}$



$f(\mathbf{r})$   
amplitude function

Describes behavior  
of the obstacle



$f(\mathbf{r}) e^{i\mathbf{k} \cdot \mathbf{r}}$

Constant  $w$  &  $\lambda$   
Different amplitudes  
**amplitude modulated wave**

## Fourier Transform & diffraction

Diffraction pattern = Fourier transform of  $f(\mathbf{r})$

$$F(\mathbf{k}) = \int_{\text{all } \mathbf{r}} f(\mathbf{r}) e^{i\mathbf{k} \cdot \mathbf{r}} d\mathbf{r}$$

Inverse transform 
$$f(\mathbf{r}) = \int_{\text{all } \mathbf{k}} F(\mathbf{k}) e^{-i\mathbf{k} \cdot \mathbf{r}} d\mathbf{k}$$

$F(\mathbf{k})$  ; contains info on the spatial distribution of diffraction pattern.

$f(\mathbf{r})$  ; contains info on the structure of obstacle.

XRD  
experiment

intensity  
 $|F(\mathbf{k})|^2$



$F(\mathbf{k})$



$f(\mathbf{r})$



Crystal  
structure

If the structure is known,

- f(**r**) is known.
- diffraction pattern F(**k**) can be computed.

If the diffraction pattern is known,

- F(**k**) is known.
- f(**r**) can be computed.

The act of diffraction = taking Fourier transform of the obstacle

Diffraction pattern of an obstacle described by f(**r**) is the Fourier transform of f(**r**), which is F(**k**).

## Experimental Limitation

➤ Information is contained in all space.

✓ It is impossible to scan all space to collect all the information.

→ some info is lost.

→ reconstruction of the obstacle from the diffraction data will be incomplete.

➤ PHASE PROBLEM

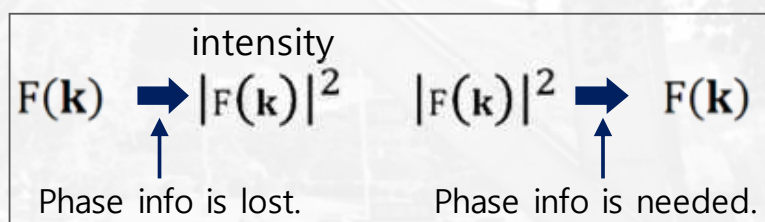
✓ Only the intensity is measured.

✓ The phase of diffraction pattern is not measured.

$$F(\vec{k}) = |F(\vec{k})| e^{i\delta}$$

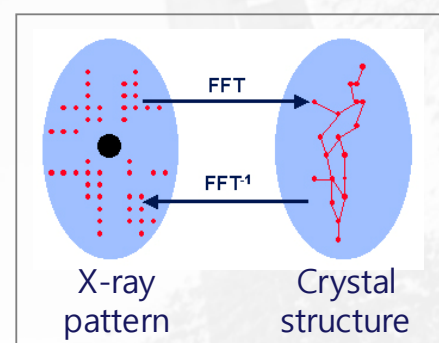
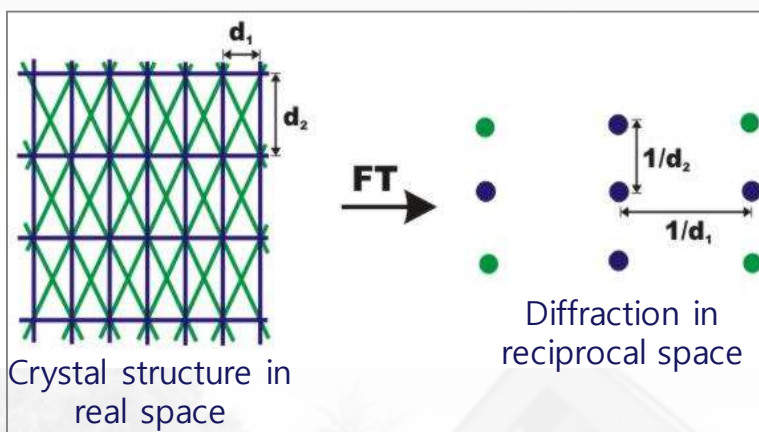
observe  $|F(\vec{k})|^2$

↓  
Info on  $\delta$  is lost.



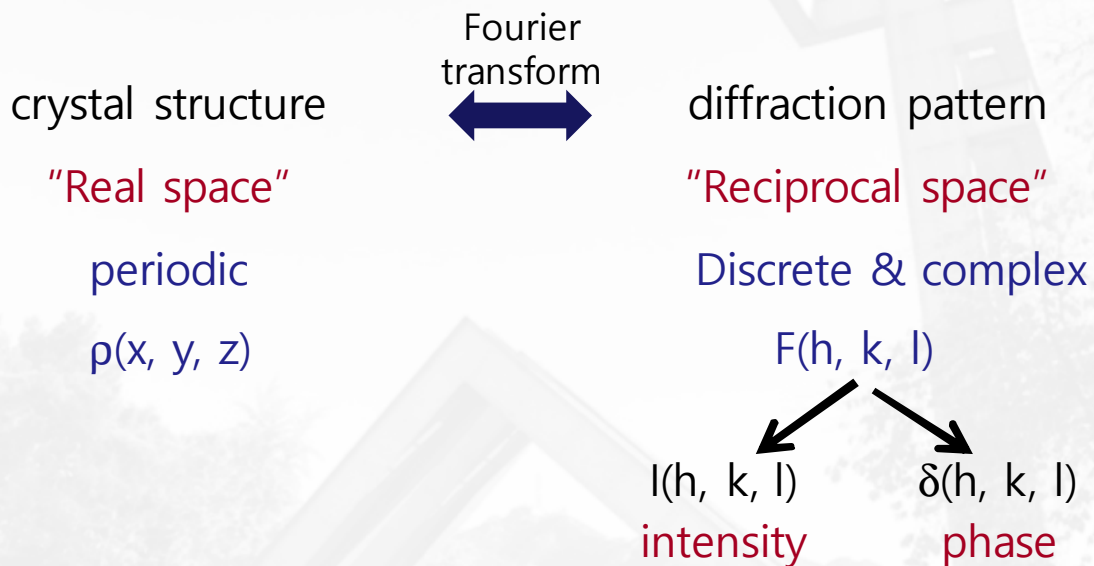
- What is measured is **INTENSITY**, not the **COMPLEX AMPLITUDE** of the diffraction pattern.
- $F(\mathbf{k})$  is complex  $F(\mathbf{k}) = |F(\mathbf{k})|e^{i\delta}$   $\delta$ ; phase factor
- What is measured is  $|F(\mathbf{k})|^2$ .
- $|F(\mathbf{k})|^2$  contains no info on  $\delta$ .
- A lot of info is lost because we can record only intensity (not the amplitude).
- The problem caused by the loss of info contained in  $\delta$ . → **PHASE PROBLEM**

## Reciprocal lattice vs Crystalline lattice in real space



The crystal structure can be deduced by performing a FT on the resultant diffraction pattern **once the phase is known**.

Phase problem can be solved by direct method, isomorphous replacement method, heavy atom method, etc.



$$F_{hkl} = \int_V \rho_{xyz} \exp[2\pi i(hx + ky + lz)] dV$$

$$\rho_{xyz} = \frac{1}{V} \sum |F_{hkl}| \exp(i\phi_{hkl}) \exp[-2\pi i(hx + ky + lz)]$$

## todos

### ➤ XRD-6 READ

- Sherwood & Cooper, Chap 4.13, 5.6, 6.1~6.7, 6.14
- Sherwood & Cooper, Chap 6.8~6.14
- Ott, Chapter 13.4
- Sherwood & Cooper, Chap 4.1~4.3, 4.11, 5.1~5.3
- Sherwood & Cooper, Chap 4.4~4.10, 5.4, 5.7~5.10

No problem solving homeworks, only reading homeworks