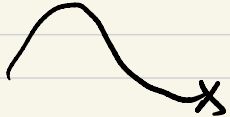
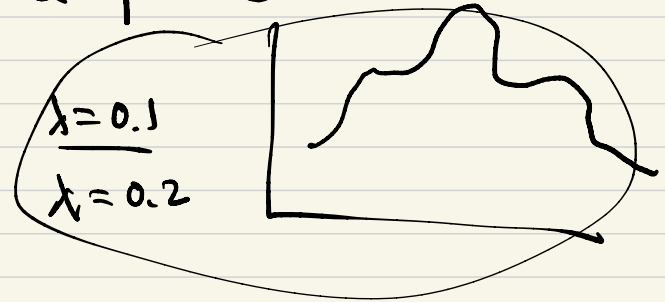


② $\begin{cases} g_0'' = 0 \\ g_n'' = 0 \end{cases}$  natural spline

③ $\begin{cases} g_0'' = \lambda g_1'' \\ g_n'' = \lambda g_{n-1}'' \end{cases} \quad 0 \leq \lambda \leq 1$



④ $\begin{cases} g_0'' = g_n'' \\ g_n'' = g_0'' \end{cases}$ periodic b.c.

$$\begin{pmatrix} 1 & 0 & \dots & 0 & -1 \\ & \parallel & & & \\ & & \parallel & & \\ & & & \parallel & \\ -1 & 0 & \dots & 0 & 1 \end{pmatrix} \begin{pmatrix} g_0'' \\ g_1'' \\ \vdots \\ g_{n-1}'' \\ g_n'' \end{pmatrix} = \begin{pmatrix} 0 \\ \vdots \\ 0 \end{pmatrix}$$

For equally spaced pts.,
 $\Delta \bar{u} = \Delta \bar{u}_1 = \Delta$

$$\frac{1}{6} [g_{i-1}'' + 4g_i'' + g_{i+1}''] = \frac{1}{\Delta^2} (y_{i-1} - 2y_i + y_{i+1})$$

$\bar{i} = 1, 2, \dots, n-1$

g_0'' & g_n''

In some cases, you may get wiggles.

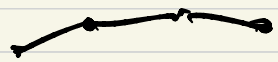
→ tension splines → pull on each ends ($y^{(iv)} - \sigma^2 y'' = 0$)

$$\rightarrow y' - \sigma^2 y = (y_{ii}'' - \sigma^2 y_{ii}) \frac{x - x_{i+1}}{x_i - x_{i+1}} + (y_{i+1}' - \sigma^2 y_{i+1}) \frac{x - x_i}{x_{i+1} - x_i}$$

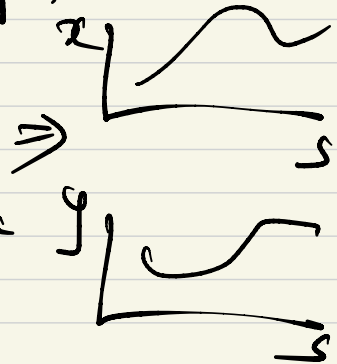
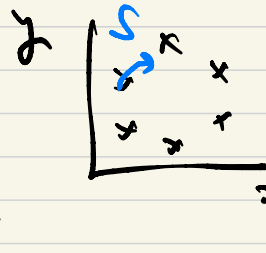
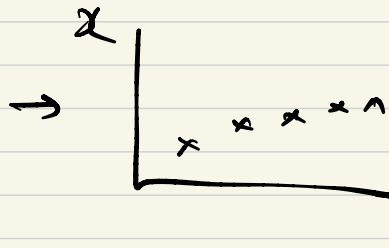
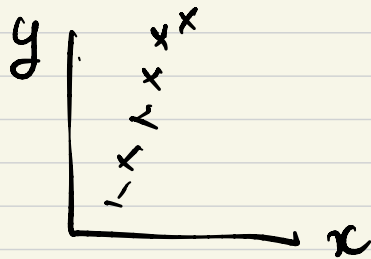
$\sigma \rightarrow \infty$: straight line interpolation

$\sigma = 0$: usual spline

$\sigma \neq 0 \Rightarrow$ wiggles in sol. → increase σ .

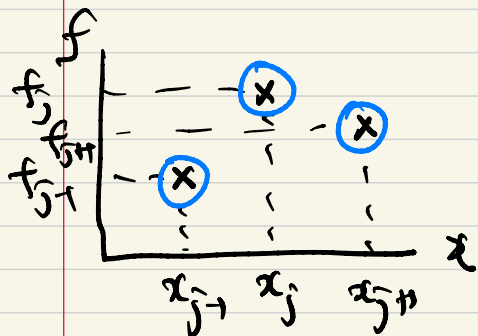


* spline has difficulties with large slopes.



Ch. 2. Numerical differentiation - finite difference

2.1 Construction of difference formulae using Taylor series.



$$(x_i, f_i) \rightarrow f_i'$$
$$\frac{df}{dx} = \lim_{\Delta x \rightarrow 0} \frac{f(x+\Delta x) - f(x)}{\Delta x}$$

Finite difference using Taylor series expansion

$$f(x_j + \Delta x) = f(x_j) + \Delta x f'(x_j) + \frac{1}{2} \Delta x^2 f''(x_j) + \frac{1}{6} \Delta x^3 f'''(x_j) + \dots$$
$$\rightarrow f'(x_j) = \frac{f(x_j + \Delta x) - f(x_j)}{\Delta x} - \frac{1}{2} \Delta x f''(x_j) - \frac{1}{6} \Delta x^2 f'''(x_j) + \dots$$

1st-order accurate
finite difference (FD)
method

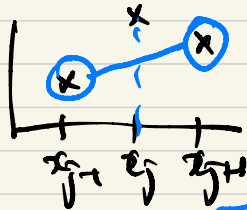
leading truncation error $O(\Delta x)$

$$\rightarrow f'_j = \frac{f_{j+1} - f_j}{\Delta x} + O(\Delta x) \quad \text{forward 1st-order FD method.}$$

Similarly, $f(x_j - \Delta x) = f(x_j) - \Delta x f'(x_j) + \frac{1}{2} \Delta x^2 f''(x_j) - \frac{1}{6} \Delta x^3 f'''(x_j) + \dots$

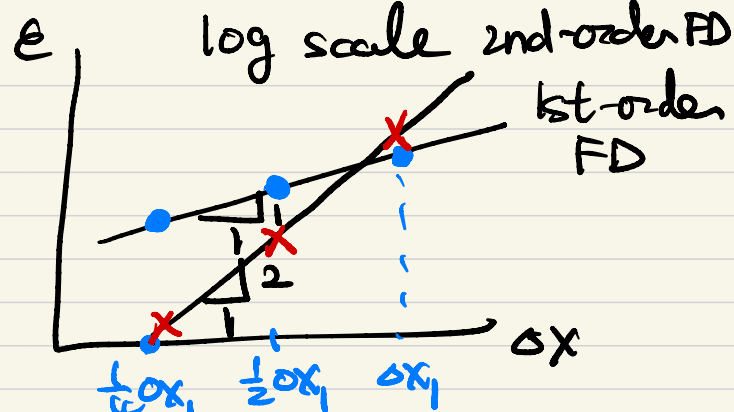
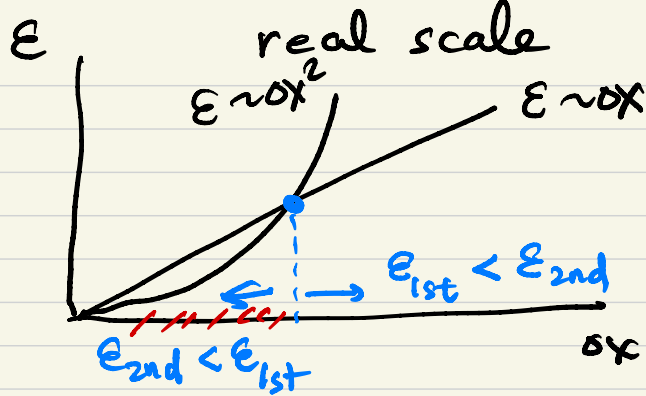
→ $f'_j = \frac{f_j - f_{j-1}}{\Delta x} + O(\Delta x)$ backward 1st-order FD method

$f(x_j + \Delta x) = \dots$
 $f(x_j - \Delta x) = \dots$ ⇒ $f'_j = \frac{f_{j+1} - f_{j-1}}{2\Delta x} - \frac{1}{6} \Delta x^2 f''_j + \dots$



2nd-order central FD method (CP2)

$$\left\{ \begin{aligned} f'_j &= \frac{f_{j+1} - f_j}{\Delta x} - \frac{1}{2} \Delta x f''_j + \dots \\ f'_j &= \frac{f_j - f_{j-1}}{\Delta x} + \frac{1}{2} \Delta x f''_j + \dots \\ f'_j &= \frac{f_{j+1} - f_{j-1}}{2\Delta x} - \frac{1}{6} \Delta x^2 f''_j + \dots \end{aligned} \right. \begin{aligned} &\rightarrow \epsilon \sim \Delta x \rightarrow \ln \epsilon \sim \ln \Delta x \\ &\rightarrow \epsilon \sim \Delta x^2 \Rightarrow \ln \epsilon \sim 2 \ln \Delta x \end{aligned}$$

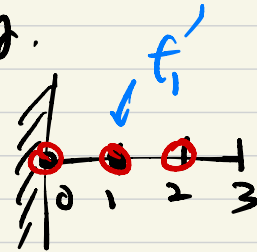


$$f'_j = \frac{f_{j-2} - 8f_{j-1} + 8f_{j+1} - f_{j+2}}{12 \Delta x} + O(\Delta x^4)$$

CD4

more grid pts. \rightarrow higher accuracy.
 How about near/on boundary?

lower-order accurate formula.



② $j=1$, CD2 $\rightarrow O(\Delta x^2)$

f'_1 using f_0, f_1, f_2, f_3, f_4 .

2.2 A general technique for construction of FD schemes

• f'_j using function values at $j, j+1, j+2$.

Q: what is the most accurate formula?

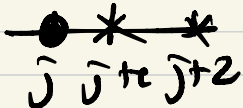
$$f'_j + a_0 f_j + a_1 f_{j+1} + a_2 f_{j+2} = O(h^?) \quad (h=Ox)$$

Taylor table	f_j	f'_j	f''_j	f'''_j	...
f'_j	0	1	0	0	0
$a_0 f_j$	a_0	0	0	0	...
$a_1 f_{j+1}$	$a_1 \cdot 1$	$a_1 \cdot h$	$a_1 \cdot \frac{1}{2} h^2$	$a_1 \cdot \frac{1}{6} h^3$...
$+ a_2 f_{j+2}$	$a_2 \cdot 1$	$a_2 \cdot (2h)$	$a_2 \cdot \frac{1}{2} (2h)^2$	$a_2 \cdot \frac{1}{6} (2h)^3$...

$$\begin{aligned} \rightarrow f'_j + a_0 f_j + a_1 f_{j+1} + a_2 f_{j+2} &= \underbrace{(a_0 + a_1 + a_2)}_{\text{3 unknowns}} f_j \\ &+ \underbrace{(1 + a_1 h + 2a_2 h)}_{\text{set as many lower-order}} f'_j + \underbrace{(\frac{1}{2} a_1 h^2 + 2a_2 h^2)}_{\text{coeffs to zero as possible}} f''_j \\ &+ (\frac{1}{6} a_1 h^3 + \frac{4}{3} a_2 h^3) f'''_j + \dots \end{aligned}$$

$$\begin{aligned} \rightarrow a_0 + a_1 + a_2 &= 0 \\ 1 + a_1 h + 2a_2 h &= 0 \\ \frac{1}{2} a_1 h^2 + 2a_2 h^2 &= 0 \end{aligned} \Rightarrow a_0 = \frac{3}{2h}, a_1 = -\frac{2}{h}, a_2 = \frac{1}{2h}$$

$$\rightarrow f_j' + \frac{3}{2h} f_j - \frac{2}{h} f_{j+1} + \frac{1}{2h} f_{j+2} = \frac{1}{3} h^2 f_j''' + \dots$$

$$\rightarrow f_j' = \frac{-3f_j + 4f_{j+1} - f_{j+2}}{2h} + \frac{1}{3} h^2 f_j''' + \dots$$


2nd-order FD
one-side difference

2.3 An alternative measure for the accuracy of FD.

HW 1

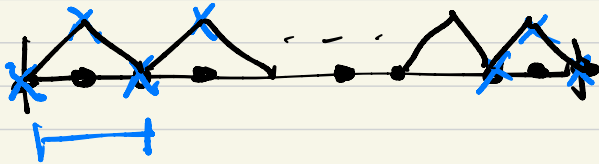
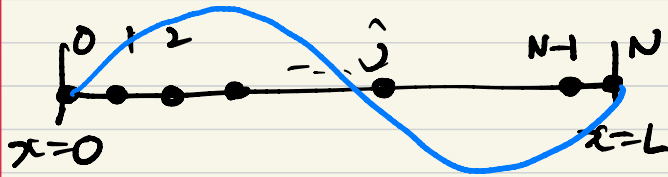
$\mathcal{O}(h^2)$ $\mathcal{O}(h^2)$

2.3 An alternative measure for the accuracy of FD

* modified wavenumber approach for measuring accuracy

$f(x) = e^{ikx} = \cos kx + i \sin kx$: pure harmonic ft. of period L

k : wavenumber



uniform mesh

$\Delta x = h = L/N$: grid spacing

smallest wavenumber Δk

$$\Delta k \cdot L = 2\pi \rightarrow \Delta k = \frac{2\pi}{L}$$

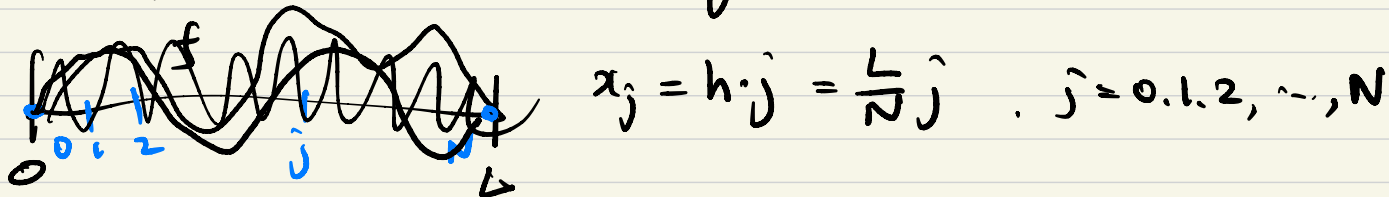
largest wavenumber K

$$K \cdot (2\Delta x) = 2\pi$$

$$K = \frac{\pi}{\Delta x} = \frac{\pi}{h} = \frac{\pi N}{L}$$

$$f(x) = e^{ikx} \rightarrow f' = ike^{ikx} = ikf : \text{exact sol.}$$

Q: How accurately CD2 computes the derivatives of f for different values of k ? $\mathcal{O}(h^2)$

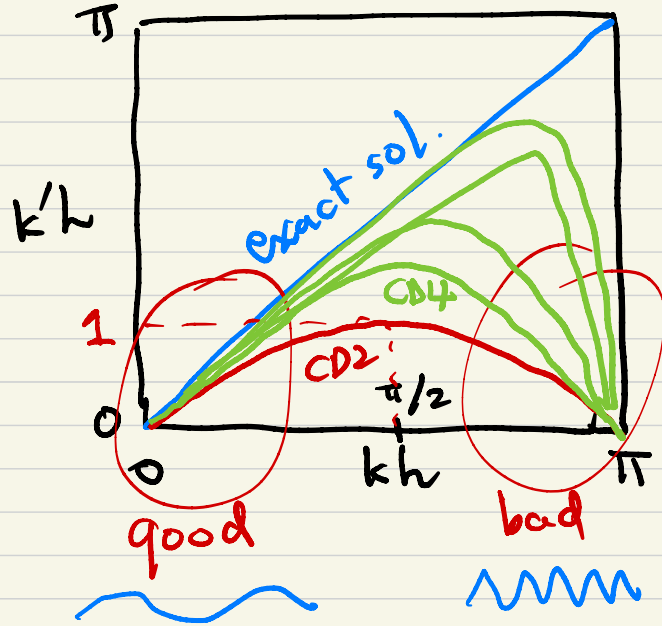


On this grid, e^{ikx} ranges from a constant for $k=0$ to a highly oscillatory ft. of period equal to mesh width ($2\pi x$) for $k = \pi N/h$.

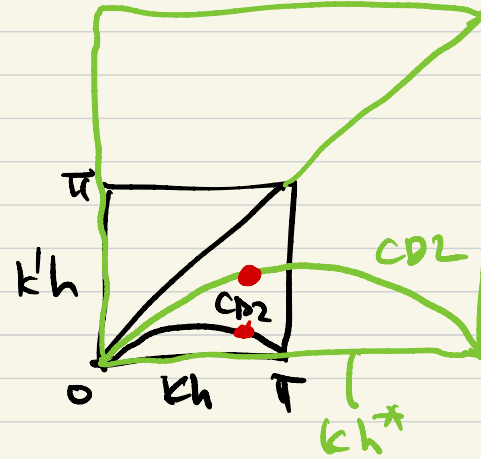
$$\begin{aligned} \text{CD2: } f'_j &= \frac{f_{j+1} - f_{j-1}}{2h} = \frac{1}{2h} (e^{ikx_{j+1}} - e^{ikx_{j-1}}) \\ &= \frac{1}{2h} e^{ikx_j} (e^{ikh} - e^{-ikh}) = i \frac{\sin kh}{h} e^{ikx_j} = i \frac{\sin kh}{h} f_j \\ &= i - 2 \sin kh \end{aligned}$$

k' : modified wavenumber

$\rightarrow k'h = \sin kh \quad (CD2)$



$N \rightarrow 2N \quad h \rightarrow \frac{h}{2} = h^*$



• 1st order FD $f'_j = \frac{f_{j+1} - f_j}{h} + O(h)$ $f = e^{ikx}$
 $f' = ikf$

$$= \frac{1}{h} (e^{ikx_{j+1}} - e^{ikx_j})$$

$$= \frac{1}{h} e^{ikx_j} (\cos kh + i \sin kh - 1)$$

$$= i \left[\frac{\sin kh}{h} + i \frac{1 - \cos kh}{h} \right] f_j$$

$e^{ik'x}$

k'_{D2} \swarrow k' : modified wavenumber
 complex number
 \searrow dispersive error \searrow dissipative error



Each FD scheme has its own modified wavenumber,

2.4 Padé approximations

$$f'_j + a f_j + b f_{j+1} + c f_{j-1} + d f_{j+2} = O(h^2)$$

Include derivatives too in the formula.

ex) Find the most accurate formula of f'_j that involves $f_{j+1}, f_{j-1}, f_j, f'_{j+1}, f'_{j-1}$.

$$\rightarrow f'_j + a_0 f'_{j+1} + a_1 f'_{j-1} + b_0 f_j + b_1 f_{j-1} + b_2 f_{j+1} = O(h^2)$$

Taylor table	f_j	f'_j	f''_j	f'''_j	$f^{(iv)}_j$	$f^{(v)}_j$...
f'_j	0	1	0	0	0	0	
$a_0 f'_{j+1}$	$a_0 \cdot 0$	$a_0 \cdot 1$	$a_0 \cdot h$	$a_0 \cdot \frac{1}{2} h^2$	$a_0 \cdot \frac{1}{6} h^3$	$a_0 \cdot \frac{1}{24} h^4$...
$a_1 f'_{j-1}$	$a_1 \cdot 0$	$a_1 \cdot 1$	$a_1 \cdot (-h)$	$a_1 \cdot \frac{1}{2} (-h)^2$	$a_1 \cdot \frac{1}{6} (-h)^3$	$a_1 \cdot \frac{1}{24} (-h)^4$...
$b_0 f_j$	b_0	0	0	0	0	0	
$b_1 f_{j-1}$	$b_1 \cdot 1$	$b_1 \cdot (-h)$	$b_1 \cdot \frac{1}{2} h^2$	$b_1 \cdot \frac{1}{6} (-h)^3$	$b_1 \cdot \frac{1}{24} h^4$...	
$b_2 f_{j+1}$	$b_2 \cdot 1$	$b_2 \cdot h$	$b_2 \cdot \frac{1}{2} h^2$	$b_2 \cdot \frac{1}{6} h^3$	$b_2 \cdot \frac{1}{24} h^4$...	

$$f_j' + a_0 f_{j+1}' + a_1 f_{j-1}' + b_0 f_j + b_1 f_{j-1} + b_2 f_{j+1} \quad \text{5 unknowns}$$

$$= (b_0 + b_1 + b_2) f_j + (1 + a_0 - a_1 - b_1 h + b_2 h) f_j'$$

$$+ (a_0 h + a_1 h + \frac{1}{2} b_1 h^2 + \frac{1}{2} b_2 h^2) f_j'' + \dots f_j''' + \dots f_j^{(4)} + \dots$$

$$\rightarrow a_0 = \dots, a_1 = \dots, b_0 = \dots, b_1 = \dots, b_2 = \dots$$

$$\rightarrow \underline{f_j'} + \frac{1}{4} \underline{f_{j+1}'} + \frac{1}{4} \underline{f_{j-1}'} - \frac{3}{4h} \underline{f_{j+1}} + \frac{3}{4h} \underline{f_{j-1}} = \frac{h^4}{120} f_j^{(4)} + \dots$$

$$\rightarrow \boxed{f_{j-1}' + 4 f_j' + f_{j+1}'} = \frac{3}{h} (f_{j+1} - f_{j-1}) + \frac{1}{30} h^4 f_j^{(4)} + \dots$$

4th order accurate FD $\mathcal{O}(h^4)$

$j = 1, 2, \dots, M-1$

3 grid pts! \uparrow

\downarrow
compact scheme

$$\begin{bmatrix} \diagdown & & & \\ & \diagdown & & \\ & & \diagdown & \\ & & & \diagdown \end{bmatrix} \begin{bmatrix} \\ \\ \\ \end{bmatrix} = \begin{bmatrix} \\ \\ \\ \end{bmatrix}$$

tri-diagonal matrix sys.
 \rightarrow easy to solve.

$n \rightarrow$ eqs. for $n+1$ unknowns $f_0', f_1', \dots, f_{n-1}', f_n'$

① $\bar{j}=0$: $f_0' = \frac{1}{2h} (-3f_0 + 4f_1 - f_2) + \mathcal{O}(h^2)$

② $\textcircled{a} \bar{j}=1$: $f_2' + 4f_1' = \frac{1}{2h} (7f_2 - 3f_0 - 4f_1) + \mathcal{O}(h^2)$ $\mathcal{O}(h^2)$

$$\begin{pmatrix} 4 & 1 & 0 & \dots & 0 \\ 1 & 4 & 1 & & \\ & & \ddots & \ddots & \\ & & & & 1 & 4 \end{pmatrix} \begin{pmatrix} f_1' \\ \vdots \\ \vdots \\ f_{n-1}' \end{pmatrix} = \begin{pmatrix} \frac{1}{2h} (7f_2 - 3f_0 - 4f_1) \\ \frac{3}{h} (f_3 - f_1) \\ \vdots \\ \vdots \end{pmatrix}$$

$\mathcal{O}(h^2)$
 $\mathcal{O}(h^k)$
 $\mathcal{O}(h^2)$

③ $\textcircled{a} \bar{j}=0$: $f_0' + 2f_1' = \frac{1}{h} \left(-\frac{5}{2}f_0 + 2f_1 + \frac{1}{2}f_2 \right) + \mathcal{O}(h^3)$ overall accuracy $\mathcal{O}(h^2 \sim h^3)$

$$\begin{pmatrix} 1 & 2 & 0 & \dots & 0 \\ 1 & 4 & 1 & & \\ 0 & & \ddots & \ddots & \\ & & & & 1 & 4 \end{pmatrix} \begin{pmatrix} \vdots \\ \vdots \\ \vdots \end{pmatrix} = \begin{pmatrix} \vdots \\ \vdots \\ \vdots \end{pmatrix}$$

Pade approx. can be easily extended to higher derivatives.

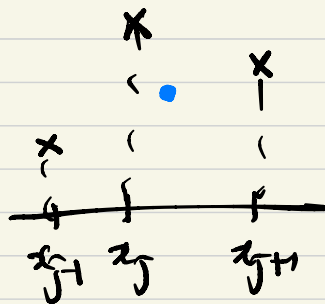
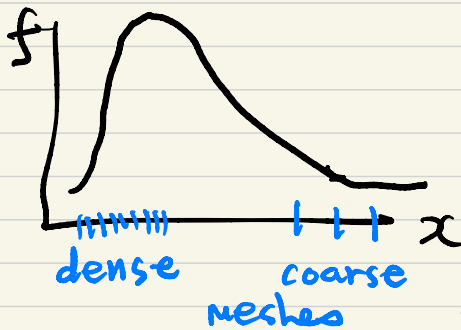
$$O(x) \quad \frac{1}{12} f_{j-1}'' + \frac{10}{12} f_j'' + \frac{1}{12} f_{j+1}'' = \frac{1}{h^2} (f_{j+1} - 2f_j + f_{j-1}) + O(h^4)$$

3 pts. $\rightarrow O(h^4)$ 4th-order accurate compact scheme

S. Lele (J. Comput. Phys, 1992) - compact schemes

7267 citations

2.5 Non-uniform grids



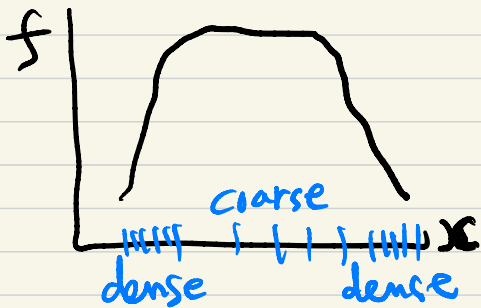
$$f_j' = \frac{f_{j+1} - f_{j-1}}{x_{j+1} - x_{j-1}} + O(h^2)$$

strictly 1st-order accurate formula

$$f_j'' = 2 \left[\frac{f_{j-1}}{h_j(h_j + h_{j+1})} - \frac{f_j}{h_j h_{j+1}} + \frac{f_{j+1}}{h_{j+1}(h_j + h_{j+1})} \right] + O(h^2)$$

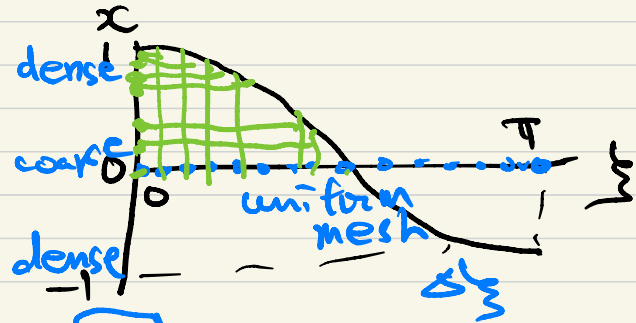
FD formula for non-uniform mesh generally has a lower order of accuracy than their counterpart w/ the same stencil for uniform mesh.

- Use a coordinate transformation



$$x \rightarrow \xi$$

ex) $x = \cos \xi$



$$\xi = g(x)$$

$$\frac{df}{dx} = \frac{df}{d\xi} \frac{d\xi}{dx} = g' \frac{df}{d\xi} \quad \leftarrow \text{FD on uniform mesh.}$$

$$\frac{d^2f}{dx^2} = \frac{d}{dx} \left(\frac{df}{dx} \right) = \frac{d}{dx} \left(g' \frac{df}{d\xi} \right) = g'' \frac{df}{d\xi} + g' \frac{d^2f}{d\xi^2} \frac{d\xi}{dx}$$

$$= g'' \frac{df}{d\xi} + g'^2 \frac{d^2f}{d\xi^2}$$