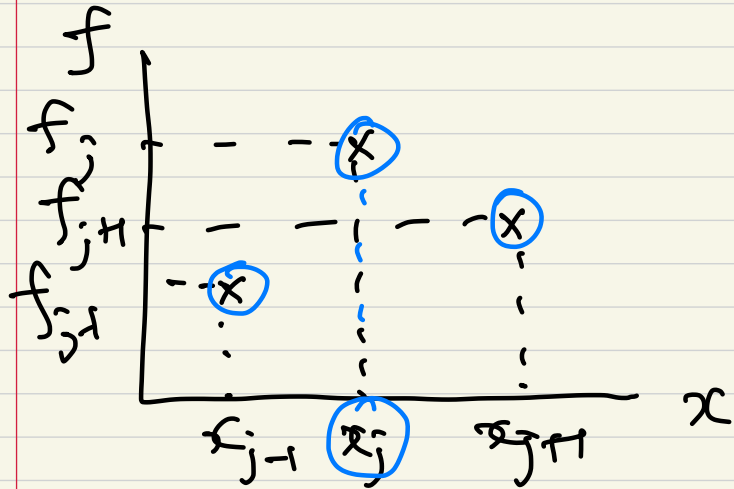


Ch 2. Numerical differentiation - finite difference

2.1 Construction of difference formulae using Taylor series



$$(x_i, f_i) \rightarrow f_i'$$

$$\frac{df}{dx} = \lim_{\Delta x \rightarrow 0} \frac{f(x+\Delta x) - f(x)}{\Delta x}$$

Finite difference using Taylor series expansion

$$f(x_j + \Delta x) = f(x_j) + \Delta x f'(x_j) + \frac{1}{2} \Delta x^2 f''(x_j) + \frac{1}{6} \Delta x^3 f'''(x_j) + \dots$$

$$\rightarrow f'(x_j) = \frac{f(x_j + \Delta x) - f(x_j)}{\Delta x} - \frac{1}{2} \Delta x f''(x_j) - \frac{1}{6} \Delta x^2 f'''(x_j) + \dots$$

1st-order accurate

finite difference (FD)

leading

truncation error
 $\mathcal{O}(\Delta x)$

$$\rightarrow f'_j = \frac{f_{j+1} - f_j}{\Delta x} + \mathcal{O}(\Delta x)$$

forward 1st-order
FD method

Similarly, $f(x_j - \Delta x) = f(x_j) - \Delta x f'(x_j) + \frac{1}{2} \Delta x^2 f''(x_j) - \dots$

→ $f'_j = \frac{f_j - f_{j-1}}{\Delta x} + O(\Delta x)$ backward 1st-order FD method

↳ $+\frac{1}{2} \Delta x f''(x_j) - \frac{1}{6} \Delta x^2 f'''(x_j) + \dots$

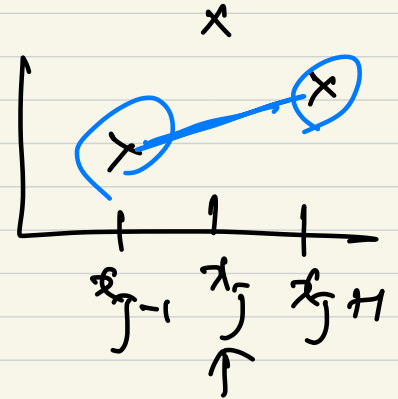
$f(x_j + \Delta x) = \dots$

+ $f(x_j - \Delta x) = \dots$

→ $f'_j = \frac{f_{j+1} - f_{j-1}}{2\Delta x} - \frac{1}{6} \Delta x^2 f'''_j + \dots$

$O(\Delta x^2)$

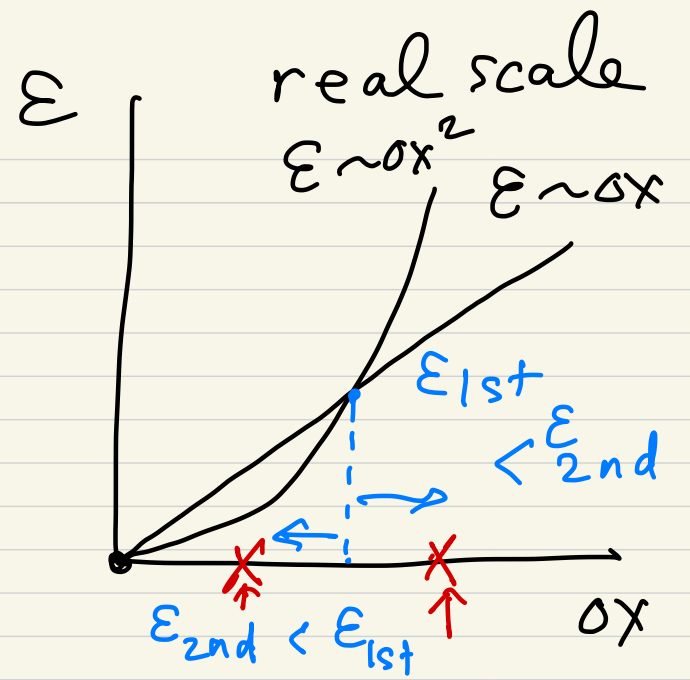
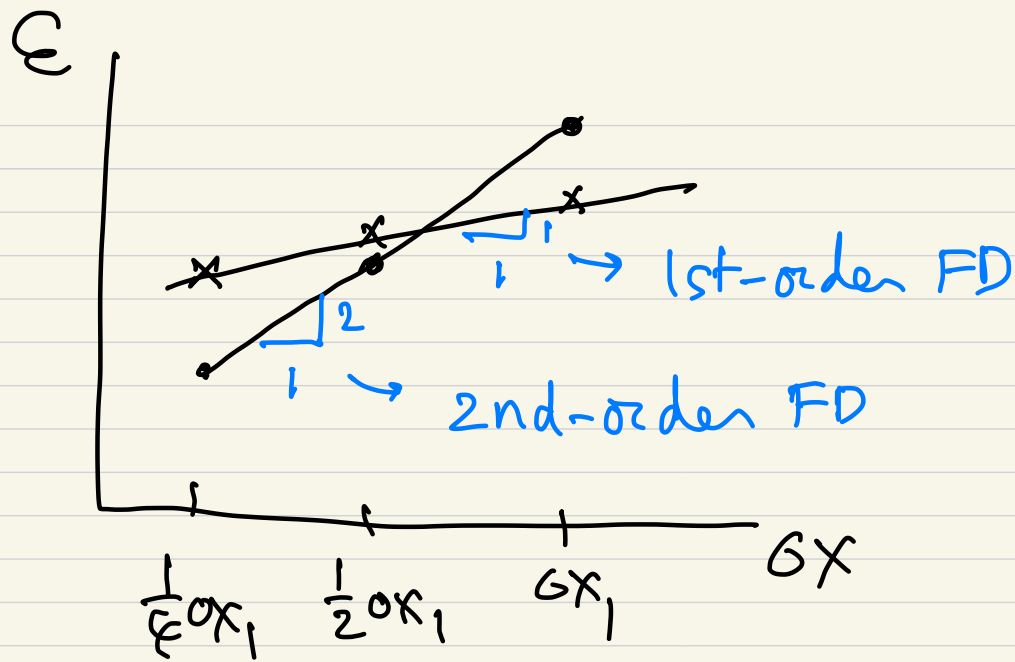
2nd-order central FD method (CØ2)



$f'_j = \frac{f_{j+1} - f_j}{\Delta x} - \frac{1}{2} \Delta x f''_j + \dots$
 $f'_j = \frac{f_j - f_{j-1}}{\Delta x} + \frac{1}{2} \Delta x f''_j + \dots$
 $f'_j = \frac{f_{j+1} - f_{j-1}}{2\Delta x} - \frac{1}{6} \Delta x^2 f'''_j + \dots$

$\epsilon \sim \Delta x \rightarrow \ln \epsilon \sim \ln \Delta x$
 $\epsilon \sim \Delta x^2 \rightarrow \ln \epsilon \sim 2 \ln \Delta x$

log scale



$$f'_j = \frac{f_{j-2} - 8f_{j-1} + 8f_{j+1} - f_{j+2}}{12\Delta x} + O(\Delta x^4) \quad \underline{\text{CD4}}$$

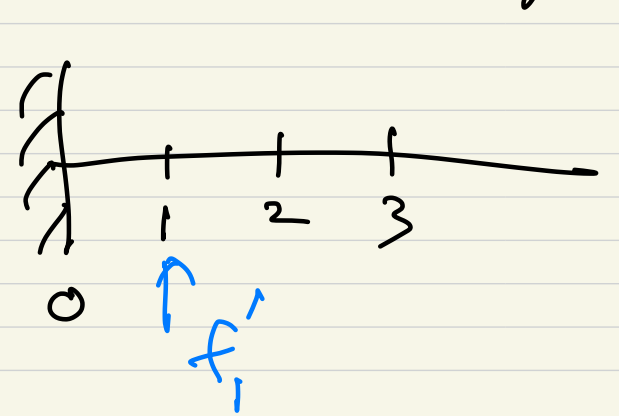
more grid pts. \rightarrow higher accuracy

problem? \rightarrow near/on boundary

lower-order accurate formula \rightarrow

① $j=1$, CD2 $\rightarrow O(\Delta x^2)$

f'_1 using f_0, f_1, f_2, f_3, f_4



2.2 A general technique for construction of FD schemes

• f_j' using function values at $j, j+1, j+2$.

Q: what is the most accurate formula?

$$f_j' + a_0 f_j + a_1 f_{j+1} + a_2 f_{j+2} = O(h^?) \quad (h \rightarrow 0)$$

<u>Taylor table</u>	f_j	f_j'	f_j''	f_j'''
f_j'	0	1	0	0	...
$a_0 f_j$	a_0	0	0	0	...
$a_1 f_{j+1}$	$a_1 \cdot 1$	$a_1 \cdot h$	$a_1 \cdot \frac{1}{2} h^2$	$a_1 \cdot \frac{1}{6} h^3$...
$+ [a_2 f_{j+2}]$	$a_2 \cdot 1$	$a_2 \cdot (2h)$	$a_2 \cdot \frac{1}{2} (2h)^2$	$a_2 \cdot \frac{1}{6} (2h)^3$...

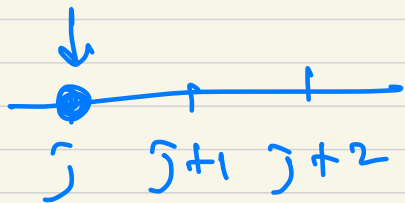
$$\begin{aligned} \rightarrow f_j' + a_0 f_j + a_1 f_{j+1} + a_2 f_{j+2} &= (a_0 + a_1 + a_2) f_j \\ &+ (1 + a_1 h + 2a_2 h) f_j' + \left(\frac{1}{2} a_1 h^2 + 2a_2 h^2\right) f_j'' + \left(\frac{1}{6} a_1 h^3 + \frac{4}{3} a_2 h^3\right) f_j''' \\ &+ \dots \end{aligned}$$

Set as many lower-order coeffs to zero as possible.

$$\begin{aligned} \rightarrow a_0 + a_1 + a_2 &= 0 \\ 1 + a_1 h + 2a_2 h &= 0 \\ \frac{1}{2} a_1 h^2 + 2a_2 h^2 &= 0 \end{aligned} \rightarrow a_0 = \frac{3}{2h}, a_1 = -\frac{2}{h}, a_2 = \frac{1}{2h}$$

$$\rightarrow f_j' + \frac{3}{2h} f_j - \frac{2}{h} f_{j+1} + \frac{1}{2h} f_{j+2} = \frac{1}{3} h^2 f_j'' + \dots$$

$$\rightarrow f_j' = \frac{-3f_j + 4f_{j+1} - f_{j+2}}{2h} + \frac{1}{3} h^2 f_j'' + \dots$$



2nd-order $\neq D$

← one-side difference

2.3 An alternative measure for the accuracy of FD,

* modified wavenumber approach for measuring accuracy.

$$f(x) = e^{ikx} = \cos kx + i \sin kx : \text{pure harmonic ft. of period } L.$$

k : wave number

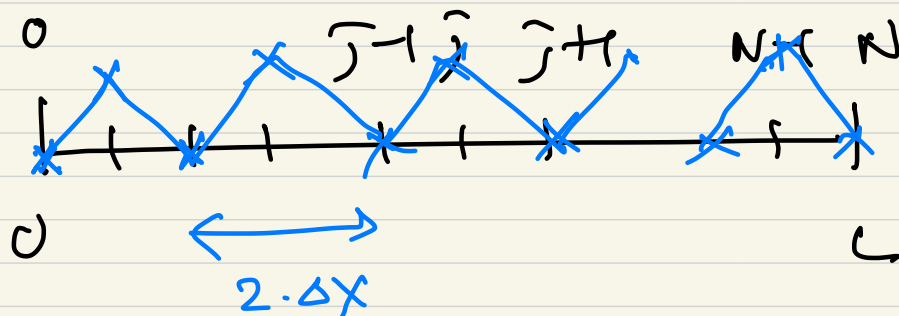


$\Delta x = h = L/N$: grid spacing
uniform mesh
largest wavelength e^{ikx}

$$\Delta k \cdot L = 2\pi$$

$$\rightarrow \boxed{\Delta k = \frac{2\pi}{L}} : \text{smallest wavenumber}$$

largest wavenumber?
smallest wave length



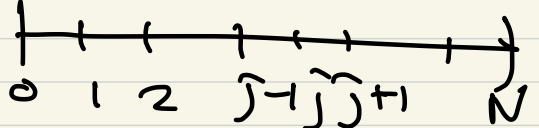
$$K \cdot (2\Delta x) = 2\pi$$

$$\rightarrow \boxed{K = \frac{\pi}{\Delta x} = \frac{\pi}{h} = \frac{\pi N}{L}}$$

largest wavenumber

$$f(x) = e^{ikx} \rightarrow f' = ik e^{ikx} = ik f \quad \text{: exact sol.}$$

Q: how accurately $\mathcal{O}2$ computes the derivatives of f for different values of k ? 0 L

$$x_j = h \cdot j = \frac{L}{N} j, \quad j=0, 1, 2, \dots, N$$


On this grid, e^{ikx} ranges from a constant for $k=0$ to a highly oscillatory ft. of period equal to mesh width ($2\Delta x$) for $k = \pi N/L (= \pi/h)$.

$$f'_j = \frac{f_{j+1} - f_{j-1}}{2h} \quad (\mathcal{O}2)$$

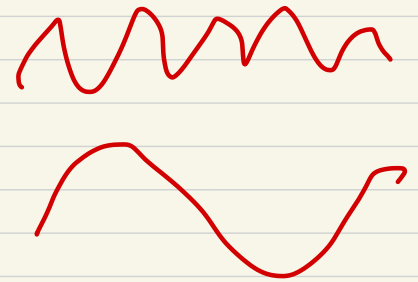
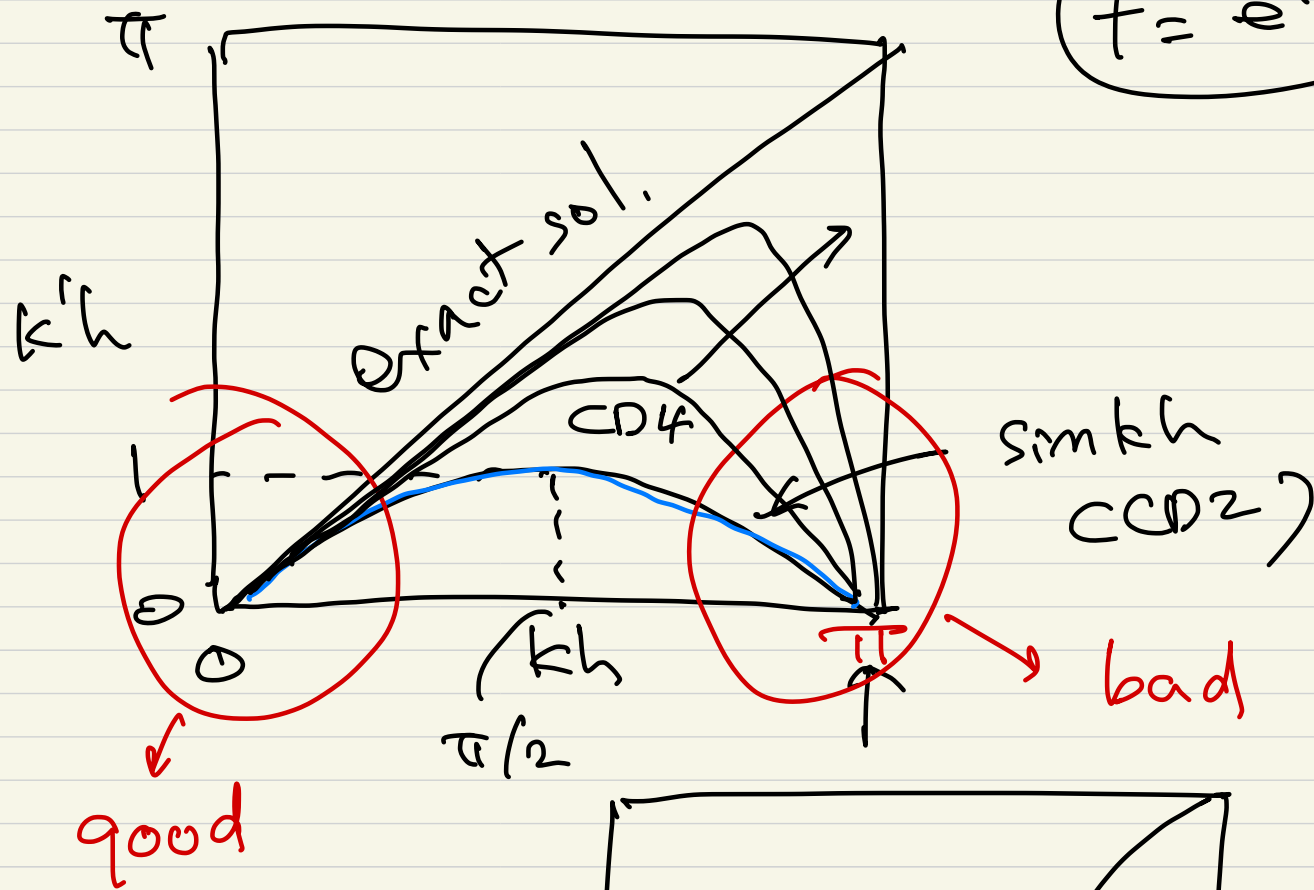
$$= \frac{1}{2h} (e^{ikx_{j+1}} - e^{ikx_{j-1}}) = \frac{1}{2h} e^{ikx_j} (e^{ikh} - e^{-ikh})$$

$$= i \cdot \frac{\sin kh}{h} e^{ikx_j} = i \frac{\sin kh}{h} f_j$$

k' : modified wavenumber

$\rightarrow k'h = \sin kh \quad (\text{CD2})$

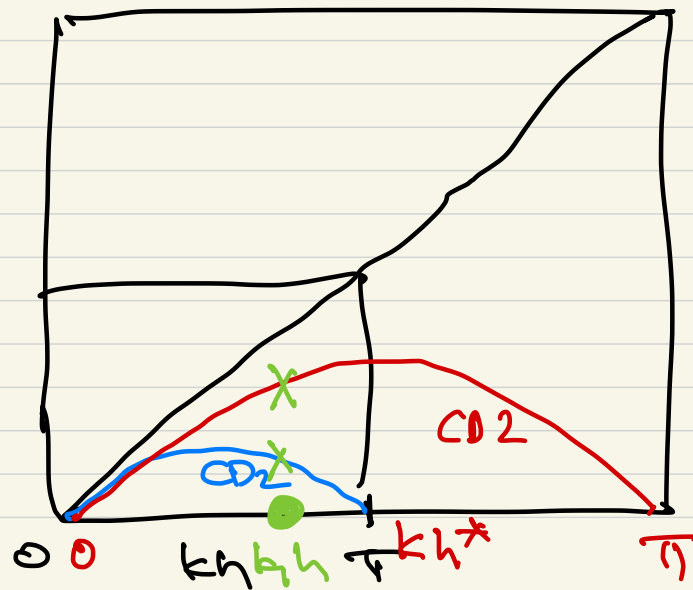
$f = e^{ikx}$
 $\rightarrow f' = ik e^{ikx}$



$f = e^{ikx}$
 \uparrow
 (k, h)

$N \rightarrow 2N$

$h \rightarrow \frac{h}{2} = h^*$



1st-order FD : $f_j' = \frac{f_{j+1} - f_j}{h} + O(h)$ $f' = i k f$
exact sol.

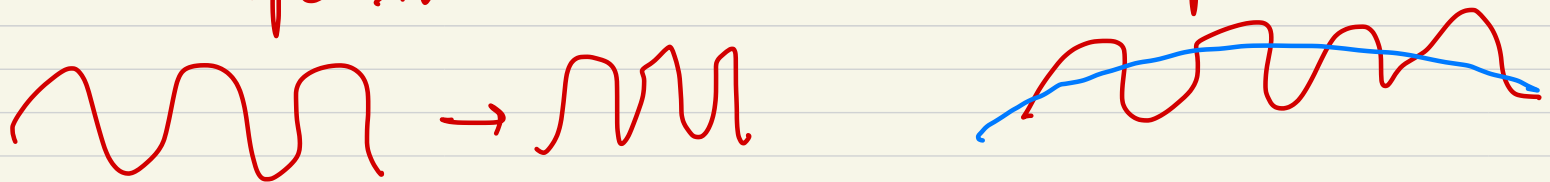
$$= \frac{1}{h} (e^{i k x_{j+1}} - e^{i k x_j})$$

$$= \frac{1}{h} e^{i k x_j} (e^{i k h} - 1) = \frac{1}{h} (-1 + \cos kh + i \sin kh) \times e^{i k x_j}$$

$$\rightarrow f_j' = i \left[\frac{\sin kh}{h} + i \frac{1 - \cos kh}{h} \right] f_j$$

k'_{CO2} \rightarrow dispersive error

k' : modified wavenumber complex number \rightarrow dissipative error



each FD method has its own modified wavenumber.

2.4 Padé approximations

$$f'_j + af_j + bf_{j-1} + cf_{j+1} + df_{j+2} = \mathcal{O}(h^2)$$

Include derivatives too in the formula.

ex) Find the most accurate formula of f'_j that involves

$$f_{j+1}, f_{j-1}, f_j, f'_{j+1}, f'_{j-1}$$

$$\rightarrow f'_j + a_0 \underline{f'_{j+1}} + a_1 \underline{f'_{j-1}} + b_0 f_j + b_1 f_{j-1} + b_2 f_{j+1} = \mathcal{O}(h^2)$$

Taylor table

	f_j	f'_j	f''_j	f'''_j	$f^{(iv)}_j$	$f^{(v)}_j \dots$
f'_j	0	1	0	0	0	0 ...
$a_0 f'_{j+1}$	$a_0 \cdot 0$	$a_0 \cdot 1$	$a_0 \cdot h$	$a_0 \cdot \frac{1}{2} h^2$	$a_0 \cdot \frac{1}{6} h^3$	$a_0 \cdot \frac{1}{24} h^4 \dots$
$a_1 f'_{j-1}$	0	$a_1 \cdot (-1)$	$a_1 \cdot h$	$a_1 \cdot \frac{1}{2} h^2$	$a_1 \cdot \frac{1}{6} (-h)^3$	$a_1 \cdot \frac{1}{24} h^4 \dots$
$b_0 f_j$	b_0	0	0	0	0	0
$b_1 f_{j-1}$	$b_1 \cdot 1$	$b_1 \cdot (-h)$	$b_1 \cdot \frac{1}{2} h^2$	$b_1 \cdot \frac{1}{6} h^3$	$b_1 \cdot \frac{1}{24} h^4$...
$b_2 f_{j+1}$	$b_2 \cdot 1$	$b_2 \cdot h$	$b_2 \cdot \frac{1}{2} h^2$	$b_2 \cdot \frac{1}{6} (-h)^3$	$b_2 \cdot \frac{1}{24} h^4$...

$$f_j' + a_0 f_{j+1}' + a_1 f_{j-1}' + b_0 f_j + b_1 f_{j-1} + b_2 f_{j+1}$$

$$= \underbrace{(b_0 + b_1 + b_2)}_{=0} f_j + \underbrace{(1 + a_0 - a_1 - b_1 h + b_2 h)}_{=0} f_j' + \underbrace{(a_0 h + a_1 h + \frac{1}{2} b_1 h^2 + \frac{1}{2} b_2 h^2)}_{=0} f_j'' + \underbrace{(\dots)}_{=0} f_j''' + \underbrace{(\dots)}_{=0} f_j^{(iv)} + \dots$$

5 unknowns

$$\hookrightarrow a_0 = \dots, a_1 = \dots, b_0 = \dots, b_1 = \dots, b_2 = \dots$$

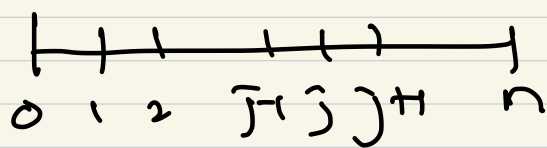
$$\rightarrow f_j' + \frac{1}{4} f_{j+1}' + \frac{1}{4} f_{j-1}' - \frac{3}{4h} f_{j+1} + \frac{3}{4h} f_{j-1} = \frac{h^4}{120} f_j^{(v)} + \dots$$

$$\rightarrow \boxed{f_{j-1}' + 4f_j' + f_{j+1}' = \frac{3}{h} (f_{j+1} - f_{j-1}) + \frac{1}{30} h^4 f_j^{(v)} + \dots}$$

3 grid pts. \rightarrow 4th-order accuracy

\Rightarrow "compact" scheme

$\mathcal{O}(h^4) \rightarrow$ 4th-order accuracy



$$j = 1, 2, \dots, n-1$$

$$\begin{bmatrix} \text{ } & \text{ } & \text{ } & \text{ } & \text{ } \\ \text{ } & \text{ } & \text{ } & \text{ } & \text{ } \\ \text{ } & \text{ } & \text{ } & \text{ } & \text{ } \\ \text{ } & \text{ } & \text{ } & \text{ } & \text{ } \\ \text{ } & \text{ } & \text{ } & \text{ } & \text{ } \end{bmatrix} \begin{bmatrix} f_j' \end{bmatrix} = \begin{bmatrix} \text{ } \\ \text{ } \\ \text{ } \\ \text{ } \\ \text{ } \end{bmatrix}$$

tri-diagonal matrix system \rightarrow easy to solve

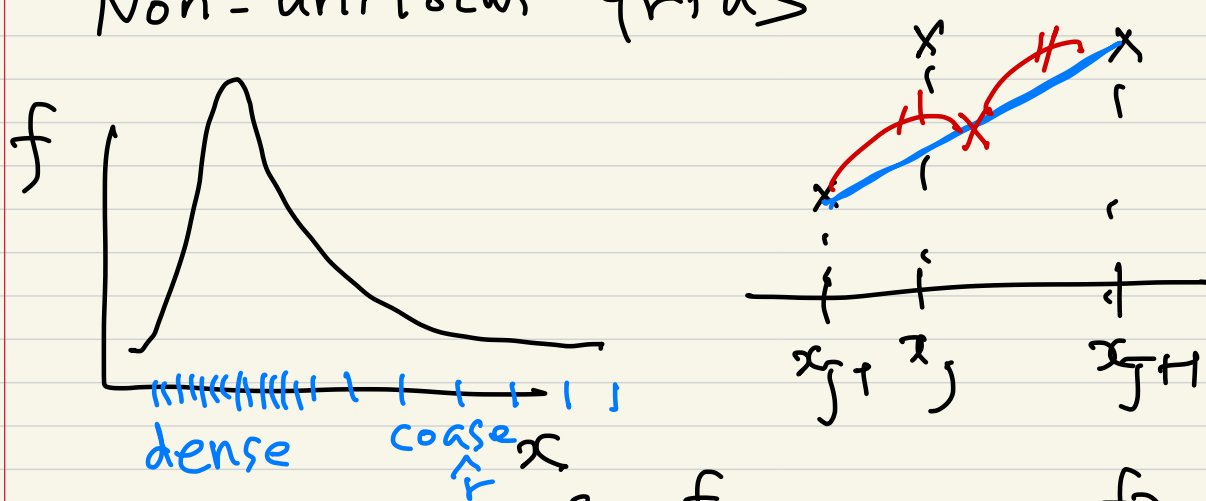
Padé approximations can be easily extended to higher derivatives.

$$\text{ex)} \quad \frac{1}{12} f_{j-1}'' + \frac{10}{12} f_j'' + \frac{1}{12} f_{j+1}'' = \frac{1}{h^2} (f_{j+1} - 2f_j + f_{j-1}) + O(h^4)$$

3pts $\rightarrow O(h^4)$ 4th-order accurate compact scheme

S. Lele (JCP, 1992) - compact schemes \rightarrow #6246 citations
 \downarrow
 J. of Computational Physics

2.5 Non-uniform grids



$$f_j' = \frac{f_{j+1} - f_{j-1}}{x_{j+1} - x_{j-1}} + O(h^2)$$

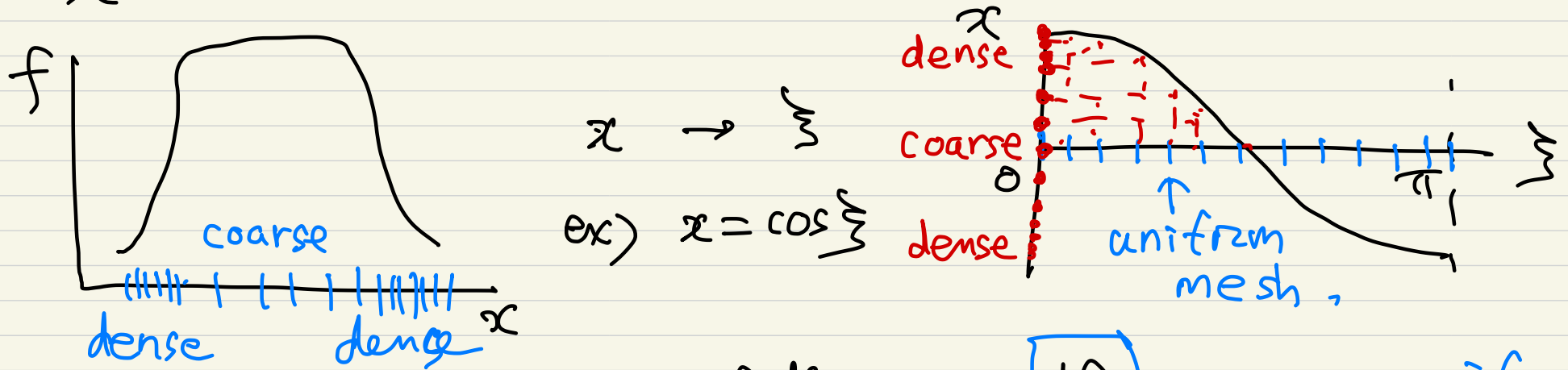
\downarrow strictly 1st-order accurate

$$h_j = x_{j+1} - x_j$$

$$f_j'' = 2 \left[\frac{f_{j-1}}{h_j (h_j + h_{j+1})} - \frac{f_j}{h_j h_{j+1}} + \frac{f_{j+1}}{h_{j+1} (h_j + h_{j+1})} \right] + O(h^2)$$

FD formula for non-uniform mesh generally has a lower order of accuracy than their counterpart with the same stencil for uniform mesh.

Use a coordinate transformation

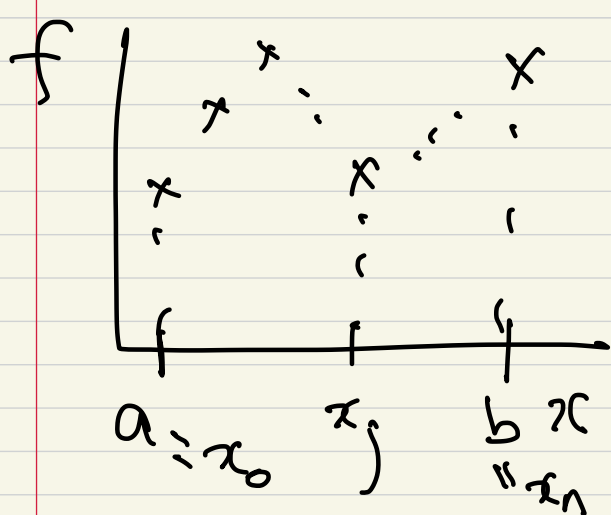


$$\xi = g(x)$$

$$\frac{df}{dx} = \frac{df}{d\xi} \frac{d\xi}{dx} = g' \left[\frac{df}{d\xi} \right] \leftarrow \text{FD on uniform mesh in } \xi$$

$$\begin{aligned} \frac{d^2f}{dx^2} &= \frac{d}{dx} \left(\frac{df}{dx} \right) = \frac{d}{dx} \left(g' \frac{df}{d\xi} \right) = g'' \frac{df}{d\xi} + g' \frac{d^2f}{d\xi^2} \frac{d\xi}{dx} \\ &= g'' \left[\frac{df}{d\xi} \right] + g'^2 \left[\frac{d^2f}{d\xi^2} \right] \end{aligned}$$

Ch.3 Numerical integration (quadrature)



$$I = \int_a^b f(x) dx = \sum_{j=0}^n f_j w_j$$

weight

3.1 Trapezoidal and Simpson's rule

Lagrange polynomial $P(x) = \sum_{j=0}^n f_j L_j(x)$

$$I = \int_a^b f(x) dx = \int_a^b P(x) dx = \int_a^b \sum_{j=0}^n f_j L_j dx = \sum_{j=0}^n f_j \int_a^b L_j(x) dx$$

$$= (b-a) \sum_{j=0}^n c_j^n f_j \quad \text{where } c_j^n = \frac{1}{b-a} \int_a^b L_j(x) dx$$

Newton-Cotes formula \swarrow cotes number

For $n=1$, $x_0 = a$, $x_1 = b$

$$L_0(x) = \frac{x-b}{a-b}, \quad L_1(x) = \frac{x-a}{b-a}$$

$$c_0^1 = \frac{1}{b-a} \int_a^b L_0(x) dx = \frac{1}{2}, \quad c_1^1 = \frac{1}{b-a} \int_a^b L_1(x) dx = \frac{1}{2}$$

$$I = (b-a) \left(\frac{1}{2} f_0 + \frac{1}{2} f_1 \right) = (b-a) \left[\frac{1}{2} f(a) + \frac{1}{2} f(b) \right]$$

trapezoidal rule

For $n=2$, $x_0 = a$, $x_1 = \frac{1}{2}(a+b)$, $x_2 = b \rightarrow c_0^2, c_1^2, c_2^2$

$$\rightarrow I = \frac{b-a}{6} \left[f(a) + 4f\left(\frac{a+b}{2}\right) + f(b) \right] \quad L_0, L_1, L_2$$

Simpson's rule

n	N	NC_0^n	NC_1^n	NC_2^n	NC_3^n	NC_4^n	NC_5^n	NC_6^n	error
1	2	1	1						$8.3 \times 10^{-2} \Delta^3 f^{III}$
2	6	1	4	1					$3.5 \times 10^{-4} \Delta^5 f^{IV}$
3	8	1	3	3	1				$1.6 \times 10^{-4} \Delta^5 f^{IV}$
4	90	7	32	12	32	7			$5.2 \times 10^{-7} \Delta^7 f^{VII}$
5	288	19	25	50	50	25	19		$3.6 \times 10^{-7} \Delta^7 f^{VII}$
6	860	41	216	27	272	27	216	41	$6.4 \times 10^{-10} \Delta^9 f^{VIII}$
⋮									

