

Chapter 11

Kinetic Theory of Gases (1)

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11.1 Basic Assumption

Basic assumptions of the kinetic theory

- 1) Large number of molecules (Avogadro's number)

$$N_A = 6.02 \times 10^{26} \text{ molecules per kilomole}$$

- 2) Identical molecules which behave like hard spheres
- 3) No intermolecular forces except when in collision
- 4) Collisions are perfectly elastic
- 5) Uniform distribution throughout the container

$$n = \frac{N}{V} \quad dN = n dV$$

n: The average number of molecules per unit volume

11.1 Basic Assumption

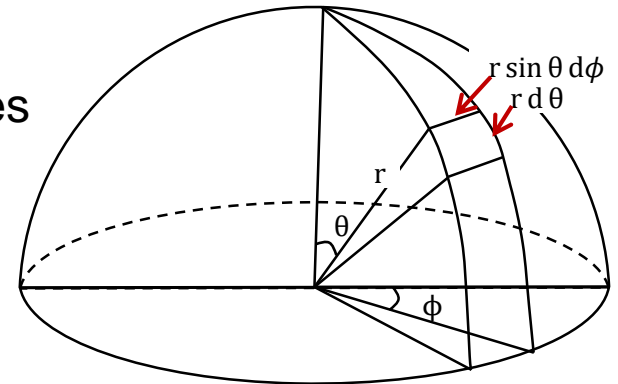
- 6) Equal probability on the direction of molecular velocity average number of intersections of velocity vectors per unit area; $\frac{N}{4\pi r^2}$

the number of intersections in dA

$$d^2 N_{\theta\phi} = \frac{N}{4\pi r^2} dA = \frac{N \sin \theta d\theta d\phi}{4\pi} \quad \text{Where } dA = r^2 \sin \theta d\theta d\phi$$

$$d^2 n_{\theta\phi} = \frac{n \sin \theta d\theta d\phi}{4\pi}$$

$N_{\theta\phi}$: The number of molecules having velocities in a direction ($\theta \sim \theta + d\theta$) and ($\phi \sim \phi + d\phi$)



11.1 Basic Assumption

7) Magnitude of molecular velocity : $0 \sim \infty$



c (speed of light)

dN_v : The number of molecules with specified speed ($v \sim v+dv$)

11.1 Basic Assumption

- Let dN_v as the number of molecules with specified speed ($v \sim v+dv$)
- $\int_0^{\infty} dN_v = N$
- Mean speed is $\bar{v} = \frac{1}{N} \int_0^{\infty} v dN_v$
- Mean square speed is $\overline{v^2} = \frac{1}{N} \int_0^{\infty} v^2 dN_v$
- Square root of $\overline{v^2}$ is called the root mean square or rms speed:

$$v_{rms} = \sqrt{\overline{v^2}} = \sqrt{\frac{1}{N} \int_0^{\infty} v^2 dN_v}$$

- The n-th moment of distribution is defined as

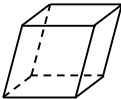
$$\overline{v^n} = \frac{1}{N} \int_0^{\infty} v^n dN_v$$

11.2 Molecular Flux

- The number of gas molecules that strike a surface per unit area and unit time
- Molecules coming from particular direction θ, ϕ with specified speed v in time dt

$$\rightarrow \theta\phi v \text{ collision } \begin{cases} \theta \sim \theta + d\theta \\ \phi \sim \phi + d\phi \\ v \sim v + dv \end{cases}$$

- The number of $\theta\phi v$ collisions with dA

= $\theta\phi v$ molecules in 

= $\theta\phi$ molecules with speed v

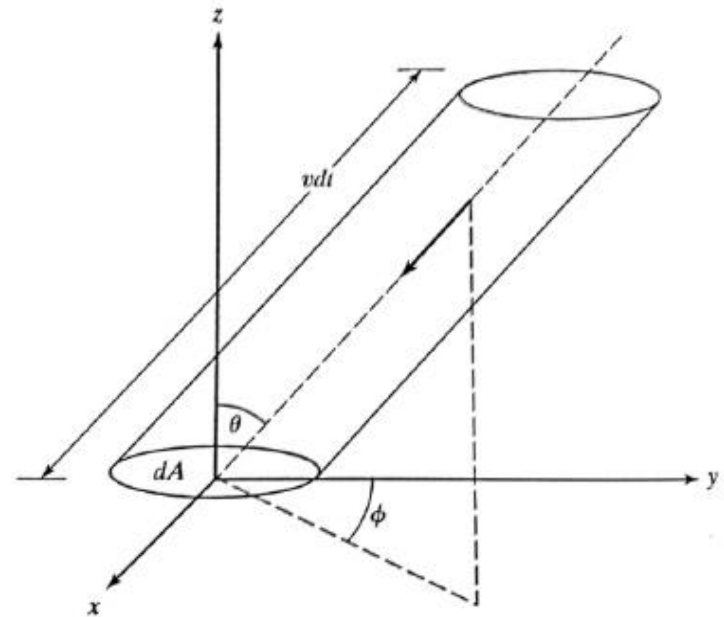
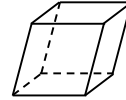


Fig. Slant cylinder geometry used to calculate the number of molecules that strike the area dA in time dt .

11.2 Molecular Flux

- How many molecules in unit volume



dn_v : Density between speed ($v \sim v+dv$)

dA : Surface of spherical shell of radius v and thickness dv (i.e., θ, ϕ molecules)

$$d^3n_{\theta\phi v} = dn_v \cdot \frac{dA}{A} = dn_v \frac{v^2 \sin \theta d\theta d\phi}{4\pi v^2}$$

- The number of $\theta\phi v$ molecules in the cylinder toward dA

Volume of cylinder: $dV = dA (v dt \cos \theta)$

$$d^3n_{\theta\phi v} dV = (dA v dt \cos \theta) dn_v \frac{\sin \theta d\theta d\phi}{4\pi}$$

11.2 Molecular Flux

- The number of collisions per unit area and time (i.e., particle flux)

$$\frac{d^3 n_{\theta\phi v} dV}{dA dt} = \frac{1}{4\pi} v dn_v \sin\theta \cos\theta d\theta d\phi$$

- Total number of collisions per unit area and time by molecules having all speed

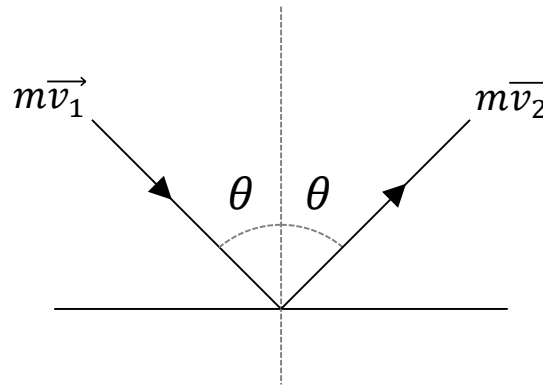
$$\int \frac{d^3 n_{\theta\phi v} dV}{dA dt} = \int_0^{2\pi} d\phi \int_0^{\pi/2} \sin\theta \cos\theta d\theta \cdot \frac{1}{4\pi} \int_0^{\infty} v dn_v = \frac{1}{4} n \bar{v} \quad (\int_0^{\infty} v dn_v = n \bar{v})$$

Cf. average speed $\bar{v} = \frac{\sum \bar{v}}{N} = \frac{\sum N_i v_i}{N} = \frac{\sum n_i v_i}{\sum n_i} = \frac{\int v dn_v}{n}$

11.3 Gas Pressure and Ideal Gas Law

- Gas pressure in Kinetic theory

Gas pressure is interpreted as impulse flux of particles striking a surface



11.3 Gas Pressure and Ideal Gas Law

- Perfect elastic $v = v'$

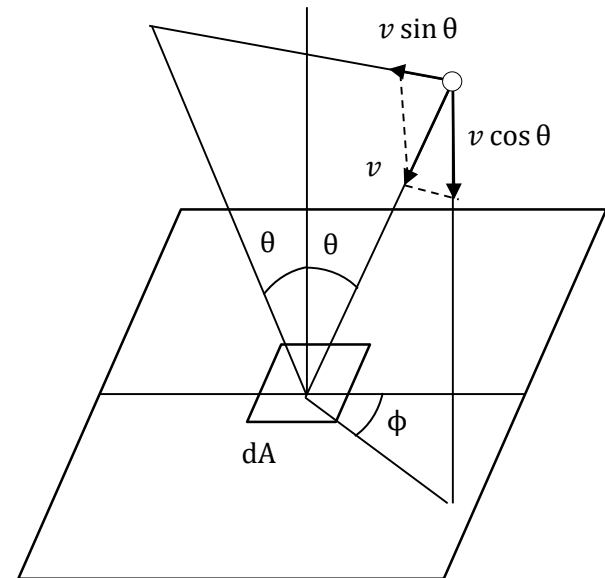
- Average force exerted by molecules $F = \frac{d(m\vec{v})}{dt} = m\vec{a} + \dot{m}\vec{v}$

- Momentum change of one molecule (normal component only)

$$mv\cos\theta - (-mv\cos\theta) = 2mv\cos\theta$$

- The number of $\theta\phi v$ collisions for dA, dt

$$\frac{d^3n_{\theta\phi v}dV}{dA dt} = \frac{1}{4\pi} v dn_v \sin\theta \cos\theta d\theta d\phi$$



11.3 Gas Pressure and Ideal Gas Law

- Change in momentum due to $\theta\phi v$ collisions in time dt

$$2mv\cos\theta \times \frac{1}{4\pi} v dn_v \sin\theta \cos\theta d\theta d\phi = \frac{1}{2\pi} mv^2 dn_v \sin\theta \cos^2\theta d\theta d\phi dA dt$$

- Change in momentum p in all v collisions $0 < \theta \leq \frac{\pi}{2}, 0 < \phi \leq 2\pi$ at all speed

$$dp = \int_0^\infty \int_0^{\pi/2} \int_0^{2\pi} \frac{1}{2\pi} mv^2 dn_v \sin\theta \cos^2\theta d\theta d\phi \cdot dA dt = \frac{1}{3} mn \overline{v^2} dA dt$$

- Change in momentum from collisions of molecules with unit time

$$\frac{dp}{dt} = d\vec{F} = \frac{1}{3} mn \overline{v^2} dA$$

$$\text{cf. } \overline{v^2} = \frac{\sum v^2}{N} = \frac{\int v^2 dn_v}{n}$$

- Average pressure $\bar{P} = \frac{d\vec{F}}{dA}$

$$\bar{P} = \frac{1}{3} mn \overline{v^2}$$

11.3 Gas Pressure and Ideal Gas Law

Since $n = \frac{N}{V}$ then pressure $P = \frac{1}{3} \frac{N}{V} m \overline{v^2} \quad \therefore PV = \frac{1}{3} N m \overline{v^2}$

EOS of an ideal gas: $PV = n \bar{R} T = m R T = \frac{N}{N_A} \bar{R} T = N k T$

N_A : Avogadro's number : 6.02×10^{26} molecules/kmole

k_B : Boltzmann constant : $k_B = \frac{\bar{R}}{N_A} = 1.38 \times 10^{-23}$ J/K

$$PV = \frac{1}{3} N m \overline{v^2} = N k T$$

$$\therefore \frac{1}{2} m \overline{v^2} = \frac{3}{2} k T$$

The temperature is proportional to the average kinetic energy of molecule

11.4 Equipartition of Energy

- Equipartition of energy

Because of even distribution of velocity of particles,

$$\overline{v^2} = \overline{v_x^2} + \overline{v_y^2} + \overline{v_z^2},$$

By assumption, no preferred direction

$$\overline{v_x^2} = \overline{v_y^2} = \overline{v_z^2} = \frac{1}{3}\overline{v^2} \quad \rightarrow \quad \frac{1}{2}m\overline{v_x^2} = \frac{1}{6}m\overline{v^2} = \frac{1}{2}kT$$

It can be interpreted that a degree of freedom allocate energy of $\frac{1}{2}kT$

11.5 Specific Heat

Total energy of a molecule in Cartesian coordinate

$$\bar{\epsilon} = \bar{\epsilon}_x + \bar{\epsilon}_y + \bar{\epsilon}_z = \frac{1}{2} m \overline{v_x^2} + \frac{1}{2} m \overline{v_y^2} + \frac{1}{2} m \overline{v_z^2} = \left(\frac{kT}{2} + \frac{kT}{2} + \frac{kT}{2} \right) = \frac{3}{2} kT$$

General expression of total energy of molecules for f -DOF (Degree of Freedom)





$$U = N\bar{\epsilon} = \frac{f}{2} NkT = \frac{f}{2} nRT \quad \leftrightarrow \quad u = \frac{U}{n} = \frac{f}{2} RT$$

$$c_v = \left(\frac{\partial u}{\partial T} \right)_v = \frac{f}{2} R \quad \text{from the above equation}$$

$$c_p = \left(\frac{\partial h}{\partial T} \right)_p = \frac{f}{2} R + R = \frac{(f+2)}{2} R \quad \text{cf) } c_p = c_v + R$$

The ratio of specific heat: $\gamma = \frac{c_p}{c_v} = \frac{f+2}{f}$

11.5 Specific Heat

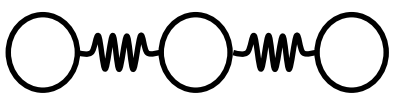
<p>Monatomic gas</p>		$\frac{1}{2}mv_x^2, \frac{1}{2}mv_y^2, \frac{1}{2}mv_z^2$  DOF
<p>Diatomic gas</p>		$\frac{1}{2}mv_x^2, \frac{1}{2}mv_y^2, \frac{1}{2}mv_z^2$ <u>Translational</u> $\frac{1}{2}I\omega_x^2, \frac{1}{2}I\omega_y^2, \frac{1}{2}I\omega_z^2$ <u>Rotational</u> negligible $\frac{1}{2}kx^2, \frac{1}{2}m\dot{x}^2$ no y,z vibration <u>Vibrational</u>  DOF

$$\frac{c_p}{c_v} = \frac{3 + 2}{3} = 1.67$$

$$\frac{c_p}{c_v} = \frac{5 + 2}{5} = 1.4$$

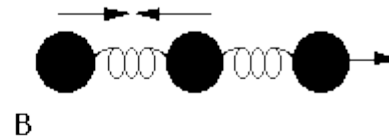
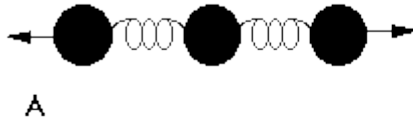
Near room temperature, rotational or vibrational DOF are excited, but not both. DOF: 7 → 5

11.5 Specific Heat

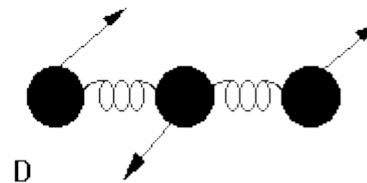
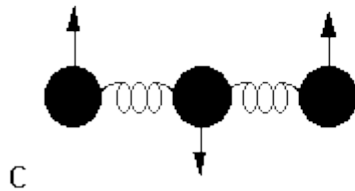
<p>Triatomic gas</p>	<p>CO₂</p> 	<p>translational 3 rotational 2 vibrational 4</p> <p>9 DOF</p>
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$$\frac{c_p}{c_v} = \frac{7 + 2}{7} = 1.28$$

- Vibration modes of CO₂

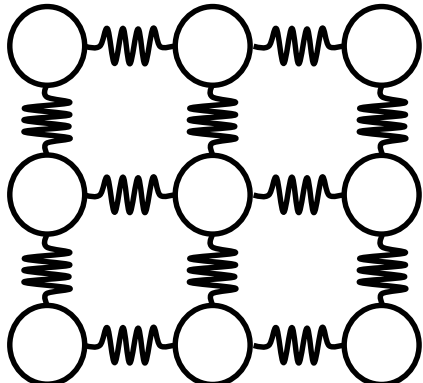


Stretch



Bending

11.5 Specific Heat

<p>Solid</p>		<p>$\frac{kT}{2}$ (kinetic) ← x,y,z direction</p> <p>$\frac{kT}{2}$ (potential) ↙</p> $U = \frac{3kT}{2} + \frac{3kT}{2} = 3NkT$ $c_v = 3R \text{ (Dulong-Petit Law)}$
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