

# Intrinsic Semiconductor

mistake

$$n_i^2 = np = N_C N_V \exp(-E_G/kT)$$

$$\frac{N_C}{N_V} = \frac{e^{2(E_C - E_F)/kT}}{e^{E_G/kT}}$$

$$n_i = N_C e^{-(E_C - E_F)/kT} = N_C \left( \frac{m_h^*}{m_e^*} \right)^{3/4} e^{-E_G/2kT}$$

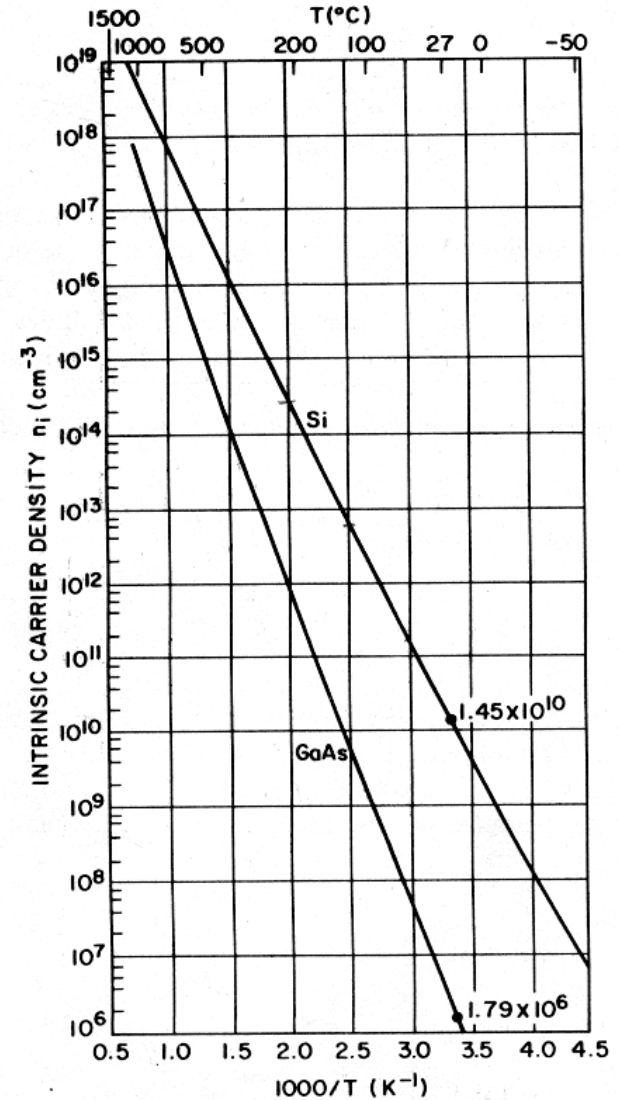
$\propto T^{3/2}$

a plot of  $\ln(n_i T^{-3/2})$  vs.  $1/T$

→ straight line with slope of  $-E_G/2k$

$$\sigma_i = N_C q \left( \frac{m_h^*}{m_e^*} \right)^{3/4} (\mu_n + \mu_p) e^{-E_G/2kT}$$

Note: T-dependence of  $N_C$  is ignored.



# Donor and acceptor imperfections (Extrinsic Semiconductor)

- The electrical conductivity of semiconductor is controlled by thermal excitations across the band gap of the material and thermal excitation from **localized imperfections**.
- Substitution of atoms with different valence into a host crystal leads to **shallow energy level** close to either the conduction band or valence band.
  - Impurity atom with +1 valence compared to the host atom: donate electron (donor), creating an electron
  - Impurity atom with -1 valence compared to the host atom: accept electron (acceptor), creating a hole
- The wave function corresponding to an electron or a hole bound to a shallow donor or acceptor extends over many atomic dimensions so that the indeterminacy of  **$\Delta x$  is large** and the indeterminacy of  **$\Delta k$  is small**.

# Donor and acceptor imperfections (Extrinsic Semiconductor)

- Ionization of a donor



$$\text{Ionization energy} = (E_C - E_D)$$

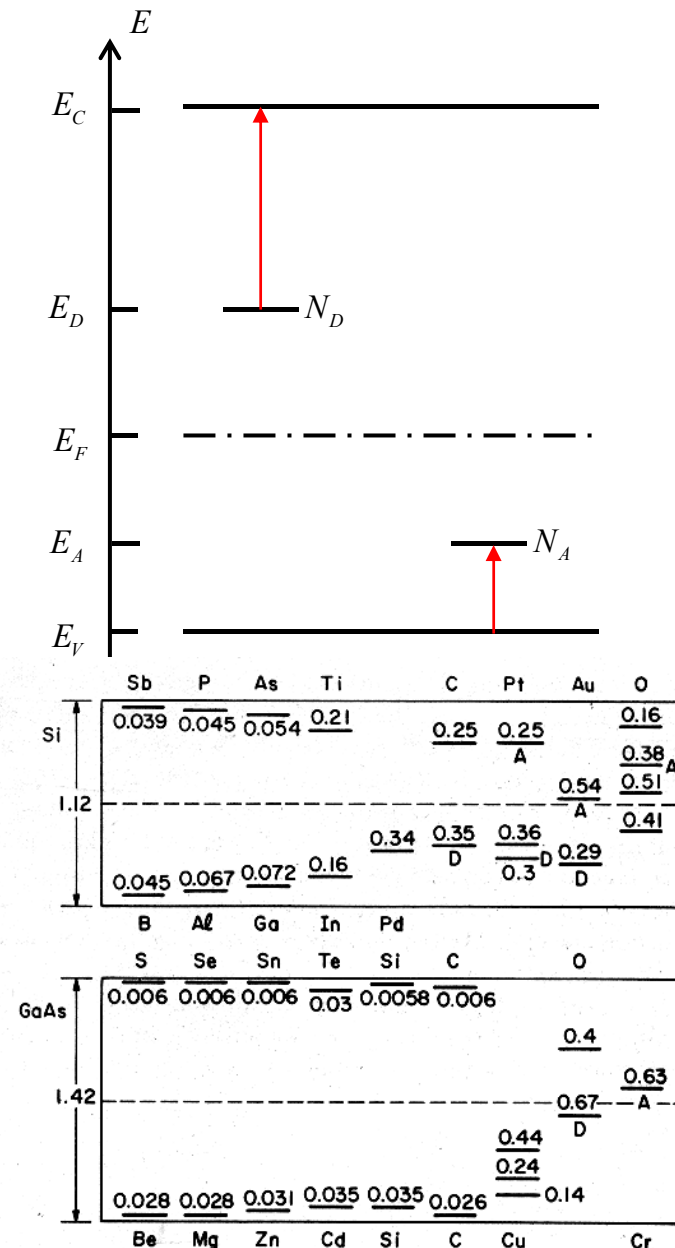
- Ionization of an acceptor



$$\text{Ionization energy} = (E_A - E_V)$$

Examples:

- donor: Si on a Ga site in GaAs, P on a Si site in Si,
- Acceptor: Si on a As site in GaAs, Ga on a Si site in Si

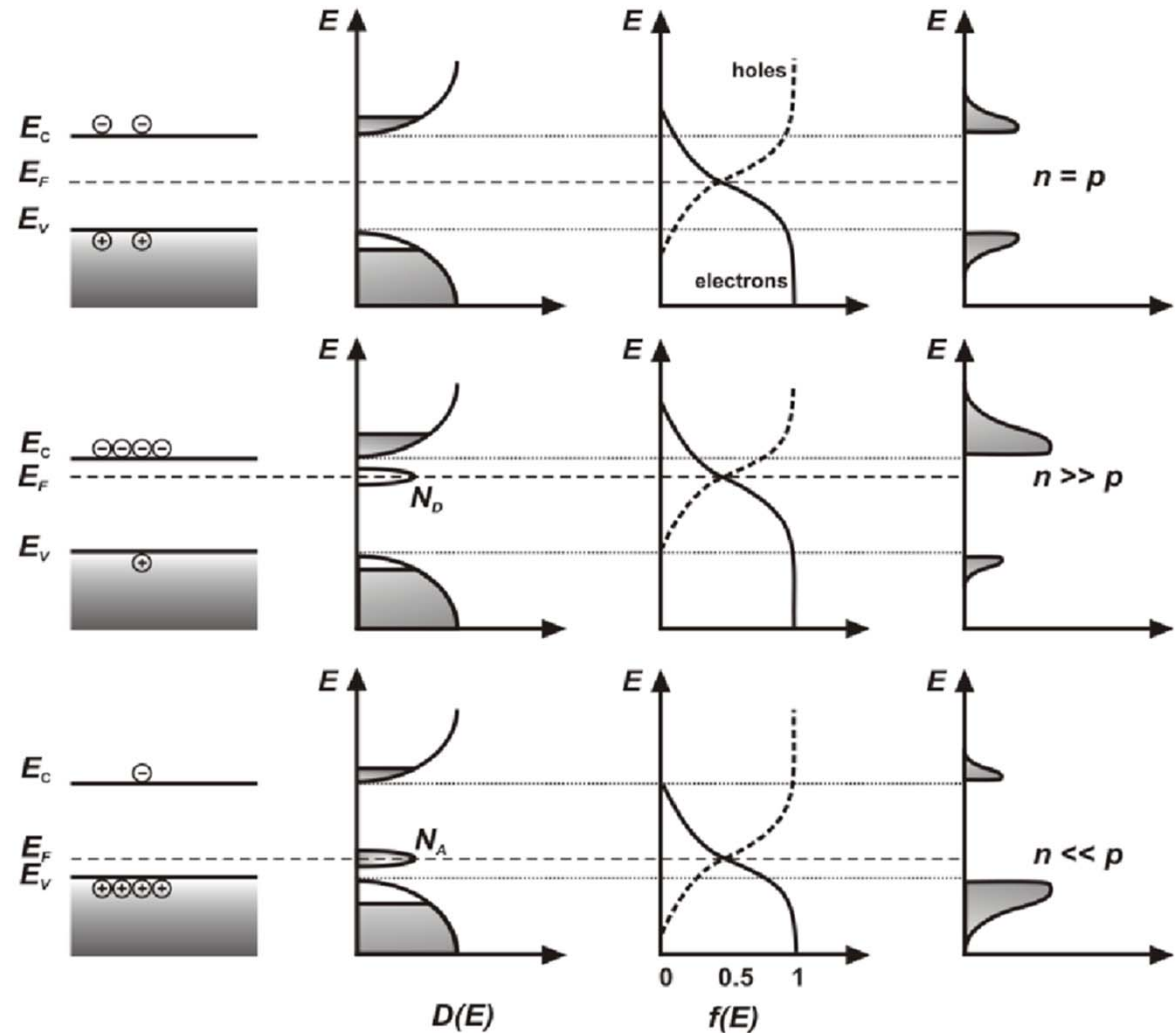


# Donor and acceptor imperfections (Extrinsic Semiconductor)

Complete ionization  
 $n = N_D$

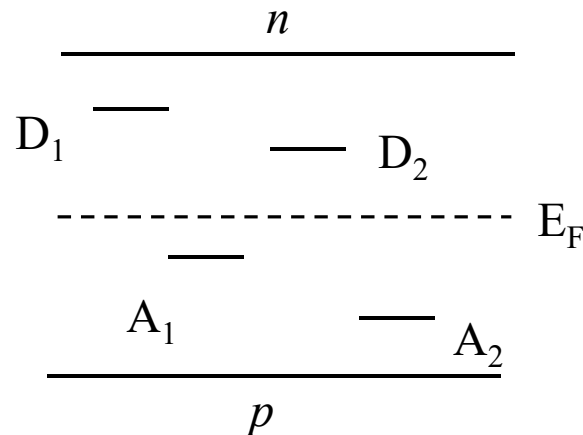
$$E_C - E_F = kT \ln\left(\frac{N_C}{N_D}\right)$$

The higher donor concentration,  
 the smaller energy difference  
 → Fermi level moves toward the  
 bottom of the conduction band



# Electrical Conductivity in Extrinsic Semiconductor

- Various donors and acceptors in extrinsic semiconductor



$N_D$ : the total density of donor impurity

$n_D$ : the density of electron-occupied donor (non-ionized donor)

$N_D^+$ : the density of the ionized donor

$N_A$ : the total density of acceptor impurity

$n_A$ : the density of electron-occupied acceptor (ionized acceptor)

$N_A^-$ : the density of the ionized acceptor

- From T-dependence of  $E_F$ , T-dependence of free carrier density ( $n(T)$ ) can be obtained.
- From charge neutrality condition

$$\sum_i (\text{negative charge})_i = \sum_i (\text{positive charge})_i$$

$$n + \sum_i (\text{electron occupied acceptors}) = p + \sum_i (\text{electron unoccupied donors})$$

$$n + n_A = (N_D - n_D) + p$$

# Electrical Conductivity in Extrinsic Semiconductor

$$(1) \quad n_D = N_D \left\{ \left( \frac{1}{g} \right) \exp \left[ \frac{(E_D - E_F)}{kT} \right] + 1 \right\}^{-1} \quad (g: \text{degeneracy factor})$$

ex)  $g = 1$  for electrons in conduction and holes in valence  
 $g = 2$  for donor and acceptor due to two spin orientations

: when  $E_F$  decreases from donor level,  $n_D$  approaches zero.  
 when  $E_F$  increases toward donor level,  $n_D$  approaches  $N_D$ .

$$(2) \quad \frac{(N_A - n_A)}{\underbrace{\hspace{1.5cm}}} = N_A \left\{ \left( \frac{1}{g} \right) \exp \left[ \frac{(E_F - E_A)}{kT} \right] + 1 \right\}^{-1} \quad (g: \text{degeneracy factor})$$

 Hole-occupied acceptors

: when  $E_F$  increases from acceptor level,  $(N_A - n_A)$  approaches zero.  
 when  $E_F$  decreases toward acceptor level,  $(N_A - n_A)$  approaches  $N_A$ .

Positively charged donors:  $N_D^+ = (N_D - n_D) = N_D \left\{ 2 \exp \left[ \frac{(E_F - E_D)}{kT} \right] + 1 \right\}^{-1}$

Negatively charged acceptors:  $N_A^- = n_A = N_A \left\{ 2 \exp \left[ \frac{(E_A - E_F)}{kT} \right] + 1 \right\}^{-1}$

# Electrical Conductivity in Extrinsic Semiconductor

From  $n + n_A = (N_D - n_D) + p$

- For intrinsic material,  $n_A = N_D - n_D = 0$

$$n = p$$

- For donor only,  $n_A = 0, p \ll n$

$$n = N_D - n_D : \text{At high T, all of donors are ionized, } n \rightarrow N_D$$

- For acceptor only,  $N_D - n_D = 0, n \ll p$

$$n_A = p : \text{At high T, all of acceptors are ionized, } p \rightarrow N_A$$

- For donors and acceptors,

(1)  $N_D \gg N_A, p \ll n$  (assumption: all acceptors are ionized,  $n_A \approx N_A$ )

$$\rightarrow n + n_A = (N_D - n_D)$$

$\rightarrow$  As T increases, all donors are ionized,  $n \rightarrow (N_D - N_A)$

(2)  $N_D \ll N_A, n \ll p$  (assumption: all donors are ionized,  $(N_D - n_D) \approx N_D$ )

$$\rightarrow n_A = N_D + p$$

$\rightarrow$  As T increases, all acceptors are ionized,  $p \rightarrow (N_A - N_D)$

(3)  $N_D \approx N_A, p \approx n \approx 0$  ( $\because$  At low T, donor and acceptor are dominant)

$$\rightarrow n_A = (N_D - n_D)$$

# Electrical Conductivity in Extrinsic Semiconductor

Temperature dependence of Fermi level for donor only,

$$n_A = 0, \quad p \ll n$$

$$n = N_D - n_D \rightarrow N_C \exp \left[ -\frac{(E_C - E_F)}{kT} \right] = N_D \left\{ 2 \exp \left[ \frac{(E_F - E_D)}{kT} \right] + 1 \right\}^{-1}$$

At high T extrinsic range, complete ionization of donors

$$\left( \frac{N_D g}{N_C} \right) \exp \left[ \frac{(E_C - E_D)}{kT} \right] \ll 1$$

$$(E_C - E_F) = kT \ln \left( \frac{N_C}{N_D} \right) \quad (n = N_D)$$

$$\sigma = N_D q \mu_n$$

At low T extrinsic range, most donors un-ionized

$$\left( \frac{N_D g}{N_C} \right) \exp \left[ \frac{(E_C - E_D)}{kT} \right] \gg 1$$

$$(E_C - E_F) = \frac{E_C - E_D}{2} + \left( \frac{kT}{2} \right) \ln \left( \frac{N_C g}{N_D} \right) \quad (n \ll N_D)$$

$$n = \left( \frac{N_C N_D}{g} \right)^{1/2} \exp \left[ -\frac{(E_C - E_D)}{2kT} \right]$$

$$\sigma = \left( \frac{N_C N_D}{g} \right)^{1/2} q \mu_n \exp \left[ -\frac{(E_C - E_D)}{2kT} \right]$$



# Electrical Conductivity in Extrinsic Semiconductor

For intrinsic semiconductor

$$n_i = N_C \left( \frac{m_h^*}{m_e^*} \right)^{3/4} \exp \left[ -\frac{E_G}{2kT} \right] \text{ and } \sigma_i = N_C q \left( \frac{m_h^*}{m_e^*} \right)^{3/4} (\mu_n + \mu_p) \exp \left[ -\frac{E_G}{2kT} \right]$$

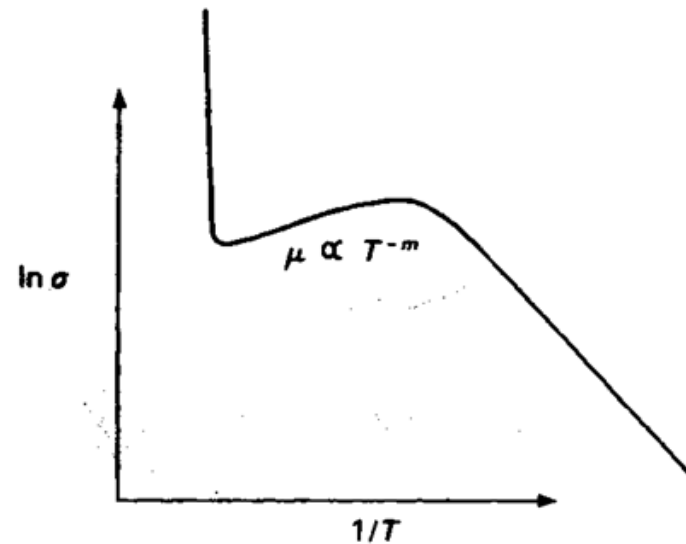
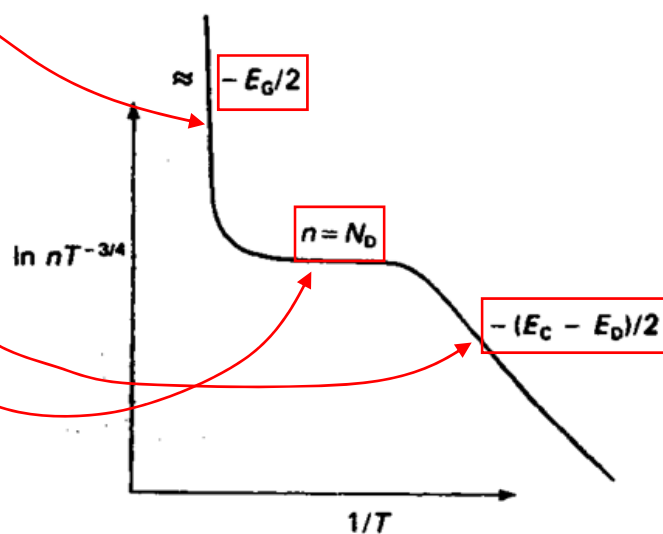
For extrinsic semiconductor with donors only,

At high T extrinsic range,

$$n = N_D \text{ and } \sigma = N_D q \mu_n$$

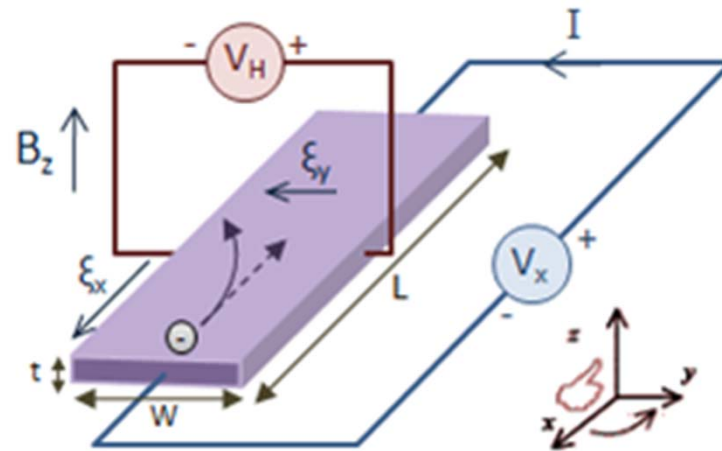
At low T extrinsic range,

$$n = \left( \frac{N_C N_D}{g} \right)^{1/2} \exp \left[ -\frac{(E_C - E_D)}{2kT} \right] \text{ and } \sigma = \left( \frac{N_C N_D}{g} \right)^{1/2} q \mu_n \exp \left[ -\frac{(E_C - E_D)}{2kT} \right]$$



# Hall Effect

The Hall effect is one of the most commonly used of the technique to distinguish the effects of carrier density and carrier mobility on conductivity.



By measuring  $V_H$

$$V_H = \frac{I_x B_z}{nte}$$
$$n = \frac{I_x B_z}{V_H t e}$$

Hall coefficient ( $R_H$ ) =  $\frac{1}{ne}$

$$\mu_{Hall} = \frac{R_H}{\rho}$$

# Other galvanomagnetothermoelectric effect

- Magnetoresistance

: change in electrical conductivity associated with an applied electric fields when magnetic field is applied.

- Seebeck effect

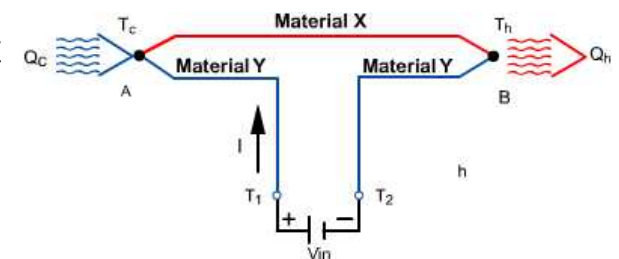
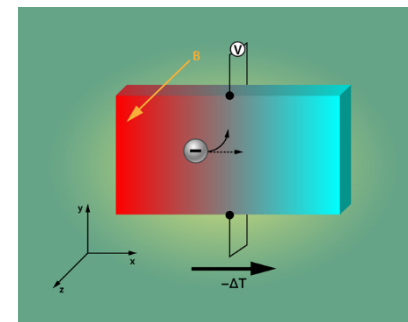
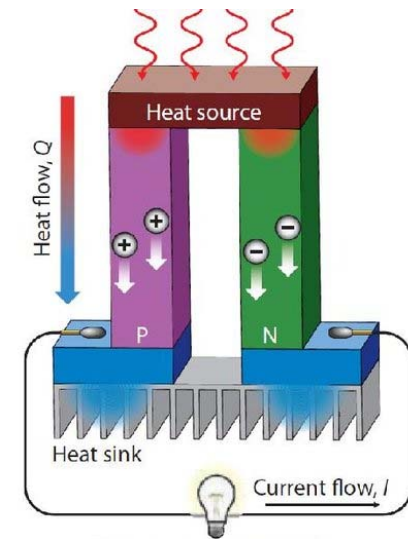
: When two ends of a conductor are at different temps, diffusion associated with this concentration gradient is counteracted by the buildup of an electric field.

- Nernst effect

: When a material is exposed to temp gradient and magnetic field, mobile charges move from the warm to the cold side and this motion induces a force coming from magnetic field, resulting in an electric field.

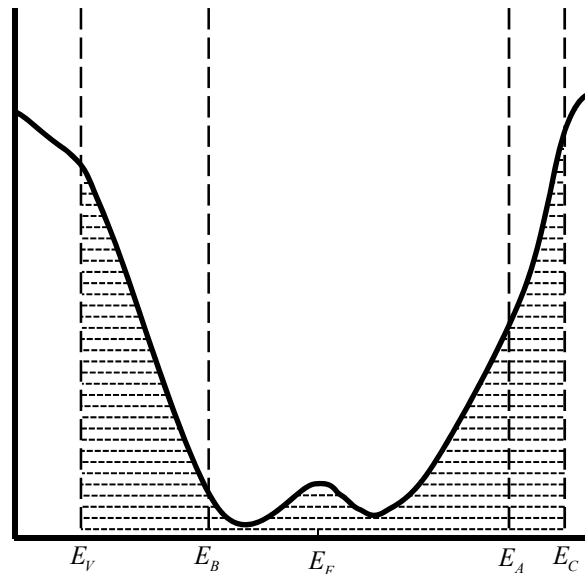
- Peltier effect

: If a voltage is applied to two terminals, electrical current will flow in the circuit. A slight cooling effect occurs at thermocouple junction *A* where heat is absorbed and a heating effect occurs at junction *B* where heat is expelled.

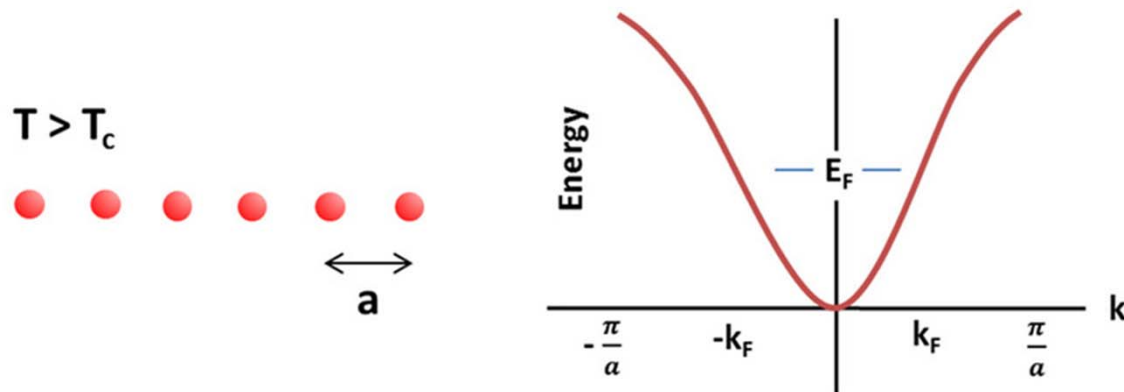


# Amorphous semiconductors

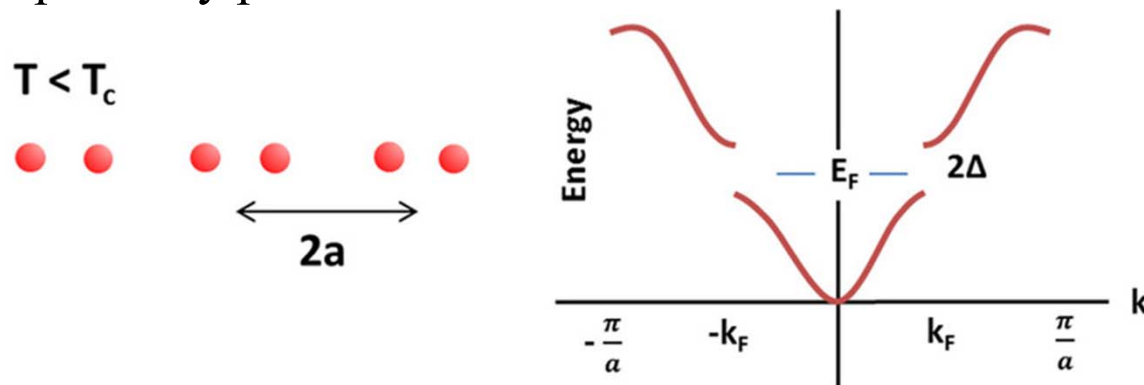
- Do not show long range order
- Larger potential for thin film transistors, large area uniform devices, solar cells,



# Charge density waves



In a metal, the electron density is highly uniform. The equilibrium positions of the ions form a perfectly periodic lattice.



In quasi-one- and two-dimensional metals, however, static modulations of the density are stable under certain conditions. The electron gas and the ion lattice spontaneously develop a periodic modulation when the temperature decreases below a critical transition temperature  $T_p$  ('p' for Peierls). The modulation of the electron density is called a charge density wave (CDW).

# Conducting polymers

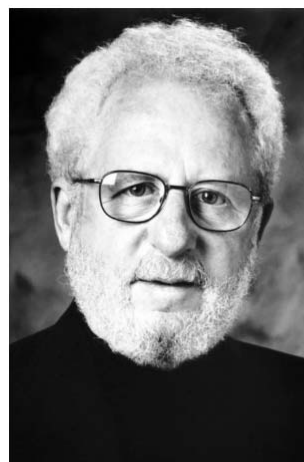
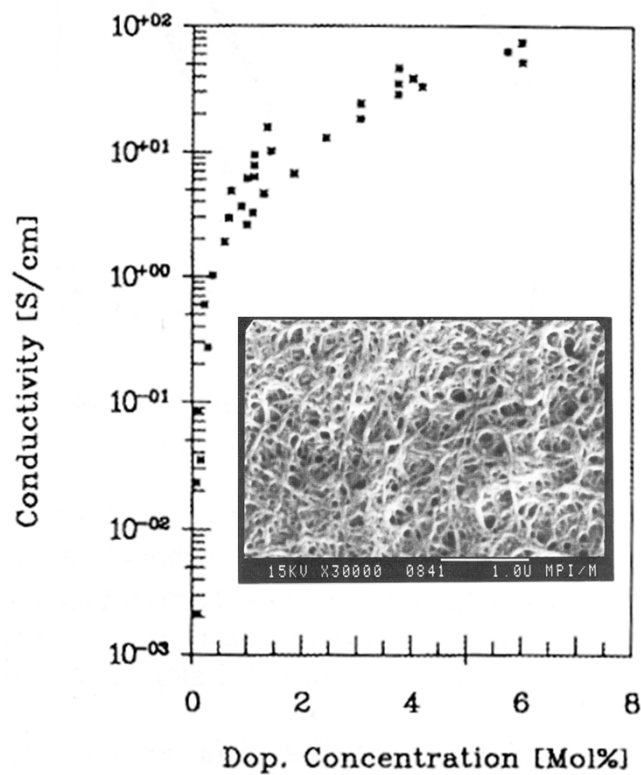
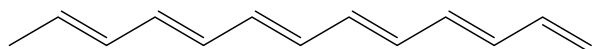


Photo from the Nobel Foundation archive.  
Alan J. Heeger



Photo from the Nobel Foundation archive.  
Alan G. MacDiarmid



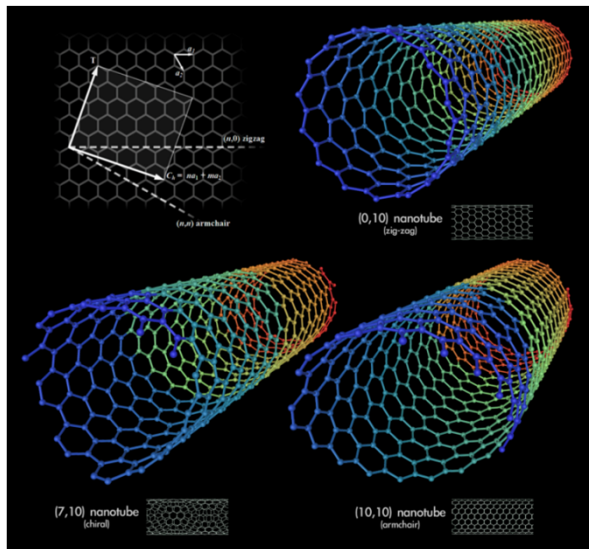
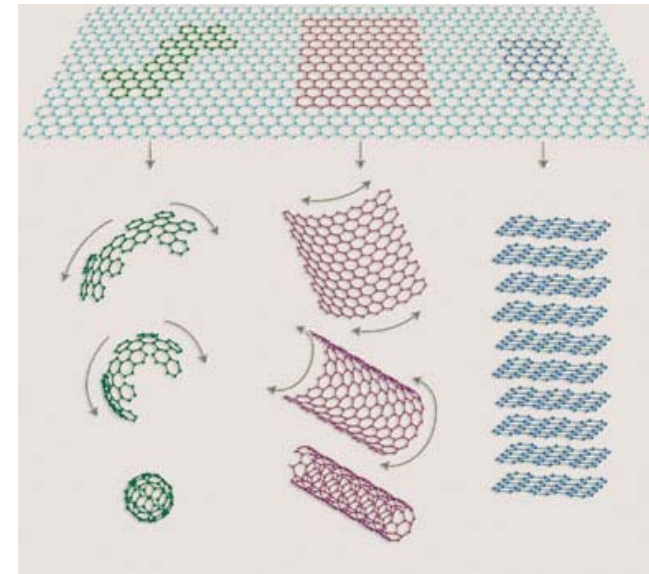
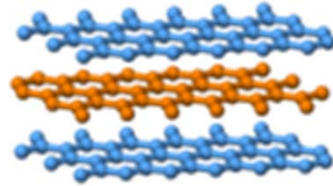
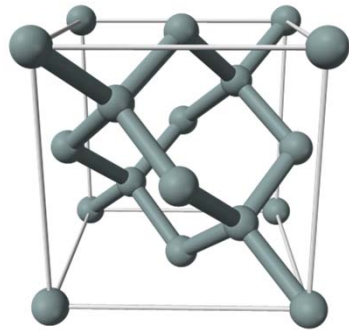
Photo from the Nobel Foundation archive.  
Hideki Shirakawa

The Nobel Prize in Chemistry 2000 was awarded for the discovery and development of conductive polymers.

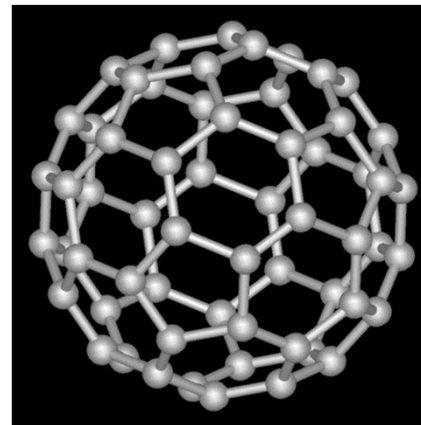
# Carbon-based Materials



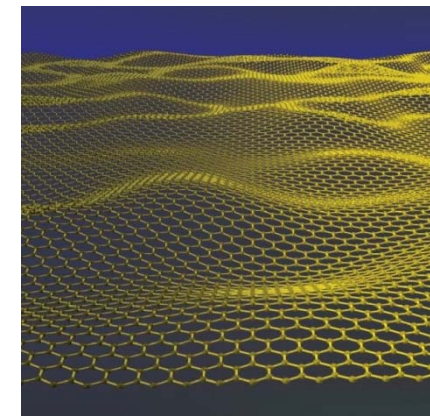
Diamond



CNT



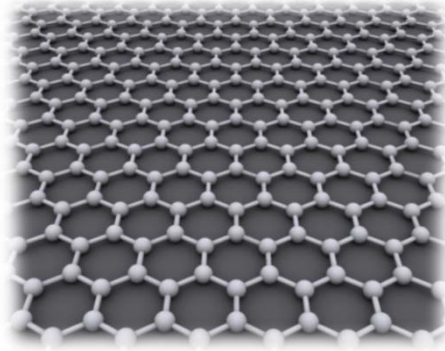
C<sub>60</sub>



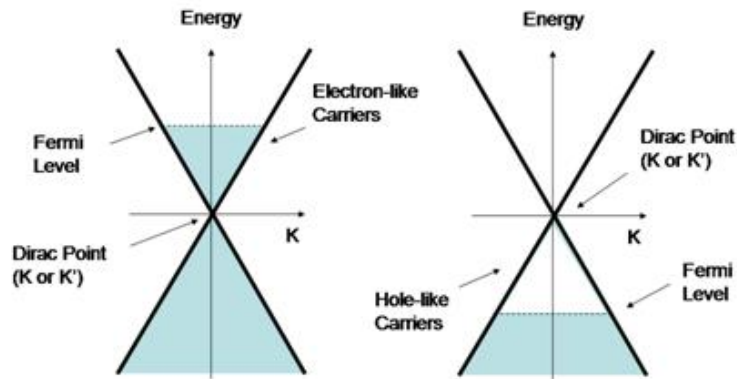
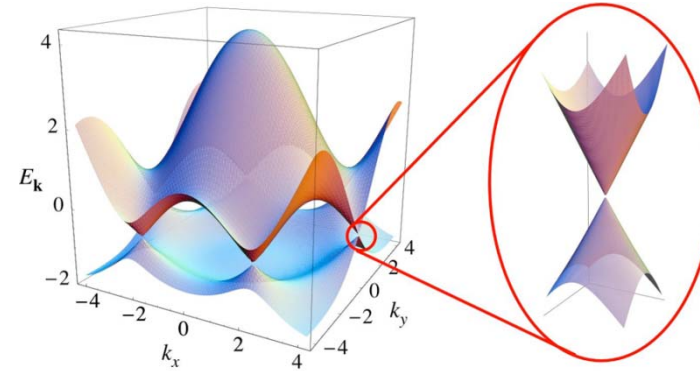
Graphene

# Carbon-based Materials: Graphene

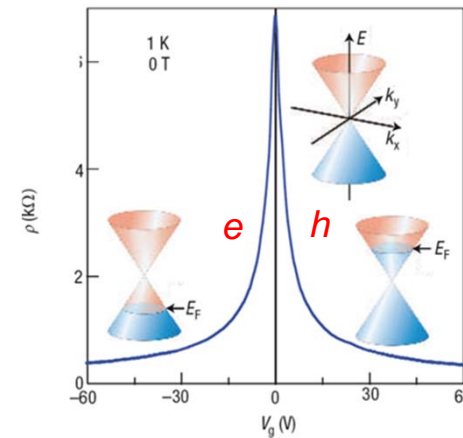
Crystal structure of graphene



Band structure of graphene



Linear dispersion



Ambipolar conductance