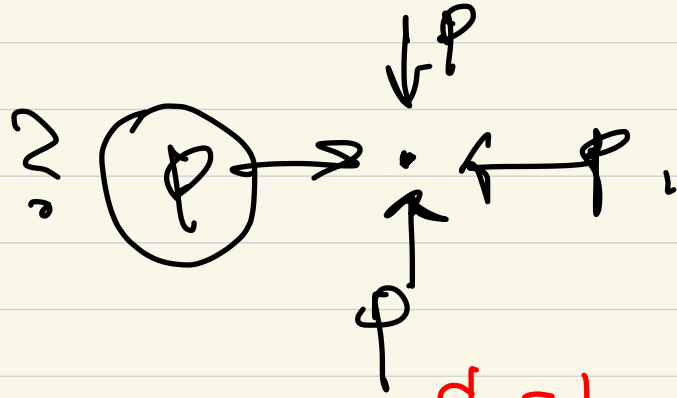


2.1. Pressure and pressure gradient.

at rest \rightarrow no shear stress, only normal stress.
(fluid pressure)

\downarrow
pressure at a point "in a static fluid" is
indep. of orientation.



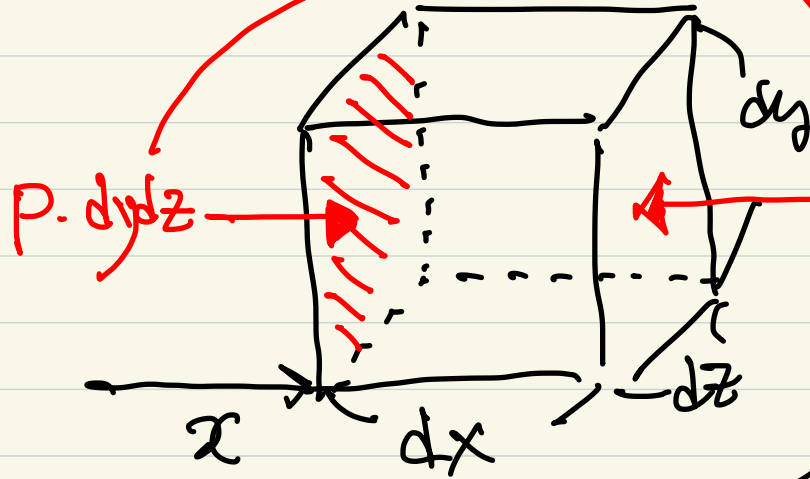
\hookrightarrow in a moving fluid, \rightarrow shear & normal stresses.

: pressure is the average of three
normal stresses on the element.

$$P = -\frac{1}{3} (\sigma_{xx} + \sigma_{yy} + \sigma_{zz}) \leftarrow \underline{\text{Ch. 4}}.$$

* pressure force on a fluid element.

Taylor series expansion.



$$(P + \frac{\partial P}{\partial x} dx) dydz.$$

$$dF_x = P \cdot dydz - (P + \frac{\partial P}{\partial x} dx) dydz.$$

$$= -\frac{\partial P}{\partial x} dx dy dz.$$

likewise, $\underline{dF}_{\text{press}} = \left(-\frac{\partial P}{\partial x} \hat{i} - \frac{\partial P}{\partial y} \hat{j} - \frac{\partial P}{\partial z} \hat{k} \right) dx dy dz.$

vector

$$= -\nabla P (\underbrace{dx dy dz}_{\text{volume}}).$$

$$\rightarrow \underline{df}_{\text{press}} = -\nabla P.$$

(net force per unit volume)

2.2 Equilibrium of a fluid element.

surface force : pressure, shear stress, ...

body force : gravity, electromagnetic force, ...

$$\rightarrow d\underline{F}_{\text{grav}} = \rho \underline{g} \cdot dx dy dz,$$

$$d\underline{f}_{\text{grav}} = \rho \underline{g}.$$

viscous force. (Ch. 4) : $\underline{f}_{\text{vis}} = \mu \nabla^2 \underline{V}$.

per unit volume.

$$\begin{aligned} \underline{\Sigma f} &= \rho \underline{a} = \underline{f}_{\text{press}} + \underline{f}_{\text{grav}} + \underline{f}_{\text{vis}} \\ &= -\nabla p + \rho \underline{g} + \mu \nabla^2 \underline{V}. \end{aligned}$$

special cases

✓ ① $\underline{V} = 0$ (or $\underline{a} = 0$; constant velocity).

✓ $\Rightarrow \nabla p = \rho \underline{g}$ (hydrostatic condition)

② rigid-body movement (rotation, translation)

ex) $\underline{u}_0 = r\Omega \Rightarrow \nabla^2 \underline{v} = 0$.

$\Rightarrow \nabla p = \rho (\underline{g} - \underline{a})$.

③ Irrotational motion. ($\nabla \times \underline{v} = 0$) $\Rightarrow \nabla^2 \underline{v} = 0$.

\hookrightarrow Bernoulli eq. (will come back to this later).

2.3. Hydrostatic pressure distributions.

($\underline{v} = 0$ or $\underline{a} = 0$, $\Rightarrow \nabla^2 \underline{v} = 0$).

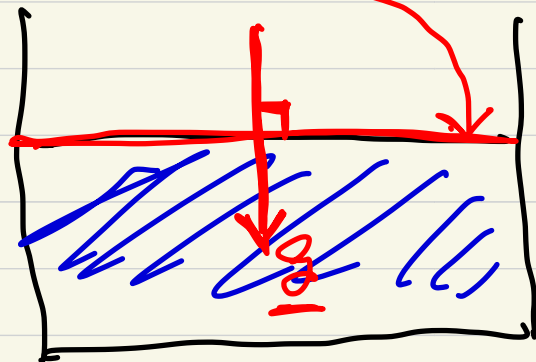
$\rho \underline{a} = -\nabla p + \rho \underline{g} + \mu \nabla^2 \underline{v} \Rightarrow \nabla p = \rho \underline{g}$ hydrostatic distribution.

magnitude and direction of the maximum spatial rate of increase in "p".

$\Rightarrow \nabla p$ is perpendicular to the surfaces of constant p .

\therefore A fluid in hydrostatic equilibrium will align its constant pressure surfaces everywhere normal to the local gravity vector.

(max. pressure increase will be in the direction of gravity, i.e., downward, for a liquid, its free surface is normal to local gravity.



• in a conventional coordinate system.

$$\underline{g} = -g \hat{k}, \quad \nabla p = \rho \underline{g} \rightarrow \frac{\partial p}{\partial x} = \frac{\partial p}{\partial y} = 0, \quad \frac{\partial p}{\partial z} = -\rho g$$

$$\therefore \frac{\partial p}{\partial z} = -\rho g \xrightarrow{\text{integrate.}}$$

$$p_2 - p_1 = - \int_1^2 \rho g dz.$$

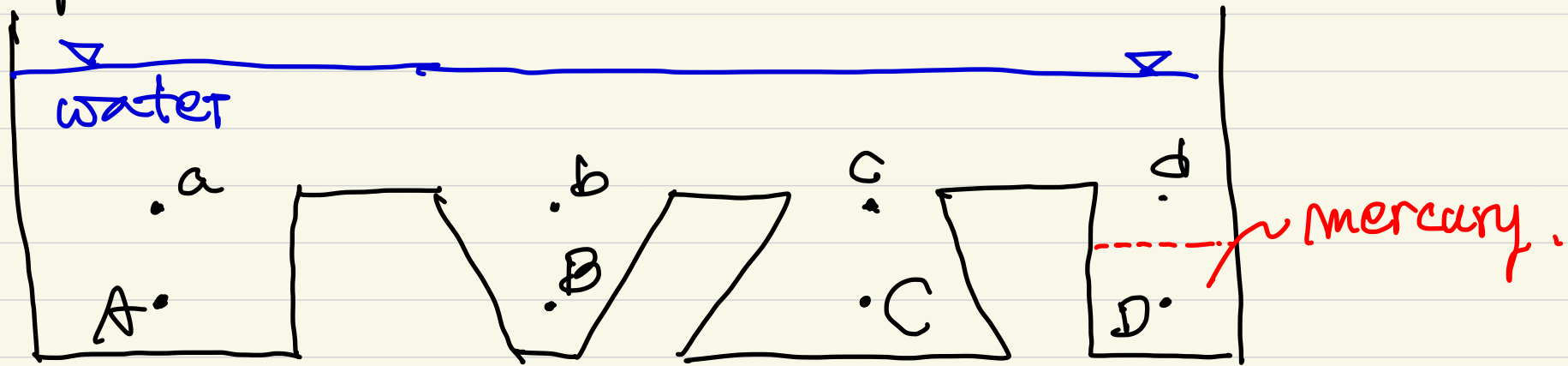
(Solution for hydrostatic problem.)

"Hydrostatic Condition" (Pascal's Law).

Pressure in a continuously distributed uniform static

fluid varies only w/ vertical distance, and is indep. of the shape of the container. The pressure is the same at all points on a given horizontal plane in the fluid. The pressure increases w/

depth in the fluid.



$$P_a = P_b = P_c = P_d, \quad P_A = P_B = P_C \neq P_D.$$

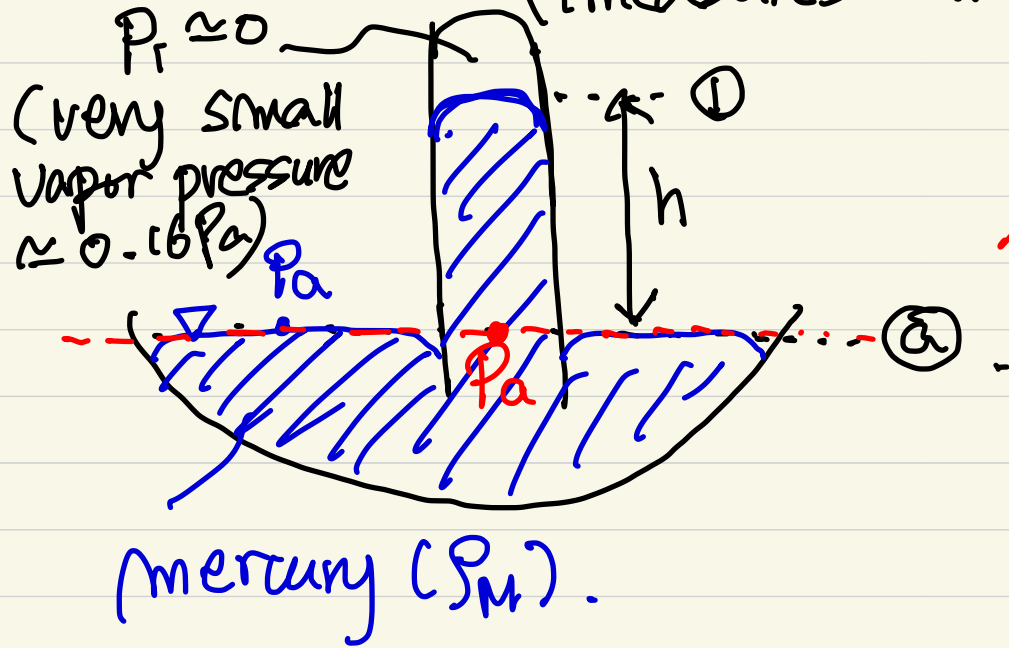
* Hydrostatic pressure in Liquids. (almost incompressible $\therefore \rho = \text{const}$)

$$P_2 - P_1 = - \int_1^2 \rho g dz = - \rho g (z_2 - z_1).$$

$$\text{or, } z_2 - z_1 = - \frac{P_2}{\rho g} + \frac{P_1}{\rho g} \quad \text{pressure head [L].}$$

* Mercury Barometer

(measures the atmospheric pressure).



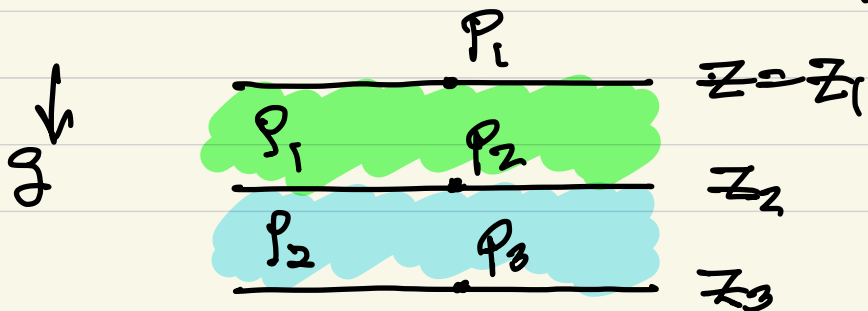
$$P_1 - P_a = -\rho_M g (z_1 - z_a)$$

$$= -\rho_M g h$$

$$\therefore h = \frac{P_a}{\rho_M g} = 761 \text{ mmHg}$$

2.4. Application to manometry.

• manometer: a device to measure pressure difference between two points.

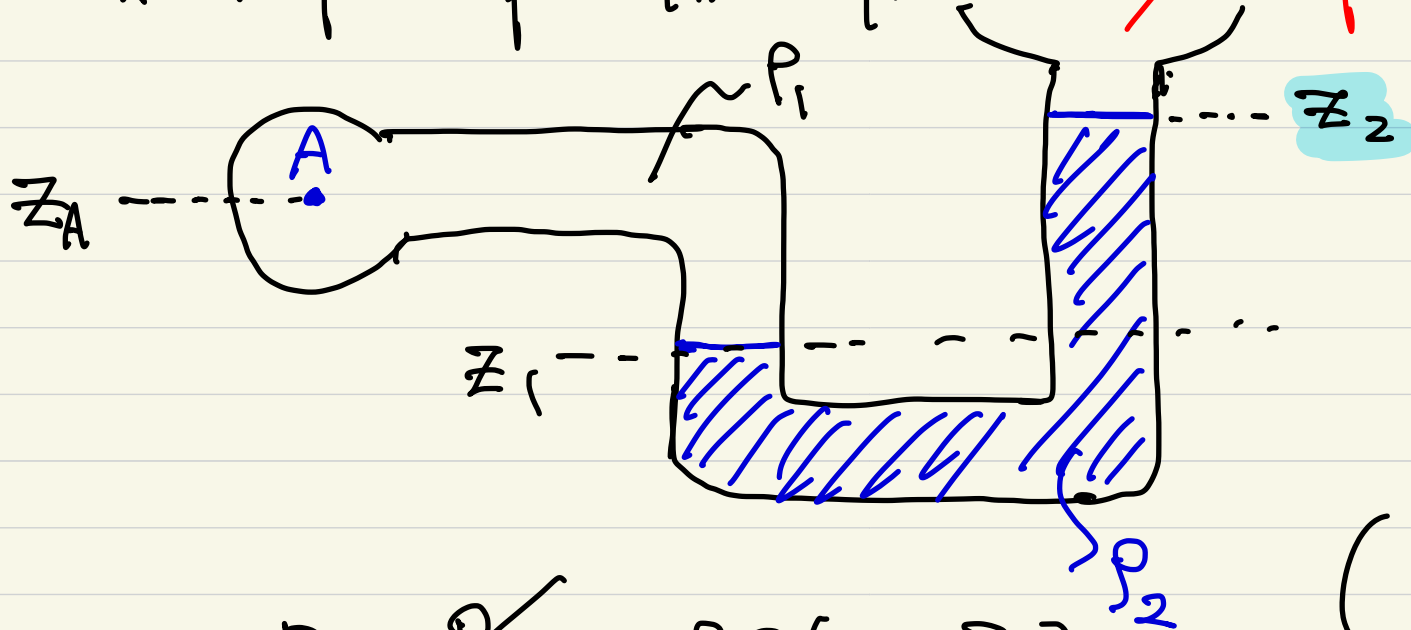


$$P_2 - P_1 = -\rho_1 g (z_2 - z_1)$$

$$P_3 - P_2 = -\rho_2 g (z_3 - z_2)$$

$$\Rightarrow P_2 - P_1 = -\rho_1 g (z_2 - z_1) - \rho_2 g (z_2 - z_1)$$

* simple open manometer. open (Pa)



$$P_A - P_1 = -\rho_1 g (z_A - z_1)$$

$$P_1 - P_a = -\rho_2 g (z_1 - z_2)$$

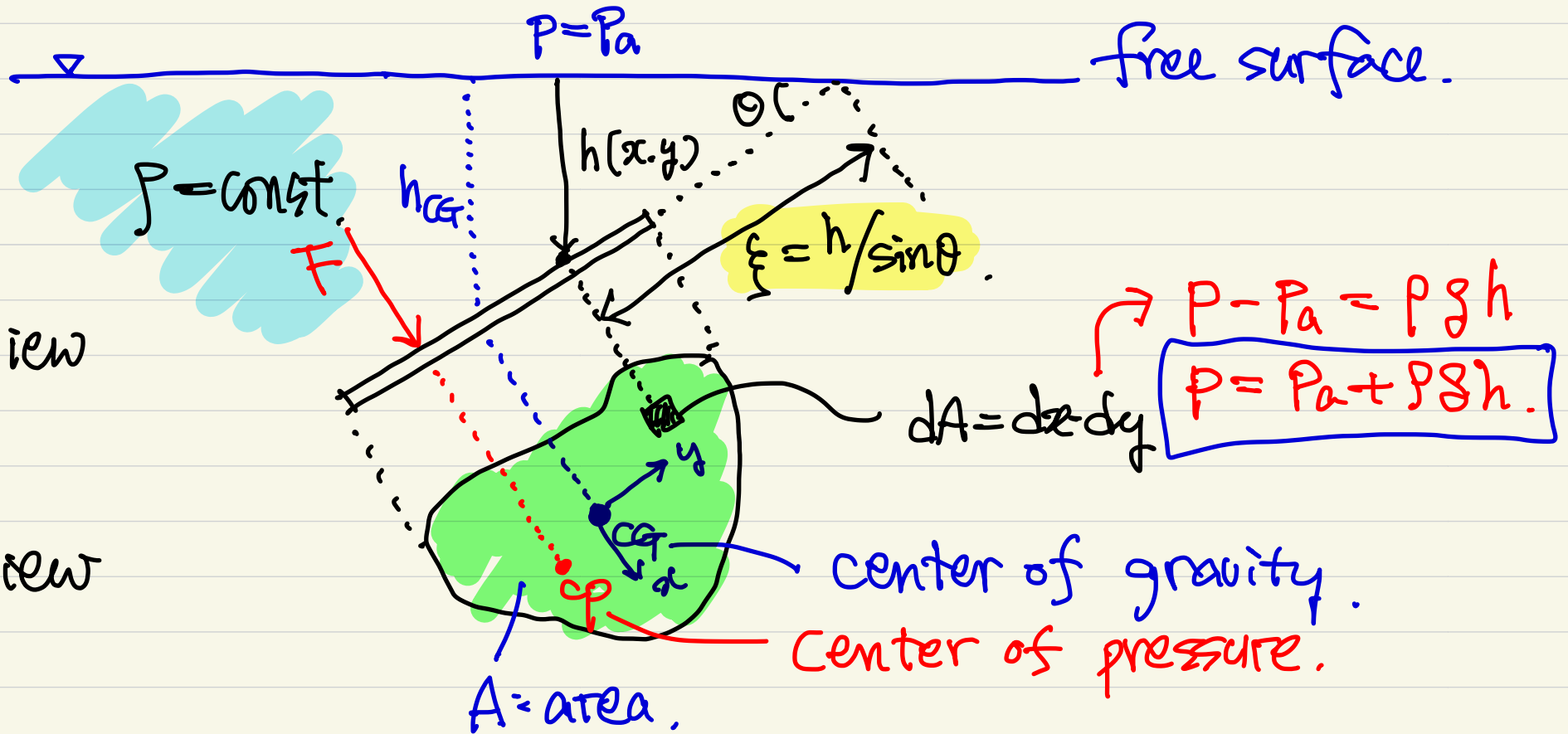
(heavier fluid (ρ_2)
 \Rightarrow better to make
 \uparrow ' $z_1 - z_2$ ' smaller)

$$P_A - P_a = -\rho_1 g (z_A - z_1) - \rho_2 g (z_1 - z_2)$$

gauge pressure.

2.5. Hydrostatic forces on plane surface.

$$(P_2 - P_1 = - \int_1^2 \rho g dz)$$



→ Total hydrostatic force on one side, F

$$: F = \int P \cdot dA = \int (P_a + \rho g h) dA$$

$$= P_a \cdot A + \rho g \int h \, dA.$$

$$(h = \xi \cdot \sin \theta)$$

$$\xi_{CG} \equiv \frac{1}{A} \int \xi \, dA \quad \text{by definition}$$

$$F = P_a \cdot A + \rho g \int \xi \cdot \sin \theta \cdot dA$$

$$= P_a \cdot A + \rho g \sin \theta \int \xi \, dA.$$

$$= P_a \cdot A + \rho g \sin \theta \cdot A \cdot \xi_{CG} = P_a \cdot A + \rho g A \cdot h_{CG}.$$

$$= (P_a + \rho g h_{CG}) A = P_{CG} \cdot A.$$

The force on one side of any plane submerged in a uniform fluid equals the pressure at the plate centroid (P_{CG}) times the plate area (A), indep. of

the plate shape or the angle (θ) at which the plate is slanted.

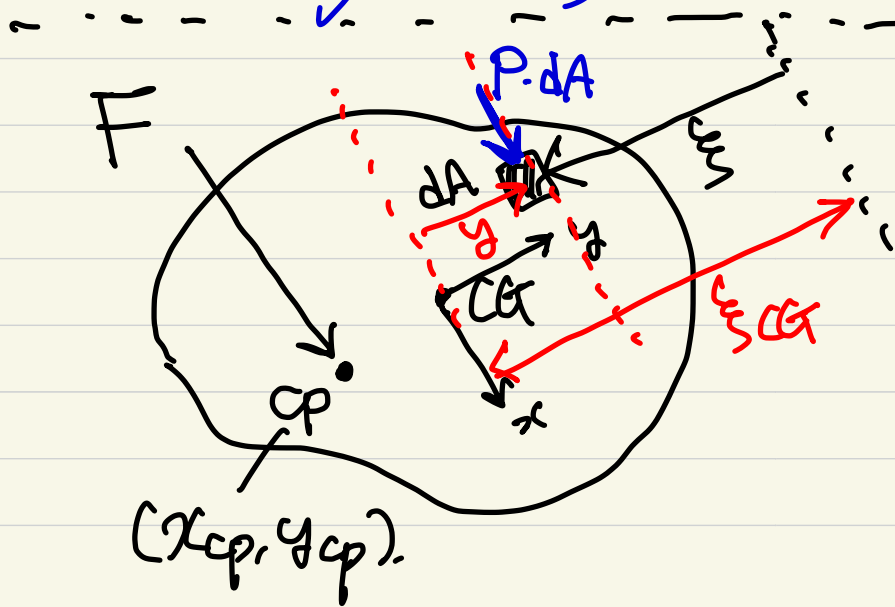
Then, where is the force acting? -

F acts not through the CG.

↓ (if it is, bending moment is zero!)

should balance the bending-moment via stress (pressure)

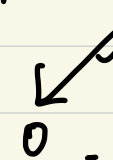
by acting at the high-pressure side.

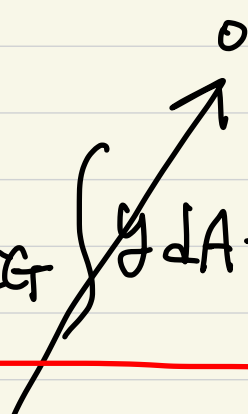


$$F \cdot y_{cp} = \int y \cdot P \, dA.$$

$$= \int y (P_a + P g h) \, dA.$$

$$= \int y (P_a + P g \xi_{CG}) \, dA.$$

$$= P_a \int y dA + \rho g \sin \theta \int y \xi dA.$$


$$= \rho g \sin \theta \int y (\xi_{CG} - y) dA. = \rho g \sin \theta \left[\xi_{CG} \int y dA - \int y^2 dA \right]$$


($I_{xx} = \int y^2 dA$, area moment of inertia)

$$\therefore \underline{F} \cdot y_{cp} = - \rho g \sin \theta \cdot I_{xx}.$$

$$= \rho g \cdot A.$$

$$\Rightarrow y_{cp} = - \rho g \sin \theta \cdot \frac{I_{xx}}{\rho g \cdot A}.$$

cp is below cg

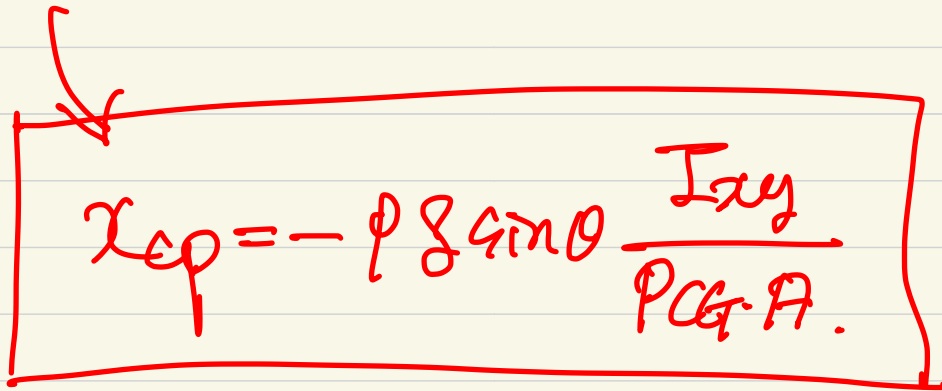
depends on θ and geo.

as the plate moves deeper, $\rho_{CG} \uparrow$.

$$\rightarrow y_{cp} \rightarrow 0. \rightarrow CP \rightarrow CG.$$

• Similarly (moment balance along x-dir).

$$F \cdot x_{cp} = \int x \cdot p dA = -\rho g \sin \theta \int xy dA.$$

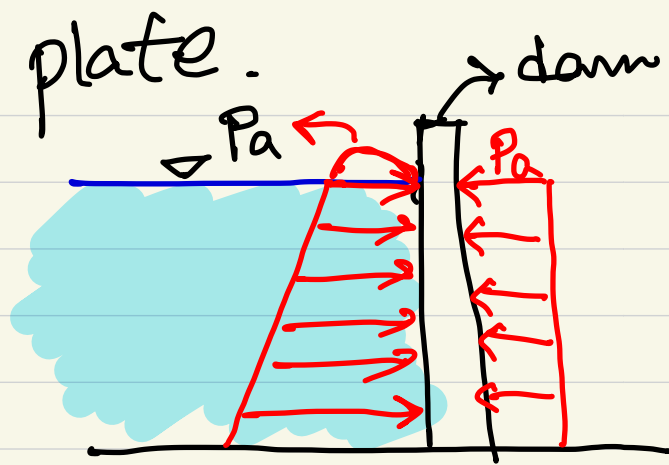

$$x_{cp} = -\rho g \sin \theta \frac{I_{xy}}{\rho g \sin \theta \cdot A}.$$

= I_{xy} (product of inertia).

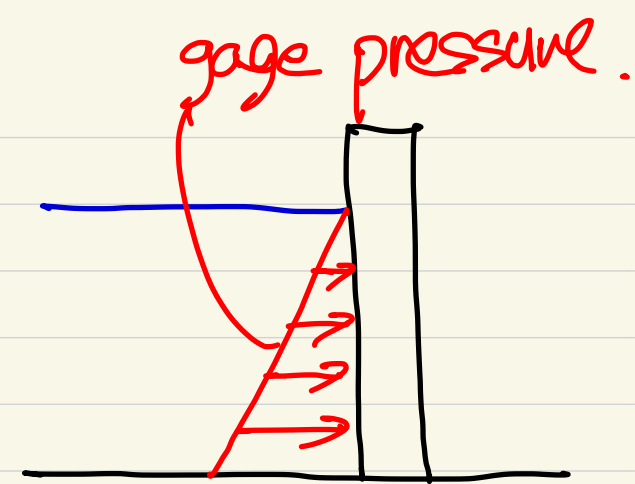
(for symmetric shape, $I_{xy} = 0$, $x_{cp} = 0$)

* Gage pressure formulae.

: P_a (ambient pressure) is neglected in most cases, because it acts on both sides of the



=



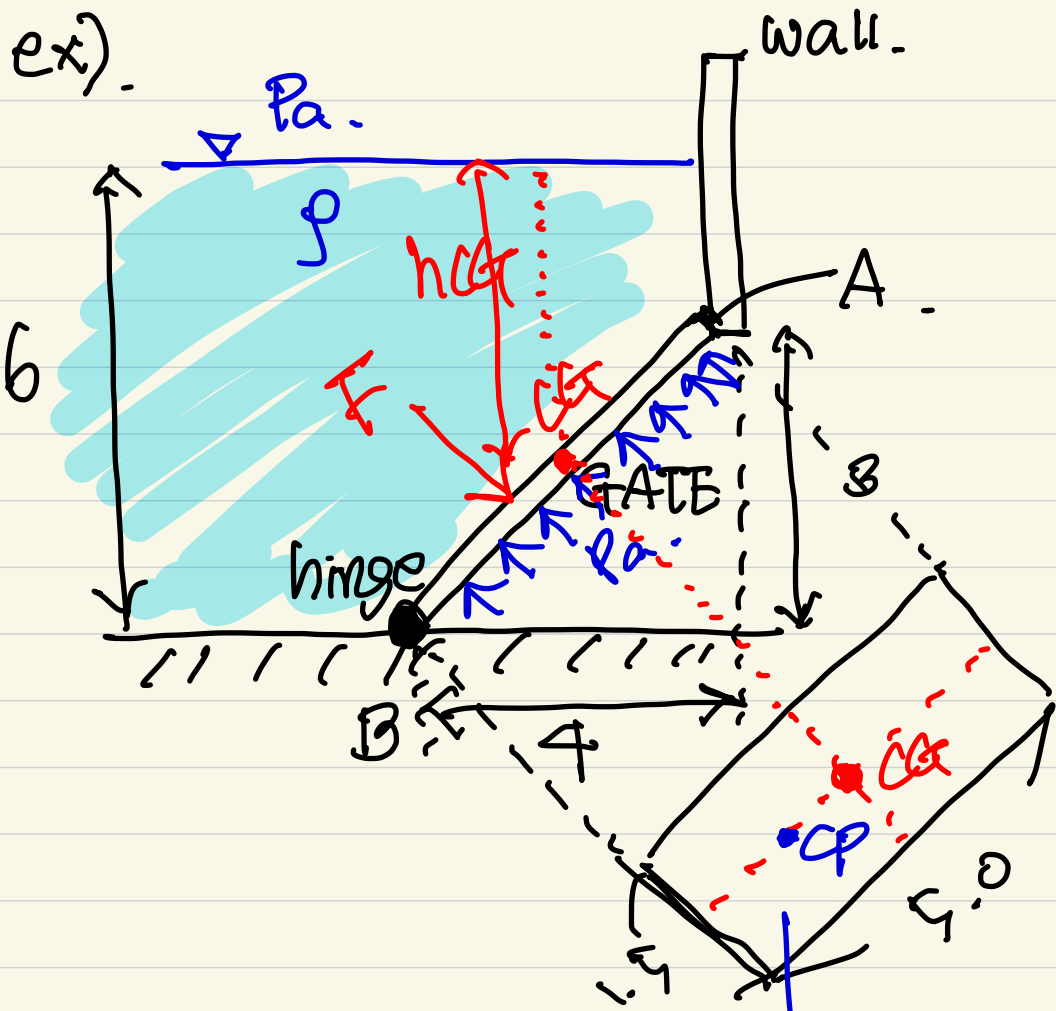
$$P_{CG} - P_a = \rho g h_{CG} \Rightarrow P_{CG} = \rho g h_{CG}$$

↓

$$F = \rho g h_{CG} \cdot A$$

$$y_{cp} = - \frac{I_{xx} \cdot \sin \theta}{h_{CG} \cdot A}$$

$$x_{cp} = - \frac{I_{xy} \cdot \sin \theta}{h_{CG} \cdot A}$$

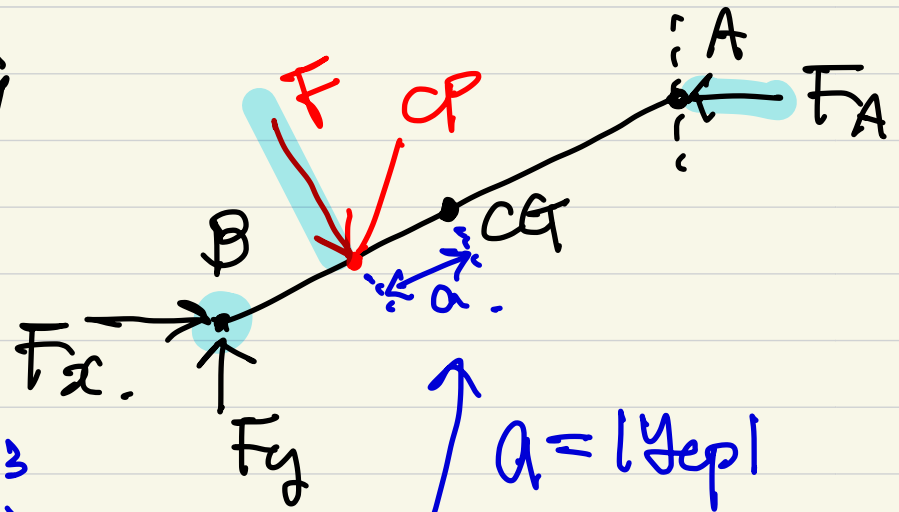


$$(a) F_G = ? = \rho \cdot g \cdot h_{CG} \cdot A$$

$$= \rho \cdot g \cdot h_{CG} \cdot A$$

↑ ↑ ↑
1.5 × 5

(b) Forces at A and B?



$$I_{xx} = \frac{1}{12} b L^3$$

$$I_{yy} = 0$$

$$a = |y_{CG}|$$

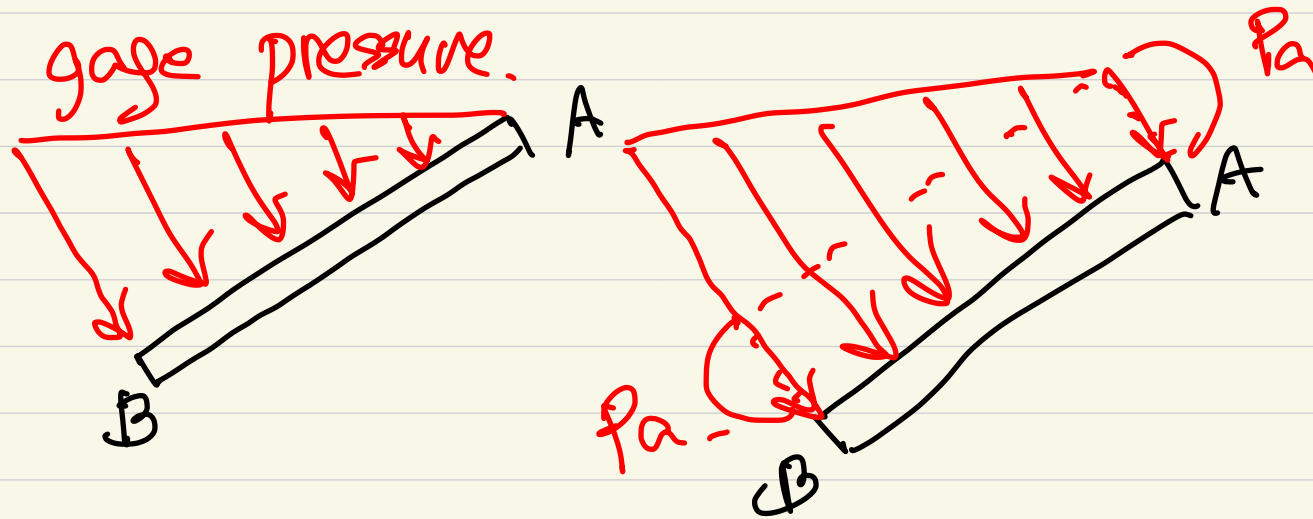
$$= \left| \frac{I_{xx} \cdot \sin \theta}{h_{CG} \cdot A} \right|$$

$$\Sigma F \rightarrow = F_x - F_A + F \sin \theta = 0$$

$$\Sigma F \uparrow = F_y - F \cos \theta = 0.$$

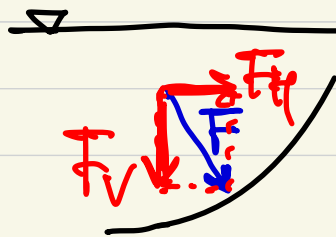
$$\textcircled{+} \Sigma M_B = F \cdot (2.5 - a) - F_A \cdot 3 = 0.$$

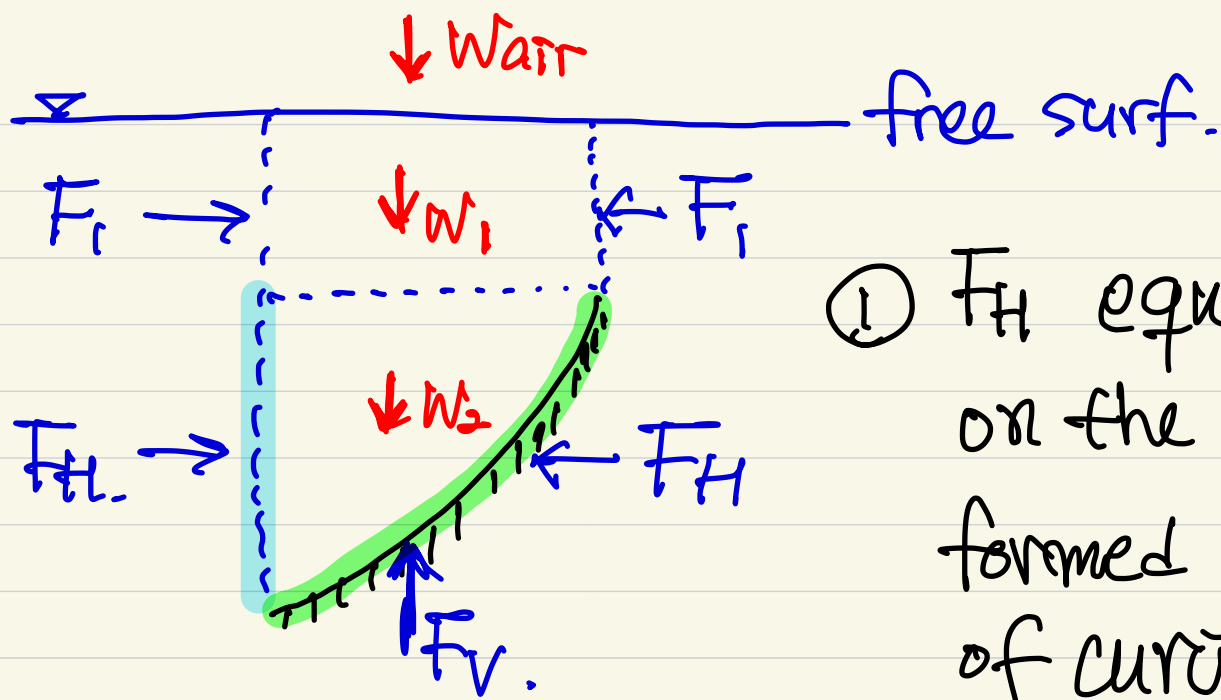
} $\rightarrow F_x, F_y, F_A.$



26. Hydrostatic forces on curved surfaces.

\rightarrow try to separate F_H and F_V .





① F_H equals the force on the plane area formed by the projection of curved surface onto a vertical plane.

② $F_v = W_{air} + W_1 + W_2$.