

$$\sum \underline{F} = \rho \frac{DV}{Dt} dx dy dz \quad (1)$$

(2) + (3)

body force :

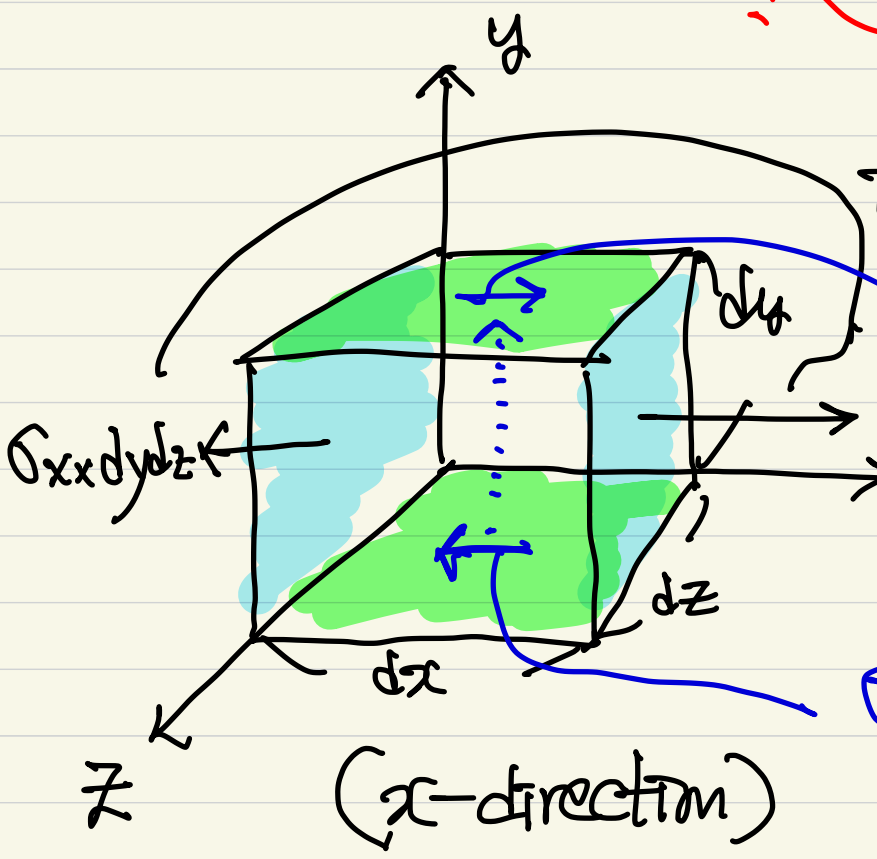
$$dF_{grav} = \rho g dx dy dz \quad (2)$$

surface (4)

$$\sigma_{ij} = -P \delta_{ij} + \tau_{ij}$$

VISCOUS STRESS

hydrostatic pressure



Taylor series expansion

$$\sigma_{xx} dy dz + \frac{\partial}{\partial x} (\sigma_{xx} \cdot dy dz) \cdot dx$$

$$\sigma_{yx} dx dz + \frac{\partial}{\partial y} (\sigma_{yx} dx dz) dy$$

$$\sigma_{yx} \cdot dx \cdot dz$$

gradient ( $\nabla P$  or  $\nabla \tau_{ij}$ ) is important!

$$\therefore dF_{x,\text{surf}} = \left[ \frac{\partial}{\partial x} \sigma_{xx} + \frac{\partial}{\partial y} \sigma_{yx} + \frac{\partial}{\partial z} \sigma_{zx} \right] dx dy dz$$

$$= \left[ -\frac{\partial P}{\partial x} + \frac{\partial \tau_{xx}}{\partial x} + \frac{\partial \tau_{yx}}{\partial y} + \frac{\partial \tau_{zx}}{\partial z} \right] dx dy dz$$

$$dF_{y,\text{surf}} = \left[ -\frac{\partial P}{\partial y} + \frac{\partial \tau_{xy}}{\partial x} + \frac{\partial \tau_{yy}}{\partial y} + \frac{\partial \tau_{zy}}{\partial z} \right] dx dy dz$$

$$+ \left[ dF_{z,\text{surf}} = \left[ -\frac{\partial P}{\partial z} + \frac{\partial \tau_{xz}}{\partial x} + \frac{\partial \tau_{yz}}{\partial y} + \frac{\partial \tau_{zz}}{\partial z} \right] dx dy dz \right]$$

$$\Rightarrow d\mathbf{F}_{\text{surf}} = \left( -\nabla P + \nabla \cdot \underline{\underline{\tau}} \right) dx dy dz \quad \text{(vector)} \quad \textcircled{3}$$

$$\text{or } dF_{i,\text{surf}} = \left( -\frac{\partial P}{\partial x_i} + \frac{\partial \tau_{ij}}{\partial x_j} \right) dx dy dz \quad \text{(tensor)}$$

$$(i,j = x,y,z)$$

① + ② + ③ : momentum conservation eq

$$\rho \underline{g} - \nabla P + \nabla \cdot \hat{\underline{\tau}} = \rho \frac{D\underline{v}}{Dt}$$

$\uparrow$  gravity       $\uparrow$  pressure force       $\downarrow$  viscous force.

$$= \rho \left( \frac{\partial \underline{v}}{\partial t} + (\underline{v} \cdot \nabla) \underline{v} \right)$$

$$= \rho \left( \frac{\partial \underline{v}}{\partial t} + u \frac{\partial \underline{v}}{\partial x} + v \frac{\partial \underline{v}}{\partial y} + w \frac{\partial \underline{v}}{\partial z} \right)$$

\* Inviscid flow.

$$\rho \frac{D\underline{v}}{Dt} = \rho \underline{g} - \nabla P \quad : \text{Euler equation.}$$

\* Newtonian Fluid.

$$\hat{\underline{\tau}} = \mu \frac{du}{dy}$$

generalization  $\rightarrow$

$$\hat{\tau}_{ij} = \mu \left( \frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right)$$

(i, j = x, y, z)

$$\hat{T}_{xx} = 2\mu \frac{\partial u}{\partial x}, \quad \hat{T}_{yy} = 2\mu \frac{\partial v}{\partial y}, \quad \hat{T}_{zz} = 2\mu \frac{\partial w}{\partial z}$$

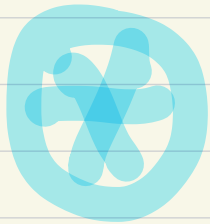
$$\hat{T}_{xy} = \hat{T}_{yx} = \mu \left( \frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} \right)$$

$$\hat{T}_{xz} = \hat{T}_{zx} = \mu \left( \frac{\partial w}{\partial x} + \frac{\partial u}{\partial z} \right)$$

$$\hat{T}_{yz} = \hat{T}_{zy} = \mu \left( \frac{\partial v}{\partial z} + \frac{\partial w}{\partial y} \right)$$

→ Continuum conservation eq.

→ Navier-Stokes eq.



$$\frac{\partial \hat{T}_{xx}}{\partial x} + \frac{\partial \hat{T}_{yx}}{\partial y} + \frac{\partial \hat{T}_{zx}}{\partial z}$$

$$= \frac{\partial}{\partial x} \left( \mu \frac{\partial u}{\partial x} + \mu \frac{\partial u}{\partial x} \right) + \frac{\partial}{\partial y} \left( \mu \frac{\partial u}{\partial y} + \mu \frac{\partial v}{\partial x} \right) + \frac{\partial}{\partial z} \left( \mu \frac{\partial w}{\partial x} + \mu \frac{\partial u}{\partial z} \right)$$

$$= \mu \frac{\partial^2 u}{\partial x^2} + \mu \frac{\partial}{\partial x} \left( \frac{\partial u}{\partial x} \right) + \mu \frac{\partial^2 u}{\partial y^2} + \mu \frac{\partial}{\partial x} \left( \frac{\partial v}{\partial y} \right) + \mu \frac{\partial^2 u}{\partial z^2} + \mu \frac{\partial}{\partial x} \left( \frac{\partial w}{\partial z} \right)$$

$$= \mu \left( \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} + \frac{\partial^2 u}{\partial z^2} \right) + \mu \frac{\partial}{\partial x} \left( \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} \right)$$

Newtonian,  
incompressible  $\nabla \cdot \underline{v} = 0$

incompressible flow  
continuity eq.

$$\rho \frac{D\underline{v}}{Dt} = \rho \underline{g} - \nabla P + \mu \nabla^2 \underline{v} \quad : \text{N-S eq.}$$

$$\rho \frac{\partial u}{\partial t} + \rho u \frac{\partial u}{\partial x} + \rho v \frac{\partial u}{\partial y} + \rho w \frac{\partial u}{\partial z} = \rho g_x - \frac{\partial P}{\partial x} + \mu \left( \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} + \frac{\partial^2 u}{\partial z^2} \right)$$

$$\rho \frac{\partial v}{\partial t} + \rho u \frac{\partial v}{\partial x} + \rho v \frac{\partial v}{\partial y} + \rho w \frac{\partial v}{\partial z} = \rho g_y - \frac{\partial P}{\partial y} + \mu \left( \frac{\partial^2 v}{\partial x^2} + \frac{\partial^2 v}{\partial y^2} + \frac{\partial^2 v}{\partial z^2} \right)$$

$$\rho \frac{\partial w}{\partial t} + \rho u \frac{\partial w}{\partial x} + \rho v \frac{\partial w}{\partial y} + \rho w \frac{\partial w}{\partial z} = \rho g_z - \frac{\partial P}{\partial z} + \mu \left( \frac{\partial^2 w}{\partial x^2} + \frac{\partial^2 w}{\partial y^2} + \frac{\partial^2 w}{\partial z^2} \right)$$

unsteady, non-linear. 2<sup>nd</sup>-order, coupled PDE.

w/ given  $\rho, \mu$ , 4 unknowns ( $u, v, w, P$ )  
(continuity + ③ N-S)

+ State equation  
 $(p = \rho RT)$ .

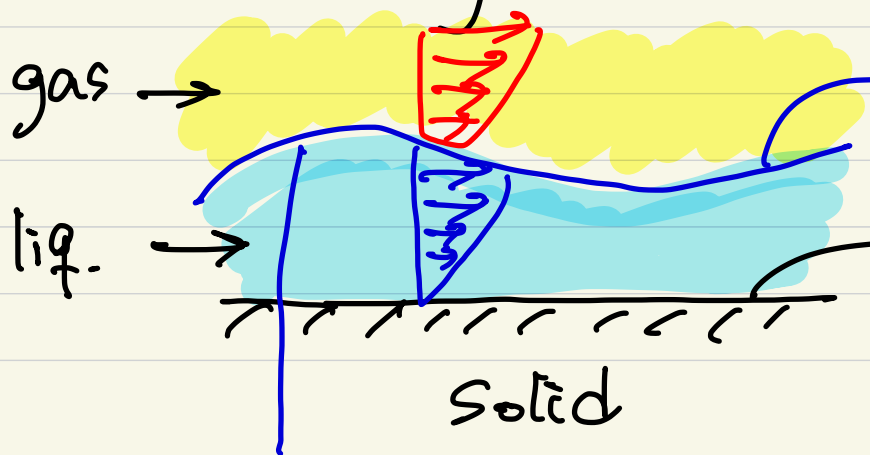
+ Energy eq (for T)

### 4.5 Differential eq. of Energy.

$$\rho \frac{D\hat{u}}{Dt} + \rho(\nabla \cdot \underline{V}) = \nabla \cdot (k \nabla T) + \Phi$$

↙ internal energy ( $d\hat{u} \approx C_v dT$ )
 ↘ thermal conductivity.
 ↙ viscous dissipation.

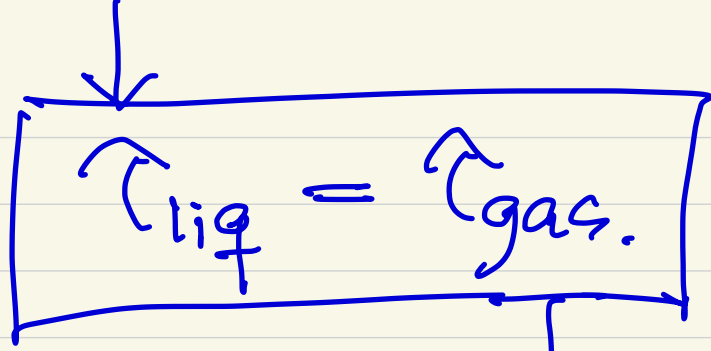
### 4.6. Boundary conditions for basic eqns.



free-surface.

no-slip condition :  $\underline{V}_{fluid} = \underline{V}_{solid}$

no-jump :  $T_{fluid} = T_{solid}$



$\Rightarrow$  simplified free-surf. BC  
(valid for river, sea, ...)

$\hat{c}_{liq} = 0$   
( $P_{liq} = P_{gas}$ )

$$dF_{i.surf} = \left( -\frac{\partial p}{\partial x_i} + \frac{\partial \tau_{ji}}{\partial x_j} \right) dx dy dz.$$

↓ dummy index. (i, j = x, y, z)

( $\rho = \text{constant}$ )  
↑

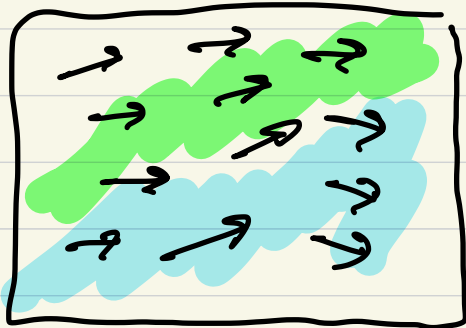
$$\frac{\partial \hat{\tau}_{xi}}{\partial x} + \frac{\partial \hat{\tau}_{yi}}{\partial y} + \frac{\partial \hat{\tau}_{zi}}{\partial z}$$

\* **incompressible** flow w/ constant properties. **Newtonian Fluid.**

$$\left. \begin{array}{l} \nabla \cdot \underline{V} = 0 \quad (\text{continuity eq.}) \end{array} \right\}$$

$$\rho \frac{\partial \underline{V}}{\partial t} + \rho (\underline{V} \cdot \nabla) \underline{V} = \rho \underline{g} - \nabla p + \mu \nabla^2 \underline{V}$$

(N-S eq.)





# \* Inviscid flow approximation.

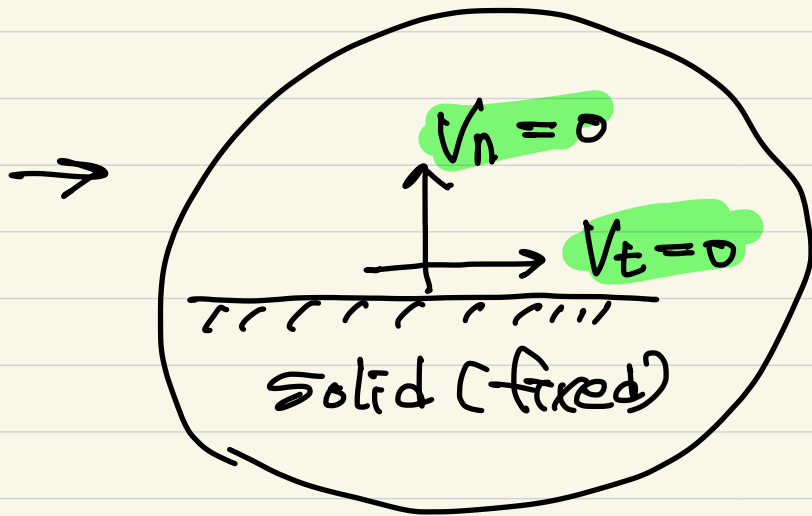
(비점성 유체).  $\mu \rightarrow 0$ .

$$\rho \frac{D\underline{V}}{Dt} = \rho \underline{g} - \nabla P \quad : \text{Euler equation.}$$

$$\rho \frac{\partial \underline{V}}{\partial t} + \rho (\underline{V} \cdot \nabla) \underline{V}$$

↓  
1<sup>st</sup>-order PDE

⇒ one BC for  $\underline{V}$ .

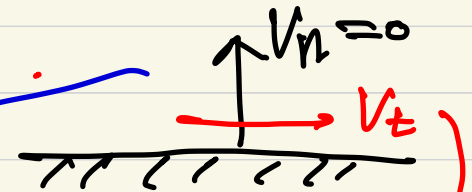


N-S eq.

no-slip condition. (w/ viscosity)

drop no-slip condition  
along the tangential  
direction.

do not allow the flow into (out of)

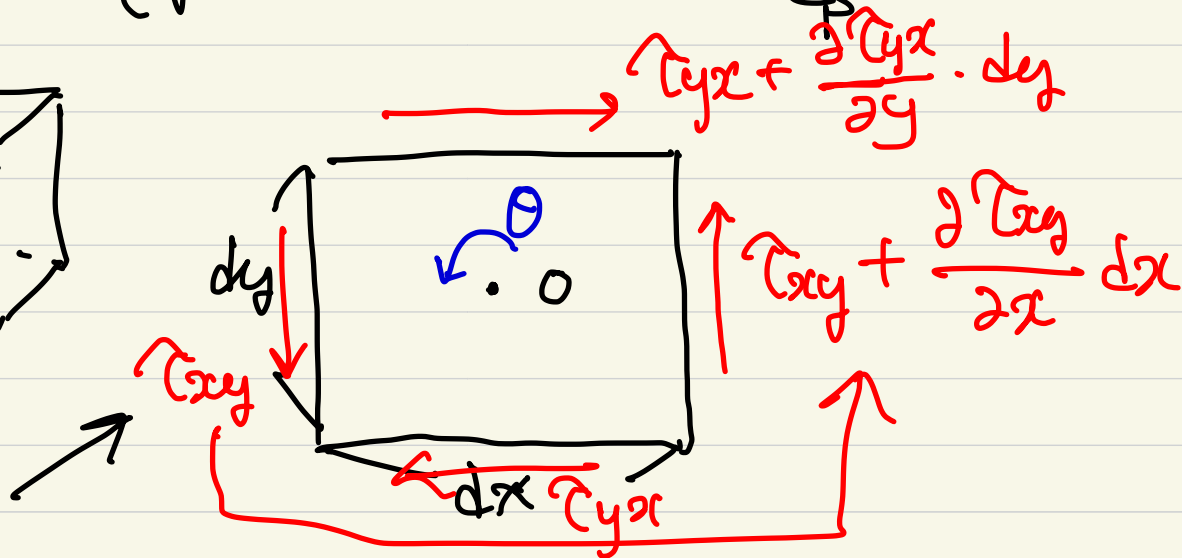
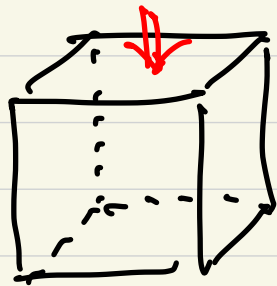


the wall.

Not BC, part of the solution. (ch. 8)

4.4. Differential eq. of angular momentum.

$$\Sigma M_o = \frac{d}{dt} \int_{CV} (\underline{r} \times \underline{v}) \rho dV + \int_{CS} (\underline{r} \times \underline{v}) \rho (\underline{v} \cdot \underline{n}) dA.$$



Taylor series exp.

$$\Sigma M_o = \left( \hat{\tau}_{xy} + \frac{\partial \hat{\tau}_{xy}}{\partial x} dx \right) dy \cdot dz \cdot \frac{1}{2} dx + \hat{\tau}_{xy} \cdot dy dz \cdot \frac{1}{2} dx$$

$$- \left( \hat{\tau}_{yx} + \frac{\partial \hat{\tau}_{yx}}{\partial y} dy \right) dx \cdot dz \cdot \frac{1}{2} dy - \hat{\tau}_{yx} \cdot dx dz \cdot \frac{1}{2} dy$$

$I_x \Rightarrow \frac{1}{12} m (w^2 + d^2)$

$\frac{1}{12} \rho dx dy dz (dx^2 + dy^2) \cdot \frac{d^2 \theta}{dt^2}$

$$\left( \hat{\tau}_{xy} - \hat{\tau}_{yx} \right) dx dy dz + \frac{1}{2} \frac{\partial \hat{\tau}_{xy}}{\partial x} dx^2 dy dz$$

$$- \frac{1}{2} \frac{\partial \hat{\tau}_{yx}}{\partial y} dx dy^2 dz$$

as  $dx, dy, dz \rightarrow 0$ .

if we neglect higher order terms.

$$\Rightarrow \boxed{\hat{\tau}_{xy} = \hat{\tau}_{yx}} \quad \left( \hat{\tau}_{xz} = \hat{\tau}_{zx}, \hat{\tau}_{yz} = \hat{\tau}_{zy}, \text{as well} \right)$$

$\hookrightarrow \vec{\zeta}$  is symmetric.

4.7 stream function (유선 함수).  
(scalar function)

for 2D incompressible flow ( $\rho = \text{const.}$ )  
continuity (mass conserv.)

$$\therefore \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0 \quad (\nabla \cdot \underline{v} = 0)$$

Let  $u \equiv \frac{\partial \psi}{\partial y}$  and  $v \equiv -\frac{\partial \psi}{\partial x}$

then,  $\frac{\partial}{\partial x} \left( \frac{\partial \psi}{\partial y} \right) + \frac{\partial}{\partial y} \left( -\frac{\partial \psi}{\partial x} \right) = 0 \rightarrow$  scalar function  
stream function.

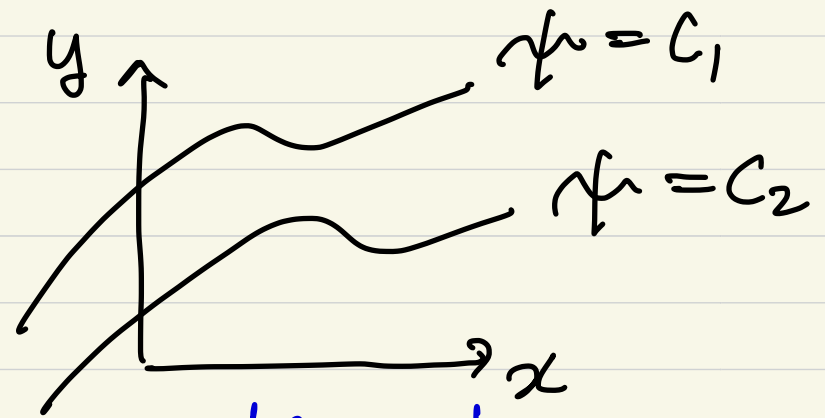
↳ automatically satisfies the continuity!

\* Geometric interpretation of  $\psi$

① Lines of constant  $\psi$

$$\psi = \psi(x, y) = C$$

Streamlines



Streamlines in 2D flow:  $\frac{dx}{u} = \frac{dy}{v}$

$$d\psi = \frac{\partial \psi}{\partial x} dx + \frac{\partial \psi}{\partial y} dy = 0$$

$\frac{-v}{u}$

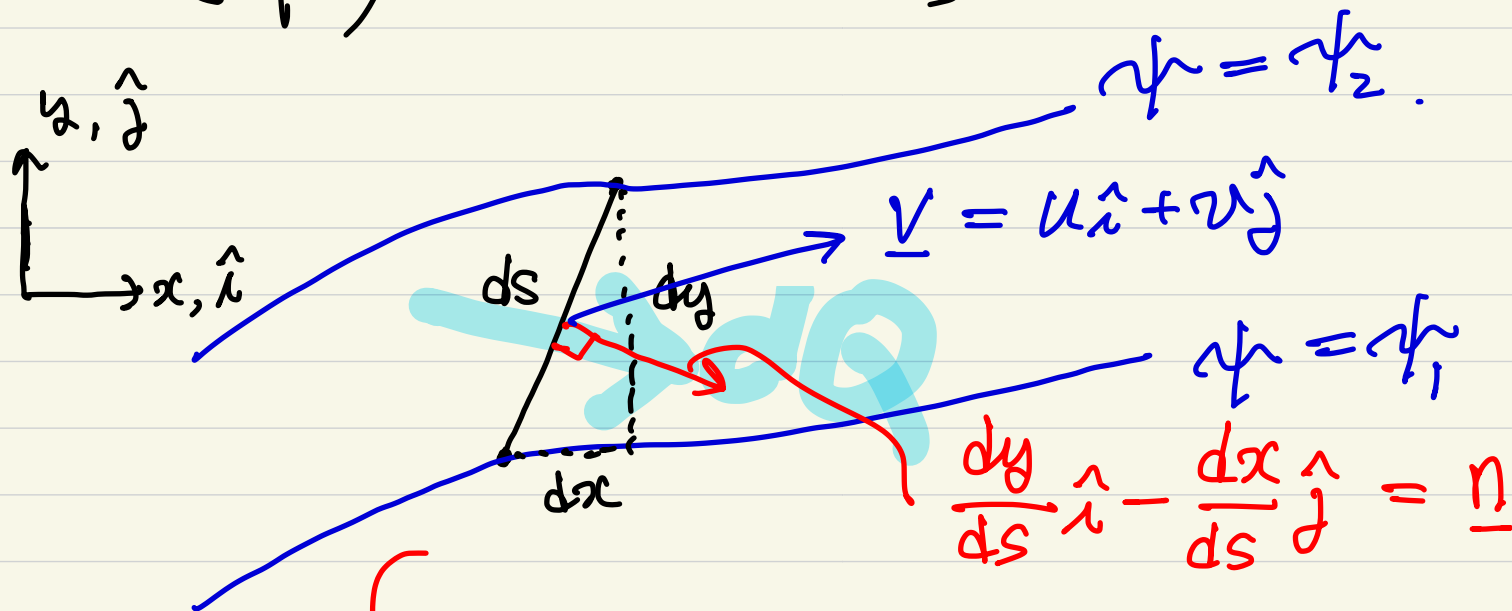
$$-v dx + u dy = 0$$

$$u dy - v dx = 0$$

along a streamline.

$\therefore \psi = C$  along a streamline.

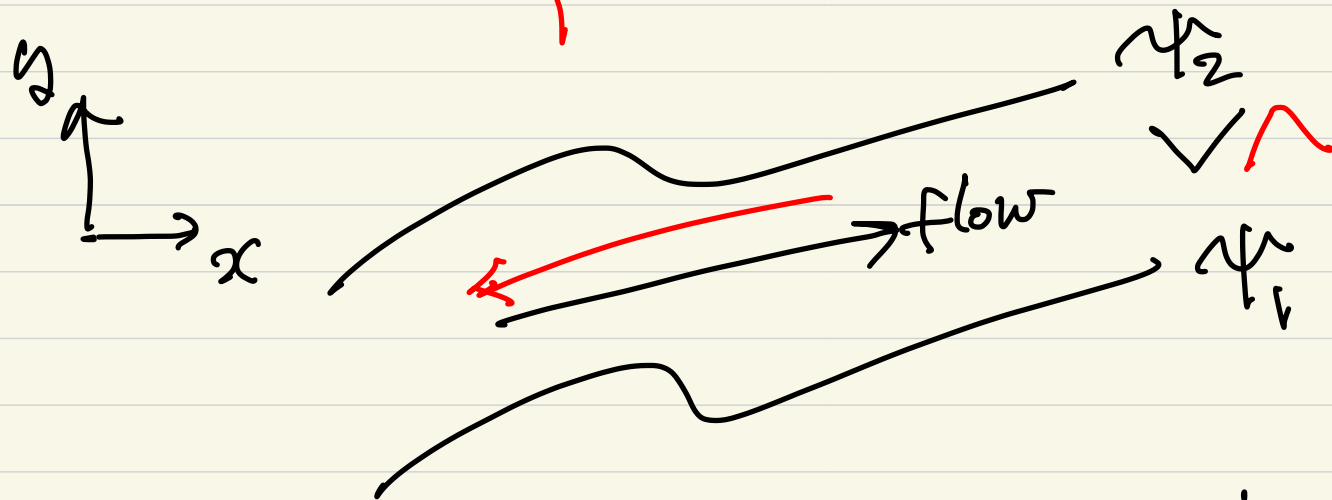
② Change in  $\psi$  across the element is equal to the volume flow through the element.  
 (physical meaning of  $\psi$  value) (2D)



$$\begin{aligned}
 dQ &= (\underline{v} \cdot \underline{n}) dA \\
 &= (u\hat{i} + v\hat{j}) \cdot \left( \frac{dy}{ds}\hat{i} - \frac{dx}{ds}\hat{j} \right) ds \\
 &= \left( \frac{\partial \psi}{\partial y}\hat{i} - \frac{\partial \psi}{\partial x}\hat{j} \right) \cdot \left( \frac{dy}{ds}\hat{i} - \frac{dx}{ds}\hat{j} \right) ds.
 \end{aligned}$$

$$= \frac{\partial \psi}{\partial x} dx + \frac{\partial \psi}{\partial y} dy = \int d\psi$$

$$\therefore Q = \int_1^2 d\psi = \psi_2 - \psi_1$$



$$u = \frac{\partial \psi}{\partial y}$$

$$v = -\frac{\partial \psi}{\partial x}$$

\* steady plane compressible flow.

(2D)

$$\text{continuity: } \frac{\partial}{\partial x} (\rho u) + \frac{\partial}{\partial y} (\rho v) = 0$$

→ compressible stream function

$$\underline{p}u \equiv \frac{\partial \psi}{\partial y}, \quad \underline{p}v \equiv -\frac{\partial \psi}{\partial x}$$

$$\Rightarrow \dot{m} = \int (\underline{v} \cdot \underline{n}) dA = d\psi \quad (\text{not } dQ = d\psi)$$

$$\dot{m}_{1 \rightarrow 2} = \psi_2 - \psi_1$$

\* In incompressible plane flow in polar coordinate.  
(2D)

$$\text{continuity: } \frac{1}{r} \frac{\partial}{\partial r} (r \underline{v}_r) + \frac{1}{r} \frac{\partial}{\partial \theta} (\underline{v}_\theta) = 0$$

$$\rightarrow r \underline{v}_r = \frac{\partial \psi}{\partial \theta}, \quad \underline{v}_\theta = -\frac{\partial \psi}{\partial r}, \quad dQ = d\psi$$

\* In general, stream function is NOT defined in 3D flows, except the axisymmetric flow. ( $\underline{v}_\theta = 0, \frac{\partial}{\partial \theta}(\cdot) = 0$ )



$$\text{continuity eq} = \frac{1}{r} \frac{\partial}{\partial r} (r v_r) + \frac{\partial}{\partial z} (v_z) = 0.$$

↓

$$\frac{\partial}{\partial r} (r v_r) + \frac{\partial}{\partial z} (r v_z) = 0.$$

$$\Rightarrow r v_r = - \frac{\partial \psi}{\partial z}, \quad r v_z = \frac{\partial \psi}{\partial r}.$$

$$dQ = 2\pi \cdot d\psi. \quad \rightarrow \quad Q_{1 \rightarrow 2} = 2\pi (\psi_2 - \psi_1)$$